

Article

A Deep Neural Network Regression of the Cahn–Hilliard Single-Particle Thermal Model for LiFePO₄ Batteries

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Abstract: Lithium-ion batteries serve as the primary sources of power for electric vehicles (EVs) and hybrid electric vehicles (HEVs). For vehicle applications, battery management systems (BMSs) are necessary to protect lithium-ion batteries from overheating and to ensure optimum vehicle performance. Our approach to developing a BMS was based on recent advances in the application of phase field models for lithium-ion batteries. In particular, our reduced-order model (ROM) utilized a dataset generated from the COMSOL® Multiphysics simulation of the Cahn–Hilliard equation for a single particle of a lithium iron phosphate (LiFePO₄) cathode: an example of using a reduced-order model (ROM) based on a single-particle model (SPM). The main innovation of our ROM is that the SPM is fully coupled to a heat transfer model at the battery cell level. We utilized principal component analysis to identify a lower-order model that could reproduce the battery's voltage and temperature response for ambient temperatures ranging from 253 to 298 K and for discharge rates ranging from 1 C to 20.5 C. The reduced-order dataset was then fitted to the experimental data for an A123 Systems 26650 2.3 Ah cylindrical battery using deep neural network (DNN) regression. The entire BMS is realized in conjunction with a digital-twin (DT) configuration with an offboard COMSOL® Multiphysics SPM simulation and training of the DNN, allowing the ROM to be periodically updated by retraining the DNN for aging batteries and actual operating conditions. In this configuration, only the trained DNN predictor function is onboard and in real time.

Keywords: lithium-ion batteries; LIB; LiFePO₄; electric vehicles; Cahn–Hilliard; principal component analysis, neural networks

1. Introduction

Global oil resources will eventually dry up and fossil-fueled vehicles are a major source of greenhouse emissions. There is an ongoing global trend of electrifying transportation, and as a result, intensive research is being conducted into electric vehicles and hybrid electric vehicles (EVs and HEVs) [1,2]. Lithium-ion batteries have become the most promising choice for EVs due to their high energy density and long cycle life [3,4]. For EV applications, battery management systems (BMSs) are necessary to prevent Li-ion batteries from overheating and overcharging and avoid potential thermal runaway. Currently, BMSs cannot store and process large amounts of data while managing the battery's state of charge (SOC), voltage, and temperature in real time [5–7]. To address this shortcoming, additional research is needed to develop reduced-order models (ROM) that can both model complex battery mechanisms and provide real-time management data. Applications of

machine learning (ML) and artificial intelligence (AI) are fertile areas for eventual solutions to this problem.

Models that describe battery dynamic processes at all levels are not feasible. Therefore, mechanistic battery models, referred to as full-order models (FOMs), need to be tailored to specific purposes that require a deep understanding of a particular aspect of the battery's operation and performance. Many mechanistic models have been developed specifically to describe battery thermal behavior [8–12]. In general, models for analysis and diagnosis purposes employ detailed simulations of the battery's physics and, thus, are often multidimensional, multiphysics models and are computationally slow. Models for control and optimization applications are usually computationally fast, but provide a limited description of the underlying physics. The most commonly used mechanistic model for single-battery cells is a lumped parameter simplification of a FOM called the pseudo-two-dimensional (P2D) model. Even given P2D simplifications of battery homogeneity and a constant electrode thickness, P2D models still use more than fifteen parameters specific to a particular battery and their computational cost is still too high for control and optimization applications. On the other hand, P2M parameters are too lumped to provide much insight into the underlying phenomena [13,14]. Decades of research have been conducted to develop reduced-order models (ROMs) that adequately retain the robustness of FOMs without excessive computational complexity. The state of the art for the various methods of achieving Li-ion battery ROMs has recently been reviewed in the literature [15,16]. We limit our following discussion of such methods to equivalent circuit models (ECMs), their model-based extensions, and the single-particle model approximation. These methods are most pertinent since our ROM, developed in the following sections, is a single-particle application and is most likely to be used instead of or alongside an ECM.

ECMs that do not consider fundamental physics have been extensively used to imitate the relationships between battery input and output systems while offering real-time computation [17–22]. ECMs use electrical circuits to simulate lithium-ion cells utilizing capacitors to shape the battery capacity, while variable resistors and controlled-voltage sources shape the temperature effect or SOC variations. Black-box modeling is another extensively applied method to provide real-time computation. It relies on developing an equivalent transform function with different inputs and outputs. Like ECMs, this method depends on experimental data for a specified battery [23]. ECMs and transform functions are generally implemented in a “mixed approach” along with thermal and aging models, as shown in Figure 1 [24].

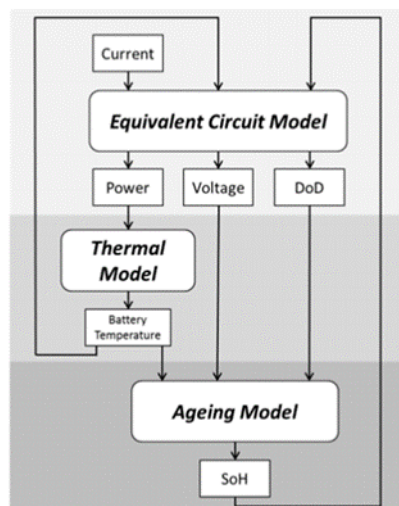


Figure 1. Equivalent-Circuit-Model-Based Battery Management System.

