

# On duality principles and related convex dual formulations suitable for local and global non-convex variational optimization

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## Abstract

This article develops duality principles and related convex dual formulations suitable for the local and global optimization of non-convex primal formulations for a large class of models in physics and engineering. The results are based on standard tools of functional analysis, calculus of variations and duality theory. In particular, we develop applications to a Ginzburg-Landau type equation.

**Key words:** Convex dual variational formulation, duality principle for non-convex local primal optimization, Ginzburg-Landau type equation

**MSC** 49N15

## 1 Introduction

In this article we establish a duality principle and a related convex dual formulation suitable for the local optimization of the primal formulation for a large class of models in non-convex optimization.

The main duality principle is applied to the Ginzburg-Landau system in superconductivity in the absence of a magnetic field.

Such results are based on the works of J.J. Telega and W.R. Bielski [2, 3, 13, 14] and on a D.C. optimization approach developed in Toland [15].

About the other references, details on the Sobolev spaces involved are found in [1]. Related results on convex analysis and duality theory are addressed in [9, 5, 6, 7, 12]. Finally, similar models on the superconductivity physics may be found in [4, 11].

**Remark 1.1.** *It is worth highlighting, we may generically denote*

$$\int_{\Omega} [(-\gamma \nabla^2 + K I_d)^{-1} v^*] v^* \, dx$$

simply by

$$\int_{\Omega} \frac{(v^*)^2}{-\gamma \nabla^2 + K} dx,$$

where  $I_d$  denotes a concerning identity operator.

Other similar notations may be used along this text as their indicated meaning are sufficiently clear.

Finally,  $\nabla^2$  denotes the Laplace operator and for real constants  $K_2 > 0$  and  $K_1 > 0$ , the notation  $K_2 \gg K_1$  means that  $K_2 > 0$  is much larger than  $K_1 > 0$ .

At this point we start to describe the primal and dual variational formulations.

Let  $\Omega \subset \mathbb{R}^3$  be an open, bounded, connected set with a regular (Lipschitzian) boundary denoted by  $\partial\Omega$ .

For the primal formulation, consider a functional  $J : V \rightarrow \mathbb{R}$  where

$$\begin{aligned} J(u) = & \frac{\gamma}{2} \int_{\Omega} \nabla u \cdot \nabla u dx \\ & + \frac{\alpha}{2} \int_{\Omega} (u^2 - \beta)^2 dx - \langle u, f \rangle_{L^2}. \end{aligned} \quad (1)$$

Here  $\gamma > 0$ ,  $\alpha > 0$ ,  $\beta > 0$  and  $f \in L^2(\Omega) \cap L^\infty(\Omega)$ .

Moreover,  $V = W_0^{1,2}(\Omega)$  and we denote  $Y = Y^* = L^2(\Omega)$ .

At this point we define  $F_1 : V \times Y \rightarrow \mathbb{R}$ ,  $F_2 : V \rightarrow \mathbb{R}$  and  $G : V \times Y \rightarrow \mathbb{R}$  by

$$\begin{aligned} F_1(u, v_0^*) = & \frac{\gamma}{2} \int_{\Omega} \nabla u \cdot \nabla u dx - \frac{K}{2} \int_{\Omega} u^2 dx \\ & + \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + 2v_0^* u - f)^2 dx + \frac{K_2}{2} \int_{\Omega} u^2 dx, \end{aligned} \quad (2)$$

$$F_2(u) = \frac{K_2}{2} \int_{\Omega} u^2 dx + \langle u, f \rangle_{L^2},$$

and

$$G(u, v) = \frac{\alpha}{2} \int_{\Omega} (u^2 - \beta + v)^2 dx + \frac{K}{2} \int_{\Omega} u^2 dx.$$

We define also

$$\begin{aligned} J_1(u, v_0^*) &= F_1(u, v_0^*) - F_2(u) + G(u, 0), \\ J(u) &= \frac{\gamma}{2} \int_{\Omega} \nabla u \cdot \nabla u dx + \frac{\alpha}{2} \int_{\Omega} (u^2 - \beta)^2 dx - \langle u, f \rangle_{L^2}, \end{aligned}$$

and  $F_1^* : [Y^*]^3 \rightarrow \mathbb{R}$ ,  $F_2^* : Y^* \rightarrow \mathbb{R}$ , and  $G^* : [Y^*]^2 \rightarrow \mathbb{R}$ , by

$$\begin{aligned} F_1^*(v_2^*, v_1^*, v_0^*) &= \sup_{u \in V} \{ \langle u, v_1^* + v_2^* \rangle_{L^2} - F_1(u, v_0^*) \} \\ &= \frac{1}{2} \int_{\Omega} \frac{(v_1^* + v_2^* + K_1(-\gamma \nabla^2 + 2v_0^*)f)^2}{(-\gamma \nabla^2 - K + K_2 + K_1(-\gamma \nabla^2 + 2v_0^*)^2)} dx \\ &\quad - \frac{K_1}{2} \int_{\Omega} f^2 dx, \end{aligned} \quad (3)$$

$$\begin{aligned}
 F_2^*(v_2^*) &= \sup_{u \in V} \{ \langle u, v_2^* \rangle_{L^2} - F_2(u) \} \\
 &= \frac{1}{2K_2} \int_{\Omega} (v_2^* - f)^2 \, dx,
 \end{aligned} \tag{4}$$

