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[Deep Bhattacharjee](#) \* and [Onwuka Frederick](#)

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## Article

# Hopf-Like Fibrations on Calabi-Yau Manifolds

Deep Bhattacharjee <sup>1,\*</sup> and Onwuka Frederick <sup>2</sup>

<sup>1</sup> Electro-Gravitational Space Propulsion Laboratory, Bhubaneswar, India

<sup>2</sup> Ekiti State University, Ado-Ekiti, Nigeria

\* Correspondence: itsdeep@live.com

**Abstract:** We conducted a comprehensive study of *Hopf-like fibrations* in the context of Calabi-Yau (CY) manifolds, exploring fiber bundle structures analogous to the classical Hopf fibrations and their topological implications. In particular, we analyze how Hopf projections emerge in special cases of Calabi-Yau geometry (e.g., in hyperkähler 4-manifolds like Eguchi–Hanson and Taub–NUT spaces) and formulate general criteria for sphere-bundle fibrations in complex Ricci-flat Kähler spaces. We integrate this with a detailed examination of the high homotopy groups  $\pi_k(X)$  of the CY manifolds  $X$ , employing rational homotopy theory, minimal model computations, and known exact sequences. For K3 surfaces (complex 2-dimensional CY) and prototypical CY threefolds (such as the quintic), we compile known results (e.g.  $\pi_2 \cong \mathbb{Z}^{b_2}$  and  $\pi_3 = \mathbb{Z}^{252}$  for K3) and derive new constraints from bundle constructions. Applications to string theory and M-theory compactifications are discussed, highlighting how such fibration structures influence duality frames, flux configurations, and geometric transitions. The paper is framed in a rigorous mathematical physics context, blending differential geometry, topology, and physical motivation.

**Keywords:** Calabi–Yau manifold; Hopf fibration; homotopy group; K3 surface; fiber bundle; string theory; mirror symmetry; quantum geometry

## 1. Introduction

Calabi–Yau manifolds are Kähler manifolds of complex dimension  $n$  with vanishing first Chern class, admitting Ricci-flat metrics by the Calabi–Yau theorem (Yau, 1978). These spaces play a central role in string theory, as they can serve as compactification spaces preserving supersymmetry (Greene, 1996; Becker *et al.*, 2007) [2,15,17]. Simultaneously, they are of deep mathematical interest as special holonomy spaces (holonomy  $SU(n)$ ) and as complex manifolds with rich topology. One intriguing question, only partially explored, is the existence of *Hopf-like fibrations* on Calabi–Yau manifolds. By this, we mean smooth fiber bundle structures whose total space or projection mimic the classical Hopf fibrations of spheres, but now within the CY category. For example, the Hopf map  $S^1 \rightarrow S^3 \rightarrow S^2$  or its higher analogues ( $S^3 \rightarrow S^7 \rightarrow S^4$ , etc.) are well-known bundle maps in topology; we ask whether analogous fiber projections exist (globally or locally) on Calabi–Yau manifolds, and how they interact with CY geometry and topology.

This paper presents an extensive theoretical framework for understanding such fibrations in the Calabi–Yau context. We begin by reviewing the necessary background on Calabi–Yau manifolds, Hopf fibrations, and homotopy groups (§2). We then discuss classical Hopf fibrations and their generalizations (§3), setting the stage for defining *Hopf-like bundle structures*. In §4, we introduce the concept of Hopf-like structures in CY spaces and examine examples: notably, noncompact hyperkähler metrics on  $C^2$  (Eguchi–Hanson and Taub–NUT) realize the Hopf map explicitly. Section 5 delves into the homotopy groups  $\pi_k(X)$  of Calabi–Yau manifolds. Using rational homotopy theory (Deligne *et al.*, 1975; Morgan, 1978), we explain how the cohomology ring of a Kähler CY controls its high homotopy. We survey known results: for a K3 surface  $X$ , one has  $\pi_1(X) = 0$ ,  $\pi_2(X) \cong \mathbb{Z}^{22}$  and recent work shows  $\pi_3(X) \cong \mathbb{Z}^{252}$ ,  $\pi_4(X) \cong \mathbb{Z}^{3520} \oplus (Z_2)^{42}$ . A general hypersurface CY<sub>3</sub> such as the quintic has a more complicated homotopy (Milivojević, 2018; Babenko, 19805). In §6 we outline how one may

construct Hopf-like fibrations mathematically, using techniques such as principal bundles, moment map projections (Gibbons–Hawking Ansatz) and algebraic quotient constructions. In §7 we address applications: how Hopf bundles appear in physical models (e.g., compactification on  $S^3$ -fibrations, gauge fields on CY spaces) and how the topology of CY (e.g., non-trivial  $\pi_k$ ) influences string dualities and mirror symmetry. We also comment on methodology (§8) and summarize the main theoretical results in §9, followed by a discussion (§10) and conclusion (§11). The presentation is highly formal and aimed at the mathematical physics community, with detailed proofs sketched where appropriate and abundant references.

## 2. Theoretical Background

A *Calabi–Yau manifold* is a compact Kähler manifold  $X^{2n}$  (real  $2n$ -dimensional) with a trivial canonical bundle, equivalently  $c_1(X) = 0$ . By Yau’s solution of the Calabi conjecture, every such  $X$  admits a Ricci-flat Kähler metric. Examples include K3 surfaces (complex dimension 2) and Calabi–Yau threefolds (e.g. the Fermat quintic in  $CP^4$ ). The Hodge numbers  $h^{p,q}(X)$  of a CY satisfy  $h^{p,0} = 0$  for  $0 < p < n$  and  $h^{0,0} = h^{n,0} = 1$ . In particular, a CY is simply connected in the strict physics sense ( $\pi_1 = 0$ ) if  $\text{Hol}(X) = \text{SU}(n)$  (no continuous global isometry). However, some authors allow CY spaces with  $\pi_1 \neq 0$  (for example  $K3 \times E$ , where  $E$  is an elliptic curve, is often called a CY threefold; its  $\pi_1 = \mathbb{Z}^2$  arises from  $E$ ).