Model-based methods have been developed that have greatly enhanced the basic ECM. Some of the predominant model-based methods have been recently reviewed in the literature including the Luenberger, sliding mode, Kalman filter, and proportional integral (PIO) methods [25–30]. Model-based methods have a battery model at their core that uses measured current and voltage signals to provide a closed-loop estimation method. These methods typically use static models where model

parameters are applied offline and assumed to not change over time. However, these models often do not ensure sufficient accuracy over a broad range of operation conditions and long time frames. In light of this, efforts have been made towards online model adaptation methods, including population-based optimization, dual filtering, and least squares (LS)-based methods such as moving window LS, continuous-time LS, and recursive least squares (RLS) [31–33].

Lithium iron phosphate (LiFePO_4) is the most frequently used phosphate-based cathode material in Li-ion batteries. LiFePO_4 has a strong tendency to separate into solid high-Li⁺-concentration and low-Li⁺-concentration phases, leading to the battery's characteristic broad voltage plateau at room temperature [34–39]. Traditionally, mathematical models of intercalation dynamics in LiFePO_4 cathodes were based on spherical diffusion or the shrinking core concept [40–42]. However, recent experimental and theoretical progress suggests that a more realistic SPM should encompass a phase field model for equilibrium and nonequilibrium solid-solution transformations [43–47]. A phase-field model is a computational method for modeling morphological and microstructure evolution in materials. They have been proposed for solid-state phase transformations, grain growth and coarsening, microstructure evolution in thin films, and crack propagation [48–51]. Research is currently underway that seeks to bridge the gap between phase-field mesoscale models and macro battery properties. For example, (Yuan et al. [52], 2021) recently used a phase field modeling approach to develop a physics-based, fully coupled model that bridges dendrite and crack propagation at the micro level with macrostate battery charging and discharging [53]. Another application (Zeng & Bazant [54], 2013) demonstrated a method for estimating the voltage plateau of LiFePO_4 batteries based on the Cahn–Hilliard phase field model solution for a single cathode particle. Our ROM was motivated by (Zeng & Bazant, 2013), and the spatial mass transfer elements of our 3-D COMSOL® Multiphysics finite element solution are similar to their 1-D, isotropic, and isothermal solution. The major enhancement of our single-particle model (SPM) is that it is a multiphysics thermal model that fully couples the battery cell's heat transfer model. Statistically, we relate the SPM simulation to battery cell property estimation. Specifically, the plateauing effect of the battery's voltage response at higher ambient temperatures and the apparent diffusion-controlled behavior at lower temperatures are related to the SPM by statistical inference. Our SPM is a specific example of a widely used method for order reduction in P2D based on a single-particle model that aims to enhance computational run time while retaining elements of the underlying physics, as opposed to an ESM [55,56]. In the single-particle thermal model preparation, the local potential and concentration gradients in the electrolyte phase were ignored and accounted for by utilizing a lumped solution resistance term. Likewise, the potential gradient in the solid phase of the electrodes was dismissed, and the porous electrode was considered as a large number of individual particles, all subjected to the same conditions. These assumptions are generally only valid under relatively low current rates and the SPM is not recommended for high-power applications, such as fast EV charging and operations involving high-power pulses, but it is well-suited to daily EV driving, where the operating ranges are less extreme. These shortcomings of the SPM are not prohibitive for our ROM since we did not attempt to model the battery voltage response by directly scaling up the SPM, but sought to only retain enough information from the SPM to make statistical inferences concerning the macro battery properties. The ROM was realized by subjecting the raw simulation results from the COMSOL® Multiphysics simulation data to principal component analysis (PCA) to determine the lowest-order simulation dataset capable of fitting the experimental data using deep neural network (DNN) regression. We validated our SPM based on available experimental data for the A123 Systems 26650 2.3 Ah battery [57].

2. Methods and Materials

The scope of our modelling effort is as follows

1. Use experimental charge/discharge and heat transfer data for the A123 Systems 26650, 2.3 Ah lithium-ion battery to estimate the parameters of our SPM.
2. Use the simulation results from the SPM and experimental data to train and test a DNN.

3. Train and test the same DNN with a similar data set except it is derived from real time operation data for the extreme road test: Up Mount Sano in Huntsville, AL
- We are stating this here to avoid confusion by our discussion of the implementation of the SPM and DNN in an actual BMS in the following sections.