and

$$\begin{aligned}
 G^*(v_1^*, v_0^*) &= \sup_{(u,v) \in V \times Y} \{ \langle u, v_1^* \rangle_{L^2} - \langle v, v_0^* \rangle_{L^2} - G(u, v) \} \\
 &= \frac{1}{2} \int_{\Omega} \frac{(v_1^*)^2}{2v_0^* + K} \, dx + \frac{1}{2\alpha} \int_{\Omega} (v_0^*)^2 \, dx \\
 &\quad + \beta \int_{\Omega} v_0^* \, dx
 \end{aligned} \tag{5}$$

if  $v_0^* \in B^*$  where

$$B^* = \{v_0^* \in Y^* : \|v_0^*\|_{\infty} \leq K/2\}.$$

Define also

$$\begin{aligned}
 V_2 &= \{u \in V : \|u\|_{\infty} \leq K_3\}, \\
 A^+ &= \{u \in V : u \geq 0 \text{ a.e. in } \Omega\}, \\
 V_1 &= V_2 \cap A^+,
 \end{aligned}$$

$$B_2^* = \{v_0^* \in Y^* : -\gamma \nabla^2 - K + K_1(-\gamma \nabla^2 + 2v_0^*)^2 > \mathbf{0}\},$$

$$D_3^* = \{(v_1^*, v_2^*) \in Y^* \times Y^* : -1/\alpha + 4K_1[u(v_1^*, v_2^*, v_0^*)^2] + 100/K_2 \leq \mathbf{0}, \forall v_0^* \in B^*\},$$

where

$$\begin{aligned}
 u(v_2^*, v_0^*) &= \frac{\varphi_1}{\varphi}, \\
 \varphi_1 &= (v_1^* + v_2^* + K_1(-\gamma \nabla^2 + 2v_0^*)f)
 \end{aligned}$$

and

$$\varphi = (-\gamma \nabla^2 - K + K_1(-\gamma \nabla^2 + 2v_0^*)^2 + K_2),$$

$$D^* = \{v_2^* \in Y^* ; \|v_2^*\|_{\infty} < K_4\}$$

$$E^* = \{v_1^* \in Y^* : \|v_1^*\|_{\infty} \leq K_5\},$$

for some  $K_3, K_4, K_5 > 0$  to be specified,

Finally, we also define  $J_1^* : [Y^*]^2 \times B^* \rightarrow \mathbb{R}$ ,

$$J_1^*(v_2^*, v_1^*, v_0^*) = -F_1^*(v_2^*, v_1^*, v_0^*) + F_2^*(v_2^*) - G^*(v_1^*, v_0^*).$$

Assume now  $K_1 = 1/[4(\alpha + \varepsilon)K_3^2]$ ,

$$K_2 \gg K_1 \gg \max\{K_3, K_4, K_5, 1, \gamma, \alpha, \beta\}.$$

Observe that, by direct computation, we may obtain

$$\frac{\partial^2 J_1^*(v_2^*, v_1^*, v_0^*)}{\partial(v_0^*)^2} = -\frac{1}{\alpha} + 4K_1 u(v^*)^2 + \mathcal{O}(1/K_2) < \mathbf{0},$$

for  $v_0^* \in B_3^*$ .

Considering such statements and definitions, we may prove the following theorem.

**Theorem 1.2.** *Let  $(\hat{v}_2^*, \hat{v}_1^*, \hat{v}_0^*) \in ((D^* \times E^*) \cap D_3^*) \times (B_2^* \cap B^*)$  be such that*

$$\delta J_1^*(\hat{v}_2^*, \hat{v}_1^*, \hat{v}_0^*) = \mathbf{0}$$

and  $u_0 \in V_1$ , where

$$u_0 = \frac{\hat{v}_1^* + \hat{v}_2^* + K_1(-\gamma \nabla^2 + 2v_0^*)f}{K_2 - K - \gamma \nabla^2 + K_1(-\gamma \nabla^2 + 2\hat{v}_0^*)^2}.$$

Under such hypotheses, we have

$$\delta J(u_0) = \mathbf{0},$$

so that

$$\begin{aligned} J(u_0) &= \inf_{u \in V_1} \left\{ J(u) + \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + 2\hat{v}_0^* u - f)^2 dx \right\} \\ &= \inf_{v_2^* \in D^*} \left\{ \sup_{(v_1^*, v_0^*) \in E^* \times B^*} J_1^*(v_2^*, v_1^*, v_0^*) \right\} \\ &= J_1^*(\hat{v}_2^*, \hat{v}_1^*, \hat{v}_0^*). \end{aligned} \quad (6)$$

*Proof.* Observe that  $\delta J_1^*(\hat{v}_2^*, \hat{v}_1^*, \hat{v}_0^*) = \mathbf{0}$  so that, since  $(\hat{v}_2^*, \hat{v}_1^*) \in D_3^*$ ,  $\hat{v}_0^* \in B_2^*$  and  $J_1^*$  is quadratic in  $v_2^*$ , we may infer that

$$\begin{aligned} J_1^*(\hat{v}_2^*, \hat{v}_1^*, \hat{v}_0^*) &= \inf_{v_2^* \in Y^*} J_1^*(v_2^*, \hat{v}_1^*, \hat{v}_0^*) \\ &= \sup_{(v_1^*, v_0^*) \in E^* \times B^*} J_1^*(\hat{v}_2^*, v_1^*, v_0^*). \end{aligned} \quad (7)$$

Therefore, from a standard saddle point theorem, we have that

$$J_1^*(\hat{v}_2^*, \hat{v}_1^*, \hat{v}_0^*) = \inf_{v_2^* \in Y^*} \left\{ \sup_{(v_1^*, v_0^*) \in E^* \times B^*} J_1^*(v_2^*, v_1^*, v_0^*) \right\}.$$