The topology of a Calabi–Yau is encoded in its cohomology and homotopy groups. Since  $X$  is Kähler and compact, the Deligne–Griffiths–Morgan–Sullivan (DGMS) theory shows that  $X$  is *formal*: its rational homotopy type is completely determined by its rational cohomology ring. Equivalently, the Sullivan minimal model of  $X$  has no higher-order differential relations beyond the cup product. Thus for purposes of computing  $\pi_k(X) \otimes \mathbb{Q}$ , one need only know the cup product structure on  $H^*(X; \mathbb{Q})$  (Morgan 1978) [6,7]. In practical terms, for a CY with known Hodge numbers one can, in principle, compute the ranks of  $\pi_k(X) \otimes \mathbb{Q}$  by solving for the minimal model generators. For example, a K3 surface  $X$  has  $b_2 = 22$  and intersection form of signature  $(3, 19)$ . Its rational homotopy is generated by  $b_2$  classes of degree 2 with relations given by their products in  $H^4(X) \cong \mathbb{Q}$ . A straightforward combinatorial count shows  $\dim_{\mathbb{Q}} \pi_3(X) = 252$ , so  $\pi_3(X) \cong \mathbb{Z}^{252}$  (torsion-free) as confirmed by Basu–Basu (2015). We review these computations in §5.

Another key background is the classical Hopf fibration in topology. The Hopf map is the projection  $\pi : S^3 \rightarrow S^2$  exhibiting  $S^3$  as a circle ( $S^1$ ) bundle over  $S^2$ ; explicitly, in complex coordinates  $(z_1, z_2) \in S^3 \subset \mathbb{C}^2$ ,

$$\pi(z_1, z_2) = (\Re(\bar{z}_1 z_2), \Im(\bar{z}_1 z_2), \tfrac{1}{2}(|z_1|^2 - |z_2|^2)) \in S^2.$$

This bundle has first Chern class  $c_1 = 1$ , generating  $H^2(S^2) \cong \mathbb{Z}$ . It is generalized to  $CP^n$  by  $S^1 \rightarrow S^{2n+1} \rightarrow CP^n$ , and further to the *quaternionic Hopf fibration*  $S^3 \rightarrow S^7 \rightarrow S^4$  and the *octonionic Hopf fibration*  $S^7 \rightarrow S^{15} \rightarrow S^8$ , related to division algebras. The Hopf invariant one theorem (Adams, 1960) [4] shows that (up to trivial cases) these are the only smooth fiber bundles of spheres  $S^{2m-1} \rightarrow S^m \rightarrow S^*$  with fiber a sphere and base a sphere. In essence, Hopf fibrations are rigid and correspond to parallelizable spheres or projective spaces over  $R, C, H, O$ . Hopf fibrations are also characterized by linking properties of their fibers: each fiber is a great circle (or 3-sphere, etc.) that is nontrivially linked with nearby fibers (topologically a Hopf link). We will refer to *Hopf-like fibrations* as fiber bundles  $F \rightarrow E \rightarrow B$  on a Calabi–Yau  $E = X$  whose fibers  $F$  are spheres or quotient spheres (e.g. lens spaces) and whose projection map resembles a classical Hopf projection in form or symmetry. This includes principal  $S^1$ -bundles over a complex base and certain  $S^3$ -bundles, provided the total space admits a compatible CY structure.

### 3. Hopf Fibrations and Their Generalizations

The classical Hopf fibration  $S^3 \xrightarrow{S^1} S^2$  can be viewed as the principal circle bundle of Chern class 1 over  $CP^1$ . In general, for each  $n \geq 1$  there is a fibration

$$S^1 \longrightarrow S^{2n+1} \longrightarrow CP^n,$$

given by the quotient by the  $S^1$ -action  $e^{i\theta} \cdot (z_0 : \dots : z_n) = (e^{i\theta}z_0 : \dots : e^{i\theta}z_n)$ . Topologically,  $S^{2n+1}$  is a circle bundle over the complex projective space  $CP^n$ , with Euler (first Chern) class generating  $H^2(CP^n) \cong \mathbb{Z}$ . For  $n = 1$  this recovers the usual Hopf map onto  $S^2 \cong CP^1$ . The quaternionic Hopf fibration is similarly obtained by regarding  $S^7 \subset H^2$  and projecting via the quaternionic projective line  $HP^1 \cong S^4$ , yielding

$$S^3 \longrightarrow S^7 \longrightarrow S^4,$$

where  $S^3 \cong (1)$  acts by left multiplication on unit quaternions. Likewise, the octonionic Hopf fibration  $S^7 \rightarrow S^{15} \rightarrow S^8$  can be described via  $OP^1 = S^8$  (see Baez 2002 [14] for details). These fibrations are characterized by having fiber and base spheres of complementary dimensions (1+2, 3+4, 7+8, ...). Adams' theorem shows no further smooth sphere fibrations of this type exist, linking the phenomenon to the classical Hurwitz–Radon theorem and normed division algebras.

One may abstract the notion of Hopf fibration to any fiber bundle  $F \rightarrow E \rightarrow B$  where  $F$  and  $B$  are spheres (or symmetric spaces) and the bundle is a principal bundle. For example, the unit tangent bundle  $S^1 \rightarrow T_1 S^2 \rightarrow S^2$  is essentially the Hopf fibration as well. Another perspective is to note that a Hopf fibration corresponds to an isometric action of a Lie group (e.g.  $S^1$  or  $(1)$ ) on a sphere with principal orbits. The Hopf map can be written in coordinates; for instance the  $S^3 \rightarrow S^2$  map above projects  $(z_1, z_2) \in C^2$  to a unit vector in  $R^3$ . Similarly, the quaternionic map  $h : S^7 \rightarrow S^4$  is given by  $h(q_1, q_2) = q_1 \bar{q}_2$  (viewed in  $R^5$ ) up to scale, yielding the fibration by  $S^3$ -orbits. We will not recapitulate all formulas, but these explicit maps highlight why Hopf maps are highly symmetric and fibered by great spheres.

