Communication

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Correction of Multispectral Singlet Oxygen Luminescent Dosimetry (MSOLD) for Tissue Optical Properties in Photofrin-Mediated Photodynamic Therapy

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Abstract: Direct detection of singlet-state oxygen ($^1O_2$) constitutes the holy grail dosimetric method for type-II photodynamic therapy (PDT), a goal that can be quantified using multispectral singlet oxygen near-infrared luminescence dosimetry (MSOLD). The optical properties of tissues, specifically their scattering and absorption coefficients, play a crucial role in determining how the treatment and luminescence light are attenuated. Variations in these properties can significantly impact the spatial distribution of the treatment light and hence the generation of singlet oxygen, and the detection of singlet oxygen luminescence signals. In this study, we investigated the impact of varying optical properties on the detection of $^1O_2$ luminescence signals during Photofrin-mediated PDT in tissue-mimicking phantoms. For comparison, we also conducted Monte Carlo (MC) simulations under the same conditions. The experimental and simulations are substantially equivalent. This study advances understanding of MSOLD dosimetry during PDT.

Keywords: photodynamic therapy (PDT); singlet oxygen luminescence; tissue optical properties

1. Introduction

Multispectral singlet oxygen ($^1O_2$) luminescence dosimetry (MSOLD) is under development for direct detection of singlet-state oxygen near-infrared (~1270 nm) for photodynamic therapy (PDT) dosimetry [1–3] which is considered the holy grail dosimetric method for type-II PDT. The weak nature of the singlet oxygen ($^1O_2$) signal arises from its extremely short lifetime in biological media due to its high reactivity with biomolecules [4–9]. Several techniques have been developed to detect the SO luminescence, including time- and wavelength-resolved spectroscopy [10–12]. Despite significant technological advances, applying these techniques in clinical settings remains a considerable challenge. Nevertheless, several research groups have successfully achieved direct detection of $^1O_2$ luminescence, both in vitro and in vivo using InGaAs photodetectors [1–4]. Moreover, it has been demonstrated that the $^1O_2$ signal can be detected in clinical environments using an optical fiber probe combined with a continuous wave (CW) laser [13,14]. While overcoming the challenges of signal detection is crucial, it is equally important to consider the effects of tissue optical properties to ensure accurate and absolute PDT dosimetry using the $^1O_2$ luminescence signal, since they significantly influence the propagation and intensity of the detected signal [15–17].

Spatial heterogeneities of optical properties in different tissue sites and variations in blood content and pigmentation are common within individual organs or tumor in the same patient and

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across different patients [18]. and the distribution of treatment laser light within the target tumor tissue and $^1O_2$ detection are affected by variations in these properties [19]. This may be misinterpreted as non-uniformity of dose deposition. Moreover, variations in optical properties lead to significant differences in the volume treated by the laser during PDT [20].

Photofrin (Pinnacle Biologics Inc., USA) is the first clinically-approved photosensitizer and one of the most widely studied and utilized agents in photodynamic therapy (PDT) [21]. Derived from hematoporphyrin derivative (HpD), Photofrin is a mixture of oligomers formed by ether and ester linkages of porphyrin units [22]. Upon activation by light at a wavelength of approximately 630 nm, Photofrin generates reactive oxygen species, $^1O_2$, which induces cytotoxic effects on targeted cancer cells. Photofrin's ability to localize in tumors is due to its preferential uptake compared to normal tissues [23].

In this study, we used both simulations and experimental approaches to evaluate the effect of tissue optical properties Photofrin-PDT and $^1O_2$ luminescence detection through optic fibers. The wavelength of the treatment laser was 632 nm and $^1O_2$ luminescence was detected at 1270 nm. Tissue optical properties at 632 nm within the clinically-relevant range were: absorption coefficient (μ_a) = 0.1 – 1.0 cm^{-1} and reduced scattering coefficients (μ_s') = 5 – 40 cm^{-1} [18]. The additional μ_a due to 100 mg/kg of Photofrin was measured as 0.504 cm^{-1}. μ_s' at wavelengths (λ) of 632 nm and 1270 nm was approximated by Mie scattering theory, using $\mu_s'(\lambda) = A\lambda^{-b}$ [24]. We simulated the spatial distribution of the treatment laser and the detection of the $^1O_2$ signal using a virtual optical fiber with the same core diameter and numerical aperture (NA) as the experimental fiber. Additionally, we compared the $^1O_2$ signal from the Monte Carlo simulation with the experimentally detected signal from Photofrin in phantoms.