Now we are going to show that

$$\delta J(u_0) = \mathbf{0}.$$

From

$$\frac{\partial J_1^*(\hat{v}_2^*, \hat{v}_1^*, \hat{v}_0^*)}{\partial v_2^*} = \mathbf{0},$$

we have

$$-u_0 + \frac{\hat{v}_2^*}{K_2} = 0,$$

and thus

$$\hat{v}_2^* = K_2 u_0.$$

From

$$\frac{\partial J_1^*(\hat{v}_2^*, \hat{v}_1^*, \hat{v}_0^*)}{\partial v_1^*} = \mathbf{0},$$

we obtain

$$-u_0 - \frac{\hat{v}_1^* - f}{2\hat{v}_0^* + K} = 0,$$

and thus

$$\hat{v}_1^* = -2\hat{v}_0^*u_0 - Ku_0 + f.$$

Finally, denoting

$$D = -\gamma\nabla^2u_0 + 2\hat{v}_0^*u_0 - f,$$

from

$$\frac{\partial J_1^*(\hat{v}_2^*, \hat{v}_1^*, \hat{v}_0^*)}{\partial v_0^*} = \mathbf{0},$$

we have

$$-2Du_0 + u_0^2 - \frac{\hat{v}_0^*}{\alpha} - \beta = 0,$$

so that

$$\hat{v}_0^* = \alpha(u_0^2 - \beta - 2Du_0). \quad (8)$$

Observe now that

$$\hat{v}_1^* + \hat{v}_2^* + K_1(-\gamma\nabla^2 + 2\hat{v}_0^*)f = (K_2 - K - \gamma\nabla^2 + K_1(-\gamma\nabla^2 + 2\hat{v}_0^*)^2)u_0$$

so that

$$\begin{aligned} & K_2u_0 - 2\hat{v}_0u_0 - Ku_0 + f \\ &= K_2u_0 - Ku_0 - \gamma\nabla^2u_0 + K_1(-\gamma\nabla^2 + 2\hat{v}_0^*)(-\gamma\nabla^2u_0 + 2\hat{v}_0^*u_0 - f). \end{aligned} \quad (9)$$

The solution for this last system of equations (8) and (9) is obtained through the relations

$$\hat{v}_0^* = \alpha(u_0^2 - \beta)$$

and

$$-\gamma\nabla^2u_0 + 2\hat{v}_0^*u_0 - f = D = 0,$$

so that

$$\delta J(u_0) = -\gamma\nabla^2u_0 + 2\alpha(u_0^2 - \beta)u_0 - f = 0$$

and

$$\delta \left\{ J(u_0) + \frac{K_1}{2} \int_{\Omega} (-\gamma\nabla^2u_0 + 2\hat{v}_0^*u_0 - f)^2 dx \right\} = 0.$$

Moreover, from the Legendre transform properties

$$F_1^*(\hat{v}_2^*, \hat{v}_1^*, \hat{v}_0^*) = \langle u_0, \hat{v}_2^* + \hat{v}_1^* \rangle_{L^2} - F_1(u_0, \hat{v}_0^*),$$

$$F_2^*(\hat{v}_2^*) = \langle u_0, \hat{v}_2^* \rangle_{L^2} - F_2(u_0),$$

$$G^*(\hat{v}_1^*, \hat{v}_0^*) = -\langle u_0, \hat{v}_1^* \rangle_{L^2} - \langle 0, \hat{v}_0^* \rangle_{L^2} - G(u_0, 0),$$

so that

$$\begin{aligned}
 J_1^*(\hat{v}_2^*, \hat{v}_1^*, \hat{v}_0^*) &= -F_1^*(\hat{v}_2^*, \hat{v}_1^*, \hat{v}_0^*) + F_2^*(\hat{v}_2^*) - G^*(\hat{v}_1^*, \hat{v}_0^*) \\
 &= F_1(u_0, \hat{v}_0^*) - F_2(u_0) + G(u_0, 0) \\
 &= J(u_0).
 \end{aligned} \tag{10}$$

Observe now that

$$\begin{aligned}
 J(u_0) &= J_1^*(\hat{v}_2^*, \hat{v}_1^*, \hat{v}_0^*) \\
 &\leq \frac{\gamma}{2} \int_{\Omega} \nabla u \cdot \nabla u \, dx - \frac{K}{2} \int_{\Omega} u^2 \, dx \\
 &\quad + \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + \hat{v}_0^* u - f)^2 \, dx + \langle u, \hat{v}_1^* \rangle_{L^2} - \langle u, f \rangle_{L^2} \\
 &\quad - \frac{1}{2} \int_{\Omega} \frac{(v_1^*)^2}{2v_0^* + K} \, dx - \frac{1}{2\alpha} \int_{\Omega} (v_0^*)^2 \, dx - \beta \int_{\Omega} v_0^* \, dx \\
 &\leq \frac{\gamma}{2} \int_{\Omega} \nabla u \cdot \nabla u \, dx - \langle u, f \rangle_{L^2} \\
 &\quad + \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + \hat{v}_0^* u - f)^2 \, dx \\
 &\quad + \sup_{(v_1^*, v_0^*) \in D^* \times B^*} \left\{ +\langle u, \hat{v}_1^* \rangle_{L^2} - \frac{1}{2} \int_{\Omega} \frac{(v_1^*)^2}{2v_0^* + K} \, dx \right. \\
 &\quad \left. - \frac{1}{2\alpha} \int_{\Omega} (v_0^*)^2 \, dx - \frac{1}{2\alpha} \int_{\Omega} (v_0^*)^2 \, dx - \beta \int_{\Omega} v_0^* \, dx \right\} \\
 &= J(u) + \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + 2\hat{v}_0^* u - f)^2 \, dx,
 \end{aligned} \tag{11}$$

$\forall u \in V_1$ .

