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# **Nuclear Physics Opportunities at European Small-Scale Facilities**

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**Abstract:** Small-scale facilities play a significant role in the landscape of nuclear physics research in Europe. They address a wide range of fundamental questions and are essential for teaching and training personnel in accelerator technology and science, providing them with diverse skill sets, complementary to large projects. The current status and perspectives of nuclear physics research at small-scale facilities in Europe will be given.

**Keywords:** small-scale facilities; nuclear instrumentation; detector arrays; accelerators; nuclear structure; nuclear reactions; nuclear astrophysics; accelerator-based nuclear science; education and training

## 1. Introduction

In 1929, Van de Graaff developed one of the first accelerators and the era of accelerator-based nuclear physics began. Almost one hundred years later, the pursuit continues using the complex, high-energy machines operating at large nuclear research centers like GSI/FAIR, GANIL and others. Meanwhile, the small-scale accelerator facilities still play a major role in keeping and providing knowledge in nuclear physics. The small-scale facilities in Europe encompass a large set of setups offering the European nuclear physics community a large variety of stable beams as well as a number of radioactive beams. They provide the possibility to carry out state-of-the-art research using different instrumentation available. They have a significant role in high-level education and training as well as the formation of the early career stage researchers in nuclear physics. Small-scale national facilities have an important role in providing know-how and expertise in developing novel experimental techniques and instrumentation for the big-scale flagship facilities from the European Strategy Forum on Research Infrastructures (ESFRI) roadmap [1]. The Nuclear Physics European Collaboration Committee (NuPECC) has strongly recommended the continuation and expansion of the training initiatives at the small-scale facilities and endorsed the support of the small-scale facilities' nuclear physics programs in its latest Long Range Plan (LRP) [2]. The current status, available instrumentation, as well as perspectives of nuclear physics research at small-scale facilities, are given. The list of reviewed facilities can be found in Table 1.

**Table 1.** The list of European small-scale facilities reviewed in this paper. The facility, country, accelerators available as well as the beams that can be provided are shown.

Facility	Country	Accelerators available	Beams
Cologne Accelerator Laboratory	Germany	10 MV FN-Tandem	Any specie that can form negative ions
		6 MV Tandetron	AMS*
Heavy Ion Laboratory	Poland	Isochronous heavy ion cyclotron (K <sub>max</sub> =160)	He-Xe
IJC Lab		15 MV Tandem (ALTO)	Proton to aggregates
	France	50 MV electron accelerator (ALTO)	e <sup>-</sup>
		4 MeV NEC pelletron accelerator (ANDROMEDE)	p-Au, He, Ne, Ar, Kr, Xe

		Licorne neutron source	n
Rudjer Bošković Institute	Croatia	6.0 MV EN Tandem	p to Au
Accelerator Facility		1 MV Tandetron	p to Au
Tandetron laboratory in Piešt'any	Slovakia	2 MV Tandetron	p, d, α
The Oslo Cyclotron Laboratory	Norway	MC-35 cyclotron	<sup>1</sup> H+, <sup>2</sup> H+, <sup>3</sup> He <sup>2+</sup> and <sup>4</sup> He <sup>2+</sup>
S-DALINAC	Germany	Superconducting electron accelerator	e <sup>-</sup>
The INFN Laboratori		13 MV TANDEM - HVEC MP	p to Au
Nazionali del Sud	Italy	Superconducting Cyclotron K800	H to U, also exotic beams see Sec. 2.8.
The Atomki Accelerator Center		MGC-20E Cyclotron	p, d, <sup>3</sup> He <sup>2+</sup> and <sup>4</sup> He <sup>2+</sup>
	Hungary	1 MV Van de Graaff	p, d, He
		5 MV Van de Graaff	Inactive
		2 MV Tandetron	P, He, B, C, O, S, Si, Cu etc.
		TR-24 cyclotron	р
CANAM RI	Czech R.	U-120 M cyclotron	$H^{+},H^{-},D^{+},D^{-},^{3}He^{2+},^{4}He^{2+}H^{+},D^{+},\\ ^{3}He^{2+},^{4}He^{2+}$
		MT25 microtron	e-
		3 MV Tandetron	Almost all elements of the periodic system
Bronowice Cyclotron Center	Poland	Proteus C-235 cyclotron	р
CNA, Seville		18/9 MeV cyclotron	p, d
	Spain	3 MV Van de Graaff Tandem 1 MV Tandetron	virtually all types of stable ions AMS*
	3	HiSPANoS neutron source	n
	Portugal	2.5 MV Van de Graaff	H, <sup>3</sup> He, <sup>4</sup> He and heavier ions
LATR, Lisbon	0	3 MV Tandem	H, <sup>3</sup> He, <sup>4</sup> He and heavier ions
INFN-HH		9 MV FN Pelletron Tandem	p to Au
		3 MV HVEE Tandetron	p to Au
		1 MV HVEE Tandetron	AMS*
		400 kV HVEE electrostatic	Ut Hot
		accelerator	H <sup>+</sup> , He <sup>+</sup>
LUNA	Italy	3.5 MV Cockroft-Walton accelerator	p, α, ¹²C
Felsenkeller underground accelerator laboratory	Germany	5 MV Pelletron accelerator	p, α, <sup>12</sup> C
CMAM, Madrid	Spain	5 MV HVEE Pelletron Tandem	p to Au
TAL, Demokritos	Greece	5.5 MV Van de Graaff Tandem	n (via neutron-induced reactions), stable beams
		250 keV single-stage accelerator PAPAP	p, d
		17 MeV Scanditronix Cyclotron 2.5 MV Tandetron	p, d AMS*
MIC, JSI	Slovenia	2 MV Tandetron	H, <sup>3</sup> He, <sup>4</sup> He and heavier ions

<sup>\*</sup> primarily used for accelerator mass spectroscopy measurements (AMS).

# 2. Small-Scale Facilities

Numerous facilities listed in Table 1. carry out research in both fundamental and applied nuclear physics. Most of the activities in fundamental nuclear research performed at these facilities are related to nuclear structure, reactions and nuclear astrophysics.

# 2.1. Cologne Accelerator Laboratory, Institute for Nuclear Physics (IKP), Cologne, Germany

The Cologne accelerator laboratory provides two accelerators, a 10 MV FN-Tandem and a 6 MV Tandetron accelerator [3]. The latter one serves solely for accelerator mass spectroscopy measurements (AMS) while the 10 MV FN-Tandem is used for research in nuclear structure and nuclear astrophysics. 10 MV Tandem provides beams of basically any ion species (which can form negative ions) over a beam current range between nA and  $\mu$ A. Various experimental setups equipped with several high-purity germanium (HPGe) as well as silicon particle detectors are used to perform research on nuclear reaction rates, nuclear lifetimes and the structure of excited states.

Nuclear astrophysics research is focused on low-energy charged-particle-induced reactions at low energies (of astrophysical interest). The proper knowledge of cross-sections is necessary for reaction network calculations of different stellar scenarios. For  $\gamma$ -ray spectroscopy, the high-efficiency  $\gamma$ -ray spectrometer HORUS (High-efficiency Observatory for  $\gamma$ -Ray Unique Spectroscopy) [4] is available. HORUS comprises up to fourteen HPGe detectors. In order to actively suppress the Compton background, six of them can be equipped with active BGO (bismuth germanate) shields. The detectors are placed at five different angles with respect to the beam axis, enabling the five angular distribution measurements that are required to precisely determine the absolute cross sections.

A detailed understanding of the nuclear physics processes is needed to properly describe different astrophysical scenarios. The  $\gamma$ -process [5] is assumed to be responsible for the largest contribution to the abundance of 35 neutron-deficient p nuclei. It creates a vast network of photo-disintegration reactions that include a plethora of different reactions on mainly unstable and exotic nuclei. In the absence of experimental results, theoretical cross-sections have to be considered. The cross sections at astrophysically relevant energies i.e. inside the Gamow window (the range of energies where nuclear reactions occur in stars) are extremely low (usually of the order of  $\mu B$ ), therefore the development of sensitive measurement techniques is very much needed.

Different direct methods are used to measure the radiative capture cross-sections. One of the most widely applied methods is the activation method [6]. In the activation method, unstable reaction products are first produced and then the radioactive decay, usually the  $\gamma$ -ray transitions from the daughter nucleus, is observed. The limitation of this technique is that it has to yield an unstable reaction product with a half-life that is long enough to transfer the sample to the corresponding counting setup. To overcome this limitation, two different other methods are mostly applied, namely, the in-beam  $4\pi$ -summing technique [7] and the in-beam high-resolution  $\gamma$ -ray spectroscopy technique [8]. The in-beam  $4\pi$ -summing technique uses a large scintillator crystal which covers a solid angle of almost  $4\pi$  around the target position and sums the energies of all  $\gamma$ -rays emitted in a certain time window. In the high-resolution  $\gamma$ -ray spectroscopy technique, the prompt  $\gamma$ -decays of the excited compound nucleus are observed. Absolute reaction cross-sections can be determined by measuring the angular distributions using an HPGe array (e.g. HORUS). Moreover, nuclear structure information such as spin and parity assignments can be obtained. The  $^{89}Y(p,\gamma)^{90}Zr$  commissioning experiment has demonstrated the reliability of the measured cross sections obtained using the HORUS array [9]. SONIC@HORUS is a setup for particle-γ coincidence measurements. It consists of the recently developed particle spectrometer SONIC (Silicon Identification Chamber) which can house up to  $12 \Delta E$ -E telescopes for the identification of ejectiles and measurement of their energy and the spectrometer HORUS which has been used to investigate the nuclear structure by measuring gamma-rays from fusion-evaporation reactions as well as from capture reactions. The identification of the ejectile is necessary to select non-dominating reaction channels like light-ion scattering and specific transfer reactions. By measuring the energy of the ejectile, nuclear level schemes can be studied in detail (the excitation energy of the nucleus can be calculated from the ejectile energy). The knowledge of the complete reaction kinematics also enables a precise and reliable determination of lifetimes using the Doppler-Shift Attenuation Method (DSAM) [10]. DSAM method is commonly used for the measurement of sub-ps lifetimes of excited nuclear states. The excited nucleus (obtained in the nuclear reaction) moves with a certain velocity and slows down during the time interval of the order of ps. The emission of  $\gamma$ -rays during this time interval leads to the observation of a line shape

whose characteristics are sensitive to the time evolution of the population of the level of interest and specifically, to its lifetime.

# 2.2. Heavy Ion Laboratory (HIL), Warsaw, Poland

HIL is a Polish infrastructure that hosts a K<sub>max</sub>=160 isochronous heavy ion cyclotron. The cyclotron provides beams (from He up to Xe) of energies 2-10 MeV/u and intensities up to a few hundred pnA [11]. Experiments can utilize different setups and apparatus including IGISOL [12] – a Scandinavian type on-line separator. The setups are described in the following. Coulomb Universal Detector Array Chamber (CUDAC) [13] is a small scattering chamber. CUDAC can accommodate an array of backwards hemisphere semiconductor detectors and forward-hemisphere monitoring Si counters. It is mostly utilized for Coulomb excitation research and measurements of fusion barrier distributions. The former Strasbourg-based ICARE (Identificateur de Charges a Rendement Eleve) [14] is a particle spectroscopy chamber. ICARE can accommodate up to forty-eight telescopes. Charged particle detection, identification and energy measurement are performed using them. The research performed in this setup encompasses barrier distribution measurements, reaction mechanism studies and novel detector tests. The experimental efforts are devoted to barrier-level distribution studies and experiments using 10,11B, 18O and 14,15N beams on light targets such as 6,7Li, <sup>9</sup>Be, <sup>12,13</sup>C. EAGLE [15] is a Central European Array for Gamma Levels Evaluation. It can host up to thirty HPGe detectors with anti-Compton shielding. EAGLE enables research on nuclear structure by  $\gamma$ -spectroscopy techniques such as  $\gamma$ - $\gamma$  angular correlations and lifetime measurements using methods such are the Doppler Shift Attenuation Method (see section 2.1.) and the Recoil Distance Doppler Shift technique (RDDS) [16] and complex Coulomb excitation (Coulex) experiments relevant for nuclear structure physics. The RDDS method is the standard method to measure lifetimes of excited nuclear states in the ps range. The method uses a plunger device that consists of a stretched target and stopper-foil which are mounted parallel to each other at a variable distance. Excited states in the nucleus of interest are populated in the target foil. The nucleus then recoils with a velocity of a few per cent of the speed of light in the direction of the stopper foil where it gets stopped. The lifetime of a level of interest is determined by the changing intensities of fully Doppler-shifted and stopped  $\gamma$ -ray components. Coincidences between gamma rays and internal conversion electrons enable the determination of transition multipolarities. The EAGLE array could be utilized in conjunction with the ULESE (University of Lodz Electron Spectrometer) [17] electron spectrometer to perform this type of measurement. High efficiency (up to 12% at 300 keV) and good energy resolution (1% at 1 MeV) characterize this electron spectrometer. Additionally, the ULESE electron spectrometer has good suppression of the emitted photons, positrons, and delta electrons. The elimination of delta electrons is crucial in electron conversion spectroscopy performed in "in-beam" measurements. The ULESE spectrometer coupled to the EAGLE array was successfully used in measurements performed at HIL. The main objective of these experiments was to study the violation of the K selection rule for electromagnetic transitions in nuclei. The absolute transition probabilities in 130Ba,132Ce, 134Nd and <sup>184</sup>Pt were measured [18].

