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Article

Adsorption of CuSO₄ on Anatase TiO₂ (101) Surface: A DFT Study

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Abstract: The rapid growth of industrial activities has increased environmental pollution, and solar-driven heterogeneous photocatalysis using TiO₂ has emerged as a promising solution. However, its wide bandgap limits its efficiency, prompting research into various optimization strategies. One of these approaches is surface functionalization. Thus, this study investigates the adsorption of CuSO₄ on the anatase TiO₂ (101) surface using density functional theory calculations. The adsorption process induced a magnetic moment of 0.97 μ _B and a slight reduction in overall bandwidth. A preferential adsorption geometry pattern with an energy of -4.31 eV was identified. Charge transfer analysis revealed a net transfer from the TiO₂ surface to the CuSO₄ molecule, with increased net atomic charges for atoms involved in new chemical bond formation, indicating a chemisorption process. These electronic structure modifications are expected to influence the electronic and catalytic properties of the material. The findings provide insights into the CuSO₄ adsorption mechanism on anatase TiO₂ (101) surface and its impact on the properties of the material, contributing to a deeper understanding of this system.

Keywords: TiO₂ (101); CuSO₄; DFT + *U*; vdW; chemisorption; surface

1. Introduction

The rapid growth of industrial activities and population has accelerated the dispersion of harmful chemicals into the environment, leading to the presence of persistent organic pollutants even in remote areas [1,2]. This global issue calls for cleaner, more sustainable production and consumption policies, along with the development of innovative technologies for air, soil, and water treatment [3].

In this sense, solar-driven heterogeneous photocatalysis is a promising and sustainable solution for removing organic and biological contaminants from water [4,5]. Among materials, Titanium dioxide (TiO₂) is a key material in several high-technology applications, including photocatalysis, biomaterials, and novel photovoltaic solar cells [6–8]. Its use in photoelectrochemical processes is particularly crucial for enabling efficient large-scale solar electricity production, addressing the growing global need for sustainable energy.

TiO₂, particularly in its anatase phase, has proven to be a promising material due to its unique electronic and catalytic properties, making it suitable for applications in photocatalysis, sensors, solar

cells, and environmental remediation [9–11]. The anatase TiO₂ (101) surface is known for its high stability and favorable energetic properties [6,12,13]. Studies have demonstrated that the anatase TiO₂ (101) surface is the most energetically favorable and commonly observed surface in natural and synthetic anatase crystals [14,15]. By examining how different molecules interact with this surface, researchers can develop insights into improving TiO₂-based materials for various applications [16,17].

However, the wide bandgap of anatase-phase TiO₂ (3.2 eV) significantly constrains its photocatalytic efficiency under solar irradiation [18]. To address this limitation, researchers have explored various strategies, including doping [19], morphological engineering [20,21], surface functionalization [22], and synthesis of substoichiometric titanium oxides, specifically Magnéli phases [23,24]. These approaches aim to modulate the electronic structure and optical properties of TiO₂, thereby enhancing its photocatalytic performance across a broader spectrum of light.

Photocatalysis is an advanced oxidation process (AOP) that facilitates the degradation of organic compounds by generating hydroxyl radicals (HO•). Nevertheless, optimizing the efficiency of photocatalysts remains challenging, requiring extensive experimental testing, often based on trial-and-error methods.

In this context, computational simulation of molecular adsorption on metal oxide surfaces has emerged as a key tool for understanding catalytic and adsorption mechanisms. This study investigates the adsorption of the copper (II) (CuSO₄) molecule on the anatase (101) surface of TiO₂ using first-principles methods based on Density Functional Theory (DFT).

To the best of our knowledge, no density functional theory (DFT) focused on the surface adsorption of this molecule. Therefore, provides unique insights that complement experimental findings [25,26]. This approach offers a deep atomic-level understanding of the adsorption process.