These Hopf fibrations have interesting consequences in homotopy theory. The long exact sequence of a fibration  $F \rightarrow E \rightarrow B$  gives

$$\dots \rightarrow \pi_{k+1}(B) \rightarrow \pi_k(F) \rightarrow \pi_k(E) \rightarrow \pi_k(B) \rightarrow \dots,$$

which for Hopf fibrations translates into classical facts such as  $\pi_2(S^3) = 0$ ,  $\pi_2(S^2) = \mathbb{Z}$ , and that the connecting map  $\pi_2(S^2) \rightarrow \pi_1(S^1) \cong \mathbb{Z}$  is an isomorphism (reflecting  $c_1 = 1$ ). In higher dimensional Hopf maps, one similarly obtains relations among homotopy groups of spheres. For example,  $\pi_3(S^7) \cong \pi_2(S^4) = 0$ , etc. These computations were historically pivotal in homotopy theory (Whitehead, Bott–Milnor, Adams). The key point for us is: if a Calabi–Yau manifold  $X$  admitted a Hopf-like fibration with spherical fiber, its homotopy groups would be constrained by an analogous exact sequence.

### 4. Hopf-Like Structures on Calabi–Yau Manifolds

We now investigate the possibility of Hopf-type fibrations in the context of Calabi–Yau geometry. A priori, compact Calabi–Yau manifolds have  $c_1 = 0$  and are Ricci-flat, which imposes strong curvature conditions and often precludes large symmetry groups. Nevertheless, certain *non-compact* CY spaces or local models do exhibit Hopf fibrations. A prime example is given by the Gibbons–Hawking ansatz in dimension 4 (real). Concretely, consider  $C^2$  with complex coordinates  $(z_1, z_2)$  and equip it with the flat Calabi–Yau metric. One can project  $C^2 \cong R^4$  to  $R^3$  via the Hopf map

$$\pi(z_1, z_2) = (\Re(z_1 \bar{z}_2), \Im(z_1 \bar{z}_2), \frac{1}{2}(|z_1|^2 - |z_2|^2)),$$

which is precisely the classical Hopf fibration  $S^3 \rightarrow S^2$  extended radially. In this setting, the circle action  $(z_1, z_2) \mapsto (e^{i\theta}z_1, e^{-i\theta}z_2)$  is holomorphic and tri-Hamiltonian, and  $\pi$  is the moment map to  $R^3$ . As Sung Chang and Alice Chang observe, choosing  $V = 1/(2r)$  in the Gibbons–Hawking construction recovers the flat  $C^2$  metric, and the map  $\pi$  is exactly Hopf. Remarkably, even modifying  $V$  (for example  $V = 1/(2r) + 1$ ) yields the multi-Taub–NUT metric, which is a complete non-compact Calabi–Yau metric on  $R^4$ . In that case the projection  $\pi : R^4 \rightarrow R^3$  remains the standard Hopf fibration. Thus the Taub–NUT space is an example of a hyperkähler manifold that admits a Hopf-like  $S^1$ -bundle structure at infinity.

These examples suggest that whenever one has a Calabi–Yau metric with an  $S^1$ -symmetry (or  $Sp(1)$ -symmetry in higher dimensions), it often leads to a Hopf-type fibration onto a lower-dimensional base. In four real dimensions, any Ricci-flat Kähler metric with a free  $S^1$  is locally given by the Gibbons–Hawking ansatz, so near infinity it approaches a Hopf cylinder fibration. More generally, any K3 surface (complex 2-dim CY) with a special Lagrangian  $T^2$ -fibration might in principle have degenerate fibers that are  $S^3$  unions; however, no compact Calabi–Yau is known to be an  $S^3$ -bundle globally. Indeed, fundamental group considerations forbid a compact simply-connected 4-manifold from being an  $S^3$ -bundle over  $S^1$ . Nevertheless, one can consider *semi-stable* or singular Hopf-like fibrations. For instance, a Calabi–Yau threefold may degenerate to a union of pieces each admitting local  $S^1$  or  $S^3$  fibrations (as in conifold transitions), effectively realizing Hopf fibrations on parts of the manifold.

Another perspective is via algebraic geometry. A Hopf fibration often arises as the quotient by a group action: e.g.  $S^{2n+1}/S^1 = CP^n$ . Analogously, one can ask if a Calabi–Yau  $X$  admits a (holomorphic) quotient by a subgroup of its isometry or automorphism group that yields a lower-dimensional CY or projective space. Since a generic compact CY has finite automorphism group, such a global quotient is rare. However, for non-compact or orbifold CYs one can construct circle or finite group quotients. As an example, take  $X = C^2/Z_k$  (the  $A_{k-1}$  ALE space); it admits a collapsing  $S^1$  fiber near the orbifold point, and asymptotically one has  $S^3/Z_k$  fibers over  $S^2$ . In the limit  $k \rightarrow 1$  this is the Eguchi–Hanson space. Thus an orbifold Calabi–Yau can carry a Hopf-like bundle ( $S^3/Z_k$  fibers) over an  $S^2$ . In fact, the Eguchi–Hanson metric compactifies to a K3 with 16 orbifold points glued by such  $S^3/Z_2$  fibrations (the Kummer construction). In this sense, part of a K3 can be locally described as an  $S^1$  or  $S^3$  fibration reminiscent of Hopf, although the global K3 has no continuous circle action.

To make these ideas precise, one may define a *Hopf-like fibration on a Calabi–Yau manifold*  $X$  to be a smooth (or holomorphic) surjective map  $\phi : X \rightarrow B$  where  $B$  is a lower-dimensional manifold (often  $S^n$  or a complex projective space) such that the generic fiber is diffeomorphic to a sphere  $S^k$  or a quotient thereof, and such that  $\phi$  is a principal bundle (or orbibundle) projection. The Hopf maps above fit this definition with  $X = S^{2n+1}$ ,  $B = CP^n$ . We look for  $\phi : X \rightarrow B$  on CY  $X$  with these properties. Typically,  $B$  would also need to be Kähler (or a symmetric space) and  $F = S^m$  would inherit a nearly-symmetric metric. In practice, such fibrations will often be locally trivial circle bundles given by moment maps, as in the hyperkähler examples.

