

## Article

# Surface and Electrical Characterization of Bilayers Based on BiFeO<sub>3</sub> and VO<sub>2</sub>

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## ABSTRACT

Thin films of BiFeO<sub>3</sub> (BFO), VO<sub>2</sub> and BFO/VO<sub>2</sub> were grown on SrTiO<sub>3</sub>(100) and Al<sub>2</sub>O<sub>3</sub>(0001) monocrystalline substrates using the RF and DC sputtering techniques. The surface of the films was characterized by profilometry, AFM and XPS. The heterostructures have roughnesses between 0.2 and 16 nm and a grain size between 20 nm and 67 nm. XPS measurements show a higher proportion of the V<sup>4+</sup> and Bi<sup>3+</sup> oxides. In the Fe region, a higher proportion of Fe<sup>3+</sup> is shown in the films. The homogeneous ordering, low roughnesses and the oxidation states on the obtained surface show a good coupling in these films. The I-V curves show ohmic behavior at room temperature and change with increasing temperature. The effect of coupling these materials in a thin film shows the appearance of hysteresis cycles I-V and R-T, typical of materials with high potential in applications such as resistive memories and solar cells.

**Keywords:** Thin films; BiFeO<sub>3</sub>/VO<sub>2</sub>; Solid-solid interface; Surface characterization; Electrical property

## 1. Introduction

Bismuth ferrite (BiFeO<sub>3</sub> (BFO)), is a multiferroic materials exhibiting coexistence of ferroelectricity (FE) and antiferromagnetism (AF) at room temperature [1]–[5]. In thin film form, the FE and AF order parameters of multiferroics could be affected for the phenomenon of magnetostriction [1]–[4] due to the change in the dimensions of the material as a consequence of stresses generated at the contact interface, either with the substrate or with another material. To study this mechanism of coupling at the interface with BFO, the vanadium dioxide (VO<sub>2</sub>) system was used, which has been widely analyzed due to the structural change it undergoes from VO<sub>2</sub> (monoclinic) to VO<sub>2</sub> (tetragonal) at a temperature of 340 K, showing weak ferromagnetic behavior [6], [7]. The combination of these materials in heterostructures gives access to the study of mixed properties and their control.

As far as we can determine, there are currently no reports in the literature that show the growth of BFO thin films over VO<sub>2</sub>. This is mainly due to the incompatibility between the crystalline structures analyzed. However, there are reports on the growth of bilayers with similar crystalline structures or lattice parameters close to those of BFO and VO<sub>2</sub> that presented good coupling, compatibility and interesting physical properties, such as the TiO<sub>2</sub> and de BaTiO<sub>3</sub>. Good compatibility might therefore be expected in the bilayer system based on BFO and VO<sub>2</sub> since the interface effects are predominant and detectable in analogous materials [8]. de la Venta et al, [6], [7], [9], [10] showed the change in the coercivity of thin films of nickel deposited on V<sub>2</sub>O<sub>3</sub> and VO<sub>2</sub>, a change attributed to the structural phase transition of V<sub>2</sub>O<sub>3</sub> and VO<sub>2</sub>. Saebeck et al, [11] studied and compiled information about the coupling between magnetism and the structural phase transition by interfacial tension in different systems of La<sub>0.7</sub>A<sub>0.3</sub>MnO<sub>3</sub> (A = Sr,Ca), Fe<sub>3</sub>O<sub>4</sub>, CoFe<sub>2</sub>O<sub>4</sub>, FeRh, Ni, Fe and Co, using as a base the structural change in the VO<sub>2</sub>, V<sub>2</sub>O<sub>3</sub> and BaTiO<sub>3</sub> compounds, resulting in structures with hybrid properties that manage to modify and/or control the magnetic behavior through effects of tension and inverse magnetostriction. A recent study [12] reports the progress in BFO-based superlattices, which the physical properties of these multifunctional materials have not been explored enough.

