

Article

Reflection of Light as a Mechanical Phenomenon Applied to the Michelson Interferometer with Light from the Sun

Filip Dambi Filipescu* 

Independent Researcher, Surprise AZ, USA

* Correspondence: filipdambi1@gmail.com

Abstract: The Sun is a frame at relative rest in which the light from Sun travels at the emitted speed C , and the Earth travels with the revolving speed V . The reflection of light as a mechanical phenomenon applies to the modified Michelson interferometer employed by Miller in his experiments with light from Sun. Unlike the Tomaschek experiments, which employ light from stars, the fringe shifts in the Miller experiments are predictable. This study derives the light speeds from the Sun entering the interferometer from 6 am to 6 pm and the corresponding fringe shifts confirming that Miller's observed fringe shifts are proof of the Earth's motion in the Sun's frame at relative rest and the Universe.

Keywords: geometrical optics; speed of light; reflection of light; modified Michelson interferometer

1. Introduction

Michelson derives the fringe shift in the space filled with ether; consequently, the speed of light from a source before and after reflection has the same magnitude c . Light reflection as a mechanical phenomenon is studied in a vacuum where the speed of light from the source emitted at the constant c changes after a mirror's reflection as a ball in an elastic collision with a wall. Therefore, the difference between these two approaches is the reflection of light by a mirror.

The studies of reflection of light as mechanical phenomena by Filipescu [1-5] continue by applying the reflection of light to the modified Michelson interferometer [6,7]. References 1 and 2, as well as References 4 and 5, are the same article. These studies [1,2,4,5] based on the elastic collision of balls with a wall in motion at the limit when the mass of balls converges to zero offer the equation

$$C_{ra} = C_s + v_i + v_r. \quad (1)$$

In Equation (1), C_{ra} is the speed of a reflected wavefront of a ray of light, C_s is the wavefront speed from a source or a mirror as a source, v_i is the mirror speed in the opposite direction of the incident wavefront, and v_r is the mirror speed in the direction of the reflected wavefront. All speeds are in the frame at absolute rest. The speeds v_i and v_r are equal and belong to the inertial frame of the mirror at the instance of collision.

Another form of the Equation (1) is

$$C_{ra} = C_s + v \cos a + v \cos b. \quad (2)$$

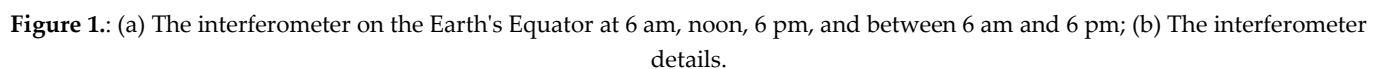
In Equation (2), the speeds $v \cos a$ and $v \cos b$ replace v_i and v_r , respectively. Angle a corresponds to the opposite direction of the incident wavefront; angle b to the direction of the reflected wavefront. These angles are measured counterclockwise

The physics phenomena are observed in the inertial frame of the Sun as in the frame at absolute rest; in other words, the inertial frame of the Sun is a frame at relative rest for Earth that revolves around the Sun with speed v . The frame at absolute rest can be replaced with the Sun's frame at relative rest for the study of the Earth's motion. Therefore, the Sun is a fixed source of light for Earth, and the Earth may be considered an inertial frame in the Sun's frame at relative rest at the time of an experiment. The light from Sun travels at the speed c in any direction in the Sun's frame at relative rest.

2. The interferometer located on the Earth's Equator

The following drawings illustrate a way to bring the light from Sun to a modified Michelson interferometer. For simplicity, the Earth's axis has no tilt.

Figure 1(b) illustrates the Michelson interferometer at 6 am, as seen in Figure 1(a) from the top of the paper. Mirrors M_1 , M_2 , M_3 , and beam splitter M belong to the instrument. Mirror M_3 replaces the instrument's source of light.



The interferometer is located in the plane Oxy and rotates counterclockwise

around Oz with an angle f measured from Ox . Thus, the initial position of the instrument is when direction OM_1 coincides with Ox for $f = 0^\circ$, as shown in Figure 1(b).

