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Article

Investigation of Variability of Flaw Strength Distributions on a Brittle SiC Ceramic

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Abstract: The present paper investigates flaw strength distributions established using various flexural tests on batches of SiC bar test specimens: 4-point bending as well as 3-point bending with different span lengths. Flaw strength is given by the elemental stress operating on the critical flaw at fracture of a test specimen. Fracture inducing flaws and their locations are identified by fractography. A single population of pores was found to dominate fracture. The construction of diagrams of p-quantile vs elemental strengths was aimed at assessing the Gaussian nature of flaw strengths. Then, empirical cumulative distributions of strengths were constructed using the normal distribution function. The Weibull distributions of strengths are then compared to the normal reference distributions. The parameters of the Weibull cumulative probability distributions are estimated using maximum likelihood and moment methods. The cumulative distributions of flexural strengths for the different bending tests are predicted from the flaw strength density function using the elemental strength model, and from the cumulative distribution of flexural strengths using the Weibull function. Flaw strength distributions that include weaker flaws potentially present in larger test pieces are extrapolated using the p-quantile diagrams. Implications are discussed about the pertinence of an intrinsically representative flaw strength distribution, with a view to failure predictions. Finally, the influence of characteristics of fracture-inducing flaw populations expressed in terms of flaw strength interval, size, dispersion, heterogeneity, reproducibility with volume change is examined.

Keywords: fracture statistics; strength; normal distribution; Weibull distribution; size effects; flaws

1. Introduction

Microstructural flaws that act as stress concentrators cause the fracture of many brittle materials including ceramics. The flaws generated during manufacturing of ceramic and glass pieces are generally distributed randomly, and they may exhibit wide variability in severity, as a result of variability in shape, nature and size. As a consequence, stress-induced fracture is a stochastic event, and fracture stress is a variate. The strength of brittle ceramics is commonly characterized by statistical distribution functions, which allows the interpretation of strength variability and the prediction of component reliability.

Statistical variability of ceramic strength is commonly characterized by the Weibull distribution function, which is based on a power law. Then, statistical-probabilistic approaches can be used to the fracture stress of specimens or components containing flaws. For this purpose, they should enable the use of laboratory data to infer the probability of fracture of a component under more complicated stress states encountered in service (nonuniform polyaxial stress states, transient stress states). Which requires considerable extrapolation to larger sizes as well as to more complex stress states than those encountered with specimens usually tested under uniaxial stresses. Therefore appropriate distribution functions and statistical-probabilistic approaches are required to describe specimen size effects on strength.

According to Weibull, « the only merit of the distribution function he proposed, is to be found in the fact that it is the simplest mathematical expression of the appropriate form, which satisfies the

necessary general conditions. » [1] However, authors raised questions about the validity of the Weibull distribution for ceramic strengths, owing to variation of estimated statistical parameters, and bias introduced by the method of estimation of Weibull parameters [2–10]. They have compared various distribution functions, and no consensus emerged on the proper form. Thus, for instance, an outcome of the detailed statistical analysis of the flexural strength data [11] using various probability models including Weibull, normal, log-normal, gamma and generalized exponential distributions, was that the gamma or log-normal distribution function, in contrast to Weibull distribution, may describe more appropriately, in certain cases, the experimentally measured strength data. The normal distribution was assumed by a few researchers on bulk ceramics and ceramic fibers [11–16]. However, it was also reported that the fitted Weibull and normal distributions of flexural strengths behave quite similarly [11,12]. The similarity of normal and Weibull distributions has also been established on fiber tows made of very large numbers (up to more than 1000) of filaments [15,16].

It is difficult to decide whether the strength data follow a Weibull distribution or not when small sample sizes and a single batch of specimens of the same dimensions are considered. The validity of a distribution function implies that the statistical parameters are constant whatever the specimen size and the stress-state. In the case of small diameter ceramic fibers tested under tensile load, the estimated Weibull parameters pertinent to flaw strengths display variation depending on various factors, including the methods of analysis of experimental strength data [15–17], the sample size and also the selection of test specimens (sampling) from a tow [15,16]. Sampling dictates the populations of critical flaws present in the tows. As reported in [18], authors did not agree on the most appropriate method of accurate determination of Weibull distribution function for brittle materials.

In a recent paper on ceramic fibers [16], it was shown that the linearity of the p-quantile vs. flaw strength relation indicates that critical flaw strength in a fiber tow is a Gaussian variable. This approach was applied to various fiber types, including SiC, carbon, glass, basalt and alumina fibers [16], to construct empirical flaw strength distributions. It was shown to avoid some of the above sources of bias for the estimation of Weibull statistical parameters [16]. This approach was applied to a monolithic ceramic in the present paper.

There are two main types of probabilistic-statistical approaches to predict the fracture stress of specimens or components containing flaws [2]:

- The phenomenological and macroscopic approaches which consider specimen strengths, like the Weibull model ;
- The approaches that consider the flaws as physical entities [2,19–25]. They are more fundamental. They are based on the flaw strength density function. The severity of a flaw is measured either using flaw strength or flaw size. The flaw strength is defined as the critical value of the local stress operating on a volume element containing a single flaw (also referred to as elemental stress) that causes extension of the flaw. Failure probability of a test piece is derived from the flaw strength density function using the following equation [21–25].

$$P = 1 - \exp \left[- \int dV \int_0^S g(S) dS \right] \quad (1)$$

where $g(S)$ is the flaw strength density function, S is the elemental strength and V is the stressed volume.

A power law equation is generally selected for $g(S)$ because it can be integrated easily for the calculation of probability using equation (1). Determination of the parameters of the power law requires the experimental distribution of flaw strengths. However, unlike fibers, the determination of this empirical distribution of flaw strengths is not straightforward for bulk ceramics since tensile tests are precluded due to practical difficulties. Bending tests are commonly used on ceramics. Therefore, most analyses and predictions of strength distributions are based on specimen strength data (macroscopic approach), neglecting the sources of fracture (the microstructural flaws). The approach proposed in the present paper, is different in essence since the strength of the fracture inducing flaws was considered. The location of critical flaws was identified and flaw strength was derived from measured data.

The present paper investigates flaw strength distributions in a bulk SiC ceramic, with the intent to propose an approach based on p-quantile diagrams to anticipate specimen size effects on the flaw strength density function with a view to predict specimen strength using the elemental strength model under various stress-states and volumes. The Weibull distributions of flaw and flexural strengths were considered for comparison purposes.