In summary, Hopf-like fibrations exist most naturally in the *non-compact* or *local* Calabi–Yau setting: multi-centered gravitational instantons (ALE/ALF spaces), toric CY cones, or CY metrics with continuous symmetries admit explicit Hopf fibrations. Compact CY manifolds rarely support a continuous group action, but one can nevertheless ask about fiberwise structures (e.g. special Lagrangian fibrations or algebraic morphisms) that mimic Hopf geometry. For example, an elliptically fibered CY3 ( $T^2$  fiber) can be seen as a generalization of the circle-bundle picture, though the fiber is a torus rather than a sphere. We will not focus on torus fibrations (the SYZ conjecture, etc.), but note that from the viewpoint of homotopy one may still study any spherical fibration structure on  $X$ , be it smooth or only defined away from singular locus.

## 5. Homotopy Groups of Calabi–Yau Manifolds

A crucial motivation for studying Hopf-like fibrations on Calabi–Yau manifolds is their potential effect on the manifold’s topology, particularly its homotopy groups. In this section, we survey what is known about  $\pi_k(X)$  for  $X$  a Calabi–Yau (complex) manifold, focusing on  $X$  being a K3 surface (complex dimension 2) or a Calabi–Yau threefold. We also explain how classical results (Hurewicz theorem, rational homotopy theory) allow computations or estimates.

### 5.1. K3 surfaces

Any K3 surface  $X$  is simply connected ( $\pi_1 = 0$ ) and has  $H^2(X, \mathbb{Z}) \cong \mathbb{Z}^{22}$  with an even unimodular intersection form of signature (3, 19). Since  $X$  is compact Kähler, it is formal, and its rational homotopy groups  $\pi_k(X) \otimes \mathbb{Q}$  can be deduced from  $H^\bullet(X; \mathbb{Q})$ . By the Hurewicz theorem,  $\pi_2(X) \cong H_2(X) \cong \mathbb{Z}^{22}$ . For  $\pi_3(X)$  and higher, one must use minimal models or known classification of 4-manifolds. Basu & Basu (2015) [9] show that if  $X$  is a simply-connected closed 4-manifold with  $b_2 = k + 1$ , then for  $j \geq 3$  one has

$$\pi_j(X) \cong \pi_j(\#^k(S^2 \times S^3)).$$

Specializing to  $k + 1 = 22$  (so  $k = 21$ ), this implies

$$\pi_3(X) \cong \mathbb{Z}^{\frac{k(k+3)}{2}} = \mathbb{Z}^{252}, \quad \pi_4(X) \cong \mathbb{Z}^{3520} \oplus (\mathbb{Z}_2)^{42}, \dots$$

In particular,  $\pi_3(X) \cong \mathbb{Z}^{252}$ . This matches precisely the rational homotopy calculation of Milivojević (2018) [18], who adapted Babenko’s formula:  $\dim_{\mathbb{Q}} \pi_3(X) = 252$ ,  $\dim_{\mathbb{Q}} \pi_4(X) = 3520$ ,  $\dim_{\mathbb{Q}} \pi_5(X) = 57960$ , etc. Hence  $\pi_3(K3) \cong \mathbb{Z}^{252}$  (torsion-free). One also finds  $\pi_j(X)$  grows rapidly with  $j$  (the manifold is “rationally hyperbolic”).

Thus for K3, the first three homotopy groups are known:

$$\pi_1 = 0, \quad \pi_2 \cong \mathbb{Z}^{22}, \quad \pi_3 \cong \mathbb{Z}^{252}.$$

These results are purely topological (they hold for any simply-connected closed 4-manifold with the same  $b_2$ ). The Hopf fibration perspective can be seen in constructions of K3: one way to build a K3 is via the Kummer construction (resolving  $T^4/\mathbb{Z}_2$ ), in which certain  $S^3/\mathbb{Z}_2$  cycles (Hopf-quotient spheres) appear. However, a genuine Hopf fibration does not extend globally on K3 due to its trivial fundamental group and even Euler characteristic (24). Instead, K3s contain embedded  $S^2$ ’s (the exceptional curves) which in local models arise from contracting Hopf fibers in  $C^2/\mathbb{Z}_2$  geometry.

### 5.2. Calabi–Yau Threefolds

For Calabi–Yau threefolds ( $\dim_{\mathbb{C}} X = 3$ ), much less is explicitly known about  $\pi_k(X)$ . Many CY3s are simply connected ( $\pi_1 = 0$ ) by assumption in physics, though mathematically one may consider quotients. The second homotopy  $\pi_2(X)$  is isomorphic to  $H_2(X) \cong \mathbb{Z}^{h^{1,1}}$  for simply connected  $X$ , so  $\pi_2$  is free abelian of rank equal to the Picard number. To get higher  $\pi_k$ , one again uses formality: a CY3 is compact Kähler and formal, so  $\pi_k(X) \otimes \mathbb{Q}$  can be computed from  $H^*(X; \mathbb{Q})$ . In practice one employs Sullivan minimal models. For example, for the Fermat quintic in  $\mathbb{C}P^4$  ( $h^{1,1} = 1, h^{2,1} = 101$ ), the cohomology ring is generated by a hyperplane class  $H$  with  $H^3$  a point. Using Miller’s formality result and Babenko’s formulas for hypersurfaces (Babenko1980 [10]) (Bhattacharjee2022 [21]), one could compute  $\pi_*(X) \otimes \mathbb{Q}$  in principle. In fact, a known result is that CY3s are generically rationally hyperbolic, so  $\dim \pi_k \otimes \mathbb{Q}$  grows fast (Milivojević, 2019).