In the Sun's frame at relative rest, a vector velocity v in the direction Oz' of the cartesian frame $Ox'y'z'$ is attached permanently to the origin O of $Oxyz$. The counterclockwise Earth's spin changes the position of $Oxyz$ while $Ox'y'z'$ rotates clockwise around the axis Oy' to keep its axes fixed.

The instrument on the Equator belongs to a local Meridian. It can be defined as the Equator's start position when the local Meridian, so the instrument, is at 6 am and frames $Oxyz$ and $Ox'y'z'$ coincide as illustrated in Figure 1(b), with the instrument at the initial position.

From 6 am to 6 pm, the center of the Sun, Earth's orbit, and the Equator's circle are in the same plane as Oxz and $Ox'z'$, even if Oxz changes its directions.

Angle g is measured from Oz to the vector velocity v at O . Figure 1(a) indicates the position of the interferometer at 6 am, which is the start position corresponding to angle $g = 0^\circ$, at a time between 6 am and 6 pm corresponding to an angle g , at noon for $g = 90^\circ$, and at 6 pm for $g = 180^\circ$.

Figure 2(a) depicts the Equator's start position at 6 am. The Earth's spin brings the interferometer at the noon position, as illustrated in Figure 2(b), at the 6 pm position, as presented in Figure 2(c), and at a time between 6 am and 6 pm position, as shown in Figure 3.

2.2. The interferometer located on the Equator at 6 am, noon, and 6 pm

Figure 2(a) illustrates the interferometer at 6 am. Point A belongs to mirror M_4 and to axis Oz . The mirror M_4 , axis Oz , mirror M_3 , and interferometer form a solid structure. M_4 rotates on an axis through point A perpendicular to the paper plane and around the axis Oz as well.

It is considering that the Sun emits parallel rays of light toward Earth. The incident ray from Sun to A and the reflected ray along the axis Oz are in the plane Oxz . The ray from Sun to A is perpendicular to velocity v , as observed from the figures below.

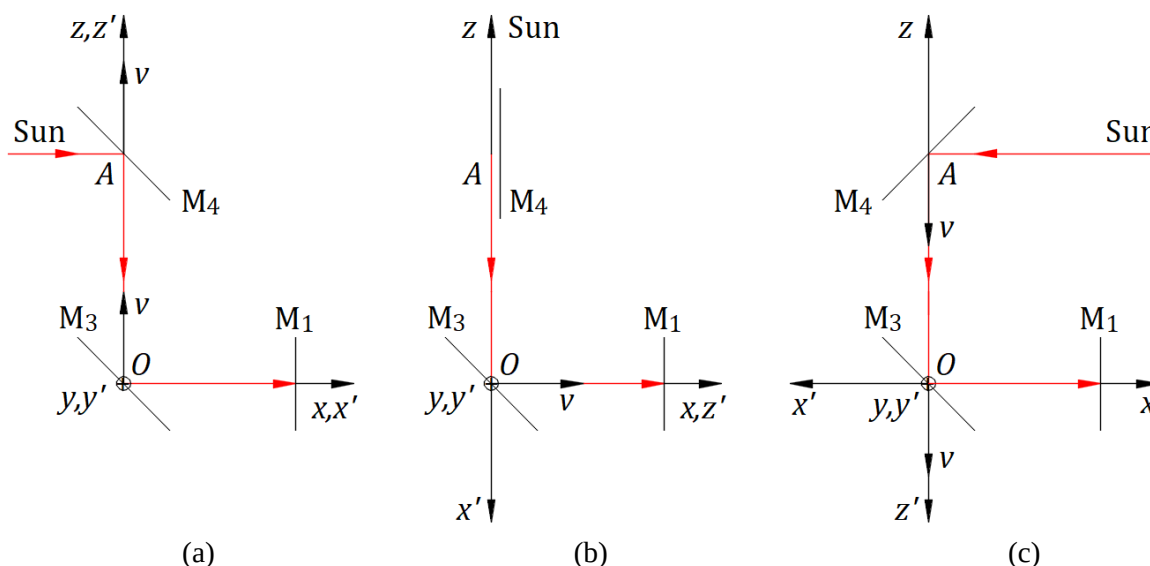


Figure 2.: (a) The interferometer located on the Equator at 6 am; (b) at noon; (c) at 6 pm.