2. Materials and Methods

Flaw strengths were determined using flexural tests on SiC bending bars and fractography. The construction of p-quantile diagrams allowed the normal distribution to be assessed, which permitted the construction of unbiased cumulative distribution functions (CDF) and probability density functions (PDF). The Weibull distribution function was then compared to the reference normal distribution. A power law for $g(S)$ was derived from 4-point bending tests. Flexural strength distributions at different 3-pt bending conditions were then predicted using equation (1) and Weibull equation. Then, broader flaw strength distributions that incorporate potential weaker flaws that were not present nor activated due to their location with respect to stress gradient in bending bars, were constructed using the p-quantile diagrams,

2.1. Generation of Flaw Strength Data

Fracture tests were performed in three and four point bending on sintered silicon carbide bars at a constant displacement rate of 1mm/mn. A servohydraulic Instron machine was used. Long and short span lengths were selected for the 3-point bending tests in order to vary the stressed volume. The long spans induce a predominantly uni-axial stress state, whereas appreciable shear stresses may exist in the section of the short span specimens.

The bar specimens were machined out of plates having dimensions of 40x40x10 mm³. The 3 batches of 70 bar test specimens comprised specimens from the different plates. After machining, all the specimens were bevelled and polished on the opposite faces using diamond disks. The specimen dimensions are given in Table 1.

Table 1. Test specimen dimensions.

	Length (mm)	Thickness $2d$ (mm)	Width b (mm)	upper span $2l_1$ (mm)	lower span $2l_2$ (mm)
3 pt long span	45	4	4	0	40
3 pt short span	20	4	4	0	10
4 pt	45	4	4	20	40

The stresses at fracture origins (flaw strengths) were derived from the maximum stress:

$$\text{- 3-point bending: } \sigma_z = \sigma_{max} \left(\frac{l-z}{l} \right) \left(\frac{x}{d} \right) \quad (2)$$

Where σ denotes the stress parallel to longitudinal axis Oz at flaw location, σ_{max} is the maximum stress in the outer surface $\sigma_{max} = \frac{3Fl}{4bd^2}$; x is the distance from the neutral plane, $2d$ is the specimen thickness, z is the distance from the loading plane, $2l$ is the span length.

$$\text{- 4-point bending: } \sigma_z = \sigma_{max} \left(\frac{x}{d} \right) \quad (3)$$

Where σ denotes the stress at flaw location parallel to longitudinal axis Oz, between the loading pins, σ_{max} is the maximum stress in the outer surface between the loading pins $\sigma_{max} = \frac{3Fl_2}{4bd^2}$, $l_2 = l-l_1$, x is the distance from the neutral plane, $2d$ is the specimen thickness.

Fracture mirrors showed the locations of fracture inducing flaws. The plane of fracture origin was always perpendicular to the ox direction. Processing flaws consisting of pores dominated fracture

(> 91%). The strengths of the other flaw types were distributed evenly. Therefore these data were kept. Figure 1 shows an example of pore fracture origin identified in the fracture surface of a 3-pt bending specimen. Fractographic analysis showed that fracture occurred always in the central loaded region of specimens, perpendicular to the direction of longitudinal tensile stress σ_z . For the 4-point bending specimens it occurred in the part of specimens located between the loading pins and close to the outer surface. Therefore, for the 4-point bending tests, the flaw strengths coincided with flexural strengths.

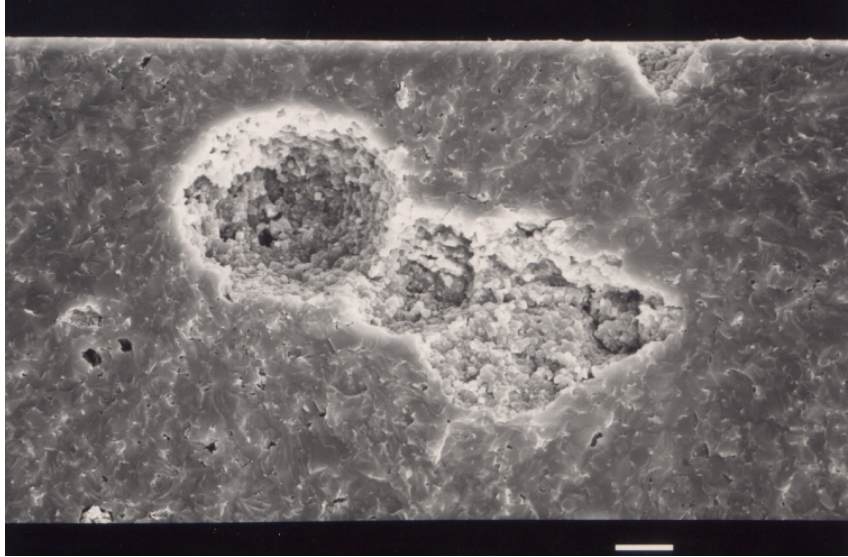


Figure 1. Pores in the fracture surface of a silicon carbide specimen. The scale bar indicates 10 micrometers.

2.2. Analysis of Flaw Strength Data

2.2.1. p-Quantile Diagrams and Normal Distribution [16].

The approach based on p-quantile diagram was proposed in a previous paper to demonstrate that filament flaw strength is a Gaussian variable [16]. It is a graphical method for comparing a Gaussian distribution to a set of data. When X is a Gaussian variable, with μ = mean and s = standard deviation, and N is a variable of the standard normal distribution, it comes:

$$P(X < x) = P\left(\frac{X-\mu}{s} < \frac{x-\mu}{s}\right) = P(N < z) = \Phi(z) \quad (4)$$

Where $P(\cdot)$ is probability, $z = \frac{x-\mu}{s}$ (5), Φ is the cumulative standard normal distribution of variable z .

Thus, when linearity of relation $z(x)$ (equation (5)) is observed for a set of x_i data, one may assume that the x_i data are occurrences of the same Gaussian variable. Then, the plot of p-quantile z_p vs. σ_p indicates whether the strength is a Gaussian variable.

$$z_p = \Phi^{-1}(p) = \frac{\sigma_p - \mu}{s} \quad (6)$$

where $p=i/(n+1)$. z_p can be extracted easily from Φ using a computer or tables available in text books.

2.2.2. Analysis Using the Normal Distribution.

The reference CDFs were calculated using the cumulative normal distribution function (equation (7)) for the values of mean μ and standard deviation s of the sets of experimental strength data:

$$P_n(\sigma) = \frac{1}{s\sqrt{2\pi}} \int_0^\sigma \exp\left(-\frac{(t-\mu)^2}{2s^2}\right) dt \quad (7)$$

The normal CDF are referred to as reference distributions because they provide unbiased graphical representation of the empirical distribution, unlike Weibull plot.

2.2.3. Analysis Using the Elemental Strength Model.

The equation of elemental strength model relates probability to flaw strength density by equation (1). This equation was used to calculate probabilities under the various bending conditions. As detailed in [2], after introduction of power-law expression of $g(S)$, equation (1) becomes:

$$P = 1 - \exp \left[-\frac{1}{V_0} \int dV \left(\frac{\sigma_1}{S_0} \right)^\lambda I \left(\lambda, \frac{\sigma_2}{\sigma_1}, \frac{\sigma_3}{\sigma_1} \right) \right] \quad (8)$$

where $\sigma_1, \sigma_2, \sigma_3$ are the principal stresses, λ and S_0 are statistical parameters. $I(\cdot)$ is a function of local and principal stresses. In the present bending tests, $\sigma_1 = \sigma_z, \sigma_2 = \sigma_3 = 0$; $I \left(\lambda, \frac{\sigma_2}{\sigma_1}, \frac{\sigma_3}{\sigma_1} \right) = I(\lambda, 0, 0)$ [2].