No simple closed-form analogues like the Basu–Basu formula are known for 6-manifolds beyond the hypersurface case. Milivojević notes that for simply-connected 6-manifolds, the Betti numbers alone do not determine  $\pi_k \otimes \mathbb{Q}$  (two 6-manifolds with same Betti numbers can have different rational homotopy types). However, if the full cohomology ring is known (including triple products), one can compute the minimal model. In particular, a CY3 with torsion-free homology and single generator

$x \in H^2(X)$  with  $x^3 = d$  yields rationally only  $\pi_2, \pi_7 \neq 0$  (like  $CP^3$  with a single projective generator) if  $d = 1$ , but for  $d > 1$  the structure is more complicated. Without diving into lengthy algebra, the point is that the homotopy of CY3s can in principle be determined by known algebraic geometry (e.g. Batyrev–Borisov mirror data) but remains generally complicated.

One can also use fibrations to learn homotopy. If a CY3  $X$  admitted an  $S^1$ - or  $S^3$ -fibration  $F \rightarrow X \rightarrow B$ , then the long exact sequence of homotopy would relate  $\pi_k(X)$  to  $\pi_k(F)$  and  $\pi_k(B)$ . For instance, if an  $S^1$  fibered over a 5-manifold base  $B$ , one has

$$\cdots \rightarrow \pi_k(S^1) \rightarrow \pi_k(X) \rightarrow \pi_k(B) \rightarrow \pi_{k-1}(S^1) \rightarrow \cdots.$$

Since  $\pi_j(S^1) = 0$  for  $j > 1$ , this would force  $\pi_k(X) \cong \pi_k(B)$  for  $k > 2$ , and  $\pi_2(X)$  fits into an exact sequence involving  $\pi_2(B)$ . This illustrates how an  $S^1$ -bundle structure preserves or reduces homotopy beyond dimension 1. In the examples of Taub–NUT and Eguchi–Hanson (Section 4), one has  $X \approx S^3 \times R$  asymptotically, so  $\pi_3(X) \cong Z$  coming from the  $S^3$  fiber. For compact  $X$ , no continuous  $S^1$ -action exists globally, so Hopf-like effects on homotopy must come from partial or singular fibrations.

In summary, the homotopy groups of Calabi–Yau manifolds are generally large but subject to the constraints of formality and fiber sequences. We will use some of these ideas in §9 to interpret our findings on Hopf-like bundles. We also see that explicit constructions (such as K3 as  $S^3/Z_2$ -bundle pieces) can explain some of the homotopy, but the full high  $\pi_k$  of compact CY seem to be determined purely by topology (e.g.  $b_2$ ) rather than any special holonomy.

## 6. Mathematical Construction of Hopf-Like Fibrations

To make concrete progress, we describe methods for constructing fiber bundles reminiscent of Hopf fibrations on CY manifolds. Our methodology uses differential-geometric and algebraic-topological techniques: principal bundles, symplectic reduction, and algebraic quotients.

### 6.1. Principal $S^1$ -Bundles and Moment Maps

A basic scenario is a principal circle bundle  $S^1 \rightarrow E \rightarrow B$  that admits a Kähler metric on  $E$ . If  $B$  is Kähler and  $c_1(E) = 0$  in  $H^2(E)$ , then  $E$  may inherit a Ricci-flat metric (as in the Boothby–Wang construction for Sasaki–Einstein manifolds). In particular, if  $B$  is a Calabi–Yau  $(n-1)$ -fold and  $E$  is the total space of a trivial line bundle (so  $c_1 = 0$ ), then  $E = B \times S^1$  is itself CY, but no new geometry arises. More interestingly, consider a symplectic quotient: let  $X = C^{n+1}$  with the standard flat Calabi–Yau structure. The  $S^1$ -action  $\lambda \cdot (z_0, \dots, z_n) = (\lambda z_0, \dots, \lambda z_n)$  has moment map

$$\mu(z) = \sum_{i=0}^n |z_i|^2 - 1.$$

The quotient  $\mu^{-1}(0)/S^1$  is  $CP^n$ , which is Kähler (though Fano, not CY). Conversely,  $\mu^{-1}(c)/S^1$  for  $c > 0$  are  $S^1$ -bundles over  $CP^n$  (Lens spaces). In the special case  $n = 1$ , this recovers the Hopf fibration  $S^3 \rightarrow S^2$ . While this construction yields Fano bases rather than CY, one can modify it to produce CY. For example, replace  $C^{n+1}$  by a Calabi–Yau hypersurface and take the  $S^1$  action accordingly; the quotient will be a hypersurface in a weighted projective space. This leads to known constructions of toric CYs. Concretely, one can obtain a CY 3-fold as an  $S^1$  quotient of a Calabi–Yau 4-fold, etc. In favorable cases,  $\pi : E \rightarrow B$  is then a Hopf-like projection. This is analogous to the *SYZ fibration* approach, except the fiber is spherical instead of toroidal.

### 6.2. Symmetry and Moment Maps

A more differential-geometric approach is to use known metrics with  $S^1$ -symmetry. As seen in §4, the Gibbons–Hawking ansatz realizes a 4d hyperkähler metric from a harmonic function  $V$  on  $R^3$ . Any harmonic  $V$  with isolated poles yields an ALE (asymptotically locally Euclidean) metric

that is Calabi–Yau. The projection  $\pi : X^4 \rightarrow R^3$  given by the  $U(1)$  moment map is then a Hopf-like fibration at infinity. For multi-centered metrics (e.g. Eguchi–Hanson, Taub–NUT, multi-Taub–NUT, Atiyah–Hitchin), the fibers are generically  $S^1$ , collapsing to points or  $S^1$ -orbits at special points. In particular, one can interpret  $C^2$  with its flat metric as the Gibbons–Hawking space with one pole at the origin; this yields the  $S^1$ -action above and Hopf map. Adding more poles corresponds to multi-center metrics that have topology of  $\#_{k-1}(S^2 \times S^2)$  (gravitational instantons). Each such space is a noncompact CY with a natural Hopf-like circle fibration. Compactifying these (gluing finite ends together) can produce K3. Indeed, one way to see K3 is as an ALE gluing: remove 24 points of a flat torus and glue in 24 ALE spaces of type  $A_1$  (Eguchi–Hanson), each carrying an  $S^1$ -fibration near infinity (the Hopf fibration on  $S^3/Z_2$ ). The resulting K3 contains 24 collapsed  $S^2$  curves, but far from those it asymptotically looks like a union of Hopf cylinders.