On the other hand, Lee et al [13] reported that the BaTiO<sub>3</sub> presents a perovskite-type crystalline structure like that of BiFeO<sub>3</sub> with magnetostriction effects. They report on the modification of the tension of SrRuO<sub>3</sub> and La<sub>0.67</sub>Sr<sub>0.33</sub>MnO<sub>3</sub> films induced by the phase transition of BaTiO<sub>3</sub> substrate, observing the effects that a biaxial stress induces on the electric and magnetic transport of these compounds.

Furthermore, Burbure et al, [14] studied the orientation and phase relationship between TiO<sub>2</sub> films and BaTiO<sub>3</sub> substrates, considering that TiO<sub>2</sub> has structure and lattice parameters of VO<sub>2</sub>, they determined that despite having a high mismatch index, bilayers that present good coupling and epitaxy were grown in different crystalline planes. Feng [15] and Sarkar et al [16], worked on the development of multifunctional nano-heterostructures of BiFeO<sub>3</sub>/TiO<sub>2</sub>, analyzing photo-ferroelectricity, transport, and non-volatile switching resistances. Zhang et al [17], used combinatorial substrate epitaxy (CSE) in 150 samples, for the determination of the phase and orientation relationships in BiFeO<sub>3</sub>/TiO<sub>2</sub> bilayers. They also proposed the configuration scheme structure whereby the coupling between the layers of these materials with mismatch is presented. However, it was shown that the structural coupling of both the anatase phase of TiO<sub>2</sub> (like VO<sub>2</sub> (M)) and the rutile phase (like VO<sub>2</sub> (R)) is ordered.

This paper deals with the growth of BFO thin films previously coated with VO<sub>2</sub>. We report on surface analysis of single material and bilayers films, their growth rate, roughness, and grain size. In addition, an electrical study was carried out.

## 2. Experimental details

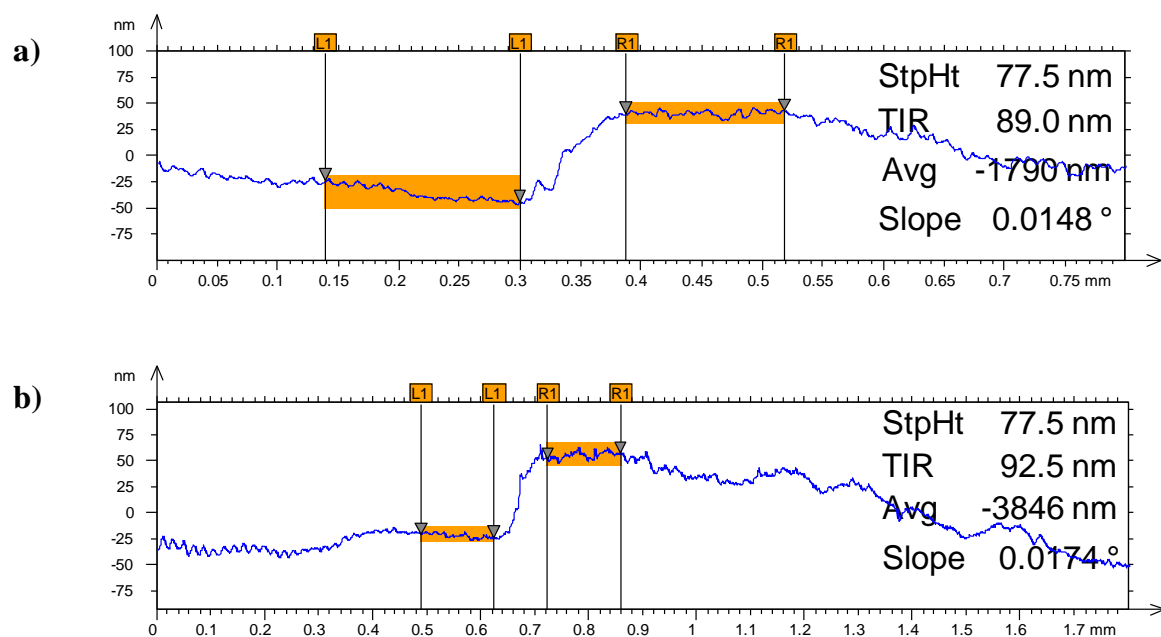
Thin films of BFO, VO<sub>2</sub> and BFO/VO<sub>2</sub> were grown on SrTiO<sub>3</sub>(100) and Al<sub>2</sub>O<sub>3</sub>(0001) substrates using RF and DC sputtering under internal pressure of  $4.5 \times 10^{-4}$  mbar and atmospheres of argon and oxygen. All materials, targets (BFO and metallic vanadium, V, with purity  $\geq 99.99$  %), and substrates, were obtained from Plasmaterials, Inc. The VO<sub>2</sub> films were grown in an O<sub>2</sub>-atmosphere at a pressure of 1.72 mbar, power of 53 W, and substrate temperature of 450° C. In turn, the top BFO layer was deposited in an argon atmosphere ( $2 \times 10^{-1}$  mbar) at a power of 100 W and a substrate temperature of 550° C. After deposition, both the single VO<sub>2</sub> films and the heterostructures were annealed in-situ for 20 minutes in 20 % of O<sub>2</sub> and 80 % of Ar mixed atmosphere ( $p = 2.4 \times 10^{-1}$  mbar) at the same growth temperature of 550 °C. Monolayers of BFO and VO<sub>2</sub> with "steps" were obtained, covering a section of the substrate with an alumina mask, thus allowing the deposition of the material only in one region of the substrate.