Coulomb interaction between the incident and target nuclei can be described by classical electrodynamics, neglecting the nuclear forces. The use of various beams and detectors which measure at broad ranges of scattering angles enable measurements of electromagnetic structure key parameters up to the high-spins. Static moments and transitional matrix elements are related to spectroscopic observables such as lifetimes, gamma-ray intensities, branching and mixing ratios. Coulomb excitation has been one of the leading experimental techniques and is established at HIL, about thirty years ago. Using the CUDAC, the first successful experimental campaigns were conducted. At first, the historical array was coupled to a set of small-size gamma detectors and was later replaced with a compact scattering chamber dedicated to work with EAGLE. A new particle array - SilCA (Silicon Coulex Array) [13] based on the DSSSD detectors is currently under development at HIL Warsaw. One of the main topics of recent studies performed at HIL was focused on the transitional region of the nuclear chart (*A*~100, *Z*~40, 50). In the transitional region of the

nuclear chart, the phenomenon of shape instabilities is relatively common and may lead to coexisting nuclear shapes. Most of the even-even nuclei in this region are also traditionally considered the best examples of vibrational nuclei. However, recent results seriously contradict this simple interpretation. Coulomb excitation studies of shape-coexisting structures in 96,98,100Mo [19] showed that triaxiality plays an important role in this region. The problem of chirality in atomic nuclei with odd numbers of protons and neutrons has attracted a lot of attention in recent years. In these nuclei, the total nuclear spin is built from the valence proton and valence neutron momenta and angular momentum of the even-even core. These three vectors can be mutually perpendicular and coupled in two manners forming left-handed and right-handed systems. Left-handed systems and righthanded systems, namely have opposite chirality in the intrinsic frame of the nucleus. On the other hand, in the laboratory frame, chirality manifests itself as the presence of two rotational bands, nearby degenerated, with the same parities. The measurements of the lifetimes of states belonging to the chiral bands are still lacking. Lifetimes carry important information on nuclear wave functions. Significant results on chiral symmetry breaking in low-energy excitations of atomic nuclei have been obtained at HIL. The gamma-spectroscopy methods, namely, DSAM for lifetime measurements [20], [21] and the time-differential perturbed angular distribution (TDPAD) to determine g-factors [22] have been employed to investigate the chiral bands. EAGLE array was employed for gamma-ray spectroscopy. The TDPAD method is based on the nuclear magnetic moment interaction with the external magnetic field. Due to that interaction, the magnetic moment of the nucleus precesses around the field axis and causes a specific angular distribution of the emitted gamma radiation. The gamma spectrometer (i.e. EAGLE) is used to obtain a good angular sensitivity necessary for this type of measurement. Recent search for a phase transition from not-chiral to chiral structure in <sup>128</sup>Cs [23] has been performed at HIL.

# 2.3. IJClab (Laboratory of the Physics of the two Infinities Irène Joliot-Curie), Orsay, France

IJClab hosts and supports ALTO (Accélérateur et Tandem d'Orsay), a French research platform that hosts two accelerators, unique in France: a 15 MV Tandem accelerator and a 50 MV electron linear accelerator. A linear accelerator is employed for the production of radioactive beams by photofission. The platform enables research in nuclear physics, astrophysics as well as multidisciplinary studies [24]. It also hosts the ANDROMEDE facility. ALTO provides rare energetic beams such as <sup>3</sup>He and <sup>14</sup>C. It is a unique facility in Europe since it produces low-energy neutron-rich beams via uranium photofission. It is also unique in its capacity to provide high-flux naturally directional neutron beam with the LICORNE neutron converter [25]. LICORNE is a high-flux directional neutron source based at the Tandem accelerator of the ALTO facility. It produces intense, kinematically focused quasi-monoenergetic beams of neutrons in the energy range of 0.5 MeV to 4 MeV. The neutron production is achieved using the capability of the ALTO Tandem accelerator to produce an intense primary beam of <sup>7</sup>Li which results in secondary beams of kinematically focused neutrons. The neutron cone allows sensitive detectors to be placed around the sample to be irradiated, the former is particularly crucial if HPGe detectors are used since they are easily damaged by the fast neutrons. LICORNE allows for both high flux and beam collimation.

Decay spectroscopy is performed on the Low Energy Branch using the BESTIOL or alias BEDO (Beta Decay Studies at Orsay) setup [26] on the High Energy Branch using the LICORNE neutron source and the NuBall array. BEDO is a movable-tape-based experiment setup. The tape's trajectory in BEDO is chosen to make the most of the space around the beam collection point, enabling the most effective deployment of various types of detectors. Significant experimental and theoretical interest has been shown in BEDO for nuclei at the Z=28 shell closure and in the mass range of N=50. The gap between the n1g<sub>9/2</sub> and n2d<sub>5/2</sub> orbitals corresponds to the N=50 shell effects and is essential to understanding the formation of this magic closure. Also, it contributes to the understanding of the origin of the strong gaps associated with spin-orbit magic numbers throughout the whole nuclear chart, in which 3-body terms of the N-N interaction are now thought to play a key role [26]. Also, BEDO has done significant research on the region close to N=82.

NuBall array [27] is an array that consists of the HPGe (for precise gamma-spectroscopy) and LaBr<sub>3</sub>(Ce) (for the fast timing measurements) detectors that coupled to the LICORNE directional neutron source enable calorimetry to select reactions and study gamma energy and multipolarities of fission products. The array's key characteristic is its ability to combine the largest peak-to-total ratio for precision gamma spectroscopy with the best time resolution conceivable.

Numerous experiments are covered under the LICORNE physics program. The research focuses on nuclear reactions and the study of the prompt gamma and neutron emission in nuclear fission as well as on the nuclear structure and the production of exotic neutron-rich nuclei. Fission is an important reaction process. It can occur either spontaneously or following a reaction or nuclear decay in the region of heavy and superheavy nuclei. Although significant research on fission has been performed for decades, fission is still rather poorly understood nowadays. One of the recent results concerns the observation of the angular momentum generation in nuclear fission [28]. The splitting of the heavy atomic nuclei is observed to produce spinning fragments. The internal generation of typically six or seven units of angular momentum in each fragment is particularly interesting for systems that start with zero (or almost zero spin). It was shown [25] that the spins of the fragment partners do not significantly correlate, which indicate the conclusion that in fission, angular momentum is created once the nucleus splits (post-scission). The comprehensive data showed that the average spin is strongly mass-dependent, varying in saw-tooth distribution. The lack of observable fragment dependency on partner mass or charge supports the spin mechanism's uncorrelated post-scission nature. The first NuBall campaign at ALTO had a diverse nuclear physics program [29], i.e., studying the superallowed beta-decay of 10C to test the weak interaction and unitarity of the CKM matrix or studying the Giant-Dipole Resonance (GDR - giant resonances are collective excitations of the atomic nucleus) feeding of low-energy structures with different deformations using NuBall coupled to the PARIS array (Photon Array for Studies with Radioactive Ion and Stable beams). The second NuBall campaign at ALTO took place last year.

Concerning the low-energy radioactive beams, a number of complementary installations are being prepared for online commissioning. POLAREX (POLARized EXotic nuclei) is a unique facility to study nuclear magnetic properties of neutron-rich nuclei produced by the ALTO facility with the technique of low-temperature nuclear orientation combined with the on-line implantation of a radioactive beam. It will also enable the studies of nuclear structure ( $\gamma$  multipolarity, nuclear magnetic moment μ, nuclear deformations) [24]. MLLTRAP is a high-precision mass measurement setup [24]. It consists of two trap electrode modules. The ions are cooled and manipulated in the first trap and then injected into the second trap where the actual mass measurement takes place by determining the cyclotron frequency of the ions. This module will be used to measure masses of fission fragments. LINO is a collinear laser spectroscopy setup [24] that will enable studies of hyperfine structure, electromagnetic moments, and mean charge radii. Collinear laser spectroscopy is a robust experimental method that enables high-precision measurement of nuclear characteristics with atomic laser excitations. The LINO setup was successfully commissioned with the stable beams in 2019. The Split-Pole spectrometer is mainly used to study two-body reactions (transfer, elastic and inelastic scattering, charge exchange). It is a very useful tool to study key nuclear astrophysics reactions that require the use of transfer reactions or inelastic scattering reactions (angular distribution measurements and excitation energy spectrum measurement with high energy resolution) to access with high precision the spectroscopic information (Ex, partial widths...) needed to calculate the reactions rates. Recently, the coupling of this spectrometer to an array of DSSSD-s in the reaction chamber has opened up opportunities to have access to the charged-particle decay branching ratios. Examples of angular distributions and correlations include studies of the reactions:  $^{70}$ Zn(d, $^{3}$ He) $^{69}$ Cu,  $^{19}$ Fe( $^{3}$ He, $^{4}$ ) $^{19}$ Ne( $\alpha$ ) $^{15}$ O [24]. Andromede is a 4 MeV NEC pelletron accelerator [30] and delivers ion beams from proton to gold, the rare gases He, Ne, Ar, K and Xe as well as molecular beams and metal clusters. One of the experiments recently performed at Andromede is the study of <sup>12</sup>C+<sup>12</sup>C fusion cross section using the STELLA (STELlar LAboratory) experimental station. The <sup>12</sup>C+<sup>12</sup>C fusion reaction is crucial for the understanding of the evolution of massive stars. A straightforward extrapolation down to the Gamow window and the energy range relevant to carbon burning in massive stars is difficult due to resonances in this reaction at energies near and below the Coulomb barrier. STELLA setup has enabled the studies of <sup>12</sup>C+<sup>12</sup>C fusion cross section at low energies, using an advanced particle-gamma coincidence technique. The method of particle-gamma coincidence is employed for background suppression. In this project, annular silicon strip detectors customized at IPHC-CNRS, Strasbourg [31] were integrated with LaBr3(Ce) detectors from the FATIMA (FAst TIMing Array) [32]. The sensitivity of the technique has effectively removed ambiguities in existing measurements made with gamma-ray or charged-particle detection alone and made it possible to obtain reliable excitation functions for the <sup>12</sup>C+<sup>12</sup>C reaction spanning eight orders of magnitude in the cross-section [33].

# 2.4. Rudjer Bošković Institute Accelerator Facility, Zagreb, Croatia

The Rudjer Bošković Institute Accelerator Facility hosts two accelerators: a 6.0 MV Tandem Van de Graaff accelerator and a 1.0 MeV Tandetron accelerator that provide a wide range of ions, from H and He to heavy ions [34]. One of the beamlines is dedicated to nuclear studies. The research is based on a large silicon detector array. The current experimental setup consists of up to four silicon detector telescopes assembled from thin single-sided ( $\Delta E$ ) SSD and thick double-sided (E) DSSSD detectors. Research topics are concentrated on the molecular and cluster structure of neutron-rich isotopes of Be, B and V and on performing measurements of the three body quasi-free reactions based on the Trojan Horse Method (THM) to accurately investigate nuclear astrophysics reactions [35]. The THM is an indirect method, unaffected by Coulomb suppression or the electron screening effect. It is used for calculating the bare nucleus astrophysical S-factor for charged particle reactions at astrophysically relevant energies. This is obtained by measuring a suitable three-body process' quasi-free crosssection. The suitable Trojan Horse (TH) nuclei must be chosen for the method's successful use; these nuclei should have a prominent cluster structure to transfer (such as nucleons, deuterons, or  $\alpha$ particles). A crucial technique for examining the behavior of nuclear forces in few-body interacting systems is the study of cluster formations in neutron-rich nuclei. Concerning the cluster and the molecular structure in neutron-rich isotopes of Be, B, and V measurement of the  $^7Li + ^7Li \rightarrow \alpha +$  $\alpha + {}^{6}He$  reaction provided the first strong indication for the molecular  $\alpha + 2n + \alpha$  structure in  ${}^{10}Be$ [36]. The measurement of the  ${}^{9}Be + {}^{4}He$  resonant scattering confirmed strong  ${}^{9}Be + {}^{4}He$ clustering in the  $^{13}C$  nucleus [37].

# 2.5. Tandetron laboratory in Piešt'any, Institute of Physics, Slovak Academy of Sciences

Tandetron laboratory in Piešt'any hosts a 2 MV Tandetron accelerator [38]. It is able to deliver proton and deuteron beams up to 4 MeV energy with an intensity of up to 25  $\mu$ A, and  $\alpha$  beams up to the energy of 6 MeV with an intensity of up to 3  $\mu$ A. Four coaxial HPGe detectors and three LaBr<sub>3</sub>(Ce) detectors for charged particles are available for use. Measurements of the angular distribution can be made with great accuracy using a precise goniometer for gamma-ray detector mounting, allowing for the determination of M1/E2 mixing ratios for  $\Delta$ J=1 transitions, which are currently poorly characterized. A gas target has been constructed for the production of quasi-monoenergetic fast neutrons through the (d,D) reaction. It is intended to conduct systematic studies of the lifetimes of excited states in stable isotopes utilizing inelastic neutron scattering and gamma-ray detection [38].

# 2.6. The Oslo Cyclotron Laboratory (OCL), University of Oslo, Norway

Research at the OCL is concentrated on spectroscopy experiments for nuclear structure and nuclear astrophysics. The laboratory hosts the K=35 cyclotron [39]. The cyclotron accelerates  $^1H^+$  in the energy range of 8 MeV to 35 Me,  $^2H^+$  in the energy range of 4 MeV to 18 MeV,  $^3He^{2+}$  in the energy range of 6 MeV to 47 MeV and  $^4He^{2+}$  in the energy range of 8 MeV to 35 MeV.