Considering the velocity vector v at A and O , the ray of light at A is reflected toward O with the speed c_{ra} obtainable with Equation (2).

$c_{ra} = c_s + v \cos a + v \cos b = c + v \cos 90^\circ + v \cos 180^\circ = c - v$. The ray is reflected at O along Ox with the entry speed into the interferometer $c_{f=0^\circ} = c_s + v \cos a + v \cos b = c_{ra} + v \cos 0^\circ + v \cos 270^\circ = (c - v) + v = c$ and is the same for any angle f , $c_f = c$. Thus, a shift of zero fringes is expected to be observed by rotating the interferometer 360° .

Within the interferometer, the light rays drift opposite to velocity v . The drift does not affect the fringe shift or may affect it insignificantly.

Figure 2(b) depicts the interferometer at noon. The light from Sun travels parallel to Oz ; therefore, no need for mirror M_4 . At this position, M_4 rotates 180° around Oz .

Considering the velocity vector v at O , the ray of light reflected at O toward M_1 at the initial position of the instrument for $f = 0^\circ$ has the entry speed $c_{f=0^\circ} = c_s + v \cos a + v \cos b = c + v \cos 90^\circ + v \cos 0^\circ = c + v$. By rotating the interferometer counterclockwise around axis Oz at $f = 180^\circ$, the entry speed $c_{f=180^\circ} = c_s + v \cos a + v \cos b = c + v \cos 90^\circ + v \cos 180^\circ = c - v$.

Figure 2(c), shows the instrument at 6 pm. Considering the velocity vector v at A and O , the ray of light at A is reflected toward O with the speed $c_{ra} = c_s + v \cos a + v \cos b = c + v \cos 90^\circ + v \cos 0^\circ = c + v$. The ray is reflected at O along Ox with the entry speed into the interferometer $c_{f=0^\circ} = c_s + v \cos a + v \cos b = c_{ra} + v \cos 180^\circ + v \cos 90^\circ = (c + v) - v = c$ and is the same for any angle f , $c_f = c$. A shift of zero fringes is expected to be observed by rotating the interferometer 360° .

3. The fringe shift in the modified Michelson interferometer located on the Earth's Equator at an angle g

Figure 3(a) presents the instrument at a time between 6 am and 6 pm when Oz' makes an angle g measured from Oz . M_4 rotates around the axis perpendicular to the paper plane. It reflects the rays from Sun at A in the direction AO . The rays belong to a semicircle on the Sun with a diameter equal to that of the Earth. Note that g has an opposite direction to the Earth's spin.

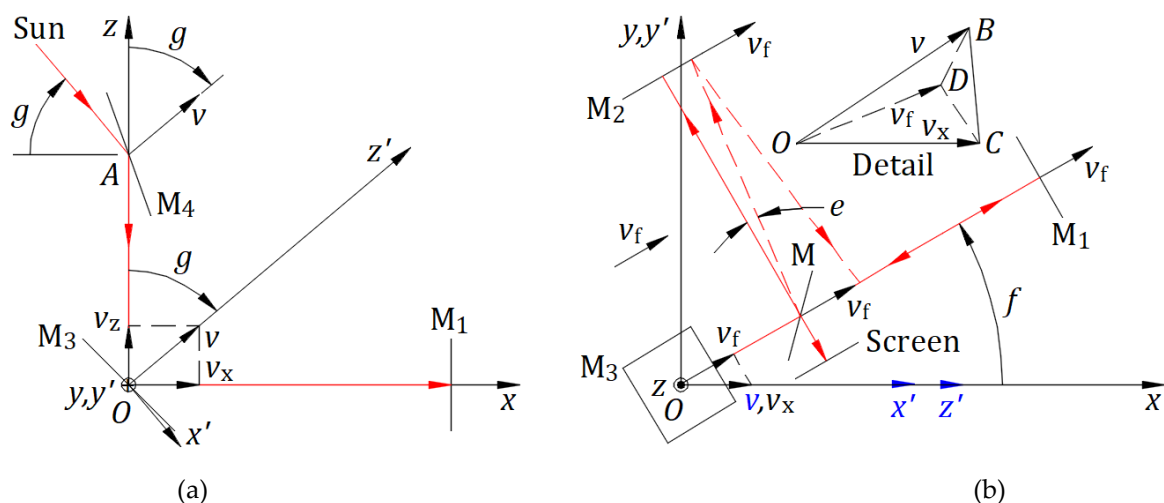


Figure 3.: (a) The interferometer at an angle g ; (b) The fringe shift in the modified Michelson interferometer.