Contact profilometry measurements were performed using a KLA Tencor D-120 equipment with a measuring range from 1 nm to 100  $\mu$ m, to determine the thickness of the thin films and their growth rate. The surface morphology of the films was examined with an atomic force microscope (AFM, Omicron SPM1) and the Gwyddion software was used to determine the statistical parameters of roughness. The images were processed with an adjusted linear plane, to eliminate the inclination before the statistical analysis. Average grain diameter was obtained using ImageJ software. This software was used to delineate the perimeter of the grain assuming a circular surface. The elemental composition of the films and the resolution spectra were studied by means of X-ray photoelectron spectroscopy (XPS) using a multi-technique scanning XPS microprobe (PHI Versa Probe II). The source of X-rays used was a monochromatic AlK $\alpha$  lamp (300 W) at a photon energy of  $2.38 \times 10^{-16}$  J. The photoelectron detection was carried out by means of a concentric hemispherical analyzer (CHA) (PHI 10360 Model) operating with an analytical aperture 0.8 m in diameter. The binding energy resolution of the instrument is  $1.12 \times 10^{-19}$  J, with K $\alpha$  radiation unfiltered from an Al ( $h\nu = 1486.6$  eV) anode, with a measurement step of 0.5 eV. The XPS spectra were adjusted using the Multipack and XPSpeak programs. The energy scale was calibrated by assigning 284.8 eV to the C1s peak corresponding to adventitious carbon. Electrical characterization with resistance curves as a function of temperature R(T) was carried out studying the metal-insulator transition behavior of the VO<sub>2</sub> compound and of the BFO/VO<sub>2</sub> system, additional measurements of current versus voltage, I-V, were made by the method of Van der Paw or method of the four points and finally the results obtained were correlated.