The tools and techniques for research of statistical properties of highly excited nuclei in the quasi-continuum region have been developed over the years in the OCL group. The Oslo Scintillator Array (OSCAR) [40] consists of 30 large-volume LaBr<sub>3</sub>(Ce) detectors. It measures high-energy  $\gamma$ -rays with excellent timing, high efficiency and good energy resolution. Most nuclear physics experiments

need the detection of high-energy  $\gamma$ -rays in coincidence with the scattered charged particles. The silicon ring (SiRi) detector array is used to detect light-ion ejectiles from transfer reactions [41]. In SiRi, 64  $\Delta$ E-E silicon telescopes are placed in eight trapezoidal pads in the lampshade geometry. For the investigation of statistical decays that require a high particle-γ coincidence rate, the OSCAR-SiRi is used. A Nuclear Instrument for Fission Fragments (NIFF) [42] can be installed within the target chamber to study actinides. NIFF consists of four individual parallel-plate avalanche counters (PPAC). Nuclear levels become increasingly close in energy as excitation energy increases, making discrete spectroscopy challenging or sometimes even impossible. In this statistical region, average properties like the nuclear level density (NLD) and  $\gamma$ -ray strength function ( $\gamma$ SF) replace the welldefined states and the decay rate between them. The Oslo method [43] allows for simultaneous determination of the functional form of the NLD density ( $\varrho$ ) and  $\gamma$ SF function (f). This method is based on a factorization of the decay probability. Thermodynamic properties in the microcanonical ensemble are obtained for 237,238,239U [44] using the Oslo method. The level density curves are exponentially increasing suggesting a linear entropy as a function of excitation energy. The almost linear entropy means that the NLD follows a close-to-constant temperature model characteristic with a critical temperature of Tc~0.4 MeV. This indicates that heating the nuclear system results in the breaking of Cooper pairs rather than an increase in temperature (phenomenological analogy to the melting ice). The study of the  $\gamma$ SF in the quasi-continuum has been quite successful. One of the most unexpected findings was that the γSF starts to increase for decreasing gamma energies below a few MeV for many light nuclei. The scissor resonance embedded in the quasi-continuum has been systematically studied. One has to note that with a large number of data obtained for a wide range of nuclei, OCL is a major contribution to the IAEA (International Atomic Energy Agency) reference database for photon strength function [45].

According to their substantially different time scales, slow (s) and rapid (r) neutron-capture processes [46] have each produced half of the heavy-element isotopes observed in the solar system while the proton-capture/photodisintegration (p) process is in charge of roughly 35 nuclides that are not produced by the s- and r-processes. In addition to the aforementioned nucleosynthesis processes an intermediate (i) neutron-capture process [47] might be important in certain stellar environments. In order to understand the heavy-element nucleosynthesis and elemental abundances, large nuclear reaction networks are applied. These networks need as an input various creation and destruction probabilities. Since both i and r processes involve very neutron-rich nuclei, there is still a lot of missing experimental information on these rates. Although the s-process mostly follows the valley of stability, there is still missing data for s-process branchings (radiative neutron-capture rates). Charged-particle reactions such as (p,  $\gamma$ ) and ( $\alpha$ ,  $\gamma$ ) for the p-process reaction network are very difficult to measure directly for sub-Coulomb energies. One relies on theoretical estimates when there are not any experimental data available. The estimates can diverge by orders of magnitude. One method of establishing experimental constraints on the astrophysical rates is to measure the NLD and  $\gamma$ SF of the residual nucleus in the radiative capture process, either (n,  $\gamma$ ) for the s, i and r-process or  $(p, \gamma)$  and  $(\alpha, \gamma)$  for the p-process. Using these experimental NLDs and  $\gamma$ SFs as input to nuclear reaction rates, significantly improved prediction of the radiative capture rates can be obtained.

Hoyle predicted (in 1954) a resonant state in  $^{12}$ C [48] and explained the production of carbon through the triple  $\alpha$ -process. An excited  $0^+$  state at the predicted excitation energy of 7.65 MeV, named the Hoyle state, was discovered soon afterwards and studied thoroughly since. The experiment at OCL used inelastic proton scattering on  $^{12}$ C to populate the Hoyle state [49]. The measured value is about 34 % higher than the currently adopted value and, if experimentally confirmed, will impact models of stellar evolution and nucleosynthesis. To confirm the discrepancy, new experiments with OSCAR are currently being performed.

2.7. The Superconducting Darmstadt Linear Electron Accelerator (S-DALINAC), Institut für Kernphysik, (IKP), Darmsatdt, Germany

IKP of the Technische Universität Darmstadt hosts a superconducting electron accelerator, the S-DALINAC. The astrophysics research at IKP [50] focuses on electron scattering studies of the form

factor and the monopole transition matrix element of the Hoyle state, the near-threshold transition in  ${}^9\text{Be}$  as well as the extraction of level densities and gamma strength functions. The most notable nuclear structure studies include [50] giant quadrupole resonance, the scissors mode (the scissors mode in nuclei refers to a pictorial image of deformed proton and neutron distributions oscillating against one another) in deformed nuclei and the competitive double gamma  $(\gamma\gamma/\gamma)$  decay.

S-DALINAC was initially constructed as a twice-recirculating accelerator with a maximum energy of 130 MeV in a continuous wave operation. In 2016 it was converted into a three-recirculating accelerator [50]. The beam is either produced in a thermionic gun or in S-DALINAC Polarized Injector (SPIN). Following both sources, the beam is firstly prepared for acceleration and then guided through the superconducting injector accelerator which is able to accelerate the beam up to 10 MeV with beam currents up to 60 µA. The injector beam can be used for the research of nuclear resonance fluorescence (NRF) or it can be deflected into the main accelerator for further acceleration and recirculation. Recirculating the beam up to three times leads to a final energy of up to 130 MeV with a beam current of up to 20 µA. Electron beams from the S-DALINAC are delivered to four major experimental setups. The Darmstadt High-Intensity Photon Setup, (DHIPS) [51] is a low-energy high-flux bremsstrahlung setup located after the linac injector. The experimental hall accommodates two magnetic spectrometers. A highly energy-resolving spectrometer (LINTOTT) [50] is used for elastic electron scattering experiments while Quadrupole CLAMshell (Q-CLAM) is used for inelastic electron scattering experiments. Q-CLAM has a large solid angle coverage and a quadrupole-dipole magnet arrangement. The facility also hosts a high-energy photon tagger NEPTUN. NEPTUN is being upgraded to considerably larger momentum acceptance. By utilizing photo-fission, photo-activation, or photon-scattering reactions, the DHIPS bremsstrahlung beams are mostly employed for nuclear structure studies. Photo-scattering reactions are frequently used in NRF experiments where the resonantly scattered  $\gamma$ -rays are detected with large-volume HPGe detectors. This method is highly selective on nuclear low-spin excitation because of the low momentum transfer induced by the photons. Experiments performed at DHIPS have significantly contributed to the studies of the fine structure of the pygmy dipole resonance in stable nuclei, studies of electromagnetic (EM) transitions between the scissors mode and intrinsic nuclear vibration, and their relation to the modelling of the 0υββ decay reactions. LINTOTT is used for the studies of spectrometry of electron-scattering reactions with low momentum transfer and high energy. It has been applied in various research studies of stable key nuclei from the deuteron, light, and medium-heavy nuclei up to 208Pb (see for example [52], [53], [54]). For the highest energy resolution, LINTOTT can be operated in a dispersionmatching mode. Four ninety-six-fold segmented silicon strip detectors and a Cherenkov counter for background suppression compose the focal plane detectors. The Q-CLAM spectrometer consists of a quadrupole magnet with a large horizontal aperture and a dipole magnet with inclined pole shoes. This allows for large solid angles essential for coincidence measurements of the type (e,e'x) where  $x=\gamma,n,p.$ . (see for example Ref. [55]). To estimate the momentum and angle of the scattered electron in the resulting complex ion optics, it is necessary to track the electron trajectory in the spectrometer. The trajectory is reconstructed using a stack of multiwire drift chambers which provide the respective position and the angle. A large Cherenkov counter is used to suppress the background. Measurements can be performed over a large angular range (25-255°) and in special chicane arrangement at 180° [52]. The latter allows for especially sensitive studies of nuclear magnetic excitation since the longitudinal electric nuclear excitation modes are kinematically suppressed at the backward angles by several orders of magnitude. For coincident γ-ray detection, GALATEA array can be placed around the target chamber. The GALATEA consists of seventeen medium-size LaBr<sub>3</sub>(Ce) detectors and was employed in the observation of the competitive double-gamma decay [56]. The double-gamma ( $\gamma\gamma$ ) decay of an excited quantum system is a fundamental second-order process of quantum electrodynamics (QED). The double-gamma decay can be distinguished from the well-known single-gamma ( $\gamma$ ) decay by the simultaneous emission of two gamma quanta, each having a continuous energy spectrum. The S-DALINAC's electron beam can be utilized to create a beam of energy-tagged bremsstrahlung photons by means of the NEPTUN setup [50]. The NEPTUN tagger setup is suitable for high-resolution studies of astrophysically relevant cross sections.

Reaction-tagged photons can be utilized to more accurately estimate the photodissociation cross-section energy dependence, particularly in the astrophysically significant region just above the threshold for  $(\gamma,n)$  reactions. Additionally, NEPTUN is employed for the extraction of the dipole polarizability as well as the observation of the complete dipole strength from well below the threshold to well above the giant resonance peak.

# 2.8. The INFN Laboratori Nazionali del Sud, (LNS), Italy

LNS hosts two particle accelerators: A Superconducting Cyclotron (SC) with a bending limit  $K_b$ =800 and focusing limit  $K_c$ =200 [57], [58] and a 13 MV HVEC MP TANDEM accelerator. The 13 MV TANDEM provides ion beams from proton to gold with a maximum energy of 25 MeV/u for masses around 200 and 90 MeV/u for fully stripped light ions, such as carbon. SC is a cyclical compact three-sector accelerator capable of accelerating ion beams from protons to uranium at energies up to 80 MeV/A. It is also possible to produce the radioactive ion beams by applying in-flight fragmentation. The FRIBS@LNS (In Fight Radioactive Ion Beams at LNS) facility produces RIBS (Radioactive Ion Beams) at intermediate energies (20-50 MeV/u). Exotic beams from  $^6$ He to  $^6$ Ni have been provided over the past years using the fragmentation of various stable beams accelerated by the SC on a Be target.

A high granularity charged-particle multi-detector called CHIMERA surrounds the target in roughly 95% of the solid angle. It comprises 1192 telescopes (300  $\mu$ m thick Si as  $\Delta E$  detector and a 3-12 cm thick CsI(Tl) crystal as residual energy (E) detector). The implemented pulse height analysis technique allows particle identification and energy measurement in a wide dynamical range including Tandem energies. Chirone is an experiment that uses the high granularity of the CHIMERA [59] detector array as well as the correlator FARCOS (Femtoscope ARray for COrrelation and Spectroscopy) [60] for reaction and spectroscopic studies in identifying the mass and charge of reaction products in order to study the effects of isospin on the reaction mechanism and the density dependence of the symmetry term of the nuclear matter equation of state (EOS). FARCOS is a modular array of telescopes arranged in a single cluster, each one consisting of three detection stages. Measurements of the angular distributions of neutron transfer reactions were obtained with the kinematical coincidence method. A study of cluster structure and levels of 10Be [61] was also performed. In the presence of additional neutrons, clustering effects could take on significantly different features from the self-conjugated nuclei's  $\alpha$ -clustering. These neutrons can take the role of covalent particles acting as a glue between the  $\alpha$  cluster centers and leading to the so-called nuclear molecules. The study of the cluster structure in exotic nuclei and the "pygmy resonance" manifestations of a particular collective nuclear motion carried out using the exotic beams produced by FRIBS have been performed. The isoscalar excitation of the Pygmy Dipole resonance in 68Ni was investigated for the first time [62] (the term "Pygmy Dipole Resonance" (PDR) has been commonly used for the E1 strength around and below the neutron-separation energy, Sn). One of the PDR's main characteristics is the behavior of isoscalar and isovector transition densities which have the same order of magnitude at the surface. The PDR population consequently employs both isoscalar and isovector probes.

Tandem and low energy Cyclotron experiments are the primary aim of MAGNEX [63] experiments. MAGNEX is a high-resolution magnetic spectrometer with a large solid angle and momentum acceptance based on the reconstruction of the ion trajectory. It consists of a vertically focusing quadrupole magnet, a bending dipole, and a number of surface coils for quadrupole and sextupole corrections. In the focal plane, a detector consisting of a gas  $\Delta E$  stage and the wall of silicon strips as the residual energy detectors (E) are positioned. Some methods for populating light neutron-rich nuclei are one and multi-neutron transfer or charge exchange reactions. Due to the spectrometer's high performance, their excitation spectrum can be determined in detail. Since the beginning of the Tandem operation, the focus of research has been the study of reaction mechanisms typical of low-energy processes, such as fission, transfer, complete and incomplete fusion. The study of cluster and molecular configuration at high spin states has also been performed.