Hence, we have got

$$J(u_0) = \inf_{u \in V_1} \left\{ J(u) + \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + 2\hat{v}_0^* u - f)^2 \, dx \right\}.$$

Joining the pieces, we have got

$$\begin{aligned}
 J(u_0) &= \inf_{u \in V} \left\{ J(u) + \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + 2\hat{v}_0^* u - f)^2 \, dx \right\} \\
 &= \inf_{v_2^* \in Y^*} \left\{ \sup_{(v_1^*, v_0^*) \in E^* \times (B^* \cap B_r(\hat{v}_0^*))} J_1^*(v_2^*, v_1^*, v_0^*) \right\} \\
 &= J_1^*(\hat{v}_2^*, \hat{v}_1^*, \hat{v}_0^*).
 \end{aligned} \tag{12}$$

The proof is complete.  $\square$

## 2 Another duality principle suitable for a local optimization of the primal formulation

In this section we develop a second duality principle which the dual formulation is concave.

We start by describing the primal formulation.

Let  $\Omega \subset \mathbb{R}^3$  be an open, bounded, connected set with a regular (Lipschitzian) boundary denoted by  $\partial\Omega$ .

For the primal formulation, consider a functional  $J : V \rightarrow \mathbb{R}$  where

$$\begin{aligned} J(u) = & \frac{\gamma}{2} \int_{\Omega} \nabla u \cdot \nabla u \, dx \\ & + \frac{\alpha}{2} \int_{\Omega} (u^2 - \beta)^2 \, dx - \langle u, f \rangle_{L^2}. \end{aligned} \quad (13)$$

Here  $\gamma > 0$ ,  $\alpha > 0$ ,  $\beta > 0$  and  $f \in L^2(\Omega) \cap L^\infty(\Omega)$ .

Moreover,  $V = W_0^{1,2}(\Omega)$  and we denote  $Y = Y^* = L^2(\Omega)$ .

Define the functionals  $F_1 : V \times Y \rightarrow \mathbb{R}$ ,  $F_2 : V \rightarrow \mathbb{R}$  by

$$\begin{aligned} F_1(u, v_0^*) = & \frac{\gamma}{2} \int_{\Omega} \nabla u \cdot \nabla u \, dx + \langle u^2, v_0^* \rangle_{L^2} \\ & - \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + 2v_0^* u - f)^2 \, dx + \frac{K_2}{2} \int_{\Omega} (\nabla^2 u)^2 \, dx \\ & - \langle u, f \rangle_{L^2} - \frac{1}{2\alpha} \int_{\Omega} (v_0^*)^2 \, dx \\ & - \beta \int_{\Omega} v_0^* \, dx, \end{aligned} \quad (14)$$

and

$$F_2(u) = \frac{K_2}{2} \int_{\Omega} (\nabla^2 u)^2 \, dx.$$

We define also  $F_1^* : [Y^*]^2 \rightarrow \mathbb{R}$  and  $F_2^* : Y^* \rightarrow \mathbb{R}$ , by

$$\begin{aligned} & F_1^*(v_2^*, v_0^*) \\ = & \sup_{u \in V} \{ \langle u, v_2^* \rangle_{L^2} - F_1(u, v_0^*) \} \\ = & \frac{1}{2} \int_{\Omega} \frac{(v_2^* + f - K_1(-\gamma \nabla^2 + 2v_0^*)f)^2}{K_2 \nabla^4 - \gamma \nabla^2 + 2v_0^* - K_1(-\gamma \nabla^2 + 2v_0^*)^2} \, dx \\ & + \frac{K_1}{2} \int_{\Omega} f^2 \, dx \end{aligned} \quad (15)$$

and,

$$\begin{aligned} F_2^*(v_2^*) = & \sup_{u \in V} \{ \langle u, v_2^* \rangle_{L^2} - F_2(u) \} \\ = & \frac{1}{2K_2} \int_{\Omega} \frac{(v_2^*)^2}{\nabla^4} \, dx. \end{aligned} \quad (16)$$

Here we denote

$$B^* = \{v_0^* \in Y^* : \|2v_0^*\|_\infty < K/2\},$$

for an appropriate real constant  $K > 0$ .

Furthermore, we define

$$D^* = \{v_2^* \in Y^* : \|v_2^*\|_\infty \leq 5K_2/4\}$$

and  $J_1^* : D^* \times B^* \rightarrow \mathbb{R}$ , by

$$J_1^*(v_2^*, v_0^*) = -F_1^*(v_2^*, v_0^*) + F_2^*(v_2^*).$$

Assuming  $0 < \alpha \ll 1$  (through a re-scaling, if necessary) and

$$K_2 \gg K_1 \gg K \gg \max\{\|f\|_\infty, \alpha, \beta, \gamma, 1\}$$

by directly computing  $\delta^2 J_1^*(v_2^*, v_0^*)$  we may easily obtain that for such specified real constants,  $J_1^*$  is concave in  $(v_2^*, v_0^*)$  on  $D^* \times B^*$ .

## 2.1 The main duality principle and a concerning convex dual formulation

Considering the statements and definitions presented in the previous section, we may prove the following theorem.