After introduction of S as a function of σ_{max} (equations (2) and (3)) into equation (1), it comes after integration:

$$P = 1 - \exp \left(-V_E(\lambda) I(\lambda, 0, 0) \left(\frac{\sigma}{S_0} \right)^\lambda \right) \quad (9)$$

$V_E(\lambda)$ results from integration of equations (2) or (3). For the 3-point bending specimens:

$$V_E(\lambda) = \frac{V}{2(\lambda+1)^2} \quad (10)$$

For the 4-point bending specimens:

$$V_E(\lambda) = \frac{V(\lambda+2)}{2(\lambda+1)^2} \quad (11)$$

$V_E(\lambda)$ is the statistically equivalent volume subject to $S = \sigma_z$ operating uniformly.

2.2.4. Analysis Using the Weibull Model

Weibull distributions of strength data were calculated using the statistical parameters derived from the experimental strength data. The Weibull distributions functions were compared to the reference normal distribution functions for assessment purpose.

The general equations of probability density function and cumulative distribution Weibull function are respectively:

$$f(\sigma) = \frac{m}{\sigma_z} \left(\frac{\sigma}{\sigma_z} \right)^{m-1} \exp \left[-\left(\frac{\sigma}{\sigma_s} \right)^m \right] \quad (12)$$

$$P(\sigma) = 1 - \exp \left[-\left(\frac{\sigma}{\sigma_s} \right)^m \right] \quad (13)$$

Where m is the Weibull parameter, σ is the strength and σ_s is the characteristic strength.

For the analysis of the influence of dimensions on strengths, the following equation of Weibull distribution was used for flexural tensile stress-state:

$$P = 1 - \exp \left[-\int \frac{1}{V_0} \left(\frac{\sigma}{\sigma_0} \right)^m dV \right] \quad (14)$$

Introducing the expressions (2) or (3) of the stress-state and integrating over the stressed volume it comes:

$$P = 1 - \exp \left[-\frac{V_E(m)}{V_0} \left(\frac{\sigma_{max}}{\sigma_0} \right)^m \right] \quad (15)$$

where $\sigma_R = \sigma_{max}$ is the flexural strength of specimens, σ is the scale factor, $V_E(m)$ is the above statistically equivalent volume for m and $\sigma_{max} = \sigma_z$ operating uniformly, V_0 is a reference volume.

The parameters of the Weibull cumulative distribution functions (CDF) were estimated using the Maximum Likelihood Estimation Method and the Moment Method [18,26,27].

(a) *Maximum Likelihood Estimation (MLE)*

The parameter estimates obtained using the maximum likelihood technique are unique (for a two-parameter Weibull distribution), and they are claimed to statistically approach the true values of the population more efficiently than other parameter estimation techniques as the size of the sample increases. They are derived from equations (16) and (17):

$$\frac{\sum_{j=1}^n (\sigma_R^j)^m \text{Ln} \sigma_R^j}{\sum_{j=1}^n (\sigma_R^j)^m} - \frac{1}{n} \sum_{j=1}^n \text{Ln} \sigma_R^j - \frac{1}{m} = 0 \quad (16)$$

$$\sigma_s = \left[\frac{1}{n} \sum_{j=1}^n (\sigma_R^j)^m \right]^{\frac{1}{m}} \quad (17)$$

where σ_R^j is the j^{th} strength in the set of data

(b) *Method of moments (MM)*

This method has the advantage of involving closed-form equation that can be handled for analysis purposes. The statistical parameters are related to the mean and variance expressions derived from the first moment of the Weibull distribution. The following expressions are used for estimating the Weibull parameters:

$$CV = \frac{s}{\mu} = \sqrt{\left[\frac{\Gamma(1+2/m)}{\Gamma^2(1+1/m)} - 1 \right]} \quad (18)$$

$$\sigma_s = \frac{\mu}{\Gamma(1+\frac{1}{m})} \quad (19)$$

where CV is the coefficient of variation, $\Gamma(\cdot)$ is the Gamma function.

3. Results

3.1. The p-Quantile Diagrams

The p-quantile vs. strength diagrams display linearity (Figure 1) with a high value of correlation coefficient (Table 2). This indicates that the flexural strengths and the flaw strengths follow normal distribution.

Table 2. Coefficients of correlation of p-quantile vs strengths, normal vs. Weibull cumulative (Pn or PW respectively) or probability distributions (PDn or PDW respectively).

	p-quantile vs strength	Pn vs PW	PDW vs PDn
4pt	0.985	0.99	0.93
3pt flex strength	0.988	0.997	
3pt flex strength short span	0.97	0.997	
3pt flaw strength	0.988	0.997	0.90
3pt flaw strength short span	0.99	0.998	0.95
Merged flaw strength	0.995	0.998	

The p-quantile vs flaw strength diagrams were distinct and depend on the bending mode. It can be noted that the difference between the distributions was larger for the flexural strengths than for the flaw strengths. The increase of the slope of the p-quantile diagrams ($=1/s$) with increasing V_E indicates that more weak flaws were present in the 4-pt bending specimens (larger equivalent volume V_E), whereas stronger flaws were present in the short span 3-pt bending specimens (smaller V_E). The long span specimens are an intermediate case. The values of V_E are given in Table 3. When the effective volume increased the high extreme values of flaw strengths decreased, whereas the low extreme values were quite similar (Table 4). The p-quantile diagrams of flaw strengths intersect at the same minimum strength (of about 250 MPa). Thus, the Figure 1 reflects shrinkage of flaw strength

distribution. This indicates that weaker critical flaws did not appear when the effective volume was increased. This may be related to the stress gradient, so that bigger flaws would not be activated by a sufficiently large stress owing to their location. This may also indicate that the spatial distribution of flaws was not reproducible or uniform.

The Weibull parameters were determined using both MM and MLE approaches except for the elemental strength model. Table 3 shows that both methods give very close parameter estimates. The statistical parameters of the elemental strength models were determined by fitting equation (9) to the reference normal distribution of 4-pt bending strength data. The statistical parameters are given in Table 3. There is a discrepancy between the estimated statistical parameters of Weibull distribution of flexural strengths and flaw strengths. This was logically expected since there is no unique relation between flaw location and flexural strength. The statistical parameters show dependence on loading mode and equivalent volume V_E . With increasing V_E , mean strengths μ , standard deviation s and characteristic strength σ decrease, whereas the scale factor and the Weibull modulus of flaw strength distribution increase.