The nuclear astrophysics program has the primary emphasis on the application of THM [35] to the processes relevant to the understanding of stellar evolution and the element nucleosynthesis. The THM can provide an indirect measurement of the astrophysical factors which in direct experiments can be inferred only by extrapolation. The method can be applied to resonant and non-resonant charged particle reactions as well as to neutron-induced reactions. Both stable and radioactive beams can be used. The method has been developed at LNS and some of the experiments at the local INFN-LNS infrastructure will be described (many that are not stated here are performed worldwide). The  ${}^{9}\text{Be}(p,\alpha){}^{6}\text{Li}$  reaction is important for the understanding of light element abundances (lithium, beryllium and boron). The measurement of the  ${}^{9}\text{Be}(p,\alpha){}^{6}\text{Li S}(E)$  factor by properly selecting the Quasi-Free (QF) contribution of the three-body reaction  ${}^{2}$ H( ${}^{9}$ Be, $\alpha$   ${}^{6}$ Li)n has been performed [64], [65]. Deuterons were used as "TH nuclei" because of the obvious p-n structure and their relative motion mainly occurring in the s-wave. The astrophysical S(E)-factor was extracted and compared to direct data. No information about the electron screening effect could be extracted due to the poor resolution of the indirect data. The  ${}^{10}$ B(p, $\alpha_0$ )<sup>7</sup>Be reaction has been measured for the first time at the Gamow peak [66] by means of the THM applied to the  ${}^{2}H({}^{10}B,\alpha {}^{7}Be)n$  QF reaction. The study of the  ${}^{10}B(p,\alpha){}^{7}Be$ reaction is of importance in nuclear astrophysics because of the challenging measurements of the corresponding S(E)-factor at low energies (Gamow window) and the astrophysical importance of the production/destroying processes of the unstable 7Be isotope [67]. The absolute value of the astrophysical S(E)-factor has been deduced, providing for the first time the measurements at the corresponding Gamow peak. Since the THM S(E) factor does not suffer from electron screening effects, it has been used to evaluate the electron screening potential value needed for the description of the low-energy direct data. This was the first independent measurement of the electron screening potential, U<sub>e</sub>, for the  ${}^{10}\text{B}(p,\alpha_0){}^{7}\text{Be}$ . The  ${}^{11}\text{B}(p,\alpha_0){}^{8}\text{Be}$  reaction was studied from 1 MeV down to astrophysical energies by applying THM to the  ${}^{2}H({}^{11}B,\alpha_{0}){}^{8}Be$  three-body reaction [68].  ${}^{27}Al(p,\alpha){}^{24}Mg$ reaction rate was measured by means of the THM applied to the  ${}^{2}H({}^{27}Al,\alpha{}^{24}Mg)n$  three-body reaction [69]. The high-accuracy measurement of the strength of the 84.3 keV resonance was obtained. The stricter constraints on the upper limits of the 71.5 keV, 193.5 keV, and 214.7 keV resonances lead to a reduction of a factor of approximately three in the reaction rate at temperatures where the MgAl cycle is especially important. Due to its high yield from 150 keV down to zero energy, the neutron capture reaction  ${}^{10}\text{B}(\text{n},\alpha)^7\text{Li}$  is becoming more and more significant for applied nuclear physics. The  $^{10}$ B $(n,\alpha)^7$ Li reaction has been investigated [70] via the Trojan Horse Method (through  $^2$ H $(^{10}$ B $,\alpha^7$ Li)H reaction) from 0 to 1 MeV. The  $\alpha_0$  and  $\alpha_1$  channels, corresponding to <sup>7</sup>Li in its ground state (g.s.) and first excited level respectively, have been analyzed and cross-sections have been measured for the two reaction channels. Understanding the fluorine abundance in the outer layers of asymptotic giant branch stars depends significantly on  ${}^{19}F(p,\alpha){}^{16}O$  reaction (AGB stars are stars that are brighter than the Sun but have lower surface temperatures). Recently, direct measurements of the  $^{19}$ F(p, $\alpha$ 0) $^{16}$ O reaction [71] at center-of-mass energies from 0.4 MeV to 0.9 MeV were performed. The LHASA (Large High-resolution Array of Silicon for Astrophysics) detector array has been used. It consists of six 300  $\mu$ m thick YY1 silicon strip detectors [72] mounted in a lamp-shade configuration. The emerging  $\alpha$ particles can be detected by LHASA over a wide angular range (from 10° to 32.

The current upgrade project of the SC is opening new perspectives and possibilities. The upgrade will enable further investigation of isospin physics in the range of Fermi energies by means of reactions with high isospin asymmetric projectiles. Additionally, astrophysically relevant reactions utilizing low-energy degraded beams will be performed. A new fragment separator called FRAISE [73] is also being built in order to fully benefit from the intensity upgrade.

# 2.9. The Atomki Accelerator Center

The Atomki Accelerator Center [74] hosts a K=20 cyclotron, two Van de Graaff accelerators with a maximum terminal voltage of 1 MV and 5 MV, respectively, and a 2 MV Tandetron accelerator installed in 2015. The MGC-20E Cyclotron is a compact isochronous cyclotron capable of accelerating the four lightest ions (p, d,  ${}^{3}$ He ${}^{2}$ + and  ${}^{4}$ He ${}^{2}$ +). At the beginning of 2020, the decision was made to stop

5 MV Van de Graaff accelerator. The Tandetron and the cyclotron have been widely used for astrophysical research.

Elements heavier than iron are mostly produced via neutron-capture (s- or r-process) reactions. However, the "p nuclei" on the proton-rich side of the valley of stability cannot be created by these processes. The  $\gamma$ -process (see section 2.1. for more details about the  $\gamma$ -process) is mainly responsible for the production of these isotopes. The synthesis of heavy, proton-rich isotopes in the  $\gamma$ -process proceeds through photodisintegration reactions. The  $(\gamma,\alpha)$  and  $(\alpha,\gamma)$  reaction cross-section calculations are highly sensitive of the selection of the  $\alpha$ -nucleus potential, which comprises a Coulomb and a nuclear part (the latter one consists of a real and an imaginary part). The knowledge of  $\alpha$ -nucleus model potential (OMP) is very important in the studies of nuclear structure, reactions, and nuclear astrophysics. Moreover, in several astrophysical applications, such as modeling the nucleosynthesis in explosive scenarios like the  $\gamma$  process, the reaction rates are taken from the Hauser-Feshbach (H-F) statistical model (see Sections 2.1, 2.6.) using global OMPs. In recent years, significant efforts have been devoted to improving the  $\alpha$ -nucleus optical potential parameterizations for astrophysical applications. The predicted cross-sections using different global  $\alpha$ -nucleus optical potential parametrizations can vary by an order of magnitude. The parameters of the global  $\alpha$ nucleus optical potentials are usually obtained from the analysis of the angular distributions of elastically scattered  $\alpha$  particles (and are adjusted to  $\alpha$ -induced cross sections if experimental data exist). The elastic scattering cross sections for the reactions  $^{110,116}$ Cd  $(\alpha,\alpha)^{110,116}$ Cd at energies above and below the Coulomb barrier were measured and provided a sensitive test for the  $\alpha$ -nucleus optical potential parameter sets [75]. The rates of the reaction network have to be known for a proper understanding of the  $\gamma$ -process. It was found, in previous rate variation studies, that the reaction  $^{128}$ Ba $(\gamma,\alpha)^{124}$ Xe influences the abundance of the p nucleus  $^{124}$ Xe [76]. Since the stellar rate for this reaction cannot be determined by direct measurement, the cross-section of the inverse  $^{124}$ Xe( $\alpha$ , $\gamma$ ) $^{128}$ Ba reaction [76] was measured. Simultaneous studies of the  $^{124}$ Xe( $\alpha$ ,n) $^{127}$ Ba reaction channel at higher energy have allowed to further identify the source of a discrepancy between data and prediction. The  $\alpha$ -beam for the irradiations was provided by the MGC-20 cyclotron of ATOMKI. An upper limit for the  $^{128}$ Ba $(\gamma,\alpha)^{124}$ Xe stellar rate was inferred from the measurements. For the first time, measurements of the elastic scattering cross sections of the  $^{113}$ In( $\alpha,\alpha$ ) $^{113}$ In [77] reaction have been performed at energies close to the astrophysically significant energy region. The high-precision experimental data were used to evaluate the predictions of the recent global and regional  $\alpha$ +nucleus optical potentials. Parameters for the local  $\alpha$ -nucleus optical potential were derived from the measured angular distributions. In addition,  $\alpha$ -induced reaction cross-section measurements on gold were measured  $(^{197}\text{Au}(\alpha,n)^{200}\text{Tl} \text{ and } ^{197}\text{Au}(\alpha,\gamma)^{201}\text{Tl}$  [78]) which could be performed almost at the astrophysically relevant energy region. The new experimental data were then facilitated to test statistical model predictions and to constrain the  $\alpha$ -nucleus optical model potential. Recent nucleosynthesis studies have demonstrated that  $(\alpha, xn)$  reactions play a particularly significant role in the production of light  $(30 \le Z \le 45)$  neutron-rich isotopes. The cross section of the  $^{100}$ Mo( $\alpha$ , n) $^{103}$ Ru [79] was measured by means of the activation method (see section 2.1 for the details of the method). Measurements of cross sections have been performed at a number of energies below the Coulomb barrier. Large discrepancies between the experimental values and statistical model prediction employing  $\alpha$ -OMP were found. However, the discrepancies could be excellently described by the Atomki-v2 potential. <sup>144</sup>Sm $(\alpha,n)$ <sup>147</sup>Gd reaction was studied [80].  $\alpha$ -beams were provided by the cyclotron accelerator of Atomki. The cross section was determined using the activation method. The  $\gamma$ -rays produced in the decay of the <sup>147</sup>Gd were measured with an HPGe. The cross section was measured from close above the  $(\alpha, n)$  threshold. The comparison of measurements with statistical model calculations using various approaches and parametrizations for the  $\alpha$ -nucleus optical potential were made and excellent agreement was obtained for two recent potentials. Nuclear reactions of astrophysical relevance are studied using the activation method.  ${}^{3}\text{He}(\alpha,\gamma){}^{7}\text{Be}$  reaction plays a crucial role in solar hydrogen burning. Its reaction cross section has been measured [81], which contributed to a better understanding of the investigated stellar processes. The  $\alpha$ -beam was provided by the MGC-20 cyclotron of ATOMKI. The p-process is one of the least known processes of nucleosynthesis, the

model calculations are not able to reproduce well the abundances of p-isotopes observed in nature.  $^{92}$ Nb is one of the isotopes present in the early Solar system which may be used to test the models since among the suggested sites is a thermonuclear supernova explosion. The cross section of the  $^{91}$ Zr(p, $\gamma$ ) $^{92m}$ Nb has been measured using the activation method (only the partial cross section leading to the isomeric state of  $^{92}$ Nb could be measured as the ground state has a half-life of  $3.47\cdot10^7$  years) [82]. The theoretical cross sections employed in previous astrophysical models and the  $^{91}$ Zr(p, $\gamma$ ) $^{92m}$ Nb measurements showed good agreement. One of the fundamental mechanisms for hydrogen burning in stars is the CNO cycle. The cycle is a sequence of reactions converting H into He on preexisting CNO nuclides and it represents the main hydrogen burning mode in stars with M  $\geq$  1.3 Mo. The first reaction of the cycle is the radiative proton capture on  $^{12}$ C and the rate of this  $^{12}$ C(p, $\gamma$ )  $^{13}$ N reaction is related to the  $^{12}$ C/ $^{13}$ C ratio observed e.g. in the Solar System. The low-energy part of the cross section was measured already measured a few times, however, the experimental data are scarce in a wide energy range, especially around the resonance energy (1.7 MeV). Therefore, the  $^{12}$ C(p, $\gamma$ )  $^{13}$ N cross section was measured in the energy range between 300 keV and 1900 keV using the activation method [83].

One of the most significant nuclear astrophysics reactions,  $^{14}N(p,\gamma)^{15}O$  affects the generation of energy in stars, stellar evolution, and nucleosynthesis. Consequently, the low-energy part of the cross section must be accurately known in order to calculate reaction rates reliably. The cross section was measured in the center-of-mass (cms) range between 550 keV and 1400 keV with the means of activation method [84]. The annihilation radiation (e-e+) following the  $\beta^+$  decay of the  $^{15}O$  was detected. This approach, which directly yields the astrophysically important total cross section, was never used before for the cross-section measurement in the investigated energy range.