2. Materials and Methods

2.1. Optical Properties Measurement for Liquid Phantom

Black India ink (Higgins Inc., IL, USA) and Nutralipid (20% B. Braun Medical Inc, PA, USA) were used, respectively, to mimic and μ_a and μ_s' of the tissue in liquid phantoms. μ_a of the ink in methanol was measured as a function of concentration in a 1cm cuvette using both a UV-Vis spectrometer (FLMS01765, Ocean Optics, FL, USA) and an InGaAs spectrometer (AvaSpec-NIR256/512-1.7-EVO, Avantes, CO, USA). $\mu_s'$was measured as a function of Nutralipid concentration using a dual-catheter technique [25] in Mie theory was applied to measurements of light intensity *versus* depth using an isotropic point detector and a isotropic point laser source.

2.2. Experimental Set-Up for Singlet Oxygen Detection

We conducted a series of experiments to detect the $^1O_2$ signal from Photofrin in methanol under 632 nm laser excitation using the Avantes spectrometer with a 200 μm slit width, as illustrated in Figure 1(a). This closely replicates the configuration used in previous mouse studies, where a laser with a 1 cm^2 spot was directed onto subcutaneous tumor grafts, and a fiber detector was positioned at the side to detect the $^1O_2$ signal [26].

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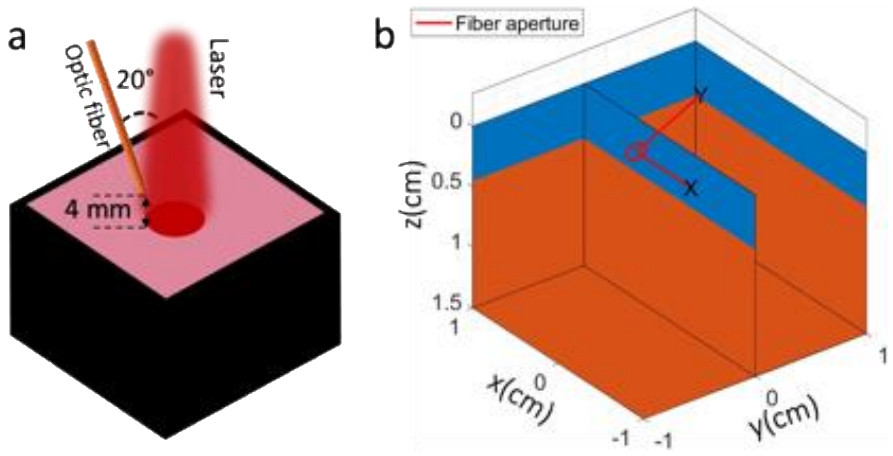

Figure 1. (a). Schematic illustration of the experimental set-up in which the laser strikes the cuvette along the z-axis and the detection fibers with NA= 0.5 are placed at 20 ± 2° to the laser beam with the fiber tip 4 ± 0.5 mm above the surface. (b) Schematic illustration of the Monte Carlo simulation configuration, mirroring the experimental set up.

To investigate how the tissue optical properties modulate the $^1O_2$ spectrum, we prepared a series phantoms with μ_a= 0.6, 0.8, 1.0, 1.2 or 1.5 cm^{-1} (ink + 0.6 cm^{-1} from 100 m/kg Photofrin) and μ_s' = 5, 10 15, 20 or 40 cm^{-1} at 632 nm from the Nutralipid. A 632 nm CW laser (A-1030: Biolitec AG., Austria) excited the photosensitizer through an optical fiber with a micro-lens tip (Pioneer Optics, CT, USA) to produce a uniform 1 cm^2 spot at the sample. A 1.5 mm core optic fiber was connected to the spectrometer to collect the $^1O_2$ luminescence.