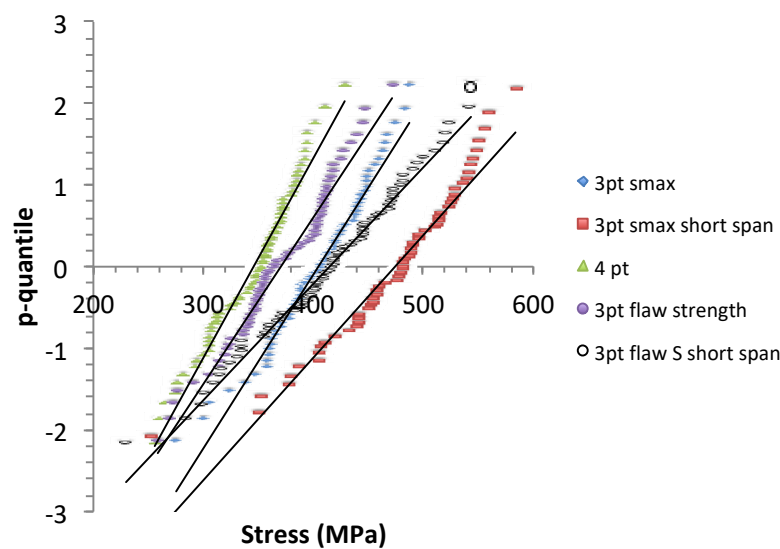


Figure 1. p-quantile diagrams obtained for the sets of flaw strength data obtained on the bending tests. The solid lines indicate the regression lines.

Table 3. Statistical parameters of the normal and Weibull distributions of flexural and flaw strengths. Also given are the statistical parameters of the elemental strength model for the 4pt flaw strengths. Also given are the statistically equivalent stressed volumes V_E . Except for the elemental strength model, the parameters of the strength distributions were determined using both MM and MLE methods.

	μ	s (MPa)	m Weibull	λ	$\sigma_s^{(1)}$ (MPa) Weibull	scale factor (MPa) Weibull	S_0 (MPa) elemental model	V_E (10^{-9} m^3)
4pt flaw strength	345.67	38.46	10.9	10.9	362.06	64.9	57	7.29
			(MM)		(MM)			
			10.8		362.44			
3pt flexural strength	404.67	44.26	(MLE)	(MM)	(MLE)	66.03	1.09	
			11.1		423.73			
			10.97		424,20			

			(MLE)	(MLE)		
			9.3	499.23		
3pt flexural strength	474.27	61.26	(MM)	(MM)		
			9.9	499.0	48.42	0.38
short span			(MLE)	(MLE)		
			9.63	389.82		
3pt flaw strength	370.3	46.2	(MM)	(MM)		
			9.35	390.1	46.99	1.42
			(MLE)	(MLE)		
3pt flaw strength	414.8	66.55	(MM)	(MM)		
			7.36	442.22		
short span			(MM)	(MM)	24.54	0.57
			7.29	442.2		
			(MLE)	(MLE)		
Merged flaw strength	377.15	59.12	(MM)	(MM)	33	
			7.56	401.65		
			(MM)	(MM)		
extrapolated flaw strength	340	76.92	(MLE)	(MLE)		
			6.64	402.93		
			(MLE)	(MLE)		
extrapolated flaw strength	295.28	94.34	(MM)	(MM)		
			5.08	369.97		
			(MLE)	(MLE)		
			7.56	401.65		
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			(MM)	(MM)		
			6.64	402.93		
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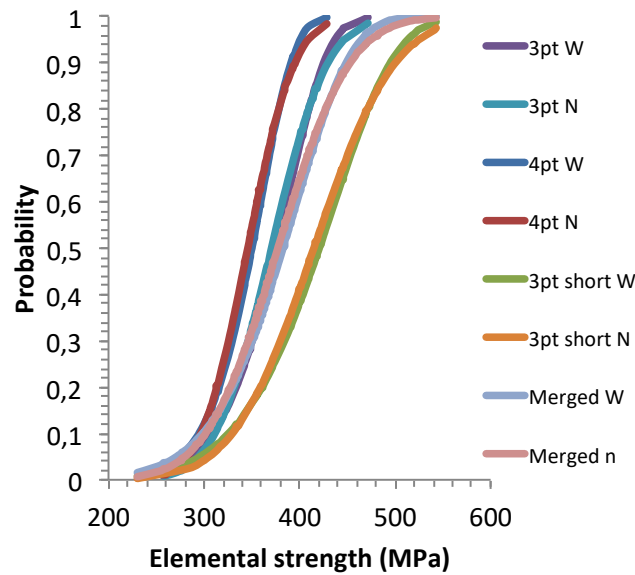


Figure 2. Comparison of Weibull (referred to as W) and normal (referred to as N) cumulative distribution functions for the elemental strengths.

Discrepancy between Gauss and Weibull probability density functions (PDF) can be noted on Figure 3. These curves amplify subtle differences between the normal and Weibull cumulative distributions that did not appear distinctly on Figure 2. This result moderates the above conclusion on correlation between normal and Weibull functions for the flaw strengths. Therefore, it can be considered that the Weibull distribution may be regarded only as an acceptable approximation of a flaw strength distribution.

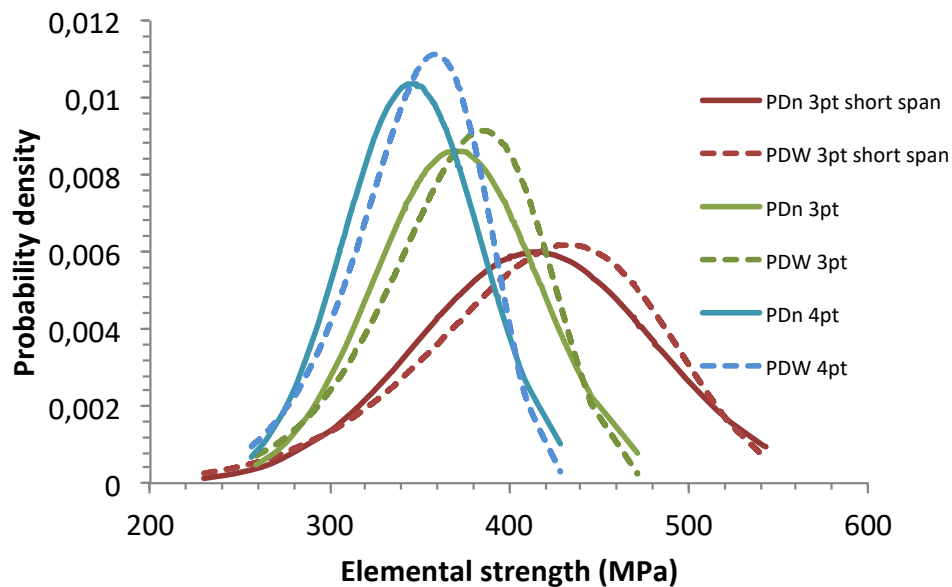


Figure 3. Comparison of Weibull (PDW) and Gauss (PDn) probability density functions.