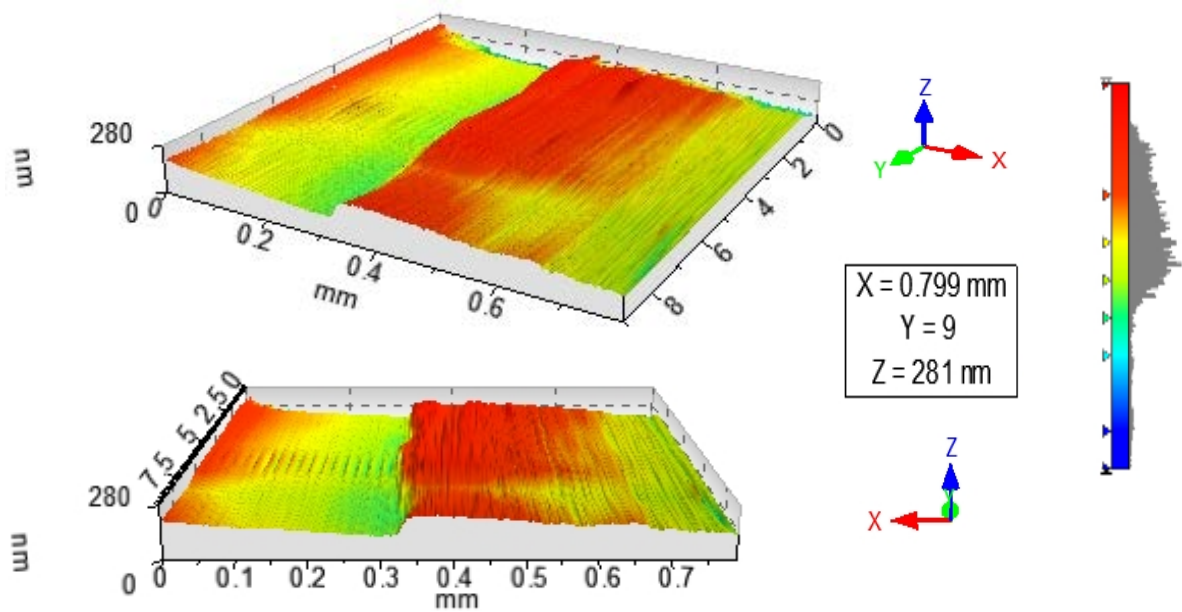
## 3. Results and discussion

### 3.1 Profilometry and AFM studies

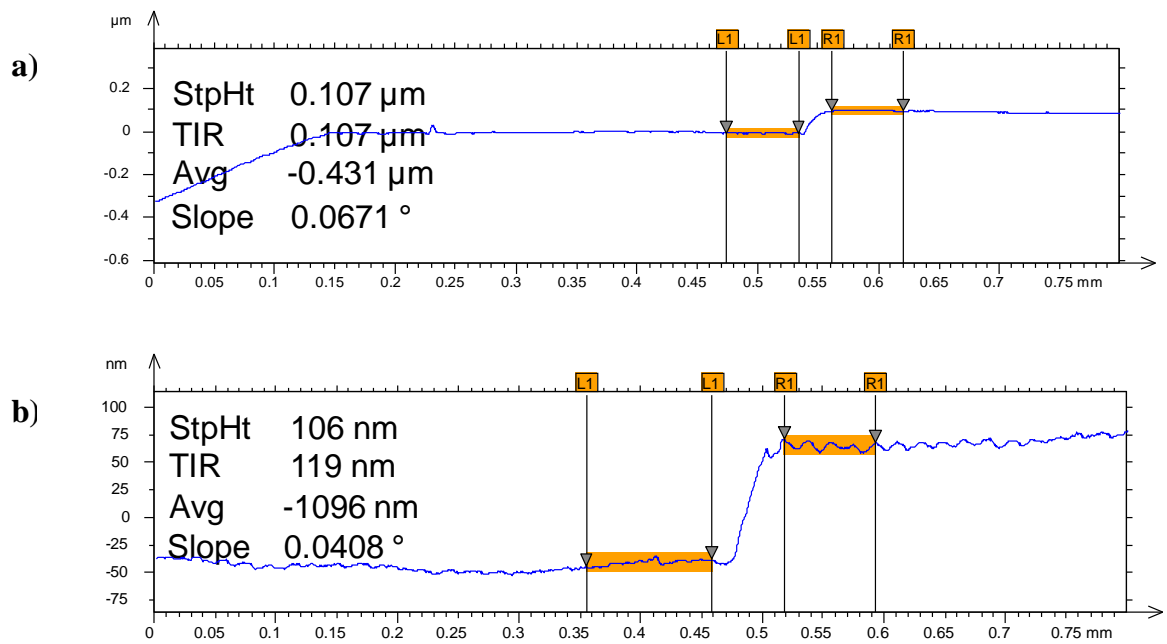
**Figures 1 and 2** show profilometry measurements of the  $\text{VO}_2/\text{Al}_2\text{O}_3$  thin films. In **Figure 1**, the approximate thickness of 77.5 nm provided by the step height parameter (StpHt) indicates a growth rate of 1.3 nm per minute. Instead, in the 3D profile uniform height regions can be seen on moving away from the step, denoting an ordered growth. In the case of  $\text{BFO}/\text{SrTiO}_3$  thin films, see **Figures 3 and 4**, the thickness of the BFO films was 106.5 nm, on average, which indicates a growth rate of 2.4 nm per minute under the BFO growth parameters. In these films the step is more marked, possibly because the mask was much better adhered and did not allow diffusion between film and mask. Although the maximum height Z is lower than in the  $\text{VO}_2$  films, there is less homogeneity and there are several zones with different heights.