Recently, several experimental anomalies were raised as possible indicators for a new light particle. According to some predictions, the experimental findings could be explained by light neutral bosons with masses between 10 MeV and 10 GeV as dark matter anomalies that couple to electrons and positrons (see Ref. [85] and references therein). The possible anomalies were studied in the  $^7\text{Li}(p,\gamma)^8\text{Be}$  reaction. The reaction was used to populate 17.6 and 18.5 MeV 1<sup>+</sup> states in  $^8\text{Be}$  using the incident proton beams from the Atomki 5 MV Van de Graaff accelerator [85]. The eterpairs were detected by five plastic  $\Delta$ E-E detector telescopes. The positions of the hits were determined by multiwire proportional counters (MWPC) placed in front of  $\Delta E$  and E detectors. Electron-positron angular correlations were measured for the isovector magnetic dipole transition from 17.6 MeV ( $J^{\pi=1+}$ , T=1) state to the ground state ( $J^{\pi=0^+}$ , T=0), and the isoscalar magnetic dipole 18.15 MeV ( $J^{\pi=1^+}$ , T=1) state to the ground state transition in 8Be. Significant enhancement relative to the internal pair creation was observed at large angles in the angular correlation for the isoscalar transition with a confidence level of  $2.5 \sigma$ . The observation could be either due to a nuclear reaction interference effect or might indicate that as an intermediate step a neutral isoscalar particle, hypothetical X17 boson, with a mass of  $16.7\pm0.35$ (stat)  $\pm0.5$ (syst) MeV/c<sup>2</sup> and J<sup> $\pi$ =1+</sup> was created. The experiment has been repeated at the ATOMKI 2 MV Tandetron accelerator using the thinner carbon backing and an increased number of telescopes (six instead of five) and DSSSDs instead of MWPCs. The anomalous angular correlation was reproduced using the new independent setup [86]. Further pieces of evidence for the X17 particle have been found in the following ATOMKI group experiments, studying two additional reactions. The angular correlation spectra of e<sup>+</sup>e<sup>-</sup> pairs produced in the <sup>3</sup>H(p,e<sup>+</sup>e<sup>-</sup>)<sup>4</sup>He nuclear reaction have been studied at the proton energies of 510, 610, and 900 keV, respectively [87]. It is possible to comprehend the key characteristics of the spectra by considering the internal and external pair creations that followed the direct proton capture by 3H. The observed peak in the angular correlation spectra at about 115° cannot be explained by these processes. This anomalous excess of e<sup>+</sup>e<sup>-</sup> pairs can be described by the creation and subsequent decay of a light particle during the direct capture process. The derived mass of particle is  $mxc^2 = 16.94 \pm 0.12$ (stat)  $\pm 0.21$ (syst) MeV. According to this mass, this is likely the same X17 particle. Moreover, the new anomaly observed in <sup>12</sup>C [88] supported the existence and the vector character of the hypothetical X17 boson. The angular correlation of e<sup>+</sup>e<sup>-</sup> pairs was investigated in the ¹¹B(p,γ)¹²C reaction and angular range of 40°≤θ≤175° for five different proton energies in the energy range of 1.50 MeV and 2.5 MeV. The measurements

were performed at the 2 MV Tandetron accelerator of ATOMKI. The Ex = 17.23 MeV,  $J^{\pi}$  =  $1^{-}$  state in  $^{12}$ C was populated by the  $^{11}$ B(p, $\gamma$ )  $^{12}$ C nuclear reaction [88]. At small angles ( $\theta \le 120^{\circ}$ ), the results can be well interpreted by the internal pair creation process of electromagnetic radiations with E1 and M1 multipolarities and by the external pair creation in the target backing. At angles greater than 120°, additional count excesses and anomalies were observed, which could be explained for by the existence of the hypothetical X17 particle. The results imply that the X17 particle was generated mainly in E1 radiation. The derived mass of the particle is  $m_x c^2 = 17.03 \pm 0.11 (\text{stat}) \pm 0.20 (\text{syst})$  MeV. The mass and derived branching ratio (Bx=3.6(3)×10<sup>-6</sup>), indicate that this is most probably the same X17 particle that was previously proposed to explain the anomaly seen in the decay of  $^8$ Be. Many experiments in the coming years are going the investigate the possibility of a new gauge boson. This will probably finally determine the existence of such a particle and constrain its properties. One of the European experiments that aims to provide a quick and reliable test of the existence of a light dark boson that couples the standard model to the dark sector is the New Judicious experiment for Dark Sector Investigations (New Jedi) project [89]. Some of the JEDI project measurements were performed at ANDROMEDE, IJC lab.

# 2.10. Centre of Accelerators and Nuclear Analytical Methods Research Infrastructure (CANAM RI), Czech Republic

CANAM RI is a Czech infrastructure for the investigation of tasks in a variety of scientific disciplines using beams of accelerated ions and neutrons. It also hosts a 4130 MC Tandetron, TR-24, and U-120M cyclotrons, and an MT25 microtron [90]. The medium-current (MC) version of the 3 MV Tandetron can accelerate ion beams in the energy range between 600 keV and 30 MeV. TR-24 cyclotron accelerates beams with energy ranging from 18 MeV to 24 MeV. Ion beams (p, H<sup>-</sup>, D<sup>+</sup>, D<sup>-</sup>,  $^{3}$ He<sup>2+</sup>,  $^{4}$ He<sup>2+</sup>) are accelerated by the isochronous cyclotron U-120M. Microtron MT25 is a cyclic electron accelerator.

The unique <sup>3</sup>He beams and deuteron beams of the U-120M cyclotron allow for research of astrophysically relevant nuclear reactions. Since it is very difficult and sometimes even not possible to measure cross sections directly at very low energies, indirect approaches offer an independent means of obtaining the relevant information at these (astrophysically significant) energies. Two indirect methods, namely, THM (see Sections 2.4. and 2.8. for more details) and ANC (the asymptotic normalization coefficients) are used. <sup>2</sup>H(d,d) and <sup>2</sup>H(d,t) reactions were studied at U-120M cyclotron by means of the THM method [91]. Data unaffected by the electron screening were obtained in the energy region from thermal energies to energies relevant to BBN. The ANC technique uses the direct transfer reactions' peripheral nature to deduce the direct component of radiative capture at zero energy. Depletion of  $^{18}$ O by (p,  $\gamma$ ) capture from the CNO cycle in AGB stars was studied [92]. The results allowed to decide between two previous contradicting findings. The  ${}^{26}$ Si(p, $\gamma$ ) capture was studied using a <sup>26</sup>Mg(d,p) reaction. Measurements were performed at the U-120M cyclotron [93]. The obtained data enabled the reaction rate to be updated, leading to new findings regarding the production of <sup>26</sup>Al in the galaxy. Precise angular distribution measurements and high-quality beams of <sup>3</sup>He and deuterium combined with a careful analysis are the foundations of the successful application of these methods.

# 2.11. Bronowice Cyclotron Center (CCB) of the Institute of Nuclear Physics PAN, Krakow, Poland.

The Bronowice Cyclotron center hosts a cyclotron Proteus C-235 [87]. Proteus C-235 is an isochronous cyclotron with a proton energy selector. It is able to deliver proton beams with energies between 70 MeV and 230 MeV and beam intensities ranging between 0.1 and 600 nA [94]. The available instrumentation include the Big Instrument for Nuclear Reaction (BINA) [95], the High Energy Gamma Ray Detector (HECTOR) [96], and the Krakow Triple Telescope Array (KRATTA) [97]. The goal of the BINA detection system is to examine breakup and elastic reactions at intermediate energies. It makes it possible to record coincidences between two charged particles over roughly a  $4\pi$  solid angle, allowing for the investigation of practically the whole phase space of elastic and breakup reactions. The detector is composed of two main parts, the forward Wall and the

backward Ball. The backward scintillator Ball is utilized to register particles scattered at higher polar angles, and the front scintillator Ball is used to precisely reconstruct charged particle momenta. Two scintillator hodoscopes for measuring energy and energy loss are present in the forward part, along with a MWPC for reconstructing particle paths. The backward part is composed of a phoswich scintillator components that serve as both the detector and at the same time form the vacuum chamber for a cryogenic deuterium target. High-efficiency array, HECTOR, measures high-energy gamma rays with energies between 2 MeV and 40 MeV. The array consists of eight large-volume BaF2 scintillators (145 mm diameter and 175 mm length). The energy resolution is ~12 % for low energy (60Co) and ~10 % for high energy (15 MeV) gamma rays. Excellent time resolution (< 1 ns) makes it possible to distinguish between gamma-rays and neutron-induced events using the time of flight method. HECTOR array can operate in coincidence with the KRATTA system to detect scattered protons. KRATTA measures the emission angle, energy and isotopic composition of light charged reaction products. It comprises thirty-eight individual modules that can be put together in any configuration. A single module consists of three identical, 500 µm thick, large area photodiodes, used also for direct detection, and of two CsI crystals. All signals are digitally processed. The middle photodiode's complex signals can be decomposed into their ionization and scintillation components via pulse shape analysis. This enables a single readout channel to obtain a satisfactory isotopic

Nuclear structure studies at Bronowice center will encompass studies of the dynamics of fewnucleon systems and the physics of nuclear clusters in order to get insight into the nucleon-nucleon interaction, measurements of collective, high-energy excitations in nuclei (e.g. giant nuclear resonances) in the yet unexplored regions of excitation energy and spin, high-resolution gamma-ray spectroscopy of nuclei produced in the process of proton-induced fission and spallation [98]. The facility will also be used to test various components of the cutting-edge detection systems being built for European ESFRI nuclear physics facilities (e.g. SPIRAL2, GSI/FAIR).

# 2.12. The Centro Nacional de Aceleradores (CNA), University of Seville, Spain

resolution. The upper energy limit for protons is about 260 MeV.

CNA hosts three particle accelerators, namely, a 3 MV van de Graaff Tandem accelerator, a cyclotron and a 1 MV Tandetron which serves as a mass spectrometer [99]. Cyclotron is able to provide 18 MeV protons and 9 MeV deuterons. Seven of the cyclotron's eight targets are used to create positron emitters, while an external beamline has been put in the eighth port to conduct nuclear physics research. The nuclear physics beamline at the 3 MV Tandem hosts a high-volume vacuum chamber, where nuclear instrumentation (detectors, electronics, DAQ), that will be used in international Nuclear Physics facilities, can be developed and tested. One of the detectors is a mini SeD: a gas ionization multiwire chamber that operates at low pressure and has proven to be very suitable for the detection of low-energy ion beams with large angular and energy straggling, thus enabling precise reconstruction of the TOF and position of beam particles at high production rates [100]. The first accelerator-based neutron source in Spain is called the HiSPANoS, short for Hispalis Neutron Source. [101]. The reactions  $p(^7\text{Li}, n)$ ,  $d(^7\text{Li}, n)$ ,  $d(^7\text{Li}, n)$ ,  $p(^9\text{Be}, n)$  at HiSPANoS create neutrons in a broad-energy range, from thermal to fast neutrons with energies up to 9 MeV.

# 2.13. The Laboratory of Accelerators and Radiation Technologies (LATR) of Instituto Superior Técnico (IST), Portugal

LATR is a Portuguese infrastructure that hosts three accelerators, namely a 2.5 MV Van de Graaff accelerator, a 3 MV Tandem [102] as well as a 210 kV ion implanter.

Elastic scattering reactions  $^{12}\text{C} + ^{16}\text{O}$  and  $^{16}\text{O} + ^{16}\text{O}$  were measured at the nuclear physics beamline of the 3MV Tandem [102]. The AMS line of the Tandem was facilitated to measure the cross section of the reaction  $^{35}\text{Cl}(n,\gamma)^{36}\text{Cl}$ .  $^{36}\text{Cl}$  has a half-live of  $3.01\times 10^5$  years. That implies that the amount of  $^{36}\text{Cl}$  existing in the earth's soil may only be explained by either continuous production in the sun wind or bombardment of the earth's surface by cosmic neutrons. The targets were obtained by the bombardment of  $^{35}\text{Cl}$  by thermal neutrons. The content of  $^{36}\text{Cl}$  was measured by AMS, leading to a cross section value that, within the uncertainties, is in agreement with the average of previously

published results [102]. Measurement of elastic scattering of protons, namely the reactions  $^6\text{Li}(p,p')^6\text{Li}$ ,  $^7\text{Li}(p,p')^7\text{Li}$ ,  $^{12}\text{C}(p,p')^{12}\text{C}$ ,  $^{19}\text{F}(p,p')^{19}\text{F}$ ,  $^{31}\text{P}(p,p')^{31}\text{P}$  at both forward and backward angles were performed at LATR facility and gave reliable information both for the purpose of acquiring precise optical model parameters and for the purpose of characterizing the excited states of nuclei.

2.14. The Horia Hulubei National Institute for R&D in Physics and Nuclear Engineering (IFIN-HH) Tandem accelerator complex, Romania

The INFN-HH tandem accelerator complex in Romania consists of 9 MV FN Pelletron tandem, 3 MV HVEE Tandetron accelerator, and 1 MV HVEE Tandetron accelerator dedicated to AMS experiments [103], [104].

The 9 MV FN Pelletron Tandem can provide beams from hydrogen to gold and can deliver currents up to  $\mu A$ . The 9 MV Tandem hosts different setups.