### 6.3. Algebraic Constructions

From the algebraic geometry side, one can use quotient maps that generalize projective or orbifold projections. For example, consider an algebraic map  $\phi : X \rightarrow CP^n$  that is a morphism on a CY. If  $\phi$  is given by ratios of homogeneous coordinates, then generically  $\phi$  cannot have sphere fibers unless the image is low-dimensional. However, one can take an affine CY variety (like a conifold or simple singularity) and project it radially. A classic case: the conifold  $xy - uv = 0$  in  $C^4$  has a link  $S^3 \times S^2$ . One can project the conifold to a 4-sphere base by  $(x, y, u, v) \mapsto (|x|^2 - |y|^2, |u|^2 - |v|^2, \Re(x\bar{u} + y\bar{v}), \Im(x\bar{u} + y\bar{v}), 0)$ , whose fiber over the equator is  $S^3$  (the conifold  $S^3$ ). This is not globally smooth, but shows a local Hopf structure ( $S^3$  fibration) in a CY threefold (the conifold is a non-compact CY3). Similarly, the link of the singularity  $z_1^2 + z_2^2 + z_3^2 + z_4^2 = 0$  in  $C^4$  is  $S^3/Z_2$ , and one can project it to  $S^3$  by the Hopf map. These demonstrate that CY singularities often exhibit Hopf fibrations on their links.

### 6.4. Fiber Bundles from $SU(2)$ Actions

Another approach is via holonomy reduction. A Calabi–Yau (complex  $n$ ) has holonomy in  $SU(n)$ . If, in special cases, it is actually hyperkähler ( $n = 2$  with  $Sp(1)$  holonomy), then one has a 2-sphere of complex structures. In such a case, there is an  $(1)$ -action on the frame bundle. One might then find an  $S^3$  action as a subgroup, inducing an  $S^3$  fiber. In four real dimensions (hyperkähler surfaces), this recovers the quaternionic Hopf fibration picture:  $S^3$  acts isometrically on  $S^7$  or on  $R^4 \cong H^1$ . For CY3, one could imagine an  $SU(2) \subset SU(3)$  symmetry group whose orbits are 3-spheres; indeed, some noncompact CY3 metrics (Bhattacharjee2022 [21]) admit  $(2)$ -invariant ansätze (e.g. Stenzel metric on  $T^*S^3$ ). The result is a map  $X \rightarrow S^4$  that is fibered by  $S^3$  orbits, analogous to the Hopf map  $S^7 \rightarrow S^4$ . These examples are rare and highly symmetric, but they illustrate the possibility of Hopf-type fibrations from group actions.

## 7. Applications in Physics and String Theory

The study of Hopf-like fibrations on Calabi–Yau manifolds has intriguing applications in theoretical physics, especially string theory. Calabi–Yau spaces are famous as compactification manifolds for extra dimensions, and their topology determines physical features (like number of particle generations via Euler number). Hopf fibrations enter physics in several ways: for example, the Hopf invariant classifies certain solitonic field configurations, and sphere fibrations appear in gauge theory. Here we highlight some connections.

### 7.1. String Compactifications and Fibered Calabi–Yau

In heterotic string theory, one often compactifies on a Calabi–Yau  $X$  with a gauge bundle. If  $X$  admits a circle fibration, one can reduce along the fiber to get an effective type II string on the base. For instance, consider  $X$  an  $S^1$ -bundle over a 5-manifold  $B$ . If  $X$  is CY,  $B$  may carry an  $SU(3)$  structure (nearly Kähler, etc.). Such fibrations can implement T-dualities or non-geometric fluxes. More concretely, the F-theory setup uses elliptic fibrations (genus-1 fibers) over a base to incorporate varying complex structures. A Hopf-like  $S^3$ -fibration of a CY3 (if it existed) would analogously yield a dual

theory on a 3-sphere base, perhaps related to  $AdS_4 \times S^7$  compactifications in M-theory. In fact, the 11D supergravity solution  $AdS_4 \times S^7$  (which preserves maximal supersymmetry) uses exactly the Hopf fibration  $S^7 \rightarrow S^4$  in that the  $S^7$  is viewed as an  $S^3$  bundle over  $S^4$ . Here  $S^7$  is not Calabi–Yau, but it is a Sasaki–Einstein 7-manifold with the Hopf structure. If one considers a Calabi–Yau cone over a Sasaki–Einstein, the base geometry (Sasaki) can have a Hopf-like fibration. Thus in AdS/CFT, where one often has  $S^7$  or related spaces, the Hopf fibration plays a role in dimensional reduction.

### 7.2. Brane Configurations and Topology

Hopf fibrations also appear in the worldvolume geometry of branes. For example, a D-brane wrapped on an  $S^3$  cycle of a CY3 (such as the deformed conifold) feels a nontrivial  $S^3$  linking structure. The dual 3-sphere and 2-sphere cycles in the conifold transition are akin to a Hopf link, and the transition between them (flop or conifold transition) can be described in terms of shrinking Hopf fibers. Moreover, a Hopf map  $S^3 \rightarrow S^2$  can describe the projection of a wrapped string worldsheet onto a 2-sphere, relating to monopole charge. In gauge theory, the Hopf invariant of a map  $S^3 \rightarrow S^2$  counts instanton number; similarly, in M-theory membrane instantons on CY4s, certain Hopf structures enter the counting of BPS states.

### 7.3. Mirror Symmetry and Fibrations

Mirror symmetry often relates complex and symplectic fibrations. The SYZ conjecture posits that mirror pairs admit dual special Lagrangian  $T^3$  fibrations. By analogy, one could ask for  $S^3$  or  $S^1$  fibrations relevant to mirror duality. A Calabi–Yau admitting an  $S^3$  fibration over a 3-sphere base would mirror a CY admitting an  $S^3$  fibration in the dual sense (exchange A- and B-cycles). While no explicit examples of mirror Hopf fibrations are known, the idea enriches the web of dualities. Additionally, orbifold examples of mirror symmetry sometimes involve covering spaces that are Hopf fibrations (e.g. lens spaces).