2.1. COMSOL Simulation of the Cahn–Hilliard Single-Particle Model

Our SPM poses the Cahn–Hilliard equation in COMSOL® Multiphysics’ standard PDF format as two coupled second-order PDEs in ion concentration and chemical potential, respectively. Details of the model variables along with all parameter magnitudes and units from the following equations are given below in the nomenclature section. We assumed the particles to be spherical and isotropic. The model equations are given below, where the overbar denotes dimensionless parameters and variables. We provide only the main points of the derivation of the Cahn–Hilliard equation and refer the reader to (Zeng & Bazant, 2013) for more details. The diffusional chemical potential based on the regular solution model and acquired from the Cahn–Hilliard free energy functional is

$$\bar{\mu} = -k_b T \ln \left[\frac{\bar{c}}{1-\bar{c}_m} \right] + \frac{\bar{\Omega}(c_m-\bar{c})}{c_m} - \frac{K V_s}{\text{cm}} \bar{V}^2 \bar{c} \tag{1}$$

The basic equation of evolution for mass conservation is

$$\frac{\partial \bar{c}}{\partial \bar{t}} = -\bar{V} \cdot \bar{q} \tag{2}$$

The ion flux is driven by the gradient of the diffusional chemical potential as

$$\bar{q} = \frac{-D_O(c_m-\bar{c})}{k_m T c_m} \bar{V} \bar{\mu} . \tag{3}$$

Voltage enters the Cahn–Hilliard SPM through Butler–Volmer kinetics obtained from transition state theory for concentrated solutions as

$$\bar{I} = \bar{I}_0 [\exp(-\alpha \bar{\eta}) i - \exp((1-\alpha)\bar{\eta})] . \tag{4}$$

where α is the electron transfer symmetry factor, $\eta = \Delta\phi - \Delta\phi_{eq}$ is the surface overpotential because of the activation polarization, $\Delta\phi$ is the local voltage drop across the interface, and $\Delta\phi_{eq}$ is the Nernst equilibrium voltage. The boundary conditions we implemented in the dimensionless form are

$$\begin{aligned} \bar{q}(o, t) &= 0, \\ \hat{n} \cdot \bar{q}(1, t) &= -\bar{q}, \\ \hat{n} \cdot \bar{V} \bar{c}(0, t) &= 0, \\ \hat{n} \cdot \bar{V} \bar{c}(1, t) &= 0. \end{aligned}$$

2.2. Coupling the Heat Transfer Model to the Cahn–Hilliard Equation

Transient temperature response and thermal power conservation were incorporated into the ROM by an enthalpy balance on the bulk battery. An A123 Systems 26650 2.3 Ah cylindrical battery was selected because it is extensively studied and property data are readily available. The battery specifications are shown in Table 1.

Table 1. Specifications for Systems A123 2.3 Ah Lithium-ion Battery.

Nominal capacity and voltage	2.3 Ah, 3.3 V
Internal impedance (1 kHz AC)	8 mΩ typical
Internal resistance 10 A, 1 s DC	10 mΩ typical
Recommended charge method	3 A to 3.6 V CCCV, 45 min
Recommended fast charge current	10 A to 3.6 CCCV

Maximum continuous discharge	70 A
Pulse discharge at 10 s	120 A
Cycle life at 10 C discharge	Over 1000 cycles
Recommended pulse charge/discharge cutoff	3.8 V to 1.6 V
Operating temperature range	243 K to 333 K
Core cell weight	70 g

A cylindrical Li-ion battery is constructed by rolling a stack of cathode/separator/anode layers. The individual layered sheets are thin, and lumped parameters are used. Therefore, material properties such as thermal conductivity, density, and specific heat capacity are presumed to be constant in a homogeneous and isotropic body. In the axial direction, the thermal conductivity is one or two orders of magnitude higher than in the radial direction, leading to a relatively uniform temperature distribution in the axial direction [58,59]. Additionally, considering natural convection, the heat transfer at the surface is much smaller than the internal heat transfer by conduction, leading to negligible temperature gradients inside the battery. Based on these assumptions, the energy balance equation in the battery can be expressed by one bulk volume-averaged temperature. To estimate the thermal response of the battery, we utilized a simplified energy balance equation for the enthalpy change for electrochemical reactions [60]. Assuming a constant system volume and pressure and neglecting heat generation because of enthalpy of mixing, the energy balance equation is presented as

$$M c_p \frac{\partial T}{\partial t} = I \left(V_{oc} - T \frac{\partial V_{oc}}{\partial T} \right) - IV + \dot{q}_{sur} \quad (5)$$

The term $T \frac{\partial V_{oc}}{\partial T}$ stands for reversible heat generation and can be calculated from the entropy of the reaction [61]. In this study, this reversible heat generation was ignored for simplicity. This simplification is justified since it is insignificant for the LiFePO₄ chemistry [62,63]. Assuming this simplification, the OCV becomes a function of SOC only, and Equation (5) was solved exactly for the battery's temperature response.

$$\bar{T}(t) = T_{surr.} - \frac{e^{-\frac{A_b h_c t}{c_p V \rho}} \left(-1 + e^{\frac{A_b h_c t}{c_p V \rho}} \right) I_b (\bar{V}(T) - \bar{V}_{oc})}{A_b h_c} \quad (6)$$

The temperature enters the Cahn–Hilliard model via Equation (1) as

$$\bar{\mu}(T) = -k_b \bar{T}(t) \ln \left[\frac{\bar{c}}{1 - \bar{c}_m} \right] + \frac{\bar{\Omega}(T)(c_m - \bar{c})}{c_m} - \frac{K V_s}{cm} \bar{\nabla}^2 \bar{c} \quad (7)$$