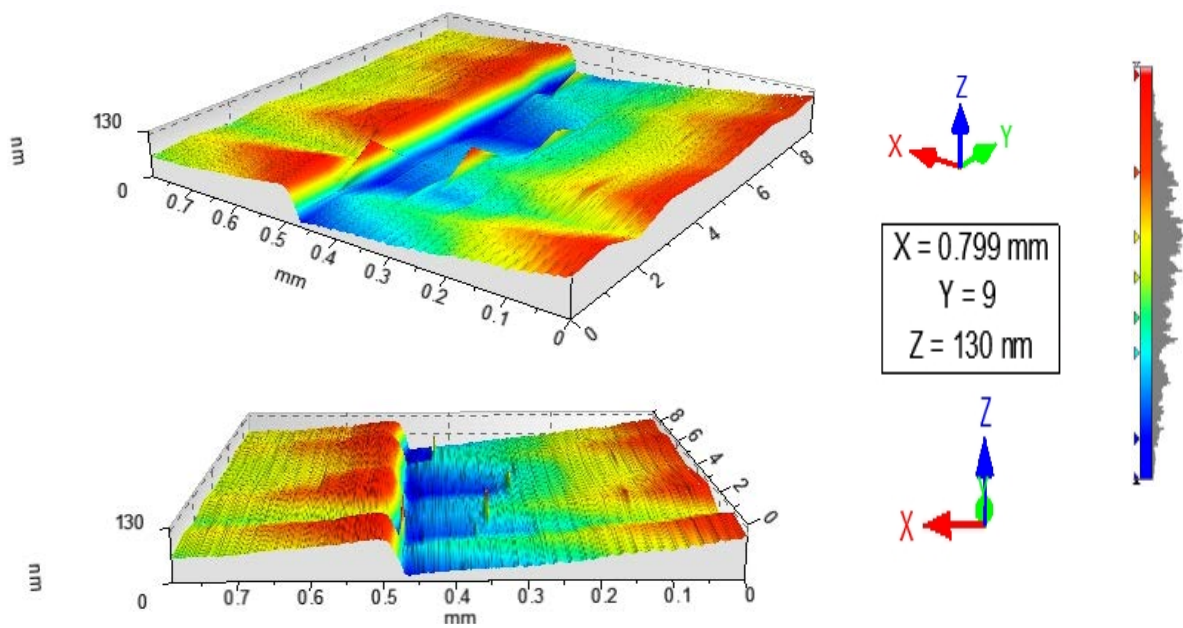
**Figure 1.** a,b) Profilometry measurements at two different points of thin films of  $\text{VO}_2$  on  $\text{Al}_2\text{O}_3$ .



**Figure 2.** 3D image obtained by profilometry of the surface of thin films of VO<sub>2</sub> on Al<sub>2</sub>O<sub>3</sub>.