2. Computational Details

The present work has been conducted using the Vienna *ab initio* Simulation Package (VASP) [27,28], which is based on first-principles Density Functional Theory (DFT) within the Generalized Gradient Approximation (GGA) method [29]. In this approach, the valence electronic states are represented using a set of periodic plane waves, and the interaction between core and valence electrons is implemented through the projector augmented wave (PAW) method [30]. The Perdew-Burke-Ernzerhof (PBE) [31] parameterized GGA functionals are used to describe the exchange-correlation interaction. The valence configuration employed in our calculations for each atom is as follows: $3p^63d^24s^2$ for Ti, $3d^{10}4s^14p^6$ for Cu, $2s^22p^4$ for O, and $3s^23p^4$ for S. Additionally, to account for the strong correlation between d electrons, an intra-site Coulomb repulsion U-term was included in the calculations, specifically, the rotationally invariant (Dudarev) approach to the GGA + U, [32] resulting in the so-called DFT + U method [33]. The utilized U parameters values were 3.5 eV for the Ti atom [34] and 7 eV for the Cu atom [35]. Furthermore, spin-polarized calculations within the DFT+U framework were applied throughout the study.

GGA and hybrid functionals have limitations in accurately capturing long-range electron correlations necessary for modelling van der Waals (dispersive) forces. This shortcoming arises because these functionals approximate the exchange-correlation energy based on local or semi-local density gradients and, in the case of hybrid functionals, include a portion of non-local Hartree-Fock exchange [36]. However, they do not inherently account for long-range dispersive interactions, which are essential for accurately modelling adsorption processes. To address this, we have employed the DFT-D3 dispersion correction method developed by Grimme, which includes a Becke-Johnson damping function [37,38]. This method enhances the accuracy of our adsorption studies by incorporating van der Waals interactions.

This study employed a cut-off kinetic energy of 550 eV which was determined by converging the total energy to less than 1 meV/atom. The Brillouin zone sampling was performed using a Γ -centered Monkhorst-Pack (MP) scheme [39]. A 3 × 2 × 1 k-point mesh was used for structural optimization of the 144-atom periodic slab, while a higher-resolution $7 \times 5 \times 1$ mesh was implemented

for subsequent static calculation to compute properties. To prevent artificial interactions between periodic slabs, a vacuum space of 20 Å and dipole correction were applied [40,41].

To accurately model the isolated CuSO₄ molecule within the context of periodic calculations, it was placed in a cubic supercell with dimensions of $15 \times 15 \times 15$ Å³. The molecule was then geometrically optimized without constraints, using the same kinetic energy cut-off as employed for the slab calculations. This approach ensures consistency in the representation of both the adsorbed and free molecular states.

3. Results and Discussion

3.1. Pristine Surface

To model the anatase (101) TiO_2 surface a slab was constructed using a (3 × 3) 12-atom primitive unit cell, resulting in a periodic slab containing 144 atoms. The slab consists of three trilayers, each containing an upper and bottom layers of O atoms bonded with a middle layer of titanium atoms. Prior to adsorption, the uppermost trilayer was fully relaxed while the undelaying layers were constrained to their bulk positions, as illustrated in Figure 1.

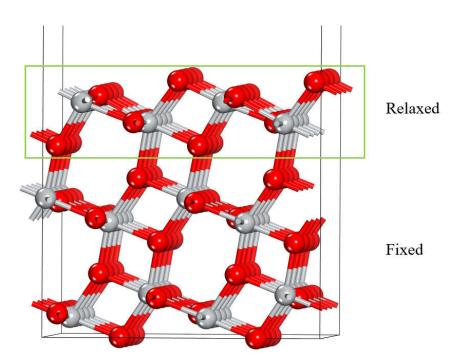


Figure 1. TiO2 anatase (101) surface structure. Red and grey spheres represent O and Ti atoms.

The computed total density of states (DOS) of the pristine surface is presented in Figure 2. It can observe hybridization between O 2p and 3d states in the valence band (VB) with mayor contributions of O 2p states. In the case of the conduction band (CB) is primarily constituted of Ti 3d states, with a small contribution from O 2p states. This composition is consistent with previous works [42,43].

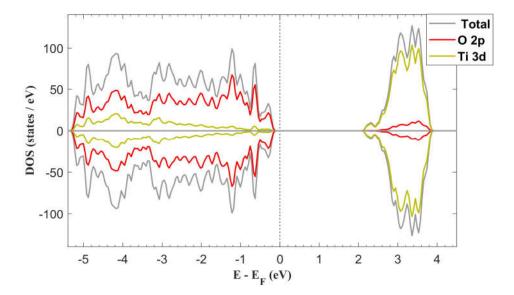


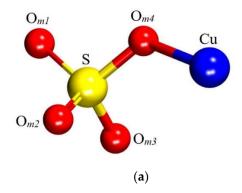
Figure 2. DFT+U total DOS and partial DOS (PDOS) of anatase TiO₂ (101) surface. The bandgap with is 2.27 eV.