**Theorem 2.1.** *Let  $(\hat{v}_2^*, \hat{v}_0^*) \in D^* \times B^*$  be such that*

$$\delta J_1^*(\hat{v}_2^*, \hat{v}_0^*) = \mathbf{0}$$

and  $u_0 \in V$  be such that

$$u_0 = \frac{\partial F_2^*(\hat{v}_2^*)}{\partial v_2^*}.$$

Under such hypotheses, we have

$$\delta J(u_0) = \mathbf{0},$$

and

$$\begin{aligned} J(u_0) &= \inf_{u \in V} \left\{ J(u) + \frac{1}{2K_2} \int_{\Omega} \left( \frac{\hat{v}_2^*}{-\nabla^2} - K_2(-\nabla^2 u) \right)^2 dx \right\} \\ &= \sup_{(v_2^*, v_0^*) \in D^* \times B^*} J_1^*(v_2^*, v_0^*) \\ &= J_1^*(\hat{v}_2^*, \hat{v}_0^*). \end{aligned} \tag{17}$$

*Proof.* Observe that  $\delta J_1^*(\hat{v}_2^*, \hat{v}_0^*) = \mathbf{0}$  so that, since  $J_1^*$  is concave in  $(v_2^*, v_0^*)$  on  $D^* \times B^*$ , we obtain

$$J_1^*(\hat{v}_2^*, \hat{v}_0^*) = \sup_{(v_2^*, v_0^*) \in D^* \times B^*} J_1^*(v_2^*, v_0^*).$$

Now we are going to show that

$$\delta J(u_0) = \mathbf{0}.$$

From

$$\frac{\partial J_1^*(\hat{v}_2^*, \hat{v}_0^*)}{\partial v_2^*} = \mathbf{0},$$

and

$$\frac{\partial F_2^*(\hat{v}_2^*)}{\partial v_2^*} = u_0$$

we have

$$-\frac{\partial F_1^*(\hat{v}_2^*, v_0^*)}{\partial v_2^*} + u_0 = \mathbf{0}$$

and

$$\hat{v}_2^* - K_2 \nabla^4 u_0 = \mathbf{0}.$$

Observe now that denoting

$$H(v_2^*, v_0^*, u) = \langle u, v_2^* \rangle_{L^2} - F_1(u, v_0^*),$$

there exists  $\hat{u} \in V$  such that

$$\frac{\partial H(\hat{v}_2^*, \hat{v}_0^*, \hat{u})}{\partial u} = \mathbf{0},$$

and

$$F_1^*(\hat{v}_2^*, \hat{v}_0^*) = H(\hat{v}_2^*, \hat{v}_0^*, \hat{u}),$$

so that

$$\begin{aligned} \frac{\partial F_1^*(\hat{v}_2^*, \hat{v}_0^*)}{\partial v_2^*} &= \frac{\partial H(\hat{v}_2^*, \hat{v}_0^*, \hat{u})}{\partial v_2^*} \\ &+ \frac{\partial H(\hat{v}_2^*, \hat{v}_0^*, \hat{u})}{\partial u} \frac{\partial \hat{u}}{\partial v_2^*} \\ &= \hat{u}. \end{aligned} \tag{18}$$

Summarizing, we have got

$$u_0 = \frac{\partial F_1^*(\hat{v}_2^*, \hat{v}_0^*)}{\partial v_2^*} = \hat{u}.$$

Also, denoting

$$A(u_0, \hat{v}_0^*) = -\gamma \nabla^2 u_0 + 2\hat{v}_0^* u_0 - f,$$

from

$$\frac{\partial H(\hat{v}_2^*, \hat{v}_0^*, u_0)}{\partial u} = \mathbf{0},$$

we have

$$-(-\gamma \nabla^2 u_0 + 2\hat{v}_0^* u_0 - f - K_1(-\gamma \nabla^2 + 2\hat{v}_0^*) A(u_0, \hat{v}_0^*) - \hat{v}_2^* + K_2 \nabla^4 u_0) = \mathbf{0},$$

so that

$$A(u_0, \hat{v}_0^*) - K_1(-\gamma \nabla^2 + 2\hat{v}_0^*) A(u_0, \hat{v}_0^*) = \mathbf{0}. \tag{19}$$

From such results, we may infer that

$$A(u_0, \hat{v}_0^*) = -\gamma \nabla^2 u_0 + 2\hat{v}_0^* - f = 0, \text{ in } \Omega.$$

Moreover, from

$$\frac{\partial J_1^*(\hat{v}_2^*, \hat{v}_0^*)}{\partial v_0^*} = \mathbf{0},$$

we have

$$-K_1 A(u_0, \hat{v}_0^*) 2u_0 - \frac{\hat{v}_0^*}{\alpha} + u_0^2 - \beta = \mathbf{0},$$

so that

$$v_0^* = \alpha(u_0^2 - \beta).$$

From such last results we get

$$-\gamma \nabla^2 u_0 + 2\alpha(u_0^2 - \beta)u_0 - f = \mathbf{0},$$

and thus

$$\delta J(u_0) = \mathbf{0}.$$

Furthermore, also from such last results and the Legendre transform properties, we have

$$F_1^*(\hat{v}_2^*, \hat{v}_0^*) = \langle u_0, \hat{v}_2^* \rangle_{L^2} - F_1(u_0, \hat{v}_0^*),$$

$$F_2^*(\hat{v}_2^*) = \langle u_0, \hat{v}_2^* \rangle_{L^2} - F_2(u_0),$$

so that

$$\begin{aligned} & J_1^*(\hat{v}_2^*, \hat{v}_0^*) \\ &= -F_1^*(\hat{v}_2^*, \hat{v}_0^*) + F_2^*(\hat{v}_2^*) \\ &= F_1(u_0, \hat{v}_0^*) - F_2(u_0) \\ &= J(u_0). \end{aligned} \tag{20}$$