Assuming that the ion activity in the electrolyte adjacent to the particle (based on the dimensionless ion concentration) is 1.0, $\Delta\phi_{eq} = -\mu/e$. This provides the voltage profile for the single-particle battery $V = V^\theta + \eta - \mu/e$, where V^θ is the standard potential defined by the half-cell voltage (3.42 V vs. Li metal). The solution for η gives the voltage response of the single-particle battery as

$$V + V^\theta = \eta = +\frac{\mu}{e} = \frac{k_b T}{e} \left(-\bar{\mu} - 2 \sinh^{-1} \left(\frac{\bar{I}}{I_0(\bar{c})} \right) \right) \quad (8)$$

2.3. Digital Twin Configuration of BMS

In this section we describe how our SPM and DNN fit into an actual BMS. This is necessary because both components must have on-board and off-board configurations. The value of our approach is realized in conjunction with a digital-twin (DT) configuration with an offboard COMSOL® Multiphysics SPM simulation and training of the DNN, allowing the ROM to be periodically updated by retraining the DNN for aging batteries and actual operating conditions. In this configuration, only the trained DNN predictor function is onboard and in real time. Our approach to implementing a DT for lithium-ion batteries follows the process of (Singh et al. [61], 2021). There is some ambiguity as to what constitutes a DT and it may be more appropriate to refer to our process as a part of the DT for the entire EV but the main focus is to designate on-board and

off-board components and the flow of information between the components. At a minimum the DT configuration for a BMS, as shown schematically in Figure 2, shows the role of the SPM and the DNN. In Figure 2 we refer to the COMSOL® Multiphysics finite element simulation of the SPM in the terminology of (Singh et al., 2021) as a heavyweight model because it has a relatively larger number of parameters that are related to the actual underlying physics of the battery. The SPM is computationally slow and must be run off-board. Also, in Figure 2 we refer to the on-board trained DNN predictor function that replaces the traditional ECM as the lightweight model because it is computationally fast and efficient and does not reflect the underlying physics. We envision the DNN predictor function as the only component of our ROM that resides on-board and in real time. The SPM calculates electrical, chemical, and electrochemical phenomena to predict the battery's performance and lifespan. Thus, providing for a deeper understanding of aging mechanisms, accurately predicting battery performance by considering the material characteristics, mass, heat and energy transfer and the electrode design.

In the initial step of developing the COMSOL® Multiphysics SPM, accurate datasets and manufactures' data was used to develop estimates of the SPM physical parameters. This initial data is not EV operational data and tends to be collected in carefully controlled experimental conditions on individual cells or packs. This initial data having served its purpose for parameter estimation cannot be treated statically because the battery's response changes over the lifetime of the battery. Battery cells and packs cannot be pulled from the EV and lab tested so updating of the SPM must rely on operational data or off-board experiments that mimic EV operation. In our BMS we assume that operational data with sufficient quality and or filtering is available for updating the SPM. The on-board DNN predictor function, however, is updated using real time operation data. As Figure 2 indicates, the DNN updates both the on-board trained DNN predictor function for on-board battery management and the SPM. During normal operation, model parameter update estimation will periodically retrain the onboard DNN predictor function and update the SPM. There is only one DNN but its training inputs come from two sources; the SPM and EV real time operation data. The methodology for obtaining the training data is given in Section 2.4. We do not address the medium by which the components of the BMS acquire the information that we prescribe. Suffice it to say that research and applications are underway into IoT devices, 4G and 5G networks, cloud computing, and artificial intelligence that can augment our model.

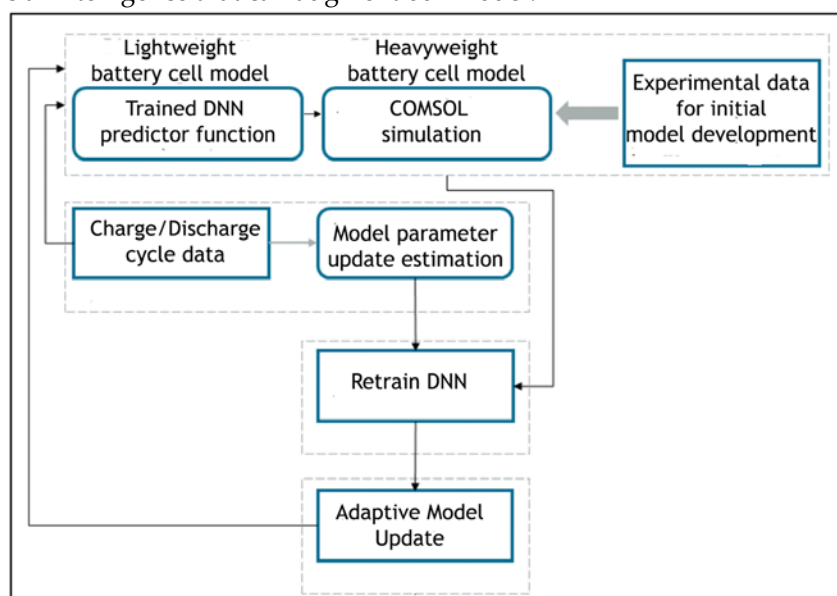


Figure 2. DT configuration showing the role of the DNN and SPM according to the process of Singh et al., 2021.