From 6 am to 6 pm for $0^\circ \leq g \leq 180^\circ$, the projection of v on Oz is $v_z = v \cos g$, and on Ox is $v_x = v \sin g$.

Considering the velocity vector v at A , the ray of light is reflected at A toward O with the speed $c_{ra} = c_s + v \cos a + v \cos b = c + v \cos 90^\circ + v \cos(g + 180^\circ) = c - v \cos g$.

At the initial position of the instrument for $f = 0^\circ$, the entry speed from O toward M_1 is $c_{f=0^\circ} = c_s + v \cos a + v \cos b = c_{ra} + v \cos g + v \cos(360^\circ - (90^\circ - g)) = (c - v \cos g) + v \cos g + v \sin g = c + v \sin g$. For $f = 180^\circ$, $c_{f=180^\circ} = c_s + v \cos a + v \cos b = c_{ra} + v \cos g + v \cos(g + 90^\circ) = (c - v \cos g) + v \cos g - v \sin g = c - v \sin g$.

The projection of velocity v on Oz is $v_z = v \cos g$, which is the speed of M_3 and the interferometer in the direction Oz for any angle g . The speed v_z does not affect the fringe shift or may affect it insignificantly.

The projection of velocity v on Ox and direction OM_1 is equal to the projection of velocity v at the initial position of the interferometer, $v_x = v_{f=0^\circ} = v \sin g$.

3.1. The fringe shift in the modified Michelson interferometer with a particular geometry located on Equator at an angle g

The top view of Figure 3(a) with the interferometer rotated by an angle f from Ox is illustrated in Figure 3(b). The Michelson instrument has a particular geometry [1,2] in which the beam splitter makes an angle of 45° , and the two opaque mirrors make the angles of 90° and 0° to the direction of light from the source. The reflected rays by beam splitter M travel perpendicular to M_2 , as illustrated by the ray in the solid red line in Figure 3(b).

In the Detail of Figure 3(b), OB , OC , and OD are along Oz' , Ox , and OM_1 , respectively. OBC and BCD planes and their intersection along BC are perpendicular to the paper plane, and OD is perpendicular to BCD . Therefore, a line from any point along BC to point D is perpendicular to OD . Thus, the projections of v and v_x on OM_1 are equal. The projection of v_x in the direction OM_1 for an angle f is $v_f = v_x \cos f = v \sin g \cos f$, which is the speed of M_3 and the interferometer in the direction OM_1 . Note that v , Ox' , and Oz' are depicted in blue to indicate that they are not in the paper plane.

Considering the velocity vector v at O , the entry speed of the ray of light reflected at O along OM_1 is $c_f = c_s + v \cos a + v \cos b = c_{ra} + v \cos g + v_x \cos f = (c - v \cos g) + v \cos g + v \sin g \cos f = c + v \sin g \cos f = c + v_f$.

Reference 2 yields from the theoretical derivation that the fringe shift is zero for any reasonable speeds c and v , therefore for any speeds $c_f = c + v_f$ and $v_f = v \sin g \cos f$.

The spin speed u of the Earth at the Equator, not shown in Fig. 3(a) and 3(b) at O , has the opposite direction to v_x . Considering the speed u , $v_x = v \sin g - u$ and $v_f = v_x \cos f$. The fringe shift is also zero for other magnitudes of speeds c_f and v_f .

3.2. The fringe shift in the modified Michelson interferometer located on the Equator at an angle g

Michelson predicts a 0.40 fringe shift in the Michelson-Morley experiment [9]. Filipescu predicts 0.40×10^{-4} [3] applying reflection of light as a mechanical phenomenon when reflected rays by beam splitter M travel as illustrated by the ray in the dashed red

line of Figure 3(b) for an angle e equal to the aberration angle, $e = \arcsin(v/c) = 0.0001$ rad, measured from the perpendicular direction to the rays from M_3 .