The PDFs show the same trends as above on the dependence of flaw strengths on the loading mode and the effective volume size. The peak probability density increased with the effective volume size, whereas the flaw strength interval decreased towards the same low strength extreme (Table 1). Like the p-quantile diagrams, this trend describes shrinkage of the flaw strength interval.

3.3. Predictions of CDFs

Failure probabilities were calculated using equations (9) and (15) and the statistical parameters derived from the 4pt-bending tests given in Table 3. The predicted CDFs were compared with the reference distributions of flexural (Figures 4 and 5) and of flaw strengths (Figure 6). The 4-pt bending CDF was well fitted by the elemental strength model for the parameters λ and S_0 given in Table 3 and for $I(\lambda, 0, 0) = 0.25$ [2] (Figure 4). The 3-pt bending CDFs were satisfactorily predicted by the elemental strength model for these latter values of λ , S_0 and $I(\lambda, 0, 0)$ (Figure 4). Figure 5 shows satisfactory Weibull predictions of CDF for the 3-point bending strengths (long span) from the 4-pt bending parameters. By contrast, the Weibull model underestimated failure probabilities for the short span flexural strengths (Figure 5) and did not describe properly the evolution of flaw strengths when the effective volume decreased (Figure 6).

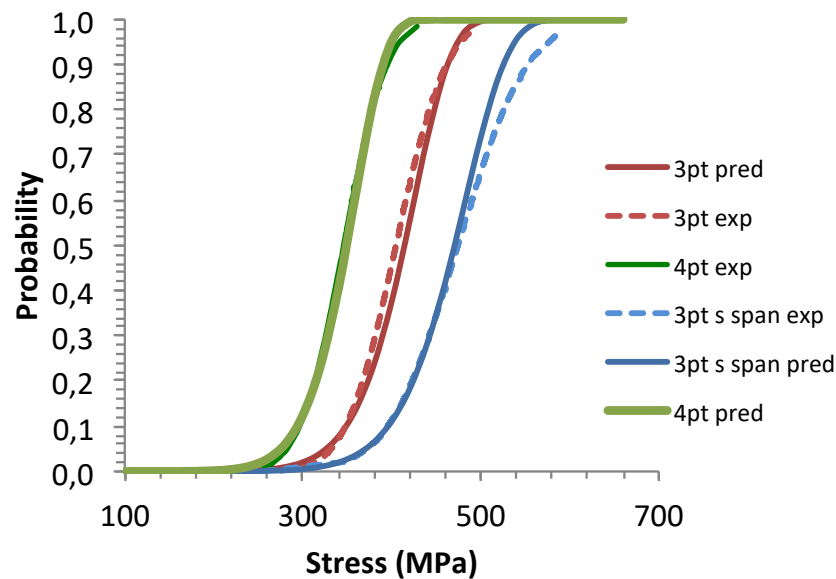


Figure 4. Size effects: Predictions of flexural strengths by the multiaxial elemental strength model equation for the flaw strength parameters derived from 4-pt bending tests. Comparison with the experimental flexural strength distributions.

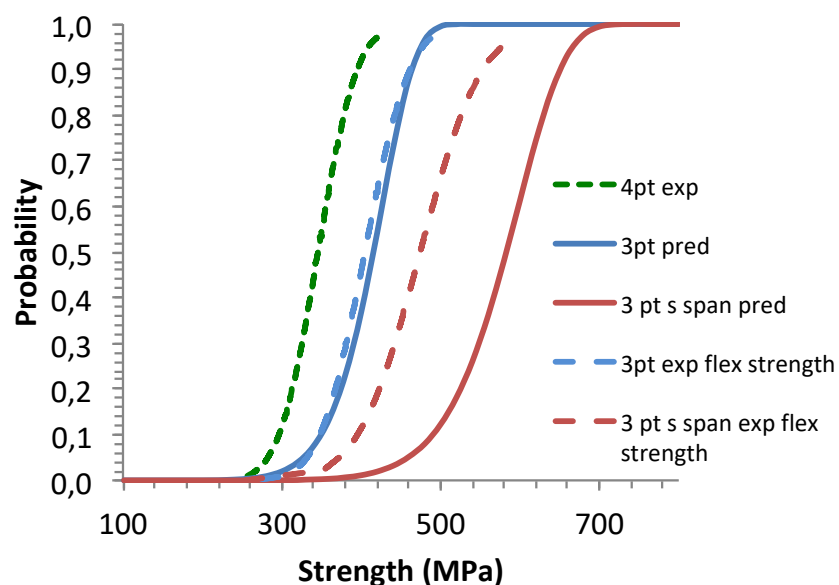


Figure 5. Size effects: Predictions of flexural strengths by the Weibull model equation for the statistical parameters derived from the 4-pt bending tests. Comparison with the experimental flexural strength distributions.

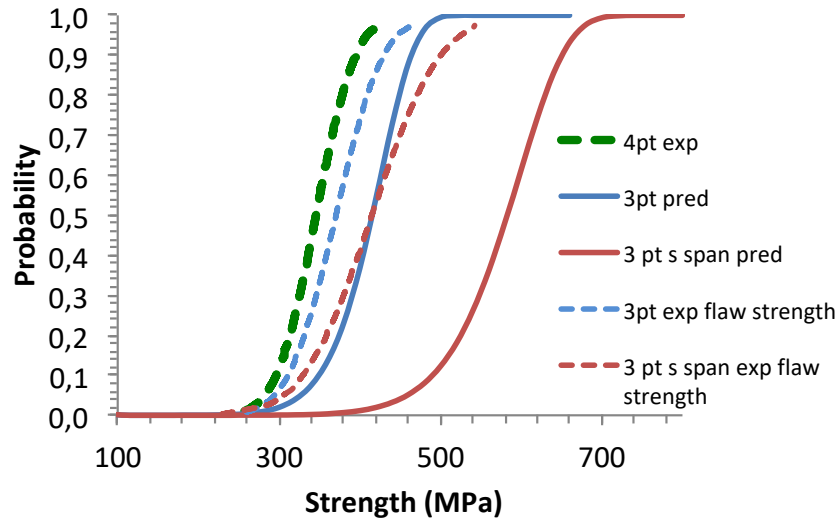


Figure 6. Size effects: predictions of flaw strengths by the Weibull model for the flaw strength parameters derived from 4 –pt bending tests. Comparison with the experimental flaw strength distributions.