2.4. Generating the Dataset from the COMSOL® Multiphysics Simulation

To create the simulation dataset over a large range of conditions, we ran the COMSOL® MULTIPHYSICS simulation for temperatures ranging from 253 to 298 K. Each run of the COMSOL® MULTIPHYSICS simulation created a vector of length 100, consisting of the right-hand side of Equation (8). All the vectors were stacked column-wise into an array to form the matrix shown in Figure 3.

$$A = \begin{bmatrix} -V^\theta + \frac{k_B T_1(t_1)}{e} [-\mu - 2\sinh^{-1}(\frac{\bar{I}}{2I_0})] & -V^\theta + \frac{k_B T_n(t_1)}{e} [(-\mu - 2\sinh^{-1}(\frac{\bar{I}}{2I_0}))] \\ -V^\theta + \frac{k_B T_1(t_2)}{e} [-\mu - 2\sinh^{-1}(\frac{\bar{I}}{2I_0})] & \dots & -V^\theta + \frac{k_B T_n(t_2)}{e} [(-\mu - 2\sinh^{-1}(\frac{\bar{I}}{2I_0}))] \\ \dots & \dots & \dots \\ -V^\theta + \frac{k_B T_1(t_{99})}{e} [-\mu - 2\sinh^{-1}(\frac{\bar{I}}{2I_0})] & \dots & -V^\theta + \frac{k_B T_n(t_{99})}{e} [(-\mu - 2\sinh^{-1}(\frac{\bar{I}}{2I_0}))] \\ -V^\theta + \frac{k_B T_1(t_{100})}{e} [-\mu - 2\sinh^{-1}(\frac{\bar{I}}{2I_0})] & -V^\theta + \frac{k_B T_n(t_{100})}{e} [(-\mu - 2\sinh^{-1}(\frac{\bar{I}}{2I_0}))] \end{bmatrix}$$

Figure 3. Schematic of the COMSOL®-simulated data for the Systems A123 2.3 Ah lithium-ion battery.

We then stacked the experimental voltage responses for the A123 Systems 26650 cylindrical 2.3 Ah cell column-wise to form the matrix shown in Figure 4.

$$V = \begin{bmatrix} V_{298_1} & V_{273_1} & V_{253_1} \\ V_{298_2} & V_{273_2} & V_{253_2} \\ \vdots & \vdots & \vdots \\ V_{298_{99}} & V_{273_{99}} & V_{253_{99}} \\ V_{298_{100}} & V_{273_{100}} & V_{273_{100}} \end{bmatrix}$$

Figure 4. Schematic of the experimental voltage data.

As shown in Figure 2, the training parameters for the DNN come from two sources, the SPM and from the EV's on-board sensors and computer. The DNN training parameters from the SPM simulation before and after PCA are shown in Figure 5. The resulting input from the SPM is a rank five array of column vectors. The input parameters for training the DNN to update the on-board DNN predictor function for battery management and updating the SPM are estimates of the same SPM parameters derived from the EV's sensor data. These inputs are also shown in Figure 5. We relate \bar{c} to SOC as a fraction the same as (Zeng & Bazant, 2013) and the temperature (T_s) and the current (I_s) are EV sensor data. The voltage plateau for the half-filled particle ($V_{c(\frac{1}{2})est}$) is scaled to I_s and $\frac{k_B T}{e} \bar{\mu}_{est}$ is estimated from Equation (7) by setting the penalty function coefficient (K) to zero. All of these parameters are defined below in the nomenclature section.

$\{\bar{c}, T_{sur}, T, I, V_{c(\frac{1}{2})}, V_{CH}, \frac{k_B T}{e} \bar{\mu}, V^\theta\} \rightarrow \{\bar{c}, T, I, V_{c(\frac{1}{2})}, \frac{k_B T}{e} \bar{\mu}\}$	$\{SOC, T_s, I_s, V_{c(\frac{1}{2})est}, \frac{k_B T}{e} \bar{\mu}_{est}\}$
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Figure 5. (Left) DNN training input from SPM before and after PCA; **(Right)** Estimated parameters and EV sensor data.

2.4. Developing the DNN

The DNN was coded in the Wolfram language as a Mathematica® notebook. Sufficient detail is given in Figure 6 to duplicate our results given a simulation dataset.