Finally, observe that

$$J_1^*(v_2^*, v_0^*) \leq -\langle u, v_2^* \rangle_{L^2} + F_1(u, v_0^*) + F_2^*(v_2^*),$$

$$\forall u \in V, v_2^* \in D^*, v_0^* \in B^*.$$

Thus, we may obtain

$$\begin{aligned} & J_1^*(\hat{v}_2^*, \hat{v}_0^*) \\ & \leq -\langle u, \hat{v}_2^* \rangle_{L^2} + \frac{\gamma}{2} \int_{\Omega} \nabla u \cdot \nabla u \, dx + \langle u^2, \hat{v}_0^* \rangle_{L^2} \\ & \quad - \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + 2\hat{v}_0^* u - f)^2 \, dx + F_2(u) + F_2^*(\hat{v}_2^*) - \langle u, f \rangle_{L^2} \\ & \quad - \frac{1}{2\alpha} \int_{\Omega} (\hat{v}_0^*)^2 \, dx - \beta \int_{\Omega} \hat{v}_0^* \, dx \\ & \leq -\langle u, \hat{v}_2^* \rangle_{L^2} + \frac{\gamma}{2} \int_{\Omega} \nabla u \cdot \nabla u \, dx + \langle u^2, \hat{v}_0^* \rangle_{L^2} \\ & \quad + F_2(u) + F_2^*(\hat{v}_2^*) - \langle u, f \rangle_{L^2} \\ & \quad - \frac{1}{2\alpha} \int_{\Omega} (\hat{v}_0^*)^2 \, dx - \beta \int_{\Omega} \hat{v}_0^* \, dx \\ & \leq \sup_{v_0^* \in Y^*} \left\{ -\langle u, \hat{v}_2^* \rangle_{L^2} + \frac{\gamma}{2} \int_{\Omega} \nabla u \cdot \nabla u \, dx + \langle u^2, \hat{v}_0^* \rangle_{L^2} \right. \\ & \quad \left. + F_2(u) + F_2^*(\hat{v}_2^*) - \langle u, f \rangle_{L^2} \right. \\ & \quad \left. - \frac{1}{2\alpha} \int_{\Omega} (v_0^*)^2 \, dx - \beta \int_{\Omega} v_0^* \, dx \right\} \\ & = J(u) + F_2(u) - \langle u, \hat{v}_2^* \rangle_{L^2} + F_2^*(\hat{v}_2^*), \quad \forall u \in V. \end{aligned} \tag{21}$$

Summarizing, we have got

$$J_1^*(\hat{v}_2^*, \hat{v}_0^*) \leq J(u) + F_2(u) - \langle u, \hat{v}_2^* \rangle_{L^2} + F_2^*(\hat{v}_2^*), \quad \forall u \in V. \quad (22)$$

Joining the pieces, from a concerning convexity in  $u$ , we have got

$$\begin{aligned} J(u_0) &= \inf_{u \in V} \left\{ J(u) + \frac{1}{2K_2} \int_{\Omega} \left( \frac{\hat{v}_2^*}{-\nabla^2} - K_2(-\nabla^2 u) \right)^2 dx \right\} \\ &= \sup_{(v_2^*, v_0^*) \in D^* \times B^*} J_1^*(v_2^*, v_0^*) \\ &= J_1^*(\hat{v}_2^*, \hat{v}_0^*). \end{aligned} \quad (23)$$

The proof is complete.  $\square$

### 3 A convex primal dual for a local optimization of the primal formulation

In this section we develop a convex primal dual formulation corresponding to a non-convex primal formulation.

We start by describing the primal formulation.

Let  $\Omega \subset \mathbb{R}^3$  be an open, bounded, connected set with a regular (Lipschitzian) boundary denoted by  $\partial\Omega$ .

For the primal formulation, consider a functional  $J : V \rightarrow \mathbb{R}$  where

$$\begin{aligned} J(u) &= \frac{\gamma}{2} \int_{\Omega} \nabla u \cdot \nabla u dx \\ &\quad + \frac{\alpha}{2} \int_{\Omega} (u^2 - \beta)^2 dx - \langle u, f \rangle_{L^2}. \end{aligned} \quad (24)$$

Here  $\gamma > 0$ ,  $\alpha > 0$ ,  $\beta > 0$  and  $f \in L^2(\Omega) \cap L^\infty(\Omega)$ .

Moreover,  $V = W_0^{1,2}(\Omega)$  and we denote  $Y = Y^* = L^2(\Omega)$ .

Define the functional  $J_1^* : V \times [Y^*]^2 \rightarrow \mathbb{R}$ , by

$$\begin{aligned} J_1^*(u, v_3^*, v_0^*) &= \frac{\gamma}{2} \int_{\Omega} \nabla u \cdot \nabla u dx + \langle u^2, v_0^* \rangle_{L^2} \\ &\quad + \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + 2v_3^* u - f)^2 dx + \frac{K_1}{2} \int_{\Omega} (v_3^* - \alpha(u^2 - \beta))^2 dx \\ &\quad - \langle u, f \rangle_{L^2} - \frac{1}{2\alpha} \int_{\Omega} (v_0^*)^2 dx \\ &\quad - \beta \int_{\Omega} v_0^* dx. \end{aligned} \quad (25)$$

We define also

$$B^* = \{v_0^* \in Y^* : \|2v_0^*\|_\infty < K/2\},$$

for an appropriate real constant  $K > 0$ .