For a modified Michelson interferometer with the geometry as in the Michelson-Morley experiment, [3,9], the fringe shift can be calculated with Reference 3 in which:

The initial, 90° , 180° , and 270° positions of the interferometer in this article replace the 180° , 270° , initial, and 90° positions in Reference 3, respectively. This is because of the different initial settings in the two articles.

The speeds v_f and $c_f = c + v_f$ corresponding to each angle g in this article replace the speeds v and c in Reference 3, respectively.

The length of the interferometer arm $L = 32$ m as used by Miller.

The fringe shift in this article at the initial position and 90° , 180° , and 270° positions, applying the numerical calculation from Reference 3, are presented in Table 1 for angles g equal to 0° , 30° , 45° , 60° , and 90° . For g equal to 120° , 135° , 150° , and 180° , the fringe shift has the same magnitudes correspondingly.

Table 1. The fringe shift for $L = 32$ m and $e = 0.0001$ at different angles g .

Positions	$g = 0^\circ$	$g = 30^\circ$	$g = 45^\circ$	$g = 60^\circ$	$g = 90^\circ$
0°	0.000E+00	-2.888E-08	0.000E+00	0.000E+00	-2.888E-08
90°	0.000E+00	5.819E-05	8.231E-05	1.008E-04	1.164E-04
180°	0.000E+00	-5.819E-05	-8.231E-05	-1.008E-04	-1.164E-04
270°	0.000E+00	2.888E-08	0.000E+00	0.000E+00	2.888E-08

For $L = 32$ m, Miller expected a 1.12 fringe shift based on Michelson’s derivation and observed 0.08 in 1921 and 0.088 in 1925 [6,7]. Based on the reflection of light as a mechanical phenomenon, Filipescu predicts the fringe shift of 1.164×10^{-4} in Reference 3, and as in Table 1, by employing Reference 3 for numerical calculation of the fringe shift in this article.

The data in Table 1 were calculated with the speeds v_f and $c_f = c + v_f$ for angles $f = 0^\circ$, 90° , 180° , and 270° and angle $e = \arcsin(v/c) = 0.0001$ rad. In Reference 3, angle e applies from zero to the aberration angle and beyond for different magnitudes of speeds v_f and $c_f = c + v_f$. The fringe shift increases with the increase of angle e . For $L = 32$ m and $e = 0.0027$ rad, the predicted fringe shift increases from Table 1 to values comparable to Miller’s observations. Angle e can have the actual magnitude of 0.0027 rad and beyond within the instrument, particularly for $L = 32$ m. The predicted fringe shift is derived with a chosen, known geometry, but the experimental result is according to the actual, unknown geometry, therefore an unknown e .

The fringe shift for $f = 0^\circ$, $f = 90^\circ$, $f = 180^\circ$, and $f = 270^\circ$ positions are presented in Table 2 for $L = 32$ m and $e = 0.0027$ rad at different angles g .

Table 2. The fringe shift for $L = 32$ m and $e = 0.0027$ rad at different angles g .

Positions	$g = 0^\circ$	$g = 30^\circ$	$g = 45^\circ$	$g = 60^\circ$	$g = 90^\circ$
0°	0.000E+00	2.888E-08	0.000E+00	0.000E+00	4.331E-08
90°	0.000E+00	4.242E-02	5.999E-02	7.348E-02	8.485E-02
180°	0.000E+00	-4.242E-02	-5.999E-02	-7.348E-02	-8.485E-02
270°	0.000E+00	-2.888E-08	0.000E+00	0.000E+00	-4.331E-08

The fringe shift in Table 2 explains the fringe shift observed by Miller at Mount Wilson [6,7] in 1921 and 1925. Therefore, Miller's experiments prove the motion of Earth in the Sun's frame at relative rest and the Universe.

Considering the speed u , $v_x = v \sin g - u$ and $v_f = v_x \cos f$. The different magnitudes of $c_f = c + v_f$ and v_f does not make a significant visual difference in the fringe shift. For example, for $g = 90^\circ$ instead of $8.485E - 02$, as in Table 2, the fringe shift is $8.353E - 02$.