The spectral range of 1200-1600 nm adequately covered the 1270 nm $^1O_2$ peak. Singular value decomposition (SVD) implemented in Matlab (The Mathworks, MA, USA) was used to fit the $^1O_2$ spectrum using basis function for the $^1O_2$ peak, the background from the laser, and near-infrared fluorescence from the Photofrin. The last was obtained by quenching the $^1O_2$ with sodium azide (NaN_3), while the $^1O_2$ basis spectrum was obtained by subtracting the quenched spectrum from the unquenched spectrum. The laser background was directly measured from the source.

2.3. Monte Carlo Modeling of Singlet Oxygen Signal

The Monte Carlo model was implemented using MCmatlab [27], which models the propagation of light in a 3D voxel space. The simulation setup is illustrated in Figure 1(b), with the dimensions matching those of the cuvette used to hold the liquid phantom (2 cm ×2 cm ×1 cm). 10^7 simulated excitation photons were launched perpendicular to the cuvette surface. The index of refraction was set at 1.4 and the anisotropy parameter (g) at 0.9. The reduced scattering coefficient in the Monte Carlo simulation matches the measured experimental values. A 1.5 mm diameter optic fiber with a numerical aperture (NA) of 0.5 was positioned at the edge of the laser beam at ~20º angle to the direction of the beam. The fiber collects both 632 nm and 1270 nm light. The singlet oxygen quantum yield is assumed to be 0.5.

Ten million 632 nm photons were launched in each Monte Carlo simulation. The spatial distribution of the 632 nm light within the phantom was simulated from the given optical properties. Generation of the 1270 nm luminescence was then simulated from the 632 nm light distribution and these photons were simulated with corresponding tissue absorption and scattering coefficients at this wavelength. A virtual optical fiber, matching the physical conditions of the experimental fiber, was used to collect the 1270 nm photons exiting the phantom.

3. Results and Discussion

Figure 2(a) illustrates the extinction spectrum of Photofrin from 600 to 800 nm, displaying the peak around 632 nm that is usually used in PDT treatments. Figure 2(b) presents the Photofrin extinction coefficient from 1200 to 1350 nm.

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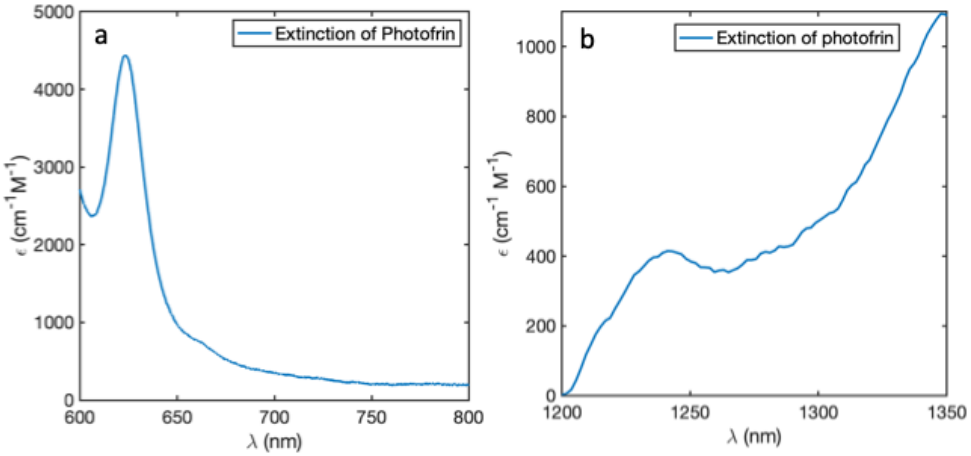

Figure 2. Extinction coefficient spectrum of Photofrin around the excitation (a) and $^1O_2$ luminescence emission (b) wavelengths. The former is consistent with literature values [28].

Figures 3(a) and (c) show the absorption coefficient of black ink for various concentrations in the range of 600-1000 nm and 1200-1350 nm. The extinction coefficient at 632 nm and 1270 nm for 100 mg/kg Photofrin are 0.504 cm^{-1} and 0.06 cm^{-1} respectively. Focusing on the excitation wavelength of 632 nm and the emission wavelength of 1270 nm, the absorption coefficient was calculated and compared with ink concentration as shown in Figures 3(b) and 3(d). This comparison presented a linear relationship between ink concentration and absorption coefficient. This allowed the manipulation of the absorption coefficient in the liquid phantom by adjusting the ink concentration in methanol.