The computed band gap width was 2.05 eV, which is smaller than the well-established value of 3.2 eV for bulk anatase TiO₂. This underestimation is a known limitation of DFT methodologies. To address this issue, we implemented the Generalized Gradient Approximation plus Hubbard U (GGA+U) approach with U = 3.5 eV. This modification resulted in a modest improvement of the band gap width, increasing from 2.05 eV to 2.27 eV, representing a 10.73% enhancement. While this adjustment partially mitigates the bandgap underestimation inherent to standard DFT calculations, it is important to note that our primary focus remains on elucidating the effects induced by the adsorption process rather than achieving perfect agreement with experimental band gap values.

3.2. Free Molecule

The dehydrated state of CuSO₄ was selected to model the molecule. The free molecule was constructed by positioning it at the center of a $15 \times 15 \times 15$ Å³ cubic supercell and subjecting it to full atomic relaxation. The optimized molecular structure is presented in Figure 3 and Table 1.

The atomic relaxation within this model led to the approximation of Cu and its nearest O atom, designated as O_{m2} . Prior to optimization, these atoms were separated by 3.25 Å and afterwards, the distance is found to be 1.85 Å, which is identical to Cu $-O_{m4}$ (Table 1). This structural rearrangement can be attributed to the breaking of bonds that the Cu atom originally formed within its crystalline structure (chalcocyanite) and subsequent relaxation that led the isolated molecule to adopt a more energetically favorable state.



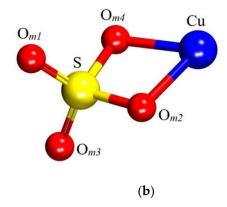


Figure 3. CuSO₄ molecule structure. (a) Experimental and (b) modelled.

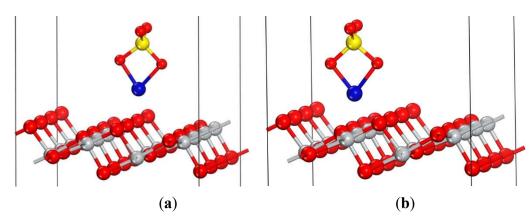
Table 1. Interatomic distances and relevant angles of the CuSO₄ molecule comparing the experimental [44] and the modelled structures.

Distance (Å)				Angle (°)	
Atoms	Exp.	Modelled	Atoms	Exp.	Modelled
S - O _{m1}	1.46	1.44	S - O _{m4} - Cu	123.44	94.04
S - O _{m2}	1.46	1.60	$O_{m1} - S - O_{m3}$	110.52	119.39
S - O _{m3}	1.46	1.44			
S - O _{m4}	1.52	1.60			
$Cu - O_{m4}$	2.04	1.85			
Cu - O _{m3}	3.71	3.46			
$Cu - O_{m2}$	3.25	1.85			
Cu - O _{m1}	4.33	3.51			

3.3. Adsorption on Anatase TiO2 (101) Surface

The anatase phase of TiO_2 is widely recognized as the most prevalent form in photovoltaic applications [45], with the (101) surface being particularly significant due to its thermodynamic stability [46]. Consequently, this study employs the anatase TiO_2 (101) surface as the substrate for investigating the adsorption behavior of CuSO₄.

To determine the optimal adsorption geometry of $CuSO_4$ on the anatase TiO_2 (101) surface, seven distinct initial orientations of the $CuSO_4$ molecule were systematically explored, each positioned over different adsorption sites on the surface. These initial configurations are illustrated in Figure 4.