The ROmanian array for SPectroscopy in HEavy ion Reactions (ROSPHERE) [98] is a multidetection array. The "mixed" configuration, which combines fast LaBr<sub>3</sub>(Ce) scintillators and largevolume HPGe detectors that have been Compton-suppressed, is the most common one used in the ROSPHERE array, which can hold up to 25 detectors. For cases requiring a high-resolution and highefficiency detection system, ROSPHERE can alternatively be used as a full HPGe array. A variety of subjects have been addressed in the physics cases, including the role of the negative-parity neutron intruder orbitals in the structure of low-spin states configurations for nuclei near the "island of inversion" [105], [106], the study of the nuclear structure near shell and/or subshell closures (see for example [107], [108], [109]), the study of neutron-rich nuclei using low energy transfer/incomplete fusion reactions induced by <sup>7</sup>Li beams ([110], [111]) as well as the interplay between collective and single-particle degrees of freedom (see Refs. [107], [108], [109], [110], [111]). The main tool to investigate collective phenomena is that of lifetime measurements. Lifetime measurement observables are closely related to the determination of the reduced matrix elements of nuclear transitions. These quantities are in turn very sensitive to details of the nuclear structure and their knowledge is necessary for the testing of different theoretical models. A broad range of experimental techniques to cover the relevant interval for lifetimes of excited states in nuclear systems was employed: the fast timing technique, the RDDS method, and the DSAM method (see Refs. [105], [106], [107], [108], [109], [110], [111] and references therein). Far from the valley of stability, deviations are observed from the conventional single-particle shell structure. Neutron-rich nuclei with  $Z \sim 10$  and N ~ 20 have unexpectedly high binding energies due to the onset of deformation [112]. The deformation occurs from the filling of the f<sub>7/2</sub> intruder orbital. In stable nuclei f<sub>7/2</sub> intruder orbital lies above the N = 20 shell closure ([113], [114]). The size of the N = 20 shell gap is reduced for neutronrich  $Z \sim 10$  nuclei, allowing excitations from the  $d_{3/2}$  to the  $f_{7/2}$  orbital to become favored, leading to the region of anomalous shell structure known as the "island of inversion." The island of inversion is known to extend from neutron numbers N = 20-22 for the Ne, Na, and Mg isotopes (Z = 10-12) [115]. The half-life of the  $I^{\pi}$  = 4-, E=2305 keV intruder state in <sup>34</sup>P has been measured as  $t_{1/2}$  = 2.0 (1) ns. The result was obtained using γ-ray coincidence and fast timing techniques with the mixed LaBr<sub>3</sub>-HPGe ROSPHERE detector array. The <sup>18</sup>O(<sup>18</sup>O,pn)<sup>34</sup>P fusion-evaporation reaction at a beam energy of 36 MeV [105] was used to populate states in <sup>34</sup>P. For small values of the mixing ratio, the B(M2) value was found to be consistent with similar transitions associated with the occupation of neutron f7/2 configurations in this mass region. In the medium-mass region of nuclei approaching the N = 40subshell closure with  $Z \sim 28$ , it is assumed from empirical observations and theoretical calculations that states with different structures such as single-particle, intruder, and collective states coexist at low and medium excitation energies ([116], [117]). The half-lives of the 9/2+, 13/2+, and 15/2+ yrast states in the neutron-rich 67Cu nucleus were determined by using the in-beam fast-timing technique [101]. The  $^{67}$ Cu nuclei were produced in an  $\alpha$ -induced reaction on a  $^{64}$ Ni target at an incident energy of 18 MeV. 9 MV tandem accelerator produced  $\alpha$ -beams. Gamma rays were detected using a setup that consisted of five HPGe detectors, four planar HPGe detectors, and eight LaBr<sub>3</sub>(Ce) scintillation detectors. The experimentally obtained E3 transition strength for the decay of the 9/2+ level to the 3/2ground state suggests that the wave function of this level might contain a collective component arising from the coupling of the odd proton p<sub>3/2</sub> with the 3- state in <sup>66</sup>Ni. The measurement of electromagnetic transition rates in nuclei in transitional regions between shell closures and permanently axially deformed nuclei can shed light on the validity of corresponding models of nuclear structure. Excited states in <sup>136</sup>Ce were populated via the <sup>124</sup>Sn(<sup>16</sup>O,4n) fusion-evaporation reaction. The gamma rays were detected using the mixed configuration of ROSPHERE array, which consisted of fourteen HPGe detectors and eleven LaBr<sub>3</sub>(Ce) scintillator detectors. Each of the HPGe detectors was surrounded by active Compton suppression shields. The half-lives of the  $I^{\pi} = 5^-$  and  $I^{\pi}$ =  $7^{-}$  yrast states with E<sub>x</sub> = 1978 keV and E<sub>x</sub> = 2307 keV in the N=78 isotone <sup>136</sup>Ce, have been measured to be 496(23) ps and 270(24) ps, respectively, using the coincident fast-timing spectroscopy technique [109] The fusion-evaporation reaction <sup>124</sup>Sn(<sup>7</sup>Li,2n) was used to populate excited states in <sup>129</sup>I [110]. The array of eight HPGe detectors and five LaBr<sub>3</sub>(Ce) scintillation detectors were used to measure inbeam γ-ray coincidences. A positive parity band structure built on the 7/2<sup>+</sup> ground state was established and the  $\pi g_{7/2}$  configuration at oblate deformation was assigned to it based on the  $\gamma\gamma$ coincidence data. The ground state collectivity of a nucleus can be accurately determined by measuring the reduced transition probability  $B(E_2; 2_1^+ \rightarrow 0_1^+)$ , and half-life measurements are anticipated to shed light on the structure of nuclei in this transitional region. The ROSPHERE array has been also used to measure the half-life of the yrast  $I^{\pi} = 2^{+}$  state in the neutron-rich nucleus <sup>188</sup>W. Fast-timing technique has been used. The resulting value of t<sub>1/2</sub> = 0.87(12) ns is equivalent to a reduced transition probability of  $B(E_2; 2_1^+ \to 0_1^+ = 85(12) W.u.)$  for this transition. Even with a rather significant uncertainty, it seems to indicate a more abrupt decrease in collectivity compared to the trend of lighter tungsten isotopes. According to the estimates for this mass region, this predicted a possibly higher softness for <sup>188</sup>W compared to stable tungsten isotopes.

One of the most important findings in the study of nuclear structure at IFIN-HH was the discovery of a shape isomer in <sup>66</sup>Ni utilizing gamma spectroscopy and heavy-ion transfer reactions at energies below the Coulomb barrier [118]. It was the discovery of the lightest-ever atomic nucleus that exhibits a photon decay hindered solely by a nuclear shape change. It is an extremely rare process involving a transition between totally distinct microscopic configurations, coexisting at similar excitation energy. This result was obtained from lifetime measurements of the first three excited 0<sup>+</sup> states pointing to the oblate, spherical and prolate nature of excitations.

The tape station for beta-decay experiments is also available at INFN-HH, using three clover detectors with 120 % relative efficiency and anti-Compton shields as well as fast LaBr<sub>3</sub>(Ce). The setup consists of multi-strip silicon detectors for particle detection that may move radially and longitudinally around the target. The setup is used for nuclear reaction and astrophysics studies.

Nuclear astrophysics research takes advantage of ion beams from a 3 MV Tandetron accelerator as well as of the ultra-low background laboratory in a salt mine at Slanic-Prahova [119]. The ultralow background radiation laboratory was built and became operational in 2006 in the former Unirea (Slanic-Prahova) salt mine at 208 m below the surface (estimated to a 560 m water equivalent (m.w.e)). In comparison to the identical spectrum recorded at the surface, in the open field, it was discovered that the overall gamma background spectrum between 40 keV and 3 MeV was 100 times smaller at the laboratory level. Given that proton and alpha-capture reactions serve as an important proxy in nucleosynthesis, two of the reactions that were studied at IFIN-HH are  $\alpha + {}^{64}$ Zn and  $\alpha + {}^{58}$ Ni [120], [121]. Both reactions are important for explosive nucleosynthesis, and in both cases, the measurements done at the IFIN-HH were able to investigate the Gamow window. However, up to now the most significant result was that obtained for the  $^{12}C$  +  $^{12}C$  fusion reaction [122], [123].  $^{12}C$  +  $^{12}C$ fusion reaction is a cornerstone case in nuclear astrophysics. There are few experimental data at energies below the Coulomb barrier, and theoretical model predictions on the fusion mechanisms often significantly disagree with the experimental results (see the references in Ref. [120]). The adjacent reaction, <sup>13</sup>C + <sup>12</sup>C was studied to obtain information on the fusion interaction at such low energies. The extra neutron in <sup>13</sup>C allowed for the production of the unstable <sup>24</sup>Na (T<sub>1/2</sub>=15.0 h) and the use of the activation method (see section 2.1.). The proton evaporation channel leads to an activity with a half-life of 15 h, suitable for samples' transfer to the salt mine. The study was able to measure the cross-section of the proton evaporation channel down to  $E_{cm}$  =2.3 MeV.

2.15. Laboratory for Underground Nuclear Astrophysics (LUNA), Gran Sasso National Laboratories (GSNL, Italy

Due to the suppression of the natural background, studying nuclear reactions that emit gamma rays is especially advantageous underground. At  $\gamma$ -ray energy exceeding 3 MeV, the 1.4 km of rock above the LUNA facility suppresses the cosmic-induced background by a factor of six. The  $\gamma$ -rays emitted in the radioactive decay of naturally occurring isotopes dominate the  $\gamma$ -ray background at lower energies. Typically, this background component is reduced by surrounding the detector with high-purity, high-Z passive shielding. The interaction of cosmic rays with such shielding is constrained by the fact that the interaction of the cosmic rays with the shielding itself produces radioactive isotopes and secondary radiation. Additionally, this issue is mitigated underground where there is considerably thicker shielding. The LUNA 400 kV accelerator provides 1 mA H<sup>+</sup> beams and 500 $\mu$ A He+ beams. It has two beamlines, one of which houses a solid target station and the other of which houses a windowless gas target system [124].

LUNA has hosted various detector setups (See Ref. [124] and references therein). Large-volume HPGe detectors are used for high-resolution spectroscopy at LUNA. To maintain the benefits of being underground, all LUNA HPGe detectors are constructed from materials with a low inherent background. Thick passive shielding constructed of lead and copper is used to suppress the environmental background caused by naturally occurring radioactive isotopes in cases where  $\gamma$ -rays with energy lower than 3 MeV need to be detected. A large-volume BGO detector covering nearly the entire solid angle surrounding the target is employed when extremely weak cross sections are measured and the sensitivity is crucial. BGO detector has also been successfully used to determine  $\gamma$ ray branching ratios exploiting γγ coincidences in segment pairs. Additionally, LUNA setup offers a set of <sup>3</sup>He counters for neutron detection as well as an array of large-area silicon detectors for charged particle detection. The  ${}^{2}H(p,\gamma){}^{3}He$  reaction (Q = 5.5 MeV), which contributes to the deuterium destruction during BBN (Big Bang Nucleosynthesis), is one of the reactions that was essential in the very beginning of the existence of the universe. The windowless gas target system and an HPGe detector installed in close geometry were used at LUNA to study the reaction  ${}^{2}$ H $(p,\gamma){}^{3}$ He at the center-of-mass energies between 30 keV and 263 keV. The reaction cross section was measured with a systematic error of less than 3% [125]. Due to its high sensitivity to the baryon density, or alternatively the baryon-to-photon ratio of the early Universe, deuterium abundance is utilized in BBN research as an indicator of cosmological parameters. The aforementioned findings enabled to substantially reduce the uncertainty on BBN predictions of the baryon density and the effective number of neutrino families.

The cosmic muon energy spectrum in a silicon semiconductor detector shows an energy maximum near zero and decreases exponentially with energy [126]. One can anticipate better signalto-background ratios in deep-underground studies aiming to detect charged particles with Si detectors since cosmic muons are greatly suppressed underground. The  $^{17,18}\text{O}(p,\alpha)^{14,15}\text{N}$  reactions which are crucial for the nucleosynthesis of key isotopes, and are used to constrain stellar models of novae, AGB, and post-AGB stars, were studied in detail [127], [128]. It was shown that enhanced background suppression was achieved. The combined effects of the background suppression underground and the generally improved experimental conditions (See ref. [129]) have resulted in the most accurate value to date for the  $E_P = 70$  keV resonance strength  $\omega \gamma$  in  $^{17}O(p,\alpha)^{14}N$ , namely  $\omega \gamma$ =  $(10.0 \pm 1.4\text{stat} \pm 0.7\text{syst})$  neV. As a result, there has been a factor of two increase in the  $^{17}\text{O}(p,\alpha)$   $^{14}\text{N}$ reaction rate and a reduced <sup>17</sup>O/<sup>16</sup>O ratio, with important consequences for the origin of some oxygenrich group II presolar grain. The improved results on the  ${}^{18}\text{O}(p,\alpha){}^{15}\text{N}$  cross sections and resonance strengths, with tighter constraints on oxygen isotopic ratios have been also obtained [128]. The sensitivity of a detector setup to the neutron background flux is closely related to the detection technique. Detectors based on neutron capture reactions (e.g. <sup>3</sup>He counters) are primarily sensitive to thermalized neutrons while organic scintillators based on elastic neutron scattering on hydrogen are only sensitive and selective to neutrons above a threshold energy. In order to modify the neutron energy spectrum (i.e., thermalize the neutrons) or decrease the neutron flux through neutron capture reactions, materials (mainly hydrogen-rich) may be used as neutron shielding around the detection

setup. The  ${}^{13}\text{C}(\alpha,n){}^{16}\text{O}$  reaction, is a main source of neutrons for the s-process. The reaction cross section was measured at LUNA [130], [131] using an array of eighteen stainless steel  ${}^{3}\text{He}$  counters to detect the produced neutrons. The counters were positioned in two concentric rings around the target chamber and embedded in a high-density polyethene moderator to thermalize the neutrons. The combination of these techniques has allowed for the measurement of  ${}^{13}\text{C}(\alpha,n){}^{16}\text{O}$  down to the astrophysically significant energies i.e. Gamow window. Activation measurements (see Section 2.1. for more details about the activation measurements) have been also performed at LUNA, some of the reactions measured include:  ${}^{3}\text{He}(\alpha,\gamma){}^{7}\text{Be}$ ,  ${}^{17}\text{O}(p,\gamma){}^{18}\text{F}$ ,  ${}^{12}\text{C}(p,\gamma){}^{13}\text{N}$  [131], [132], [133]. The LUNA collaboration has been studying low- energy nuclear reactions relevant to astrophysics for three decades. The work has led to unprecedented precision in the measurement of key reaction cross sections and to major breakthroughs in our understanding of the inner workings of stars. A vast number of measurements relevant to BBN, p-p Chain, CNO cycle, NeNa and MgAl Cycles, and s-process nucleosynthesis have been performed (See Ref. [124] and references therein).