<pre> dataset = Transpose[A]; {u, Σ, v} = SingularValueDecomposition[dataset]; u[[;;, 1 ;; 4]].Σ[[1 ;; 4]].Transpose[v]; dataset = Transpose[Join[{%[[All, 1]], %[[All, 2]], %[[All, 3]], %[[All, 4]], Vdata}]]]; Table[Histogram[dataset[[;;, i]], {i, 5}]; </pre> <p>a</p>
<pre> trainN = Floor[.8 Length[dataset]]; {dataTrain, dataTest} = TakeDrop[RandomSample@dataset, trainN]; Dimensions[dataTrain]; Dimensions[dataTest]; xTrain = dataTrain[[All, 1 ;; 4]]; yTrain = dataTrain[[All, 5]]; xTest = dataTest[[All, 1 ;; 4]]; yTest = dataTest[[All, 5]]; </pre> <p>b</p>
<pre> xTrainStandardized = Transpose@Map[Standardize, Transpose[xTrain]]; Mean[xTrainStandardized]; StandardDeviation[xTrainStandardized]; xTestStandardized = Transpose@Table[(xTest[[All, i]] - Mean[xTrain][[i]]) / StandardDeviation[xTrain][[i]], {i, 4}]; Mean[xTestStandardized] // N; StandardDeviation[xTestStandardized] // N; </pre> <p>c</p>
<pre> net1 = NetChain[{50, ElementwiseLayer["SELU"], 50, ElementwiseLayer["SELU"], 50, ElementwiseLayer["SELU"], 50, ElementwiseLayer["SELU"], 50, 1}, "Input" → 4, "Output" → "Scalar"] net1 = NetInitialize[net, Method → {"Random", "Weights" → .01, "Biases" → 0}] trained = NetTrain[net1, train, LossFunction → MeanAbsoluteLossLayer[], Method → "RMSProp", MaxTrainingRounds → 25] predicted = trained[xTestStandardized] actualPredicted = Table[{yTest[[i]], predicted[[i]], {i, Length[yTest]}} ListPlot[actualPredicted, PlotLabel → "Scatter Plot of actual versus predicted values", AxesLabel → {Actual, Predicted}] Transpose[{yTest, predicted}] </pre> <p>d</p>

Figure 6. Mathematica code for the DNN; (a) We used PCA to determine the minimum number of features for the ROM from the simulation dataset (A, Figure 3) that adequately fit the experimental data (V, Figure 4); (b) Eighty percent of the reduced-order dataset was randomly selected for training and twenty percent for testing; (c) The training data were then normalized to mean = 1.0 and standard deviation of 0.0 and the test data were standardized relative to the mean and standard deviation of the training data; (d) The DNN was constructed and trained using a mean absolute loss layer and RMS propagation for gradient descent.

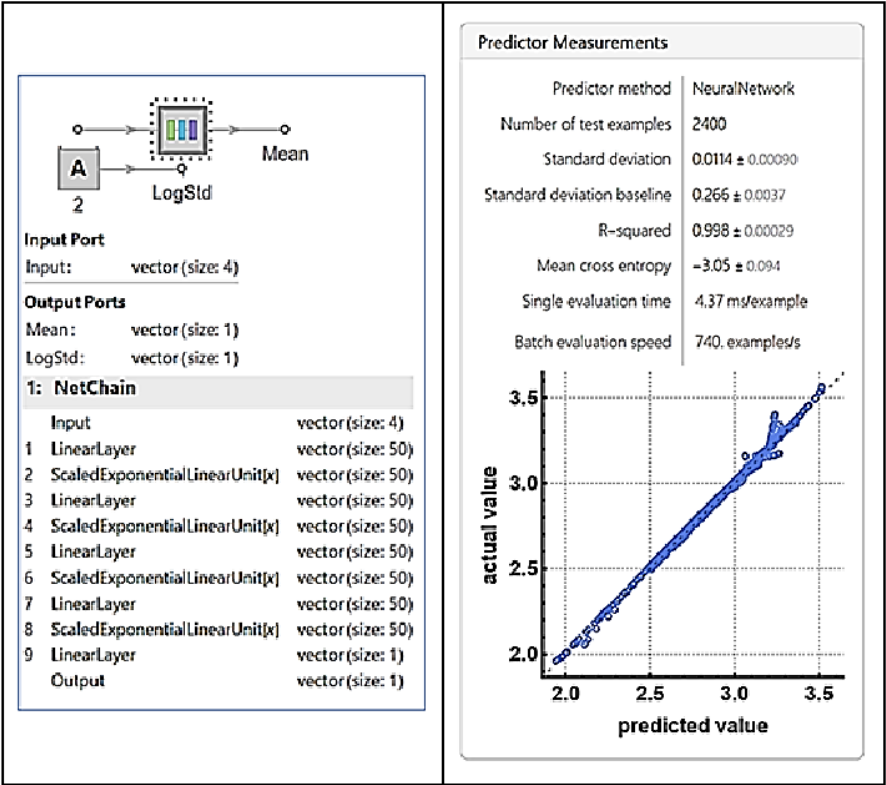


Figure 7. (Left): Layer details of the DNN. **(Right):** Error report for DNN testing and comparison plot (blue circles are actual predictions and the dashed line are perfect predictions)

3. Results and Discussion

As shown in Figure 7, the DNN is a simple feed forward neural network (FNN) with nine hidden layers, consisting of five linear layers and four elementwise scaled exponential linear layers (SELU). The DNN has a rank five vector input for training and a scalar output. The residual plot and error report for the predicted and experimental values are also shown in Figure 7. The important things to note in the error report for the trained DNN predictor function is the regression correlation coefficient and the speed of the processing of examples in this case on a desktop computer CPU. We tested and verified the DNN for a 1 C discharge rate for ambient temperatures ranging from 253 to 298 K, as shown in Figure 9. Also, in Figure 9, the model results are compared to the experimental results for discharge rates ranging from 6.8 to 20.5 C for an ambient temperature of 298 K [65].