4. Discussion : Implications

4.1. Merged Flaw Strength Distribution

The elemental strengths are independent of volume size, being strengths of flaws. Therefore, the three flaw strength data sets were merged. The outcome gives a linear p-quantile diagram (Figure 7) with a high correlation coefficient (Table 2). The result indicates that the merged diagram follows also a single normal distribution function and thus that all the flaws belong to the same family (Figure 7). This allows us to infer that the populations of flaw strengths in test specimens are subsets of a wider set of data. The merged population may characterize the whole population for the batches of test specimens of the present paper, as suggested by the same value for the lowest flaw strengths. However, it may not include more severe flaws that were not activated owing to the stress state that activates unequally the flaws.

The statistical parameters are given in Table 3. They are close to the parameters of the flaw strength distribution derived from the short span 3-point bending tests.

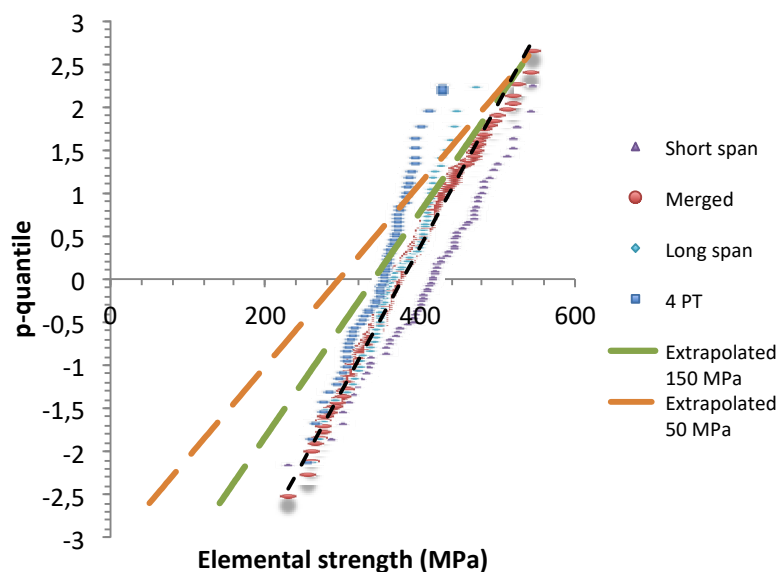


Figure 7. Merged and extrapolated p-quantile – flaw strength diagrams. Comparison with those obtained from the 3-point and 4-point bending tests.

4.2. Extrapolated Flaw Strength Distribution

The comparison of flaw strengths on PDF, or p-quantile diagrams indicates that the underlying flaw populations displayed variation. Thus, they cannot be regarded as representative of the material type. On the basis of the above results, it seems logical to think that the flaw strength distributions obtained on various batches of specimens belong to a unique distribution characteristic of the whole flaw population existing in the material at larger volume size. The existence of a material characteristic flaw strength distribution was evidenced on fiber tows in [28]. Thus, it was shown that there exists a critical tow size above which the flaw strength distribution is invariant. Therefore, an attempt was made to construct a flaw strength distribution that would be representative of the SiC ceramic of the present paper. This population would be obtained on a larger volume such that the underlying flaw population includes flaws weaker than those characterized on the bending specimens and present in the original plates. The extrapolation to distributions including weaker flaws was carried out using the p-quantile diagrams.

A linear p-quantile diagram is symmetrical to the mean strength per construction since z_{pmax} and z_{pmin} are symmetrical to $z_p=0$. These remarkable features enable to construct extrapolated p-quantile diagrams with broader flaw strength distributions.

z_{pmax} and z_{pmin} depend on the number of data through equations :

$$z_{pmax} = \Phi^{-1} \left(\frac{n}{n+1} \right) \quad (20)$$

$$z_{pmin} = \Phi^{-1} \left(\frac{1}{n+1} \right) \quad (21)$$

The slope of the p-quantile diagram is :

$$\frac{z_{pmax}-z_{pmin}}{S_{max}-S_{min}} = \frac{1}{s} \quad (22)$$

Predictions of p-quantile diagrams and CDF for a specified value of S_{min} were carried out as follows. The value of n is selected first. Then z_{pmax} and z_{pmin} are calculated using equations (20) and (21). Then, the values of S_{min} and S_{max} allow the definition of the p-quantile straight line. The value of μ corresponds to $z_p = 0$ and the standard deviation is derived from the slope ($= 1/s$) (equation (22)). From μ and s one can derive the Weibull parameters m and S_s (equations (18) and (19)).

Figure 7 shows the extrapolated p-quantile diagrams for $z_{pmax} = -z_{pmin} = 2.6$, $n=210$, $S_{min}=50$ and 140 MPa, $S_{max}=540$ MPa. Figures 8 and 9 show the corresponding normal CDFs and the PDFs. These CDFs give the probability of existence of a flaw strength in a piece of material containing flaws having strengths larger than 50 or 140 MPa. $S_{min} = 140$ MPa, is the strength of a 1mm flaw, and $S_{min} = 50$ MPa, is the strength of a 8 mm flaw, when $K_{IC} = 5 \text{ MPa}\sqrt{\text{m}}$. 8 mm is quite large with respect to the thickness of the initial plates. However, it may be considered as a conservative limit. The value of low strength extreme can be refined using non destructive testing methods to detect and eliminate those specimens with flaws larger than a given value. S_{max} was taken as the high strength extreme of the merged distribution.

Such representative function would be useful for comparison purposes of materials. It also allows anticipation of appropriate flaw strength distribution for the prediction of strength at large volumes. If the critical volume size existed above which the piece contains the whole population of flaws, there would not be size effect above the critical volume size. Below the critical volume size, representative distribution of flaw strengths depends on the reproducibility of the flaw strength distribution.

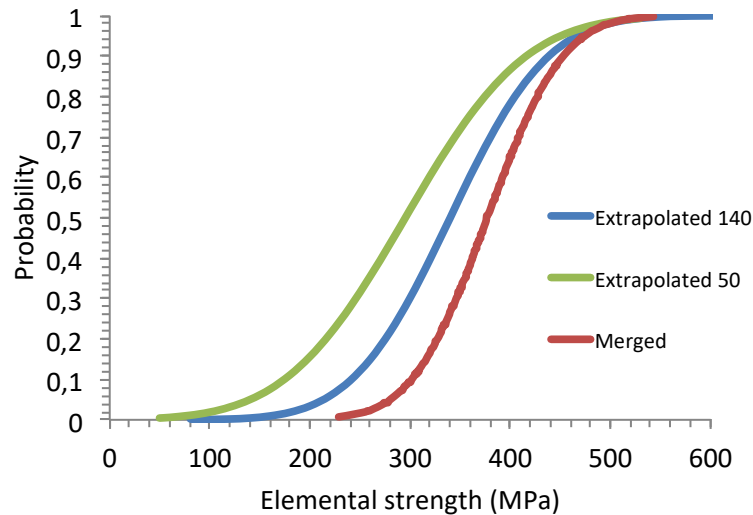


Figure 8. Normal CDFs derived from the merged and extrapolated p-quantile – flaw strength diagrams.