Cross-section measurements spanning over as wide an energy range as possible are required in situations when the Gamow window cannot be reached in order to facilitate theoretical extrapolation. Because the reactions of more complex burning processes, such as the burning of helium and carbon, occur at higher temperatures, the laboratory studies of them require higher beam energies. The installation of the new LUNA MV machine took place in 2021. The LUNA MV accelerator is a single-ended inline Cockroft-Walton accelerator with a max 3.5 MV terminal voltage. The accelerator is able to accelerate proton,  $\alpha$ , and  $^{12}$ C beams in the energy range between 300 keV and 3.5 MeV.

# 2.16. Felsenkeller underground accelerator laboratory, Helmholtz-Zentrum Dresden-Rossendorf (HZDR), Germany

The Felsenkeller underground accelerator laboratory in Dresden, Germany is a shallow underground laboratory as opposed to the deep underground laboratories (such as LUNA) [134]. It is located under 45 meters of hornblende mazonite rock overburden (140 m.w.e.). With the exception of muons, this depth is sufficient to entirely shield all cosmic ray radiation's components. In the case of Felsenkeller, the key finding was that, at  $\gamma$ -ray energies of 5-8 MeV, the background rate in detectors at the shallow underground location with active detector shielding is only 2-3 times worse than at the deep underground laboratory [135] That background is low enough for a shallow underground accelerator which can be a complementary facility to deep underground ones. The Felsenkeller underground accelerator laboratory hosts a 5 MV Pelletron accelerator. Typical ion beam currents up to 30  $\mu$ A are reached for helium and carbon ions, and similar performance is expected for proton and other beams. The Felsenkeller underground laboratory's scientific agenda comprises the investigation of a number of important astrophysical processes, including  $^3$ He( $\alpha$ , $\gamma$ ) $^7$ Be and  $^{12}$ C( $\alpha$ , $\gamma$ ) $^{16}$ O [124]. A number of HPGe detectors surrounded by active veto detectors will be used to fully utilize the background reduction capabilities of the shallow underground laboratory.

# 2.17. Centre of Micro-Analysis of Materials (CMAM) at the Autonomous University of Madrid, Spain

CMAM is a Spanish Infrastructure that hosts an HVEE (High Voltage Engineering) 5 MV tandem linear accelerator [136]. The 5 MV Tandem provides MeV ion beams of any stable element. The main nuclear physics research area is the study of relevant nuclear reactions for astrophysics, typically those with excited states of certain nuclei near the particle threshold. The nuclei and the states of interest are populated in low energy reactions, and studied by particle and gamma detection. The end of the beamline is equipped with a big versatile reaction chamber. All setups share a design that can prevent the very strong signal from Rutherford scattering that would mask the reaction channel of interest. Additionally, the detector setups are easily interchangeable. A setup of fourteen Si detectors, each divided into four subdetectors, is available for charged particle detection at forward angles [137]. At backward angles, a setup of three DSSSD detectors is positioned, covering an angle from 85° to 170°. By using flanges in various directions, several gamma-detecting system types can be incorporated. As an example of the studies of astrophysical relevance carried out in this beamline, one can stress the measurement of the cross section of the  ${}^3\text{He}({}^4\text{He},\gamma){}^7\text{Be}$  in the energy range 1-3 MeV

[138]. Also, the excited states in  $^{12}$ C were explored using the  $^{10}$ B( $^{3}$ He,p3 $\alpha$ ) reaction, in order to get additional information of the triple-alpha process as well as to study  $\alpha$  clustering in light nuclei. The current research focuses on the study of the  $^{19}$ F(p, $\alpha\gamma$ ) $^{16}$ O reaction [139]. This reaction helps to understand the radiative capture of  $\alpha$  on  $^{12}$ C which is crucial for nuclear astrophysics since it has a major role in determining the C to O ratio during the burning of stars.

2.18. The Tandem Accelerator Laboratory (TAL), Institute of Nuclear and Particle Physics (INPP) of the National Centre for Scientific Research "Demokritos" (NCSRD), Greece

TAL is currently the primary research facility of the NCSRD's Institute of Nuclear and Particle Physics (INPP). It hosts the 5.5 MV T11/25 Van de Graaff Tandem accelerator. Quasi-monoenergetic neutron beams can be produced. TAL also hosts the 250 keV single-stage accelerator PAPAP [140]. PAPAP can accelerate proton and deuteron beams with currents of hundreds of µA. CALIBRA's goal is to create, operate, and exploit a cluster of accelerator laboratories for ion-beam research and applications at "Demokritos." The CALIBRA project's funding is being used to carry out a number of infrastructural improvements at the facility. A 17 MeV Scanditronix Cyclotron and a 2.5 MV Tandetron AMS accelerator, both fully functional machines, have been donated to the TAL by the University Medical Center Groningen (UMCG), the Netherlands, and the Institute of Archaeology of the University of Oxford, UK, respectively, for the establishment of the Cyclotron and the AMS Labs. The donated cyclotron produces and accelerates protons (p) and deuterons (d) of energy 17 and 8.5 MeV, respectively. The NeoPtolemos γ -summing spectrometer, is a cylindrically shaped (14 "×14") NaI(Tl) detector with a borehole of 32 mm diameter along its axis. Its absolute efficiency for a twofold  $\gamma$  cascade is higher than 50 %. The TAL group created a method called " $4\pi$ -summing". The method makes it possible to measure angle-integrated spectra rather than many angular distributions. (see for example Ref. [141], [141], [142]). This method relies on the use of a large-volume NaI(Tl) crystal detector with the highest possible absolute  $\gamma$  -ray detection efficiency. As the name " $4\pi$ -summing" suggests, the all  $\gamma$ -rays that de-excite the entry state of a compound nucleus and form  $\gamma$  cascades are summed. This practically allows to analyze only one single peak -"sum peak" instead of multiple γ transitions. GASPAR is an acronym for "GASP for Astrophysics Research," where "GASP" stands for the GAmma SPectrometer [143] that used to be located at the INFN Laboratori Nazionali di Legnaro (INFN-LNL). GASP comprises 40 Compton-suppressed HPGe and an 80-crystal BGO calorimeter that covers 80% of the solid angle. Compared to the NeoPtolemos calorimeter, the GASPAR BGO Ball has certain advantages like its ability, not only to sum  $\gamma$  transitions forming a  $\gamma$ -cascade but also to provide the multiplicity of the γ-cascade, which determines the detector's summing efficiency. For the study of neutron-emitting processes, an array of sixteen <sup>3</sup>He-gas-filled neutron counters is also provided [144]. The multipurpose large-volume scattering chamber and the deuterium-filled gas cell for the generation of quasi-monochromatic neutrons are among the experimental tools that are most frequently used. [140].

The primary objective of the Nuclear Astrophysics program at TAL is to further clarify nuclear physics aspects of the p-process [140]. The solar system p-nuclei abundances still pose a challenge for all p-process nucleosynthesis models, which fail to reproduce them, especially in the case of the light p-nuclei. It is strongly needed that the nuclear physics uncertainties entering the astrophysical calculations are reduced. For the goals of this research program, systematic cross-section measurements of proton and  $\alpha$ -particle capture reactions are performed at the in-house Tandem (as well as abroad). More than twenty capture reactions have been investigated at NCSR Demokritos so far (see [140] and references therein, also refs. [142], [145]). The measurements of the (n, 2n), (n, p), (n,  $\alpha$ ) and (n, f) cross sections have been the focus of the successful research program. The neutron facility of TAL, which can produce monoenergetic neutron beams with fluxes of  $\approx 10^5$ – $10^6$  ncm<sup>-2</sup> s<sup>-1</sup> in the energy ranges of thermal to 450 keV, 4 MeV to 11.5 MeV, and 16 MeV to 20.5 MeV by using the <sup>7</sup>Li(p,n), <sup>2</sup>H(d, n) and <sup>3</sup>H(d, n) reactions [146], [147] , is a key tool in the implementation of this program. The laboratory has also taken part in a coordinated research program on proton-induced gamma-ray emission (PIGE) technique improvement, evaluating existing PIGE data, and measuring

new ones, funded by the IAEA. Numerous differential cross-sections have been measured in a variety of energies and angles for this purpose. [148], [149].

# 2.19. Microanalytical Centre, Jožef Stefan Institute (MIC, JSI), Slovenia

2 MV Tandetron is the only research accelerator in Slovenia [150]. Compared to similar installations in Europe, the Ljubljana 2 MV Tandetron is particularly powerful, as in addition to two standard ion sources (sputter and duoplasmatron), it also uses an extremely bright multicusp source of negative hydrogen ions. The nuclear physics research at 2 MV Tandetron focuses on reactions between light nuclei, in particular on the electron screening effect. The cross sections for reactions involving positively charged nuclei are higher in the presence of atomic electrons than they are in the absence of electrons. Unfortunately, the effect is not understood well enough to predict its role in stellar plasma, but it was noticed that the presence of electrons in the solid state may enhance reaction rates by several orders of magnitude [151]. The effect is particularly large when various hydrogen isotopes are implanted into a metallic lattice. Recently, it was discovered that electrons were emitted instead of  $\gamma$  rays in the nuclear reaction  ${}^2H(p,\gamma){}^3He$  at a significantly higher rate than would be expected from the internal conversion process [152]. This demonstrated that electrons actively participate in the reaction and do not just simply lower the Coulomb barrier. Since the nuclear physics group is also very active in detector development, a testing station for tests of detectors for GSI/FAIR is being currently commissioned

## 3. Conclusions

The nuclear physics field focuses its top research at large-scale facilities which encompass highly competitive research proposals as well as complex detector arrays and acquisition systems. On the other hand, small-scale facilities play an important role in the landscape of nuclear physics research in Europe. The current status, available instrumentation, as well as perspectives of nuclear physics research at nineteen European small-scale facilities are given. The small-scale facilities contribute to our understanding of nuclear physics through research in nuclear structure, reactions, and nuclear astrophysics. They perform a diverse research program in nuclear physics and applications and are specifically endorsed by the NuPECC LRP. In addition, small-scale facilities develop novel technologies and methodologies. They address a wide range of fundamental questions and are essential for teaching and training personnel in accelerator technology and science, providing them with diverse skill sets. Nuclear physics research at small-scale facilities has certainly thrived in the last decade. The current upgrades and new facilities such as e.g. the new fragment separator FRAISE at LNS Catania, LUNA MV, and the Felsenkeller underground accelerator laboratory are paving the road to exciting new discoveries. Nuclear physics is an international venture and we hope that this review paper will be useful for researchers working at small-scale facilities, to understand the position of their facility in the European accelerator-based nuclear physics landscape and find possible synergies.

# References

- 1. [Online]. Available: https://www.esfri.eu/esfri-roadmap.
- 2. NuPECC, "The Nuclear Physics European Collaboration Committee (NuPECC) Long Range Plan," 2017. [Online]. Available: https://www.esf.org/fileadmin/user\_upload/esf/Nupecc-LRP2017.pdf.
- 3. [Online]. Available: https://www.ikp.uni-koeln.de/en/institute/facilities/accelerators.
- 4. L. Netterdon et al, Nucl. Instrum. Methods Phys. Res. A, vol. 100, p. 94, 2014.
- 5. W. Hillebrandt et al, vol. 8, p. 116, 2013.
- 6. G. Gyürky et al, Eur. Phys. J. A, vol. 41, p. 55, 2019.
- 7. A. Spyrou et al, Phys. Rev. C, vol. 76, p. 015802, 2007.
- 8. S. Galanopoulos et al, Phys. Rev. C, vol. 67, p. 015801, 2003.
- 9. S. Pickstone et al, Nucl. Inst. and Meth. in Phys. Res. A, vol. 875, p. 104, 2017.
- 10. E. Stokstad et al, Nucl. Phys. A, vol. 1, p. 145, 1970.
- 11. [Online]. Available: https://www.slcj.uw.edu.pl/en/heavy-ion-laboratory-hil.