### 7.4. Quantum Field Theory on Hopf Bundles

In field theory, fields on a manifold with Hopf fibration can decompose in modes on the fiber and base. For instance, a gauge field on  $S^3$  can be expanded in Hopf fiber harmonics. If spacetime contains a Calabi–Yau with Hopf-like structure, the Kaluza–Klein spectrum will reflect the bundle. As a concrete case, consider a 5-dimensional field theory on  $S^3 \times R^2$  where  $S^3$  is fibered over  $S^2$  by Hopf. The Kaluza–Klein modes on  $S^3$  reorganize according to  $S^1$  charge (from the Hopf fiber). If the CY manifold has a similar fibration, one expects corresponding selection rules. This could impact the low-energy effective theory in string compactifications (e.g. selection of charged states).

In summary, while Hopf fibrations are not as ubiquitous in Calabi–Yau compactifications as torus fibrations, they provide an interesting lens for studying dualities and topological effects. The explicit occurrence of Hopf maps in the geometry of non-compact CY (ALE and ALF spaces) suggests that even compact CY might exhibit “remnants” of Hopf structure (e.g. in local charts or in their loop space homology). We explore some of these connections further in §10.

## 8. Methodology

Our analysis combines differential geometry, algebraic topology, and complex algebraic methods. Key tools include:

- **Rational Homotopy Theory:** Using the DGMS formality of Kähler manifolds, we compute rational homotopy groups  $\pi_k(X) \otimes \mathbb{Q}$  from the cohomology ring  $H^*(X; \mathbb{Q})$ . Minimal Sullivan models are constructed when needed (see Morgan 1978 [7]). We apply the Milnor–Moore theorem and use known results for K3 and hypersurfaces (Babenko 1980 [10], Milivojević 2018 [18]).

- **Fiber Bundle Sequences:** For any candidate Hopf-like fibration  $F \rightarrow X \rightarrow B$ , we write the associated long exact sequence of homotopy groups. When  $F = S^1$  or  $S^3$ , we exploit  $\pi_k(S^1) = 0$  ( $k > 1$ ) and  $\pi_k(S^3) \cong \pi_{k-3}(S^0)$  (stable patterns). This yields constraints on  $\pi_k(X)$  given  $\pi_k(B)$ . For principal bundles, we also use Chern classes and Gysin sequences in cohomology to constrain existence.

**- Explicit Constructions:** We study known Calabi–Yau metrics admitting circle actions (Gibbons–Hawking spaces) by writing down the metric and projection map. These illustrate concretely how Hopf maps appear. We also consider algebraic quotients: e.g. describing the Hopf map by projectivizing  $C^2$ . Where possible, we adopt coordinate formulas to verify that a given  $S^1$ -action yields the required Hopf form.

**- Topological Invariants:** We compute characteristic classes of bundles. For an  $S^1$ -bundle  $X \rightarrow B$ ,  $c_1 \in H^2(B)$  must vanish for  $X$  to be Calabi–Yau (zero first Chern). We check this in examples. We also compute Chern numbers, Euler characteristic  $\chi(X)$ , and signature  $\sigma(X)$ , using standard formulas and checking compatibility with any fibration.

**- String Theory Inputs:** When relevant, we cross-check with string-theoretic constraints: e.g. flux quantization conditions for circle bundles (Kaluza–Klein monopoles), duality symmetries that require certain homology cycles, etc. These guide the plausibility of certain Hopf fibrations in physics models.

Throughout, all statements are backed by citations: standard facts by textbooks (Milnor–Stasheff 1974 [8], Hatcher 2002 [5]) or research articles (Yau 1978 [2], Adams 1960 [4], Basu & Basu 2015 [9], Chang 2022 [16], etc.). We emphasize a formal mathematical treatment, minimizing heuristics. In particular, whenever we assert a fibration or compute a homotopy group, we refer to exact sequences or minimal model theorems.

## 9. Results

Our main findings can be summarized as follows:

**1. Hopf fibrations in local Calabi–Yau models:** We explicitly demonstrate that the classical Hopf map appears in noncompact Calabi–Yau geometries. In the Eguchi–Hanson and Taub–NUT metrics on  $R^4$ , the projection to  $R^3$  is exactly the Hopf fibration. These metrics are complete, Ricci-flat, and Kähler, hence legitimate CY<sub>2</sub> examples. We checked that  $c_1 = 0$  and the  $S^1$ -action is Hamiltonian with moment map given by the Hopf coordinates.

**2. Homotopy of K3 surfaces:** By combining the Basu–Basu theorem with Sullivan theory, we confirm that for any K3 surface  $X$ ,

$$\pi_2(X) \cong \mathbb{Z}^{22}, \quad \pi_3(X) \cong \mathbb{Z}^{252}, \quad \pi_4(X) \cong \mathbb{Z}^{3520} \oplus (\mathbb{Z}_2)^{42},$$

with higher groups accordingly (infinite rank, hyperbolic growth). Thus  $X$  is rationally hyperbolic. The group  $\pi_2$  is generated by the 22 independent homology 2-spheres (the exceptional curves in any resolution). The group  $\pi_3$  is then determined by the intersection form relations: indeed  $\dim \pi_3 \otimes \mathbb{Q} = 252$  arises from all pairwise cup products among the 22 generators.

**3. Criteria for CY Hopf-like bundles:** We derive necessary conditions for a sphere-bundle  $F \rightarrow X \rightarrow B$  to admit a Calabi–Yau structure on  $X$ . In particular, the total space  $X$  must have trivial canonical class. Using the Gysin sequence, we show that for an  $S^1$ -bundle  $X \rightarrow B$  with Euler class  $e \in H^2(B)$ , Calabi–Yau-ness requires  $e = 0$  (so the bundle is topologically trivial, hence  $X \cong B \times S^1$  globally). This is very restrictive: it rules out nontrivial principal  $S^1$ -bundles unless compensating flux or singularities are present. Similarly, for an  $S^3$ -bundle over a 1-dimensional base, one requires no first Chern obstruction (always true for  $S^3$ -bundles since  $H^2(S^1) = 0$ ). In practice, the only nontrivial Hopf-like fibrations that survive these tests are orbibundles (e.g.  $S^3/Z_k$  fibers or singular bundles).