4. The fringe shift in the modified Michelson interferometer located on an Earth's Latitude

If the solid structure from the Equator at 6 am, as in Figure 2(a), moves on the local Meridian to the North Pole, the observer sees the instrument as in Figure (4a). From the Equator to the North Pole, the frame $Ox'y'z'$ rotates 90° around Ox' to keep its axes fixed. After rotation, Oz' and the speed v have the direction of Oy .

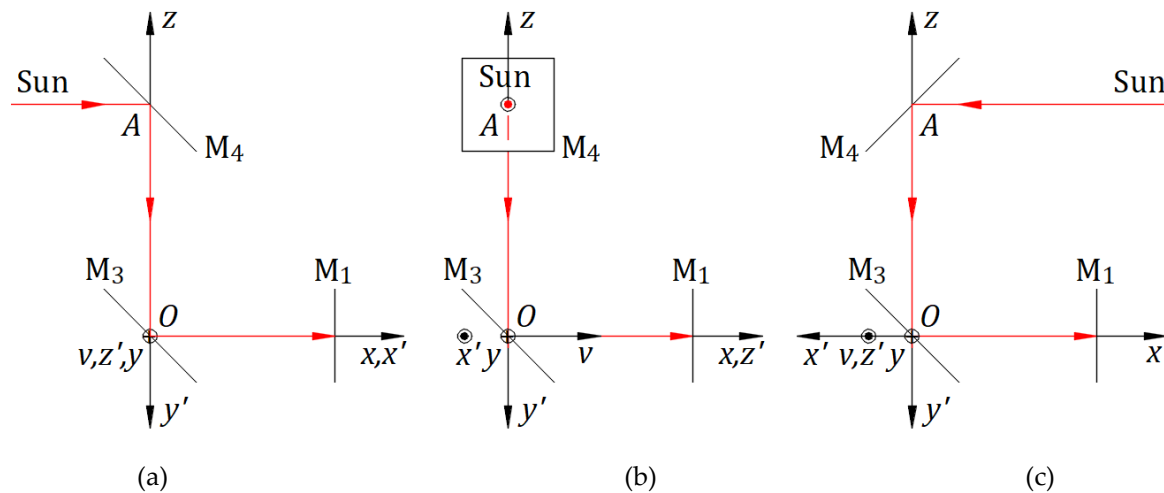


Figure 4.: (a) Interferometer located on the North Pole at 6 am, (b) at noon; (c) at 6 pm.

In the Sun's frame at relative rest, the Earth's spin changes the position of $Oxyz$ while $Ox'y'z'$ rotates around the axis Oy' to keep its axes fixed. Figure (4a) illustrates the instrument at 6 am, Figure (4b) at noon, and Figure (4c) at 6 pm. At the North Pole mirror M_4 rotates clockwise only around Oz .

The right view of Figure 2(a) is as in Figure 5(a). In the Sun's frame at relative rest, moving the solid structure from the Equator to the North Pole, the position of $Oxyz$ changes while $Ox'y'z'$ rotates around the axis Ox' to keep its axes fixed, as visualized in Figure 5(b). Therefore, the velocity v with its axis Oz' rotates in the plane Oyz around Ox' with angle h measured from the axis Oz . For $h = 0^\circ$, the interferometer is at the Equator, and for $h = 90^\circ$ at the North Pole. In this rotation from the Equator to the North Pole, Mirror M_4 stays fixed. It reflects from A to O the rays belonging to a quarter circle on the Sun with a diameter equal to that of the Earth.

Figure 5(b) depicts the interferometer at angle h . The planes Oxz and $Ox'z'$ are perpendicular to the paper plane and make the angle h between them.

The Latitude's start position can be defined when the instrument at the initial position is located at the intersection of the local Meridian with the local Latitude at 6 am and

Oz' is in the plane Oyz making an angle h measured from Oz . The projection of v on Oz is $v_{z,h} = v \cos h$.

The plane $Ox'z'$ is parallel to the Equator's plane and Oxy to the Earth's local surface.