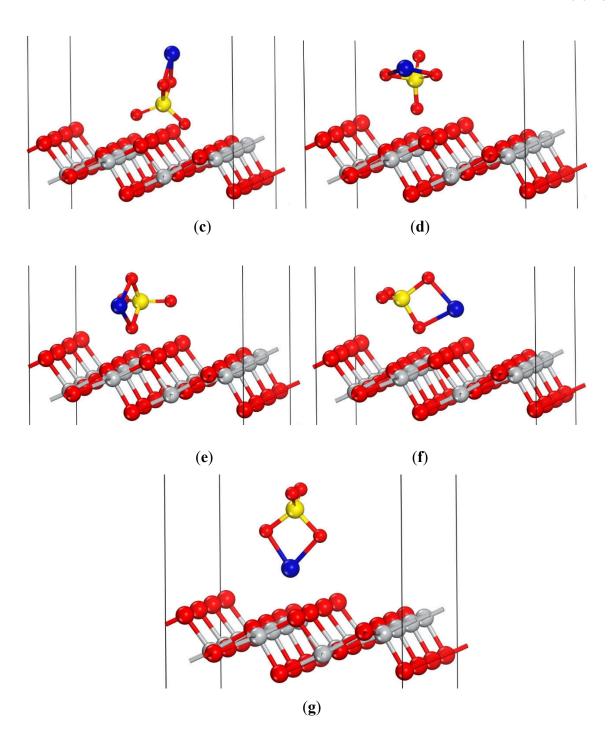


Figure 4. Initial adsorption structures of the CuSO₄ on the anatase TiO₂ (101) surface, prior geometry optimization. Atoms are represented by colored spheres: Cu (blue), S (yellow), O (red), and Ti (grey). For clarity, only the uppermost atoms of the TiO₂ surface are displayed.

Following the initial placement of the CuSO₄ molecule, a comprehensive atomic relaxation was performed. This optimization procedure was conducted without imposing any constraints on the atomic displacements of both the adsorbed molecule and the atoms within the uppermost trilayer of the TiO_2 surface. Upon completion of the geometry optimization, the adsorption energy (E_{ads}) for each configuration was calculated using the following equation:

$$E_{ads} = E_{(surf + mol)} - E_{(surf)} - E_{free(mol)}$$

Where $E_{(surf+mol)}$ refers to energy of the surface with the adsorbate, $E_{(surf)}$ is the energy of the pristine surface while $E_{free(mol)}$ represents the energy of the free molecule. Computed adsorption energies are shown in Table 2.

Table 2. Computed adsorption energies.

Configuration	Eads (eV)
a	-2.06
b	-4.31
c	-1.25
d	-4.31
e	-4.00
f	-4.31
g	-2.67

As can be observed in Table 2, configurations that achieved the lowest adsorption energies are b, d and f, all of which reach an identical adsorption structure. The second most favorable configuration was e, which is depicted in Figure 5.

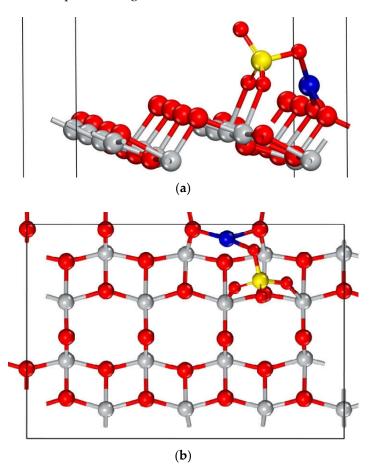


Figure 5. Adsorption geometry of configuration e: (a) Side and (b) top views. The blue, yellow, red and grey spheres represent atoms of Cu, S, O and Ti, respectively.

However, the adsorption geometry with the lowest adsorption energy corresponds to configurations b, d, and f, which have been selected for further analysis due to its higher likelihood of occurrence. Its structure closely resembles configuration e, as is illustrated in Figure 6.

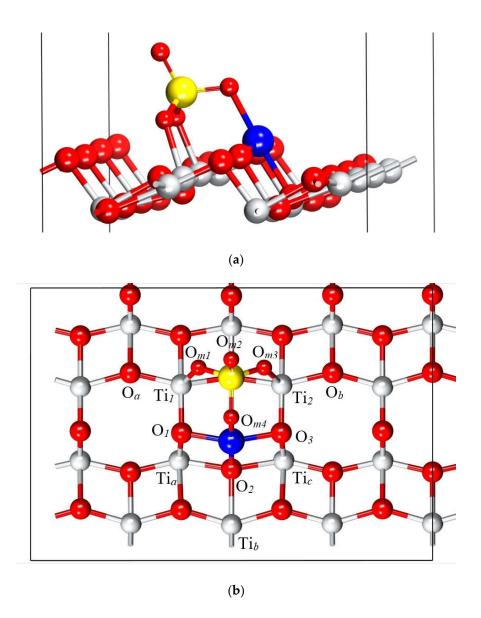


Figure 6. Most favorable adsorption geometry of CuSO₄ on the anatase TiO₂ (101) surface: (a) Side and (b) top views. The blue, yellow, red, and white spheres represent atoms of Cu, S, O, and Ti, respectively.