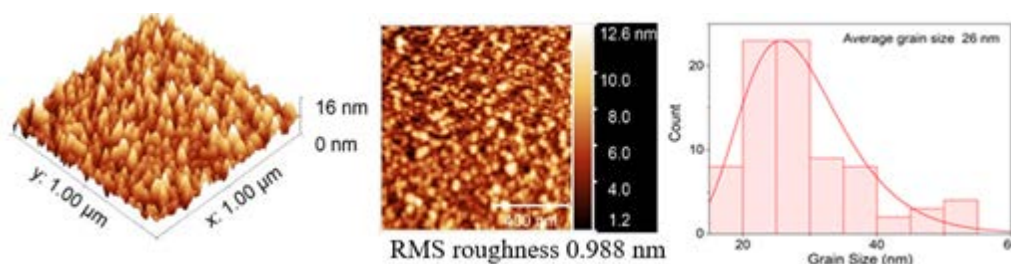
**Figure 3.** a,b) Profilometry measurements at two different points of thin films of BFO on SrTiO<sub>3</sub>.



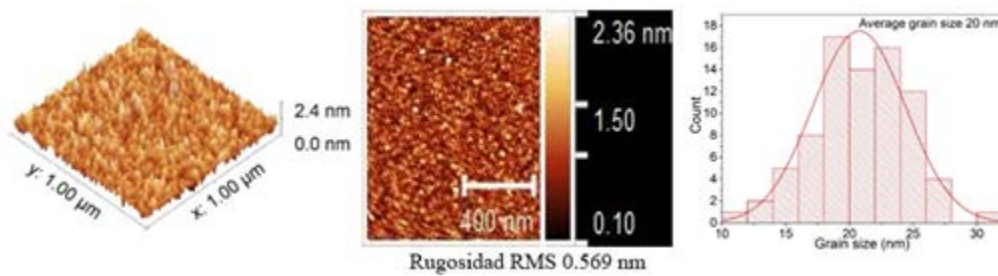
**Figure 4.** 3D image obtained by profilometry of the surface of thin films of BFO on SrTiO<sub>3</sub>.

Surface roughness plays an important role on the effects generated at the interface of the films. For example, in multiferroics systems the roughness can influence the magnetic and electrical properties of these films. Therefore, the surface of the individual layers and the BFO/VO<sub>2</sub> heterostructures is analyzed by AFM, see **Figures 5, 6, and 7** [18]. Good granular ordering and RMS roughnesses between 1 and 5 nm were observed for the individual layers and between 6 and 18 nm for the bilayers, see **Figure 8**. The films exhibited a dense morphology, a continuous surface, and a fine grain microstructure without cracks. The samples with greater presence of high peaks are due to the agglomerations of particles that appear in longer growth times, in the same way the grain size for long growth times is smaller and with more uniformity.

**Figure 5** shows the image of the surface scan of a section of  $1 \times 1 \mu\text{m}^2$  and a histogram with the average grain size, which was 26 nm for a monolayer film of VO<sub>2</sub>. In these monolayers, surface roughness and grain size increase with deposition time. However, **Figure 6** reveals that the grain size of the BFO films is highly homogeneous. No large peaks or valleys were observed. The average grain size is 20 nm and the roughness does not seem to increase significantly with the increase in thickness.