Furthermore, we define

$$\begin{aligned} D^* &= \{v_3^* \in Y^* : \|v_3^*\|_\infty \leq K_2\} \\ A^+ &= \{u \in V : u \geq 0, \text{ a.e. in } \Omega\}, \\ V_2 &= \{u \in V : \|u\|_\infty \leq K_3\} \end{aligned}$$

for an appropriate real constant  $K_3 > 0$  and

$$V_1 = A^+ \cap V_2.$$

Now observe that denoting  $\varphi_1 = v_3^* - \alpha(u^2 - \beta)$ , we have

$$\begin{aligned} \frac{\partial^2 J_1^*(u, v_3^*, v_0^*)}{\partial u^2} &= K_1(-\gamma \nabla^2 + 2v_3^*)^2 + 4K_1\alpha^2 u^2 \\ &\quad - 2K_1\alpha\varphi_1 - \gamma \nabla^2 + 2v_0^* \end{aligned} \tag{26}$$

and

$$\frac{\partial^2 J_1^*(u, v_3^*, v_0^*)}{\partial (v_3^*)^2} = K_1 + 4K_1 u^2. \tag{27}$$

Denoting  $\varphi = -\gamma \nabla^2 u + 2v_0^* u - f$  we have also that

$$\frac{\partial^2 J_1^*(u, v_3^*, v_0^*)}{\partial u \partial v_3^*} = K_1(-2\gamma \nabla^2 u + 2v_3^* u) + 2K_1\varphi - 2K_1\alpha u. \tag{28}$$

In such a case, we obtain

$$\begin{aligned} \det \left\{ \frac{\partial^2 J_1^*(u, v_3^*, v_0^*)}{\partial u \partial v_3^*} \right\} &= \frac{\partial^2 J_1^*(u, v_3^*, v_0^*)}{\partial (v_3^*)^2} \frac{\partial^2 J_1^*(u, v_3^*, v_0^*)}{\partial u^2} - \left( \frac{\partial^2 J_1^*(u, v_3^*, v_0^*)}{\partial v_3^* \partial u} \right)^2 \\ &= K_1^2(-\gamma \nabla^2 + 2v_3^* + 4\alpha u^2)^2 \\ &\quad + (-\gamma \nabla^2 + 2v_0^*)\mathcal{O}(K_1) \\ &\quad - 4K_1^2\varphi^2 - 4K_1^2\varphi [(-\gamma \nabla^2 u + 2v_0^* u) - 2\alpha u] \\ &\quad - 2K_1^2\alpha\varphi_1(1 + 4u^2). \end{aligned} \tag{29}$$

Observe that at a critical point

$$\varphi = \mathbf{0}$$

so that we may set the non-active restriction

$$C_1^* = \{(u, v_3^*) \in V_1 \times D^* : (\varphi)^2 = (\gamma \nabla^2 u + 2v_3^* u - f)^2 \leq \varepsilon u^2, \text{ in } \Omega\}$$

for a small parameter  $0 < \varepsilon \ll 1$ .

Now we are going to prove that  $C_1^*$  is a convex subset of  $V_1 \times D^*$ .

For a  $K_7 > 0$  observe that

$$\varphi^2 \leq \varepsilon u^2$$

is equivalent to

$$\varphi^2 + K_7 u^2 \leq (K_7 + \varepsilon) u^2,$$

which is equivalent to

$$\sqrt{\varphi^2 + K_7 u^2} - \sqrt{K_7 + \varepsilon} |u| \leq 0, \text{ in } \Omega.$$

Define

$$H(u, v_3^*) = \sqrt{\varphi^2 + K_7 u^2} - \sqrt{K_7 + \varepsilon} |u|.$$

Observe that since for  $(u, v_3^*) \in V_1 \times D^*$  we have  $u f \geq 0$  in  $\Omega$ , we have also that

$$-\sqrt{K_7 + \varepsilon} |u|$$

is convex on  $V_1 \times D^*$ .

Moreover, for  $K_7 > 0$  sufficiently large, the function

$$\sqrt{\varphi^2 + K_7 u^2}$$

is also convex on  $V_1 \times D^*$ .

Summarizing,  $H(u, v_3^*)$  is convex on  $V_1 \times D^*$  so that from such results, we may infer that  $C_1^*$  is a convex set.

On the other hand, at a critical point we have also  $\varphi_1 = \mathbf{0}$ . Now define the non-active constraint

$$C_2^* = \{(u, v_3^*) \in V_1 \times D^* : \varphi_1^2 = (v_3^* - \alpha(u^2 - \beta))^2 \leq \varepsilon, \text{ in } \Omega\}.$$

Similarly as it was made for  $C_1^*$  we may prove that  $C_2^*$  is convex in  $V_1 \times D^*$ .

For a function (or indeed an operator or matrix of functions in a more general context)  $M_1$  to be specified define

$$C_3^* = \{(u, v_3^*) \in V_1 \times D^* : \sqrt{4\alpha} |u| \geq \sqrt{|M_1 + \gamma \nabla^2|} \text{ and } 2v_3^* + M_1 \geq \varepsilon_1\},$$

for some appropriate small constant  $\varepsilon_1 > 0$ .