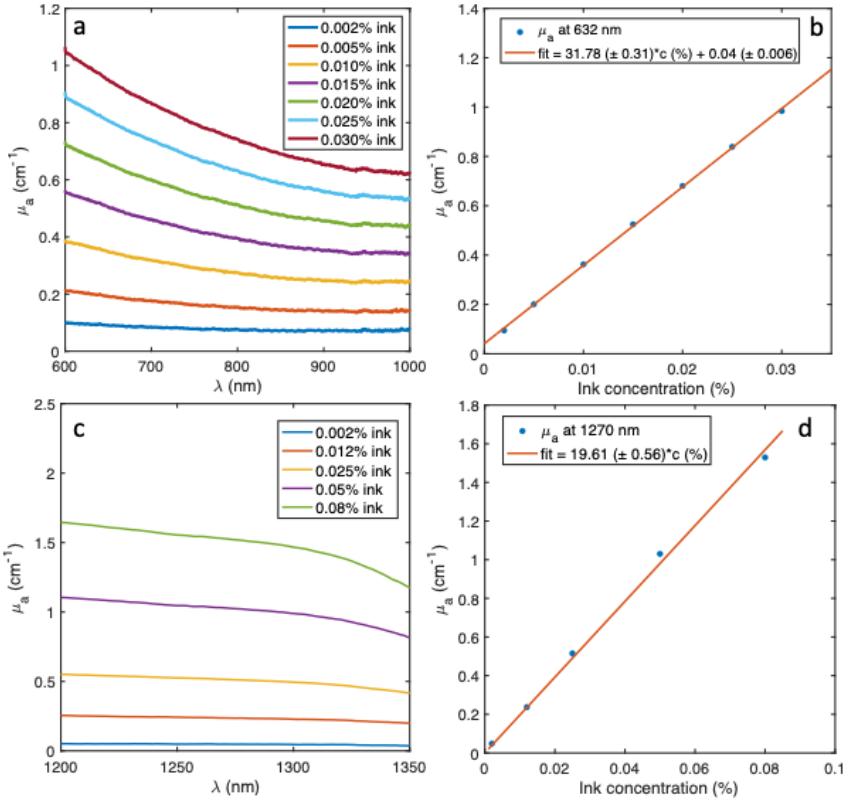

Figure 3. (a) The absorption spectrum of black ink from 600 nm to 1000 nm. (b) Calculated absorption coefficient at 632 nm as a function of concentration, together with a linear fit to the measurements. (c, d) Corresponding measurements from 1200-1350nm and at 1270 nm.

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Figure 4 (a) presents the light fluence rate per unit source strength as a function of distance along the catheter, measured using an isotropic detector and a 661 nm laser source positioned in two parallel catheters. The measured spectrum was fitted using a MATLAB program based on Mie theory, allowing the extraction of the scattering coefficients. Four phantoms were prepared with a fixed ink concentration, while the Nutralipid concentration varied from 0.3 to 0.6%. A linear relationship is seen between the measured scattering coefficient and the Nutralipid concentration in Figure 4 (b). The non-zero intercept in the fit likely stems from minor experimental imperfections or limitations.