In Figure 5(b), in the Sun's frame at relative rest, on the Latitude at the constant angle h from 6 am to 6 pm, the Earth's spin rotates the frame $Oxyz$ while the velocity v with its axis Oz' rotates around the fixed Oy' with angle g from the plane Oyz at 6 am to Ox at noon and then at 6 pm the velocity v with its axis Oz' is in the opposite of its direction at 6 am. Note that g has an opposite direction to the Earth's spin. For example, angle $g = 0^\circ$ at 6 am, $g = 90^\circ$ at noon, and $g = 180^\circ$ at 6 pm. Thus, on the same Meridian, angle g is the same when the instrument is on different Latitudes, as at the Equator.

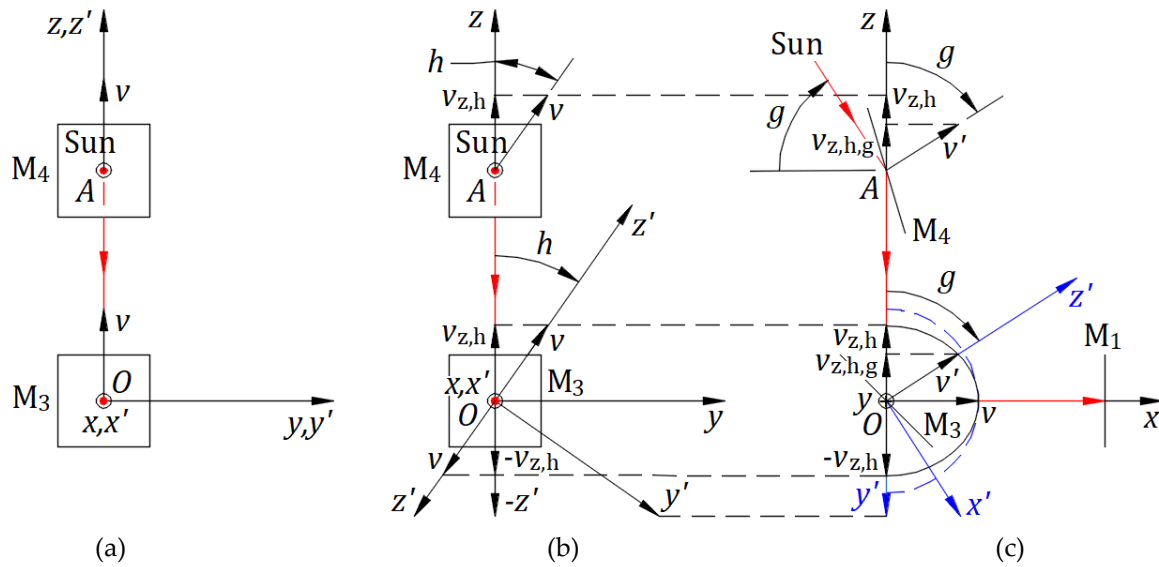


Figure 5.: (a) Interferometer located on the Equator at 6 am; (b) on a Latitude at angle h ; (c) the left view of (b).

In Figure 5(b), the view from the opposite direction of Oy' shows the vector v in the plane $Ox'z'$ from 6 am to 6 pm belonging to a semicircle with the origin at O and the diameter $2v$. The semicircle is on the front of and perpendicular to the paper plane. Thus, the projection of v on Ox is $v_x = v_{f=0^\circ} = v \sin g$, which is the speed of M_3 and the interferometer in the direction Ox , for $g = 0^\circ$ at 6 am to $g = 180^\circ$ at 6 pm, as in Section 3. The speed v_x is independent of angle h .

Suppose the plane $Ox'z'$, with its semicircle viewed from the opposite direction of Oy' , rotates clockwise around Ox' to coincide with the plane Oxz . The semicircle locks like in the plane Oxz depicted in the blue dashed line of Figure 5(c). An identical image is visualized for any angle h .

In Figure 5(b), the view from the opposite direction of Oy sees the projection of the semicircle of the vector v as a semi ellipse in the plane Oxz as illustrated in Figure 5(c). The projection of this semi ellipse on Oz represents the speed $v_{z,h,g}$ of M_3 and interferometer along Oz .