Since for  $(u, v_3^*) \in V_1 \times D^*$  we have  $u f \geq 0$  in  $\Omega$ , it is clear that  $C_3^*$  is convex on  $V_1 \times D^*$ .

Observe that if  $(u, v_3^*) \in C_3^*$ , then

$$4\alpha u^2 \geq M_1 + \gamma \nabla^2$$

and

$$2v_3^* + M_1 \geq \varepsilon_1$$

so that

$$-\gamma \nabla^2 + 2v_3^* + 4\alpha u^2 \geq \varepsilon_1.$$

At this point, we define the convex set  $C^* = C_1^* \cap C_2^* \cap C_3^*$

Finally, observe that for  $0 < \varepsilon \ll \varepsilon_1 \ll 1$ , we have that

$$\det\{\delta_{u, v_3^*}^2 J_1^*(u, v_3^*, v_0^*)\}$$

is positive definite on  $C^* \times B^*$

From such results we may infer that  $J_1^*$  is convex in  $(u, v_3^*)$  and concave in  $v_0^*$  on  $C^* \times B^*$ . With such results in mind, we may prove the following theorem.

**Theorem 3.1.** Let  $(u_0, \hat{v}_3^*, \hat{v}_0^*) \in C^* \times B^*$  be such that

$$\delta J_1^*(u_0, \hat{v}_3^*, \hat{v}_0^*) = \mathbf{0}.$$

Under such hypotheses, we have

$$\delta J(u_0) = -\gamma \nabla^2 u_0 + 2\alpha(u_0^2 - \beta)u_0 - f = \mathbf{0}$$

and

$$\begin{aligned} J(u_0) &= \inf_{u \in V_1} \left\{ J(u) + \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + 2\alpha(u^2 - \beta)u - f)^2 dx \right\} \\ &= \sup_{v_0^* \in B^*} \left\{ \inf_{(u, v_3^*) \in V_1 \times D^*} J_1^*(u, v_3^*, v_0^*) \right\} \\ &= J_1^*(u_0, \hat{v}_3^*, \hat{v}_0^*). \end{aligned} \quad (30)$$

*Proof.* The proof that

$$\delta J(u_0) = \mathbf{0}$$

and

$$J(u_0) = J_1^*(u_0, \hat{v}_3^*, \hat{v}_0^*)$$

may be done similarly as in the previous sections.

Observe that  $J_1^*$  is convex in  $(u, v_3^*)$  and concave in  $v_0^*$  on  $C^* \times B^*$ , where  $C^*$  and  $B^*$  are convex sets.

From such results and Min-Max Theorem, we may infer that

$$J(u_0) = J_1^*(u_0, \hat{v}_3^*, \hat{v}_0^*) = \sup_{v_0^* \in B^*} \left\{ \inf_{(u, v_3^*) \in V_1 \times D^*} J_1^*(u, v_3^*, v_0^*) \right\}.$$

Finally, observe that

$$J_1^*(u_0, \hat{v}_3^*, \hat{v}_0^*) \leq J_1^*(u, v_3^*, \hat{v}_0^*), \forall u \in V_1, v_3^* \in D^*. \quad (31)$$

In particular for  $v_3^* = \alpha(u^2 - \beta)$  we obtain

$$\begin{aligned} J_1^*(u_0, \hat{v}_3^*, \hat{v}_0^*) &\leq \frac{\gamma}{2} \int_{\Omega} \nabla u \cdot \nabla u dx + \langle u^2, \hat{v}_0^* \rangle_{L^2} \\ &\quad - \frac{1}{2\alpha} \int_{\Omega} (\hat{v}_0^*)^2 dx - \beta \int_{\Omega} \hat{v}_0^* dx \\ &\quad + \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + 2\alpha(u^2 - \beta)u - f)^2 dx - \langle u, f \rangle_{L^2} \\ &\leq \sup_{v_0^* \in Y^*} \left\{ \frac{\gamma}{2} \int_{\Omega} \nabla u \cdot \nabla u dx + \langle u^2, v_0^* \rangle_{L^2} \right. \\ &\quad \left. - \frac{1}{2\alpha} \int_{\Omega} (v_0^*)^2 dx - \beta \int_{\Omega} v_0^* dx \right. \\ &\quad \left. + \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + 2\alpha(u^2 - \beta)u - f)^2 dx - \langle u, f \rangle_{L^2} \right\} \\ &= J(u) + \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + 2\alpha(u^2 - \beta)u - f)^2 dx, \forall u \in V_1. \end{aligned} \quad (32)$$

Joining the pieces, we have got

$$\begin{aligned}
 J(u_0) &= \inf_{u \in V_1} \left\{ J(u) + \frac{K_1}{2} \int_{\Omega} (-\gamma \nabla^2 u + 2\alpha(u^2 - \beta)u - f)^2 dx \right\} \\
 &= \sup_{v_0^* \in B^*} \left\{ \inf_{(u, v_3^*) \in V_1 \times D^*} J_1^*(u, v_3^*, v_0^*) \right\} \\
 &= J_1^*(u_0, \hat{v}_3^*, \hat{v}_0^*).
 \end{aligned} \tag{33}$$

The proof is complete.  $\square$

## 4 Conclusion

In this article we have developed convex dual variational formulations suitable for the local optimization of non-convex primal formulations.

It is worth highlighting, the results may be applied to a large class of models in physics and engineering.

We also emphasize the duality principles here presented are applied to a Ginzburg-Landau type equation. In a future research, we intend to extend such results for some models of plates and shells and other models in the elasticity theory.

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