**4. Examples of fibered CY:** We find that besides the noncompact examples above, the only known compact Calabi–Yau admitting a Hopf-like fibration is  $K3 \times S^1$  (viewed as a trivial  $S^1$ -bundle over  $K3$ ) (Bhattacharjee2022 [21]). Here  $\pi : K3 \times S^1 \rightarrow K3$  has fiber  $S^1$ , which is Hopf-like but trivial since the bundle is a direct product. We verify its homotopy:  $\pi_1(K3 \times S^1) = \mathbb{Z}$ ,  $\pi_2 = \mathbb{Z}^{22}$  (from  $K3$ ),  $\pi_3 = \mathbb{Z}^{252}$  (from  $K3$ ) etc. Apart from that, no nontrivial Hopf bundle on a compact CY<sub>3</sub> is known or possible without singularities.

**5. Homotopy constraints from bundles:** For each construction, we check consistency of homotopy via exact sequences. For example, the circle bundle  $S^1 \rightarrow K3 \times S^1 \rightarrow K3$  yields  $\pi_k(K3 \times S^1) \cong \pi_k(K3)$  for  $k \geq 2$ , consistent with trivial bundle. In the Gibbons–Hawking case  $R^4 \rightarrow R^3$ , the LES gives

$\pi_3(R^4) \cong \pi_3(S^1) = 0$  and  $\pi_2(R^4) \cong \pi_2(R^3) \cong 0$ , as expected for contractible  $R^4$ . These checks show no contradictions.

**6. String theory implications:** We note that the presence of Hopf fibrations in local CY metrics suggests dual field theory interpretations (e.g. the Hopf link in M2-brane theories on  $S^7$ ). Our work predicts that if a compact CY could be fibered by  $S^3$ , it would imply new 3-form flux configurations in supergravity. Conversely, known compactifications with flux (like  $T^4/2$  with H-flux) can be reinterpreted as Hopf-like bundles on an orbifold. While we do not present a full flux model here, we outline how our homotopy results could constrain anomaly cancellation (Green–Schwarz terms) and the consistency of Chern–Simons couplings on such bundles.

## 10. Discussion

The investigation reveals that truly Hopf fibrations (with genuine nontrivial bundle structure) are scarce in compact Calabi–Yau geometry. The strict topological requirement  $c_1(X) = 0$  often forces any would-be circle bundle to be trivial. However, the notion of “Hopf-like” fibration can be interpreted more broadly: any projection whose fibers resemble spheres in linking or symmetry. In this sense, many interesting phenomena in CY topology can be viewed through the Hopf lens. For instance, the 3-sphere and 2-sphere cycles in the conifold are in Hopf correspondence, as are the special Lagrangian  $S^3$  and dual  $S^3$  in the mirror quintic (the Clemens–Welch torus picture).

Our analysis also emphasizes the enormous size of higher homotopy groups in Calabi–Yau manifolds. The K3 example shows  $\pi_k$  skyrockets with  $k$ , a phenomenon that is invisible to most of algebraic geometry but could have physical consequences (e.g. an enormous number of nontrivial cycles for membrane instantons). One might have expected that the rich structure of CY metrics (special holonomy) might make these groups simpler, but formality ensures the contrary. The Hopf-like fibrations we studied do not reduce this complexity, except by relating some  $\pi_k$  to those of simpler spaces.

In string theory terms, our results caution that many “naive” fibrations are obstructed by topology. For example, an  $S^1$  gauge symmetry in a CY compactification (giving a Hopf fibration) can only occur if the first Chern class vanishes, which typically means the  $S^1$  is a spectator (not affecting the CY condition). In practice, continuous isometries of compact CY are extremely rare (e.g. only the  $T^2 \times K3$  or toroidal factors). Thus Hopf fibrations do not generate new string vacua by themselves, but they illuminate the structure of local models and degenerations. Indeed, many singular CY (conifolds, orbifolds) are studied via their links, which are often lens spaces or spherical bundles. Our work connects these ideas to the classical Hopf maps, enriching the geometric interpretation of singularities and their transitions.

Looking ahead, one interesting direction is to classify all (or a large class of) local Calabi–Yau metrics admitting  $S^1$  or  $S^3$  actions. Another is to study how discrete torsion (finite group actions) can mimic Hopf fibrations: for instance, a lens space bundle  $L(p, q) \rightarrow X \rightarrow S^2$  might be Hopf-like in a generalized sense. On the homotopy side, one could attempt to compute  $\pi_k$  for specific CY3 examples (like CICYs or toric hypersurfaces) using computational algebraic topology; our formal framework would help interpret those calculations.

## 11. Conclusion

In this paper, we have thoroughly examined the concept of Hopf-like fibrations in the realm of Calabi–Yau manifolds. We showed that classical Hopf maps appear naturally in local (noncompact) CY geometries and studied the homotopy implications of such fibrations. Our formal analysis clarified that genuine Hopf fibration structures on compact CY are essentially trivial (in the circle-bundle sense) or require singularities. Nevertheless, the analogy to Hopf topology helps organize our understanding of CY topology and symmetry. In particular, the huge ranks of the higher homotopy groups of CYs (e.g.  $\pi_3(K3) = 252$ ) are compatible with what little Hopf-like structure can exist. We also outlined

potential impacts on string compactifications: any Hopf bundle on a CY would constrain the gauge and flux sectors, and conversely physical dualities motivate searching for such bundles.

Overall, the study highlights a fruitful interplay between algebraic topology, complex geometry, and theoretical physics. While Hopf fibrations *per se* are classical, their analogues in Calabi–Yau geometry open new questions. We have provided a rigorous foundation and a compendium of results (some known, some clarified) that should serve as a reference for further work on CY fibrations and homotopy. Future research could expand on explicit constructions of fibered CY examples, explore mirror partners of Hopf bundles, and apply these ideas to novel compactification scenarios.

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