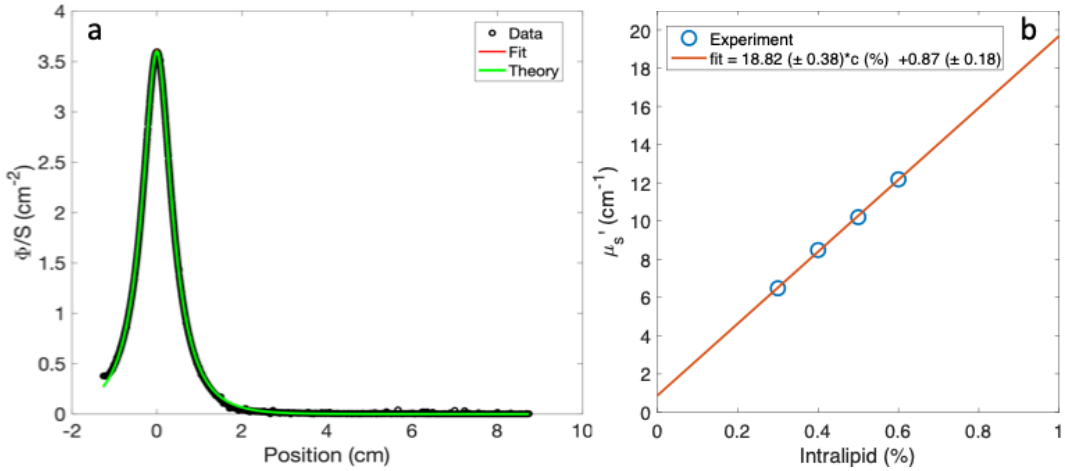

Figure 4. (a) Example of the measured light fluence rate per unit source strength (Φ/S) at 661 nm *versus* distance along the catheter from a point source: fitted $\mu_s' = 8.5\ cm^{-1}$. (b) Reduced scattering coefficient from (a) at 661 nm for Intralipid as a function of concentration and the corresponding linear fit.

Figure 5 presents the measured $^1O_2$ luminescence spectrum of Photofrin in a turbid medium with $\mu_a = 1.0\ cm^{-1}$ and $\mu_s' = 15\ cm^{-1}$. Using Singular Value Decomposition (SVD) fitting, the singlet oxygen signal, Photofrin luminescence and laser background were separated into distinct individual curves, as illustrated in Figure 5.

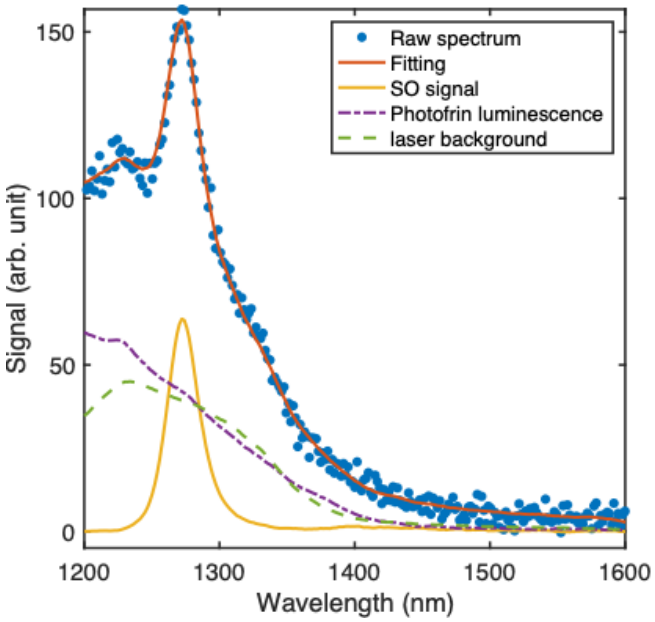

Figure 5. Singlet oxygen luminescence spectrum of Photofrin, measured in a turbid phantom with $\mu_a = 1.0\ cm^{-1}$ $\mu_s' = 15\ cm^{-1}$, together with the calculated component spectra.

Figure 6 (a) shows spectra and the extracted $^1O_2$ luminescence peak with varying absorption and fixed scattering, showing the inverse dependence. Figure 6(c) shows the corresponding normalized

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singlet oxygen signal as a function of μ_a, together with the corresponding Monte Carlo simulations, showing excellent agreement. Figure 6 (b) shows the spectra and corresponding $^1O_2$ signal for varying μ_s' with μ_a fixed at 1.1 cm^{-1}. The $^1O_2$ signal increase non-linearly with μ_s' in the range 5-20 cm^{-1} and saturates above this. Figure 7(d) shows the normalized intensity of the calculated $^1O_2$ signal in Figure 7(b) and the corresponding Monte Carlo simulation results, demonstrating very good agreement.