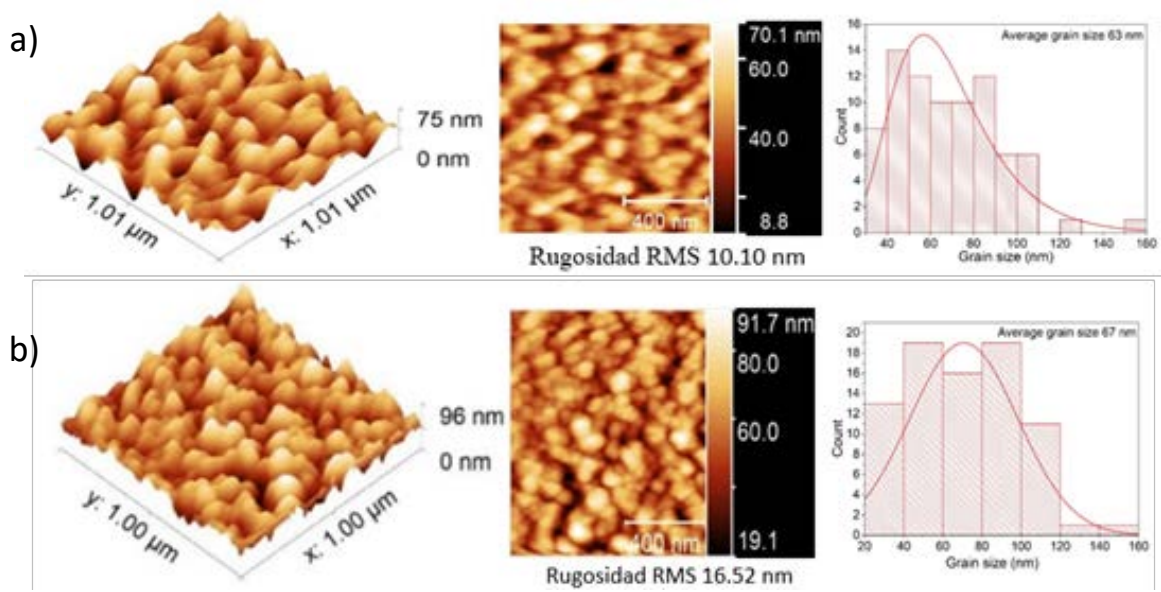
**Figure 5.** AFM measurements and average grain size for VO<sub>2</sub> thin films



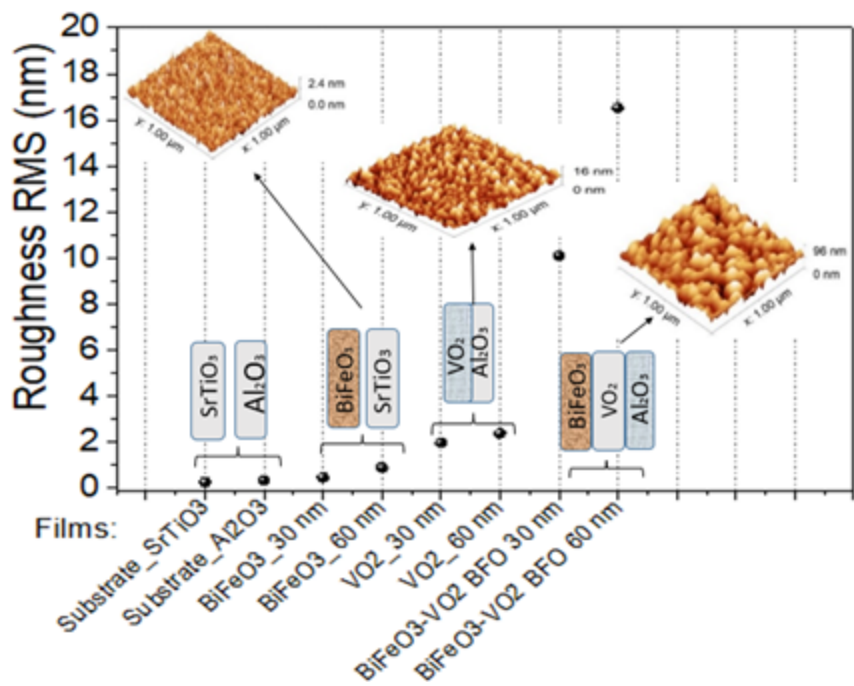
**Figure 6.** AFM measurements and average grain size for BFO thin films

**Figure 7** shows that the grain size of the BFO/VO<sub>2</sub> films is 63 nm and increases to 67 nm when the BFO layer is 30 and 60 nm, respectively. This could be because the grain geometry is more elongated in the bilayers of 30 nm and roughness 10.10 nm than with bilayer of 60 nm and roughness 16.62 nm. This shows us a more elongated grain geometry in the bilayers, especially the bilayers with greater thicknesses. This elongation in the grains influences the roughness which increases with respect to the individual layers, this translates into greater growth by islands and agglomerations of scattered particles with greater height.