- 12. [Online]. Available: https://www.slcj.uw.edu.pl/en/the-warsaw-igisol.
- 13. [Online]. Available: https://www.slcj.uw.edu.pl/en/coulomb-excitation-at-the-warsaw-cyclotron.
- 14. [Online]. Available: https://inis.iaea.org/collection/NCLCollectionStore/\_Public/38/068/38068159.pdf.
- J. Mierzejewski et al, Nucl. Instrum. Methods Phys. Res. A, vol. 659, p. 84, 2011.
- 16. P. Petkov et al, Nucl. Instr. Meth. A, vol. 431, p. 208, 1999.
- 17. J. Perkowski et al, Rev. Sci. Inst., vol. 85, p. 043303, 2014.
- 18. T. Morek et al, Phys. Rev. C, vol. 63, vol. 63, p. 034302, 2001.
- 19. K. Wrzosek et al, Phys. Rev. C, vol. 86, p. 064305, 2012.
- 20. E. Grodner et al, Phys. Rev. Let., vol. 97, p. 172501, 2006.
- 21. E. Grodner et al, Phys. Lett. B, vol. 703, p. 46, 2011.
- 22. E. Grodner et al, Phys. Rev. Let., vol. 120, p. 022502, 2018.
- 23. E. Grodner et al, Phys. Rev. C, vol. 106, p. 014318, 2022.
- 24. A. Stocchi et al, Nuclear Physics News, vol. 32 1, 2022., vol. 32 1, 2022.
- 25. J. Wilson et al, Phys. Procedia, vol. 59, p. 31, 2014.
- 26. [Online]. Available: https://alto.ijclab.in2p3.fr/en/instrumentation-en/bedo/.
- 27. M. Lebois et al, Acta Phys. Pol. B, vol. 50, p. 3, 2019.
- 28. J. Wilson et al, Nature, vol. 590, p. 566, 2021.
- 29. J. Wilson, "Report for the conseil scientifique of the the in2p3," 2019. [Online]. Available: https://www.in2p3.cnrs.fr/sites/institut\_in2p3/files/page/2019-07/5-Doc-WILSON.pdf.
- 30. [Online]. Available: https://andromede.in2p3.fr/?lang=en.
- 31. M. Heine et al, EPJ Web of Conferences, vol. 165, p. 01029, 2017.
- 32. M. Rudigier et al, Nucl. Instrum. Meth. A, vol. 969, p. 163967, 2020.
- 33. G. Fruet et al, PRL, vol. 124, p. 192701, 2020.
- 34. I. Bogdanović Radović, Nuclear Physics News, vol. 30 2, 2020.
- 35. A. Tumino et al, Annu. Rev. Nucl. Part. Sci, vol. 71, pp. 345-376, 2021.
- 36. N. Soić et al, Europhys. Lett., vol. 34, p. 7, 1996.
- 37. M. Freer et al., Phys. Rev. V, vol. 84, p. 034317, 2011.
- 38. M. Venhart et al, Nuclear Physics News, vol. 33 1, 2023.
- 39. A. Görgen et al, Eur. Phys. J. Plus, pp. 136-181, 2021.
- 40. F. Zeiser et al, Nucl. Instr. Meth. A, vol. 985, p. 164678, 2021.
- 41. M. Guttormsen et al, Nucl. Instr. Methods A, vol. 648, p. 168, 2011.
- 42. T. Tornyi et al, Nucl. Instr. Methods A, vol. 738, p. 6, 2014.
- 43. A. Schiller et al, Nucl. Instr. Methods A, vol. 447, p. 498, 2000.
- 44. M. Guttormsen et al, Phys. Rev. C, vol. 88, p. 024307, 2013.
- 45. S. Goriely et al, Eur. Phys. J. A, vol. 55, p. 172, 2019.
- 46. F.-K. Thielemann, Eur. Phys. J. A, vol. 59, p. 12, 2023.
- 47. A. Choplin et al, "A&A," vol. 667, p. A155, 2022.
- 48. F. Hoyle, Astrophys. J. Suppl. Ser., vol. 1, p. 12, 1954.
- 49. T. Kibédi et al, Phys. Rev. Lett, vol. 125, p. 182701, 2020.
- 50. N. Pietralla, Nuclear Physics News, vol. 28 2, 2018.
- 51. K. Sonnabend et al, Nucl. Instrum. Meth. Phys. Res. A, vol. 640, p. 6, 2011.
- 52. N. Ryezayeva et al, Phys. Rev. Lett, vol. 100, p. 172501, 2008.
- 53. M. Chernykh et al, Phys. Rev. Lett., vol. 105, p. 022501, 2010.
- 54. C. Kremer et al, Phys. Rev. Lett., vol. 117, p. 172503, 2016.
- 55. P. P von Neumann-Cosel et al, Phys. Rev. Lett., vol. 88, p. 202304, 2002.
- 56. C. Walz et al, Nature, vol. 405, p. 526, 2015.
- 57. "https://www.lns.infn.it/en/," [Online].
- 58. M. Lattuada, Nuclear Physics News, vol. 22 1, 2012.
- 59. A. Pagano et al, Nucl. Phys. A, vol. 734, pp. 504-511, 2004.
- 60. E. Pagano et al, EPJ Web of Conferences, p. 10008, 2016.
- 61. D. Dell'Aquila et al, EPJ Web of Conferences, vol. 117, p. 06011, 2016.
- 62. N. Mortorana et al, IL NUOVO CIMENTO, vol. 41 C, p. 199, 2018.
- 63. F. Cappuzzello et al, Eur. Phys. J. A, vol. 52, p. 167, 2016.
- 64. L. Lamia et al, ApJ, vol. 811, p. 99, 2015.

- 65. S. Romano et al, Eur. Phys. J. A, vol. 27, pp. 221-225, 2006.
- 66. C. Spitaleri et al, Phys. Rev. C, vol. 90, p. 035801, 2014.
- 67. S. Simonucci et al, ApJ, vol. 764, p. 118, 2013.
- 68. L. Lamia et al, J. Phys. G: Nucl. Part. Phys., vol. 39, p. 015106, 2012.
- 69. M. Cognata et al, Phys. Lett. B, vol. 826, p. 136917, 2022.
- 70. R. Sparta et al, Eur. Phys. J. A, vol. 170, p. 57, 2021.
- 71. G. Guardo, Eur. Phys. J A, vol. 59, p. 65, 2023.
- 72. [Online]. Available: http://www.micronsemiconductor.co.uk/silicon-detector-catalogue/.
- 73. N. Martorana et al, Front. Phys., vol. 10, p. 1058419, 2022.
- 74. I. Vajda et al, AIP Conference Proceedings, vol. 1852, p. 060002, 2017.
- 75. G. Kiss et al, J. Phys.: Conf. Ser., vol. 337, p. 012029, 2012.
- 76. Z. Halász et al, Phys. Rev. C, vol. 94, p. 045801, 2016.
- 77. G. Kiss et al, Phys. Rev. C, vol. 88, p. 045804, 2013.
- 78. T. Szücs et al, Phys. Rev. C, vol. 100, p. 065803, 2019.
- 70 T C----- di et al I Disse Comf Com --- 1 1660 -- 012041 20
- 79. T. Szegedi et al, J. Phys.: Conf. Ser., vol. 1668, p. 012041, 2020.
- 80. G. Gyürky et al, Eur. Phys. J. A, vol. 59, p. 59, 2023.
- 81. C. Bordeanu et al, Nucl .Phys. A, vol. 908, pp. 1-11, 2013.
- 82. G. Gyürky et al, J. Phys. G: Nucl. Part. Phys., vol. 48, p. 105202, 2021.
- 83. G. Gyürky et al, Eur. Phys. J A, vol. 59, p. 59, 2023.
- 84. G. Gyürky et al, Phys. Rev. C, vol. 105, p. 022801, 2022.
- 85. J. Krasznahorkay et al, Phys. Rev. Lett., vol. 116, p. 042501, 2016.
- 86. D. Firak et al, EPJ Web of Conferences, vol. 232, p. 04005, 2020.
- 87. A. Krasznahorkay et al, Phys. Rev. C, vol. 106, p. 061601.
- 88. A. Krasznahorkay at al, Phys. Rev. C, vol. 106, p. 061601, 2022.
- 89. B. Bastin et al, EPJ Web of Conferences, vol. 275, p. 01012, 2023.
- 90. A. Macková et al, Eur. Phys. J. Plus, vol. 136, p. 558, 2021.
- 91. A. Tumino et al, Astrophys J, vol. 785(2), p. 96, 2014.
- 92. V. Burjan et al, Eur. Phys. J. A, vol. 55(7), no. 114, 2019.
- 93. G. D'Agata et al, Phys. Rev. C, vol. 103(1), p. 015806, 2021.
- 94. [Online]. Available: https://ccb.ifj.edu.pl/en.home.html.
- 95. I. Ciepał et al, Phys. Rev. C, vol. 99, p. 014620, 2019.
- 96. A. Maj et al, Nucl. Phys. A, vol. 571, p. 185, 1994.
- 97. J. Łukasik et al, Nucl. Instrum. Methods Phys. Res. A, vol. 709, p. 120, 2013.
- 98. A. Maj et al, "White Book on the Future of Low-Energy Nuclear Physics in Poland and the Development of the National Research Infrastructure," 2020.
- 99. J. Gómez-Camacho et al, Eur. Phys. J. Plus, vol. 136, p. 273, 2021.
- 100. A. Garzón-Camacho et al, IEEE Trans. Inst. Meas., vol. 64, p. 318, 2015.
- 101. [Online]. Available: https://cna.us.es/index.php/en/facilities/financed-feder-funds/equipments/78-neutron-line.
- 102. E. Alves et al, Eur. Phys. J. Plus, vol. 136, p. 684, 2021.
- 103. D. Bucurescu et al, Nucl. nstr. and Meth. in Phys. Res. A, vol. 837, pp. 1-1, 2016.
- 104. [Online]. Available: https://www.nipne.ro/.
- 105. P. Mason et al, Phys. Rev. C, vol. 85, p. 064303, 2012.
- 106. T. Alharbi et al, Appl. Radiat. Isot., vol. 70, p. 1337, 2012.
- 107. C. Niță et al, Phys. Rev. C, vol. 89, p. 064314, 2014.
- 108. G. Bocchi et al, Phys. Rev. C, vol. 89, p. 054302, 2014.
- 109. T. Alharbi et al, Phys. Rev. C, vol. 91, p. 027302, 2015.
- 110. D. Deleanu et al, Phys. Rev. C, vol. 87, p. 014329, 2013.
- 111. P. Mason et al, Phys. Rev. C, vol. 88, p. 044301, 2013.
- 112. C. Thibault et al, Phys. Rev. C, vol. 12, p. 644, 1975.
- 113. B. Wildenthal et al, Phys. Rev. C, vol. 22, p. 2260, 1980.
- 114. E. Warburton et al, Phys. Rev. C, vol. 41, p. 1447, 1990.
- 115. N. Orr et al, Phys. Lett. B, vol. 258, p. 29, 1991.
- 116. F. Recchia et al, Phys. Rev. C, vol. 85, p. 064305, 2012.

- 118. S. Leoni et al, Phys. Rev. Lett., vol. 118, p. 162502, 2017.
- 119. R. Margineanu et al, Appl. Radiat. and Isot., vol. 66, p. 1501, 2008.
- 120. L. Trache et al, EPJ Web of Conferences, vol. 227, 2020.
- 121. D. Tudor et al, AIP Conf. Proc., vol. 2076, p. 060010, 2019.
- 122. N. Zhang et al, Phys. Lett. B, vol. 803, p. 135278, 2020.
- 123. N. Zhang et al, Phys. Lett. B, vol. 801, p. 135170, 2020.
- 124. M. Aliotta et al, Annu. Rev. Nucl. Part. Sci., vol. 72, pp. 177-204, 2022.
- 125. V. Mossa et al, Nature, vol. 587, pp. 210-213, 2020.
- 126. M. Misiaszek et al., Appl. Radiat. Isot., vol. 81, p. 146, 2013.
- 127. C. Bruno et al, Phys. Rev. Lett., vol. 117(14), p. 142502, 2016.
- 128. C. Bruno et al, Phys. Lett. B, vol. 790, pp. 237-242, 2019.
- 129. C. Bruno et al, Eur. Phys. J. A, vol. 51, p. 94, 2015.
- 130. L. Csedreki et al, Nucl. Instrum. Meth. Phys. Res. A, vol. 994, p. 165081, 2021.
- 131. G. Ciani et al, Phys. Rev. Lett., vol. 127(15), p. 152701, 2021.
- 132. D. Scott et al, Phys. Rev. Lett., vol. 109(20), p. 202501, 2012.
- 133. A. Di Leva et al, Phys. Rev. C, vol. 89(1), p. 015803, 2014.
- 134. [Online]. Available: https://www.chetec-infra.eu/ta/tna-labs/felsenkeller/.
- 135. T. Szücs et al, Eur. Phys. J. A, vol. 55, p. 174, 2019.
- 136. A. Redondo-Cubero et al, Eur. Phys. J. Plus, vol. 136, p. 175, 2021.
- 137. L. Fraile et al, Nucl. Instrum. Methods Phys. Res. A, vol. 513(1), p. 287, 2003.
- 138. M. Carmona-Gallardo et al, Phys. Rev. C, vol. 86(3), p. 032801, 2012.
- 139. M. Alcorta et al, Phys. Rev. C, vol. 86(6), p. 064306, 2012.
- 140. S. Harissopulos et al, Eur. Phys. J. Plus, vol. 136, p. 617, 2021.
- 141. S. Harissopulos, Eur. Phys. J. Plus, vol. 133, p. 332, 2018.
- 142. V. Foteinou et al, Eur. Phys. J. A, vol. 55, p. 67, 2019.
- 143. N. Medina et al, APH N.S. Heavy Ion Phys., vol. 2141, 1995.
- 144. S. Harissopulos et al, Phys. Rev. C, vol. 72, p. 062801, 2005.
- 145. S. Harissopulos et al, Phys. Rev. C, vol. 87, p. 025806, 2013.
- 146. R. Vlastou et al, Phys. Proc., vol. 66, p. 425, 2015.
- 147. A. Kalamara et al, Eur. Phys. J. A, vol. 55, p. 187, 2019.
- 148. K. Preketes-Siglas et al, vol. 368, p. 71, 2016.
- 149. K. Preketes-Sigalas et al, Nucl. Instrum. Methods Phys. Res. B, vol. 368, p. 71, 2016.
- 150. J. Vesic, Nuclear Physics News, Vols. 32, 3, 2022.
- 151. F. Raiola et al, Eur. Phys. J. A, vol. 19, p. 283, 2004.
- 152. M. Lipoglavšek et al, Phys. Lett. B, vol. 773, p. 553, 2017.

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