Figure 6 illustrates the resulting structure of the adsorption process, highlighting the formation of bonds between the molecule and the surface, which indicates a chemisorption mechanism. Notably, the molecule binds to the surface through its oxygen and copper atoms. This interaction significantly alters the geometry of both the molecule and the uppermost surface layers, inducing atomic displacements in these regions. To quantify these changes, the interatomic distances and angles of the atoms labelled in Figure 6b were measured and are presented in Table 3. The newly formed bonds are between O_{m1} and O_{m3} of the molecule with Ti_1 and Ti_2 on the surface, while the Cu atom interacts with surface oxygen atoms O_1 , O_2 , and O_3 . As a result of this process, the surface atoms tend to displace towards the molecule, predominantly in the upwards direction. Ti_1 and Ti_2 move about 0.13 Å, while O_1 and O_3 show a minor displacement of 0.03 Å. O_2 , located in a deeper layer, experiences a more significant displacement of 0.21 Å.

The bond length between the O_1 and O_2 atoms with the Cu atom is 1.94 Å which is consistent with the 1.95 Å observed in the monoclinic crystal structure of Copper (II) oxide (CuO) [47]. In the case of O_m —Ti bonds measure 2.02 Å, which is close to the experimental within the bulk of 1.98 Å (Table 3) [48].

Atomic charges were calculated using Bader population analysis [49] and are presented in Table 4. The data reveals that atoms that exhibited the most significant changes in atomic charge are those

belonging to the molecule. This observation is consistent with the fact that the molecule receives a net charge of 0.11 *e* from the surface. Furthermore, Table 4 also demonstrates that atoms directly involved in bond formation experience an increase in their net atomic charge, suggesting an ionic character to these chemical bonds.

Table 3. Interatomic distances and relevant angles for the pristine surface and the surface with the adsorbate. Atoms labels adhere to the notation established in Figure 6b.

Bond length (Å)				Angle (°)	
Atoms	Free molecule	Adsorbed molecule	Atoms	Free molecule	Adsorbed molecule
S - O _{m1}	1.45	1.53	$Ti_1 - O_{m1} - S$	_	132.57
$S - O_{m2}$	1.57	1.42	$S - O_{m3} - Ti_2$	_	132.62
S - O _{m3}	1.45	1.53	$S - O_{m4} - Cu$	93.82	112.29
$S - O_{m4}$	1.57	1.51	$O_{m1} - S - O_{m3}$	116.04	107.58
Cu - O _{m4}	1.91	1.92			
Atoms	Pristine	Adsorbed	Atoms	Pristine	Adsorbed
	surface	molecule		surface	molecule
O _{m1} - Ti ₁	-	2.02	O_a - Ti_1 - O_{m1}	_	102.16
Om3 - Ti2	_	2.02	O_{m3} – Ti_2 – O_b	_	102.24
Ti1 - Oa	1.98	1.97	$Ti_1 - O_1 - Ti_a$	100.96	101.53
Ti ₂ - O _b	1.98	1.97	$Ti_2 - O_3 - Ti_c$	100.96	101.52
Cu - O ₁	_	1.94	$Ti_a - O_2 - Ti_c$	155.23	152.98
$Cu - O_2$	_	2.02	$Cu - O_2 - Ti_b$	_	108.19
Cu - O ₃	_	1.94			
$O_1 - Ti_a$	1.94	1.96			
O_2 – Ti_b	2.06	2.16			
O ₃ – Ti _c	1.94	1.96			

Table 4. Atomic charges of the atoms labelled in Figure 6b.