In Figure 5(c), the projection of the velocity v that belongs to Oz' on the plane Oxz is v' . Angle g is measured from the plane Oyz to the plane perpendicular to the paper plane that includes axis Oz' .

Note that Ox' , Oy' , and Oz' are depicted in blue to indicate that they are not in the paper plane. In Figure 5(c), Oz' rotates in the back of the paper plane from $g = 0^\circ$, when it is behind Oz , to the paper plane coinciding with Ox for $g = 90^\circ$, then rotates in the front of the paper plane to above of $-Oz$ for $g = 180^\circ$. Ox' rotates in the front of paper plane from Ox for $g = 0^\circ$ to $-Ox$ for $g = 180^\circ$.

The projections of v and v' on Oz are equal for the same reason given in Detail of Figure 3(b).

At the Latitude's start position, for $g = 0^\circ$, the projection of v on Oz is $v_{z,h} = v_{z,h,g=0^\circ} = v \cos h$ for a constant h . For $0^\circ \leq g \leq 180^\circ$ from 6 am to 6 pm, the projection of v on Oz is $v_{z,h,g} = v_{z,h} \cos g = v \cos h \cos g$, which is the speed of M_3 and the interferometer in the direction Oz for any angle g . The speed $v_{z,h,g}$ does not affect the fringe shift by rotating the instrument 360° around Oz or may affect it insignificantly.

Considering the velocity vector v' at A , the ray of light is reflected at A toward O with the speed $c_{ra} = c_s + v' \cos a + v' \cos b = c + v' \cos 90^\circ + v' \cos(g + 180^\circ) = c - v' \cos g$.

Considering the velocity vector v' at O , the entry speed of the light ray reflected at O in the direction OM_1 for an angle f is $c_f = c_s + v' \cos a + v' \cos b = c_{ra} + v' \cos g + v_x \cos f = (c - v' \cos g) + v' \cos g + v_x \cos f = c + v_f$ and $v_f = v_x \cos f = v \sin g \cos f$, as in Section 3.1.

Introducing the same speeds $c_f = c + v_f$ and $v_f = v \sin g \cos f$ in References 2 and 3, the same magnitudes of the fringe shift are obtained as in Sections 3.1. and 3.2., respectively.

Therefore, the study of the fringe shift at the Equator applies to any Latitude on Earth to the Michelson interferometer with a particular geometry [1,2] and to the Michelson interferometer with a geometry, as in References 3 and 9, as well.

5. Conclusions

Mirror M_4 is adjusted to receive only the rays from Sun that are parallel in the direction from Sun to Earth. In this case, the fringe shift follows a pattern for angle $0^\circ \leq g \leq 180^\circ$, as in Tables 1 and 2. If mirror M_4 reflects rays coming from other directions, the speeds v_f and $c_f = c + v_f$ have different expressions, from these derived here, for identical angles g and f , and the fringe shift may be random for sequential values of angle g .

The adjusting of mirror M_4 , the light paths geometries, the magnitude of angle e , the drift of the light paths because of velocity v_z , even the Doppler effect, and others may affect the fringe shift more or less. We cannot master all these factors to get the fringe shift we want, but the experimental fringe shifts have to be within a reasonable theoretical expectations range.

Reference 3 offers the theoretical expectations of the fringe shift for light reflection as a mechanical phenomenon and the Michelson derivation. In this article, for $L = 32$ m, the fringe shift varies with angle e , as in Tables 1 and 2. For Michelson derivation, the fringe shift is $1.1636 \cong 1.12$ for $e = 0$ rad, 1.1637 for $e = 0.0001$ rad, and 1.1668 for $e = 0.0027$ rad, practically a constant. A different constant is obtained for each magnitude of L ; for $L = 11$ m, as in Michelson derivation, the fringe shift is 0.40 for any angle e . Therefore, for the same length L , Michelson's derivation yields a constant high fringe shift that does not vary with angle e , different from the reflection of light as a

mechanical phenomenon that offers a low fringe shift that varies with angle e . The variation of the fringe shift with angle e makes it possible to explain the results of Miller's experiments.

In conclusion, to explain Miller's observed fringe shift, a study has to offer a reasonable theoretical expectations range and the variation of the fringe shift with angle e .

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