We tested the DNN predictor function using the road test: Up Mount Sano in Huntsville, AL [66]. This is an extreme road test and we had to input a moving average of the current data. Figure 10 shows the DNN fit for the road test data for 15, 30, and 60 second update times for the moving average. Longer update time on the order of 60 seconds would be necessary for updating the SPM. If shorter than 15 seconds updating is necessary it would be necessary to modify the DNN at the expense of computation complexity perhaps requiring on-board GPU capability.

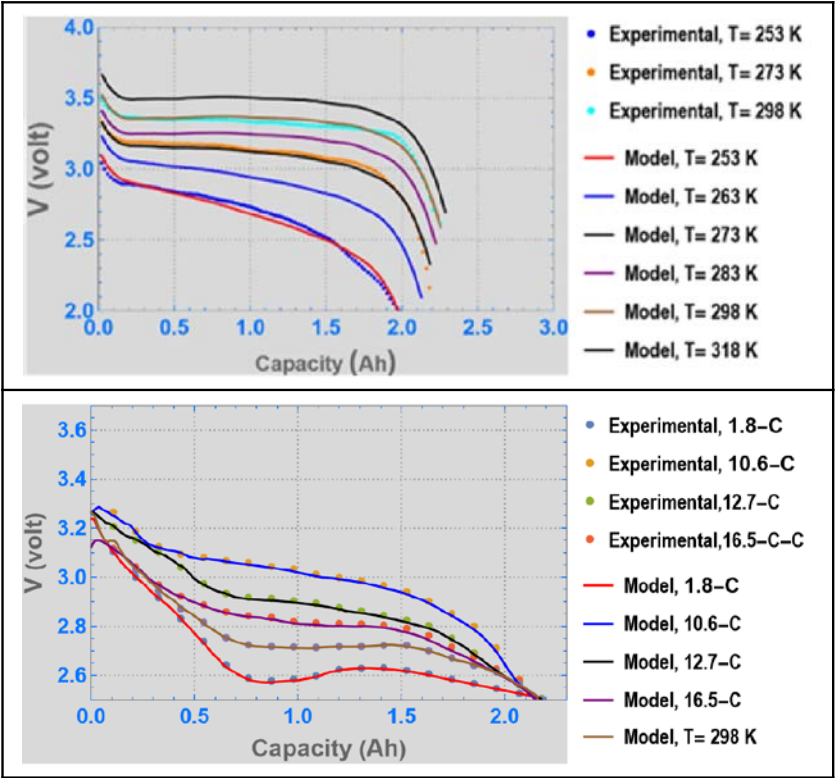


Figure 9. (Top): Voltage response of DNN for temperature range 253–298 K and 1 C discharge. **(Bottom):** Thermal model validation for fast discharge at ambient temperature of 298 K.

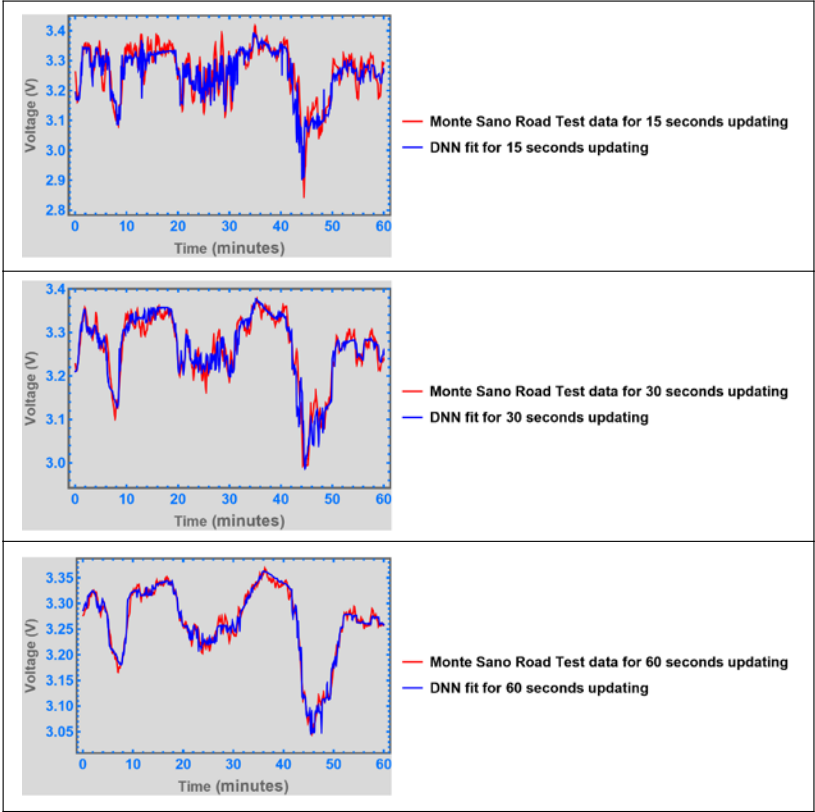


Figure 10. Onboard DNN predictor function testing for road test: “Up Mount Sano in Huntsville, AL” for 15, 30 and 60 second updating.