Atomic charge (e)					
Atom	Free molecule	Adsorbed molecule			
Cu	1.00	1.14			
S	4.20	3.77			
O_{m1}	-1.24	-1.26			
O_{m2}	-0.85	-1.27			
O_{m3}	-1.68	-1.26			
O _{m4}	-1.43	-1.23			
Atom	Pristine surface	Adsorbed molecule			
Ti ₁	2.26	2.32			
Ti_2	2.26	2.32			
Ti_a	2.30	2.30			
Ti_b	2.26	2.24			
Ti_{c}	2.30	2.30			

O_1	-1.03	-1.09	
O_2	-1.18	-1.23	
O ₃	-1.03	-1.09	
O_a	-1.20	-1.19	
O_b	-1.20	-1.19	

The computed DOS of the system with the adsorbed molecule is presented in Figure 7. Analysis of this data reveals a band gap width of 2.19 eV, which is narrower than the 2.27 eV observed for the pristine surface. Regarding the band's composition it can be noticed that the VB exhibits some minor presence of O 2*p* and Cu 3*d* states from the molecule. Sulphur states are also present at deep energy levels within the VB, but their intensity is very small and thus not distinguishable in the graph.

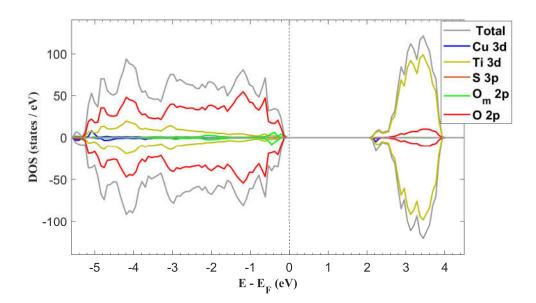


Figure 7. Total DOS and partial DOS (PDOS) of anatase TiO₂ (101) surface with one adsorbed molecule. The corresponding band gap width is 2.19 eV.

The DOS exhibits asymmetrical contributions from the molecular atomic states, resulting in an induced magnetic moment within the system. Following the adsorption process, the total magnetic moment measures 0.97 μ B, arising from both the molecule and its interacting surface atoms. The predominant contribution comes from the Cu 3d states at 0.72 μ B, while the Ti 3d and O 2p states contribute 0.02 μ B and 0.23 μ B, respectively.

The induced magnetic moment coupled with bandgap narrowing, could alter the optical absorption properties of the material and potentially enhance the catalytic performance. This phenomenon aligns with previous experimental observations, including the favorable influence of CuSO₄ on the catalytic activity observed in [25] and the reported photoexcitation under blue light irradiation [26]. While these studies attribute enhanced activity to the formation of intermediate energy levels within the bandgap, our computational analysis suggests that the effect primarily originates from bandgap narrowing.

4. Conclusions

This study presents a first-principles analysis of CuSO₄ adsorption on the anatase TiO₂ (101) surface, employing density functional theory (DFT) calculations. The key findings are as follows:

(1) Adsorption geometry: A preferential adsorption pattern was identified, characterized by the lowest adsorption energy of -4.31 eV. This configuration was consistently obtained regardless of the initial approach of the molecule to the surface, suggesting its prevalence in real samples.

- (2) Adsorption Mechanism: The substantial adsorption energy and the formation of new chemical bonds between CuSO₄ and the TiO₂ surface provide strong evidence for a chemisorption process.
- (3) Charge Transfer: Atomic charge analysis reveals a net charge transfer from the TiO₂ surface to the CuSO₄ molecule. Atoms directly involved in forming new chemical bonds exhibit increased net atomic charges, indicating the ionic nature of these bonds.
- (4) Electronic Structure Modification: The adsorption process induces a magnetic moment and a slight reduction in the bandwidth. These electronic structure alterations are expected to impact on the electronic and catalytic properties of the material.

Author Contributions: Conceptualization, F.M., D.C., A.S., S. A. and J.C.; methodology, F.M., D.C., S.A; software, F.M.; validation, S.A., J.C. and A.S.; formal analysis, D.C, F.M, S.A..; investigation, F.M., D.C., A.S., S. A. and J.C.; writing—original draft preparation, F.M., D.C., A.S., S. A. and J.C.; X.X.; writing—review and editing, F.M., D.C., A.S., S. A. and J.C.; funding acquisition, D.C. All authors have read and agreed to the published version of the manuscript.

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