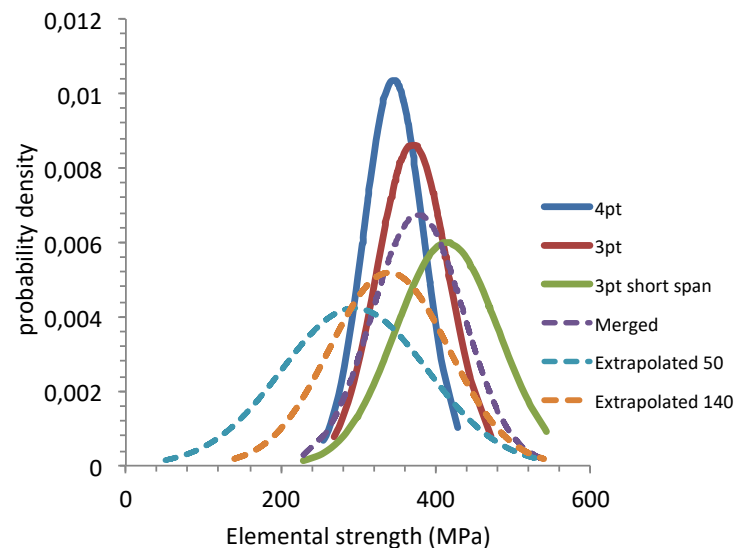


Figure 9. Gaussian PDFs derived from the merged and extrapolated p-quantile – flaw strength diagrams.

4.3. Influence of n and Size

The p-quantile diagrams depend on the following parameters: the boundaries of the z_p interval determined by n the number of strength data and the boundaries of the elemental strength interval. These values define the p-quantile diagram, the normal statistical parameters s and μ , and the Weibull parameters.

In an uniform tensile stress state, the increase of volume size may be regarded as equivalent to an increase of the number of flaw strength data n . In the presence of a stress gradient, this trend may be affected, as weak flaws may not be activated due to their location with respect to stress magnitude, whereas stronger flaws may be activated instead. A larger number of bending specimens is thus required to increase the chance that a broader spectrum of flaws is activated.

4.3.1. Influence of n on p-Quantile Maximum Values

The influence of n on $z_{pmax} = -z_{pmin}$ was calculated using equation (20). Figure 10 indicates that z_{pmax} increases slowly with n , for $n > 200$. But, there is no limit.

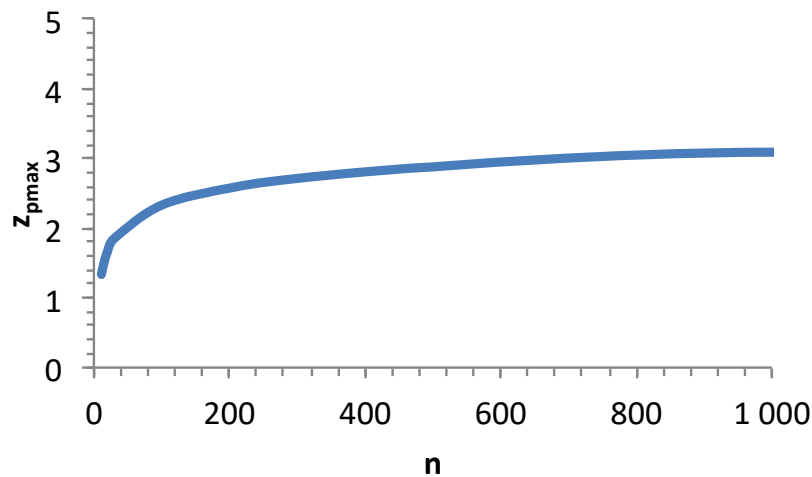


Figure 10. Influence of the number of flaw strength data n on the extreme values z_{pmax} ($=-z_{pmin}$).

4.3.2. Influence of n and Volume on Statistical Parameters

From the equation (26) of the slope of the p-quantile diagram it comes for standard deviation:

$$s = \frac{S_{max} - \mu}{z_{pmax}} = \frac{S_{max} - \mu}{\Phi^{-1}\left(\frac{n}{n+1}\right)} \quad (23)$$

Figure 11 shows that equation (24) is a good approximation of coefficient of variation ($CV = \mu/s$; equation (18)) for the range of m values observed on ceramics. It is not significantly affected by variation of constant around 1.2.

$$CV = \frac{1.2}{m} \quad (24)$$

Trends in statistical parameters variation with the value of n or of volume size were determined using equations (23) and (24). Main results are reported in Table 5.

When the flaw strength interval and μ are constant, s is a decreasing function of increasing n (Figure 12). As a consequence, the Weibull modulus increases with n as shown by Figure 13. It appears that there is no maximum value for m . Thus, it may be difficult to identify an appropriate estimate of m . For the purpose of comparison of different distributions, it is thus necessary to consider batches of identical numbers of specimens.

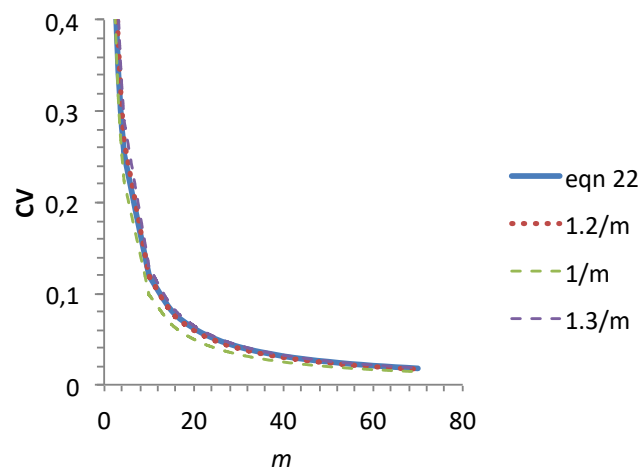


Figure 11. Variation of coefficient of variation (equation (18)) with m . Comparison with functions $1/m$, $1.2/m$ and $1.3/m$.

Table 5. Main trends in the influence of n and V changes on Weibull modulus.