4. Conclusions

Our ROM was motivated by (Zeng & Bazant, 2013) and the spatial mass transfer elements of our 3-D COMSOL® Multiphysics finite element simulation are similar to their 1-D isotropic and isothermal solution. The major enhancement of our single-particle model (SPM) is that it is a multiphysics, thermal model that fully couples the battery cell’s heat transfer model. The battery cell level thermal model is a robust macro enthalpy balance for the battery despite the simplifying assumptions. The OCV in Equation (5) provides an approximation of the voltage profile of the limiting reversible cycle and allows for an explicit estimation of the cell’s temperature profile based on the updated voltage. We relate the SPM simulation statistically to the battery cell voltage and temperature responses. Specifically, the plateauing effect of the battery’s voltage response at higher ambient temperatures and the apparent diffusion-controlled behavior at lower temperatures are related by statistical inference to the SPM. The ROM was realized by subjecting the raw simulation results from the COMSOL® Multiphysics simulation data to PCA to determine the lowest-order simulation dataset capable of fitting the experimental data. The reduced-order dataset was then fitted using DNN regression. The model was validated for discharge rates ranging from 1 C to 20.5 C and for ambient temperatures ranging from 253 to 298 K by comparison with the manufacturer’s experimental data.

Aside from the SPM providing a robust thermal model the other significant outcome of our model has to do with the statistical inference we drew from the SPM for the initial training of the DNN and then retraining the DNN to reproduce the drive test data from: “Up Mount Sano in Huntsville, AL” with a simple FNN with only nine hidden layers. We believe this is evidence of a significant correlation between the meso-scale and macro-scale properties. In particular we associate this to the strong correlation between the single particle voltage for the half-filled particle and the macro cell voltage at 50% SOC where the chemical potential is at or near zero.

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Nomenclature

A_b	Battery surface area	Parameter	$4.12 \times 10^{-2} \text{ [m}^2\text{]}$
c	Concentration	Variable	$\text{[mol m}^{-3}\text{]}$
\bar{c}	Dimensionless concentration	Variable	$\frac{c}{c_m}$
c_m	Maximum concentration	Parameter	$1.379 \times 10^{28} \text{ [m}^{-3}\text{]}$
c_p	Specific heat coefficient	Parameter	$825 \text{ J [kg}^{-1} \text{ K}^{-1}\text{]}$
h_c	Convection heat transfer coefficient	Parameter	$5.0 \text{ [W m}^{-2} \text{ K}^{-1}\text{]}$
I	Current	Variable	[A]
\bar{I}	Dimensionless current	Variable	$\frac{Rp}{c_m n e D_0} I$
I_0	Current density	Parameter	$1.6 \times 10^{-4} \text{ [A m}^{-2}\text{]}$

\bar{I}_0	Dimensionless current density	Parameter	$\frac{R_p}{c_m n e D_0} I_0$
I_s	Current from online sensors	Variable	[A]
k_b	Boltzmann constant	Constant	3.13×10^9 [eV K ⁻¹]
q	Radial flux	Variable	[s ⁻¹ m ⁻²]
\bar{q}	Dimensionless radial flux	Variable	$\frac{R_p}{c_m D_0} q$
$q_{surr.}$	Heat loss to surroundings	Variable	[W]
R_p	Particle diameter	Parameter	1×10^{-7} [m]
T	Temperature	Variable	[K]
$T_{surr.}$	Surroundings temperature	Parameter	(253–298) [K]
T_s	Temperature from on-board sensor	Variable	[K]
V	Voltage	Variable	[V]
\bar{V}	Dimensionless voltage	Variable	$\frac{e V}{k_b T}$
$V_{c(\frac{1}{2})}$	Single particle voltage for half-filled particle	Variable	[V]
$V_{c(\frac{1}{2})est}$	Cell voltage at 50 % SOC	Variable	[V]
V^θ	Reference voltage	Constant	3.42 [V]
V_{CH}	Voltage simulated by the SPM	Variable	[V]
\bar{V}^θ	Dimensionless reference voltage	Constant	$\frac{e V^\theta}{k_b T}$
α	Electron transfer symmetry factor	Parameter	0.5
η	Activation potential	Variable	-
$\bar{\eta}$	Dimensionless activation potential	Variable	$\frac{e}{k_b T} \eta$
$\Delta\phi_{eq}$	Near-equilibrium voltage	Variable	[V]
μ	Chemical potential	Variable	[eV]
$\bar{\mu}$	Dimensionless chemical potential	Variable	$\frac{\mu}{k_b T}$
$\frac{k_b T}{e} \mu$	Potential energy term from the SPM	Variable	-
$\frac{k_b T}{e} \bar{\mu}_{est}$	Potential energy estimated from Equation with 7	Variable	-
V_{oc}	Open circuit voltage		
v	Battery volume	Parameter	3.42×10^{-5} [m ³]
ϱ	Battery density	Parameter	1824 [kg m ⁻³]
Ω	Enthalpy of mixing	Parameter	0.115 [eV]
$\bar{\Omega}$	Dimensionless enthalpy of mixing	Parameter	$\frac{\Omega}{k_b T}$

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