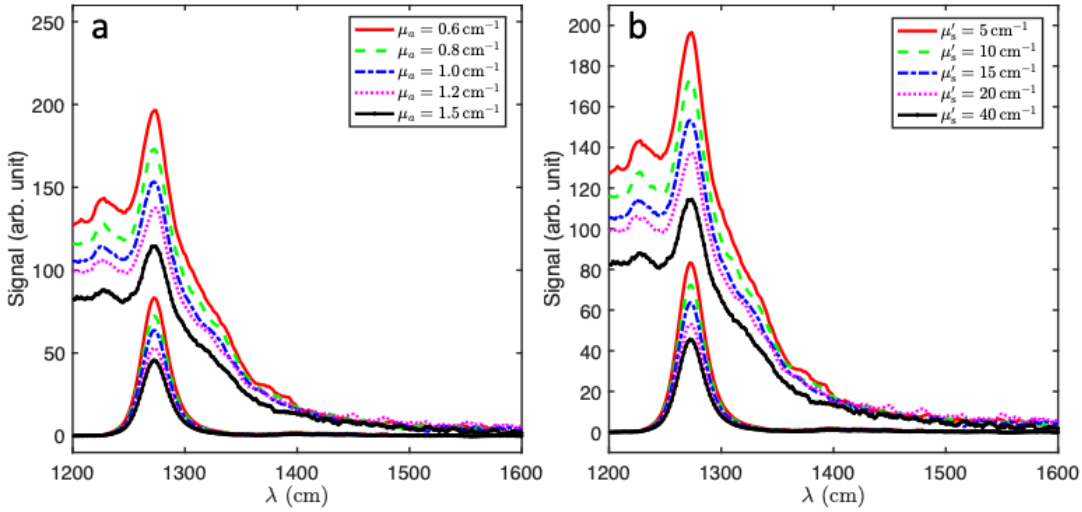

Figure 6. Measured singlet oxygen luminescence spectra with 10-point smoothing the corresponding spectra obtained by SVD fitting. (a) for μ_a ranging from 0.6 to 1.5 cm^{-1} at 632 nm (ink + Photofrin contributions) with μ_s' fixed at 15 cm^{-1}, (b) Corresponding spectra with μ_s' = 5-40 cm^{-1} and μ_a = 0.5 cm^{-1} at 632 nm.

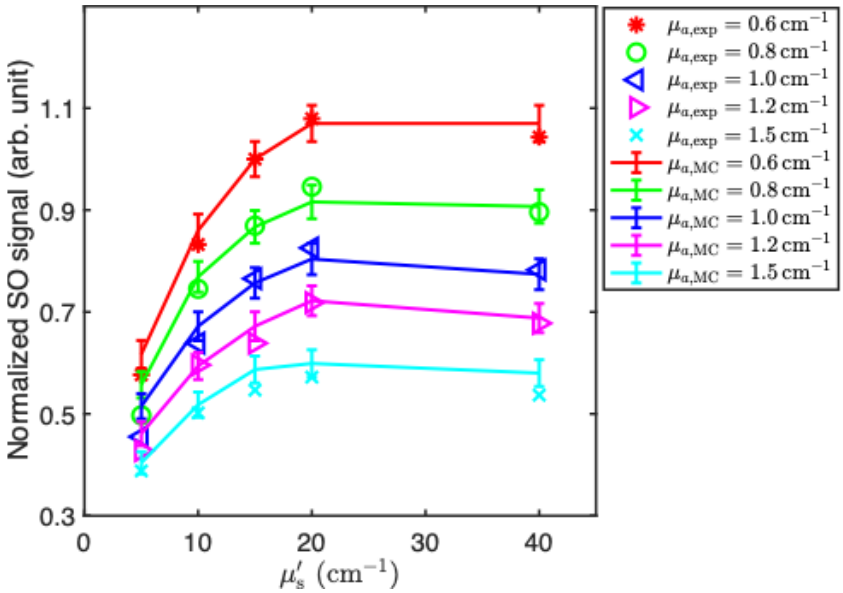

Figure 7. Dependence of the normalized $^1O_2$ signal on the scattering coefficient at 632 nm, comparing the experimental data (points) and the Monte Carlo simulations (solid lines), normalized to the value at μ_a = 0.8 cm^{-1} and μ_s'= 15 cm^{-1}.

In order to compare the experimental data and MC simulations for all 25 combinations of μ_a *and* μ_s', we normalized the results to the reference point where μ_a = 1.0 cm^{-1} and μ_s'= 5 cm^{-1} as shown in Figure 7. The excellent agreement indicates that the MC model can be used to simulate the relative $^1O_2$ signal for any other conditions.