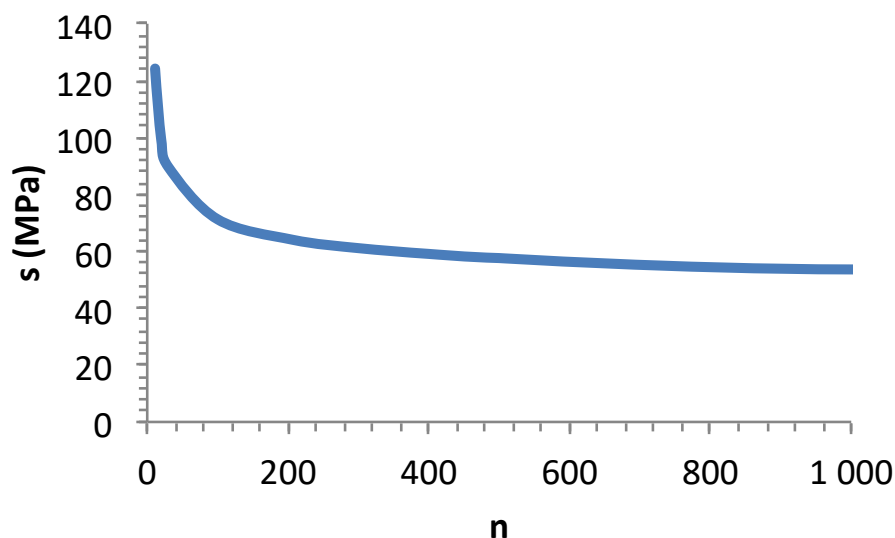
n, V	S_{min}	S_{max}	mean μ	slope $1/s$	s	m
$n \nearrow V=cst$	constant	constant	constant	\nearrow	\searrow	\nearrow
$n \nearrow V=cst$	constant	\nearrow	\nearrow	\nearrow	\searrow	\nearrow
$n=cst V \nearrow$	constant	\searrow	\searrow	\nearrow	\searrow	indeterminate
$V \nearrow n \searrow$	constant	\searrow	\searrow	\searrow	\nearrow	\searrow
$V \nearrow$	\searrow	\searrow	\searrow	\nearrow	\searrow	indeterminate
$V \nearrow$	\searrow	\searrow	\searrow	\searrow	\nearrow	\searrow

When the flaw strength interval is constant and shifted towards lower values, μ decreases and s is a decreasing function as the slope increases with n . The trend in Weibull modulus variation is undetermined on the basis of equation (24) (Table 5). m may increase, decrease or remain constant. It remains constant only if μ and s remain proportional. If n is constant, the slope and s are constant so that m decreases according to equation (24). This latter case was not reported in Table 5.

When the flaw strength interval is changed and shifted to lower values, μ decreases and s may increase or decrease with increasing n . m decreases with s increase. Then, its variation is indeterminate when s decreases (Table 5). As before, it remains constant only if μ and s remain proportional.

Both latter cases can be encountered at constant or increasing volume. Thus it appears that m does not take unique value. It depends on the variability of the underlying flaw strength population in the test specimens from batch to batch, characterized by the variation in standard deviation (Table 5). m constant requires that μ and s are proportional according to equations (18) and (24). It can be anticipated that this requirement implies a reproducible spatial distribution of the flaws.

These results are consistent with many papers that reported a big variation of m estimates from various laboratories testing different batches of specimens. Obviously, this is an issue for the prediction of the volume dependence of strength. For this reason, the determination of a characteristic flaw strength distribution as proposed in this paper, is an interesting alternative that provides the probability of presence of flaws in the material and not only in small specimens. It can be used for comparison purposes of materials or components when those components with flaws larger than a given value have been eliminated. Furthermore failure predictions on the components can be performed with the appropriate parameters instead of the parameters estimated on small specimens that may display significant variability.

**Figure 12.** Influence on n on standard deviation when the flaw strength interval and μ are constant.

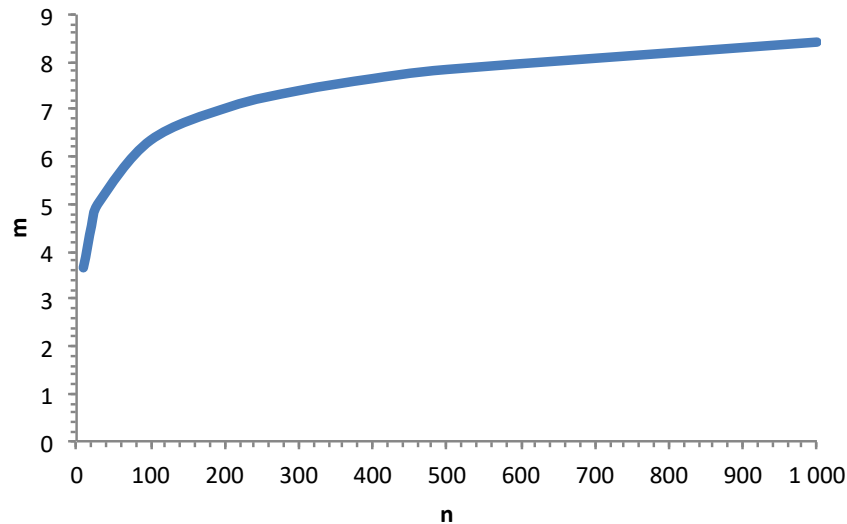


Figure 13. Influence of n on m when the strength interval and the mean strength are kept constant.

4.4. Validity of the Uniaxial Analysis for Short Span

In a previous paper [29] it was shown that the shear components can be reasonably neglected when $d/l=0.4$. At short span local perturbations can arise in the vicinity of supports and loading point. It was found that the equation of the elementary beam theory provides a satisfactory value of strength [24]. The possible effect of loading points is not detrimental since it concerns local compressive stresses. Fractography of the broken specimens clearly showed that failure occurred from the tensile part of specimens. Furthermore the plane of fracture origin was always perpendicular to the longitudinal stress. This suggests that the elemental strength was parallel to the longitudinal stress. This supports the validity of the uniaxial elemental strength model for the short span specimens of this paper.

5. Conclusion

The p-quantile diagrams showed that the flaw strength data followed normal distributions in different non uniform stress states, as well as when they were merged into a single set. The flexural strengths also shown normal distribution.

The strength data also followed Weibull distribution functions. Maximum Likelihood and Moment Method gave very close estimates of Weibull distribution parameters.

The fit of Weibull distributions to normal distributions was acceptable.

The elemental strength model permitted satisfactory predictions of flexural strength distributions from the flaw strength distribution of the 4-point bending specimens.

Weibull predictions from 4-pt bending strength distributions overestimated the 3-pt bending flaw strength distributions and the short span flexural strength distributions.

The strength distribution may characterize only the batch of test specimens used for the determination of strength data. The distribution of critical flaw strengths, may vary from batch to batch. The population of critical flaws is characterized by the size given by the number of flaw strength data (n), the weakness given by the strength interval (S_{min} , S_{max}), and the mean strength (μ), the dispersion given by the standard deviation (s), the reproducibility given by CV . The Weibull modulus variability results from the variation of these parameters. It remains constant when n or V increase only if μ and s remain proportional (constant CV).

The extrapolation of experimental p-quantile diagrams by including weaker flaw strengths is an alternative approach to size effects based on the probability of existence of critical flaw strengths. It allows also anticipating flaw strength distributions in larger specimens for failure predictions using the elemental strength model.

The critical specimen size for which the flaw strength population is representative of the total critical flaw population inherent to the material type can be approached using p-quantile diagrams and a technique of detection of the most deleterious flaws.

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