Based on the Monte Carlo simulations, we calculated a correction factor relative to μ_a = 1.0 cm^{-1} and μ_s'= 15 cm^{-1} according to the following equation:

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$$CF_{\mu_a,\mu_s'} = \frac{SO_{ref}}{SO_{\mu_a,\mu_s'}}$$

The results are summarized in Table 1, we can see the correction factor can be as large as 1.97 and as small as 0.69. This large variation of $^1O_2$ signal due to tissue optical properties indicates the necessity of using this correction factor.

Table 1. Correction factor for tissue optical properties for 100 mg/kg Photofrin (μ_a = 0.6 cm^{-1}) as a function of μ_a (ink + Photofrin) and μ_s' at 632 nm, normalized to the Monte Carlo simulation at μ_a= 1.0 cm^{-1} and μ_s' = 15 cm^{-1}.

μ_s' (cm^{-1})	μ_a (cm^{-1})				
	0.6	0.8	1.0	1.2	1.5
5	1.22	1.36	1.47	1.64	1.87
10	0.88	0.99	1.13	1.28	1.46
15	0.76	0.87	1.0	1.13	1.29
20	0.71	0.83	0.94	1.05	1.26
40	0.71	0.84	0.98	1.10	1.31

4. Discussion and Conclusions

Photodynamic therapy (PDT) is a minimally invasive medical treatment that combines a photosensitizing agent, laser light, and oxygen to generate cytotoxic reactive oxygen species (ROS), predominantly singlet oxygen ($^1O_2$). However, antioxidants present in tissues can neutralize ROS by scavenging these reactive species, thereby reducing PDT's effectiveness [29]. High antioxidant levels in tissue can diminish PDT efficacy by lowering $^1O_2$ concentration, which impacts overall therapeutic outcomes. This highlights the importance of directly detecting the near-infrared $^1O_2$ luminescence—a gold standard for accurate PDT dosimetry—since it directly measures the concentration of the primary cytotoxic agent for Type II photosensitizers.

However, variations in tissue optical properties, such as blood oxygenation levels, antioxidant concentration, and pigmentation, can markedly impact the detected signal. To investigate this, we used an Avantes spectrometer to measure the $^1O_2$ luminescence from a liquid tissue-simulating phantom containing Photofrin in methanol, while systematically varying the scattering and absorption coefficients. Employing Singular Value Decomposition (SVD), we successfully extracted the $^1O_2$ signal from the background and Photofrin luminescence in the spectrum. Consistency was observed between the experimental results and Monte Carlo simulations for the normalized $^1O_2$ luminescence. Correction factors were calculated as a function of absorption and scattering, normalized to a specific combination of absorption (μ_a = 1.0 cm^{-1}) and scattering (μ_s' = 15 cm^{-1}). The detected singlet oxygen signal exhibits an approximately linear decrease with tissue absorption coefficient, while it increases with scattering up 20 cm^{-1}, after which it begins to saturate. While these results apply to a sample and irradiation/detection geometries, they demonstrate the strong dependence of the detected singlet oxygen luminescence signal on the tissue optical properties. Hence, independently measuring the absorption and scattering coefficients of the target tissue at the treatment wavelength and around 1270 nm, for example, by diffuse reflectance spectroscopy [30] should significantly improve the accuracy of singlet oxygen luminescence dosimetry for PDT.

In addition, the close agreement between the Monte Carlo simulations and the experimental data for this particular setup indicates that the same approach may be used with some confidence to determine the corresponding correction factors for other target tissue, irradiation and detector geometries. The concordance between experiment and simulations also further validates the use of MSOLD as a singlet oxygen luminescence dosimetry technique.

Author Contributions: Conceptualization, W.Y. and T.Z.; methodology, W.Y., D.S., B.L.; formal analysis, W.Y.; investigation, W.Y.; resources, T.Z., A.D.; data curation, X.X.; writing—original draft preparation, W.Y., M.J.;

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writing—review and editing, T.Z., B.C.W., R.H.H., V.V.; supervision, T.Z.; funding acquisition, T.Z. All authors have read and agreed to the published version of the manuscript.

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Data Availability Statement: The data will be made available upon request.

Conflicts of Interest: Theresa M. Busch reports equity from Simphotek, Inc.

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