Thus, is observed in the literature reports that these films grow layer by layer following the structure of terraces and substrate steps [19]-[23]. Atoms reaching the surface of the substrate are placed in the middle of the step and, from there, the structure grows in 2D form until complete coalescence between the different levels of layer. The atoms deposited on the substrate are more strongly attached to the next layer and so on. Growth with a larger and not so homogeneous grain size can be seen for the top layer of BFO/VO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> with respect to the individual layers [19]. A degree of organization and absence of cracks is observed, indicating good diffusion and coupling between the BFO, VO<sub>2</sub> and substrate layers. These characteristics seen in the films contribute to create additional stress that increases with thickness, favoring the formation of islands in the next layer to be deposited. This stress can determine a shift in the critical transition temperature of VO<sub>2</sub> and therefore in the bilayer [20]–[22].



**Figure 7.** AFM measurements and average grain size for BFO/VO<sub>2</sub> thin films with BFO layer of a) 30 nm and b) 60 nm



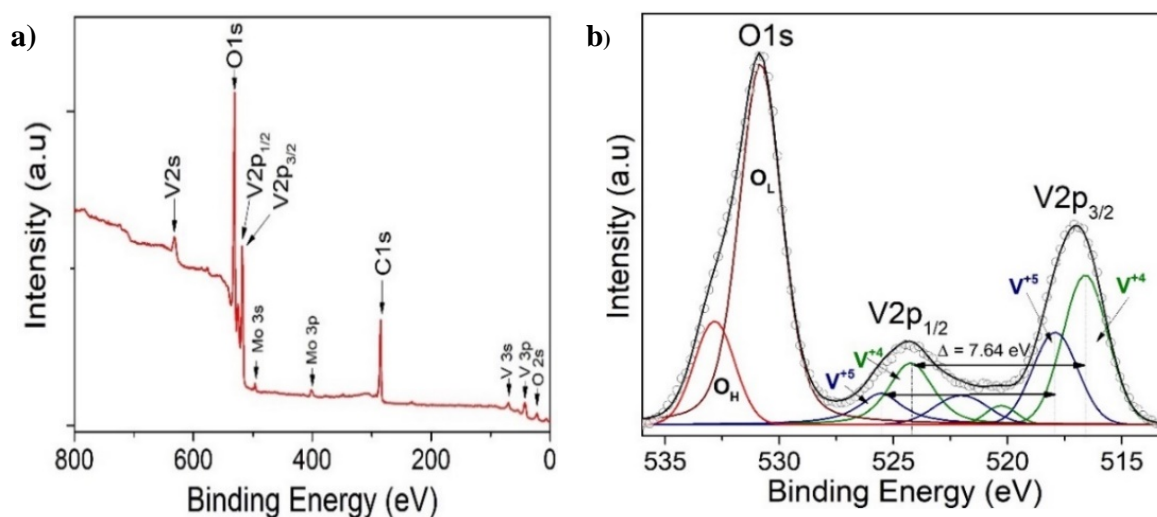
**Figure 8.** Roughness of films based on BFO and VO<sub>2</sub>

**Figure 8** presents a summary of the roughness of the grown films. A semi-exponential upward trend can be observed, showing us how roughness increases in these films as the thickness and number of layers of these two thin-film materials increase.

3.2 Chemical states

The chemical studies have been carried out through XPS analysis to detect the chemical states in these thin films. The analysis of surfaces by XPS is a very valuable resource when determining both quantitatively and qualitatively the composition and chemical state of surfaces. However, the analysis by this technique of first-row transition metals and their oxides and hydroxyls is a challenge due to the complexity of their  $2p$  spectra resulting from asymmetric peaks, the complex division of multiplets and the uncertain and overlapping union of energies [24].

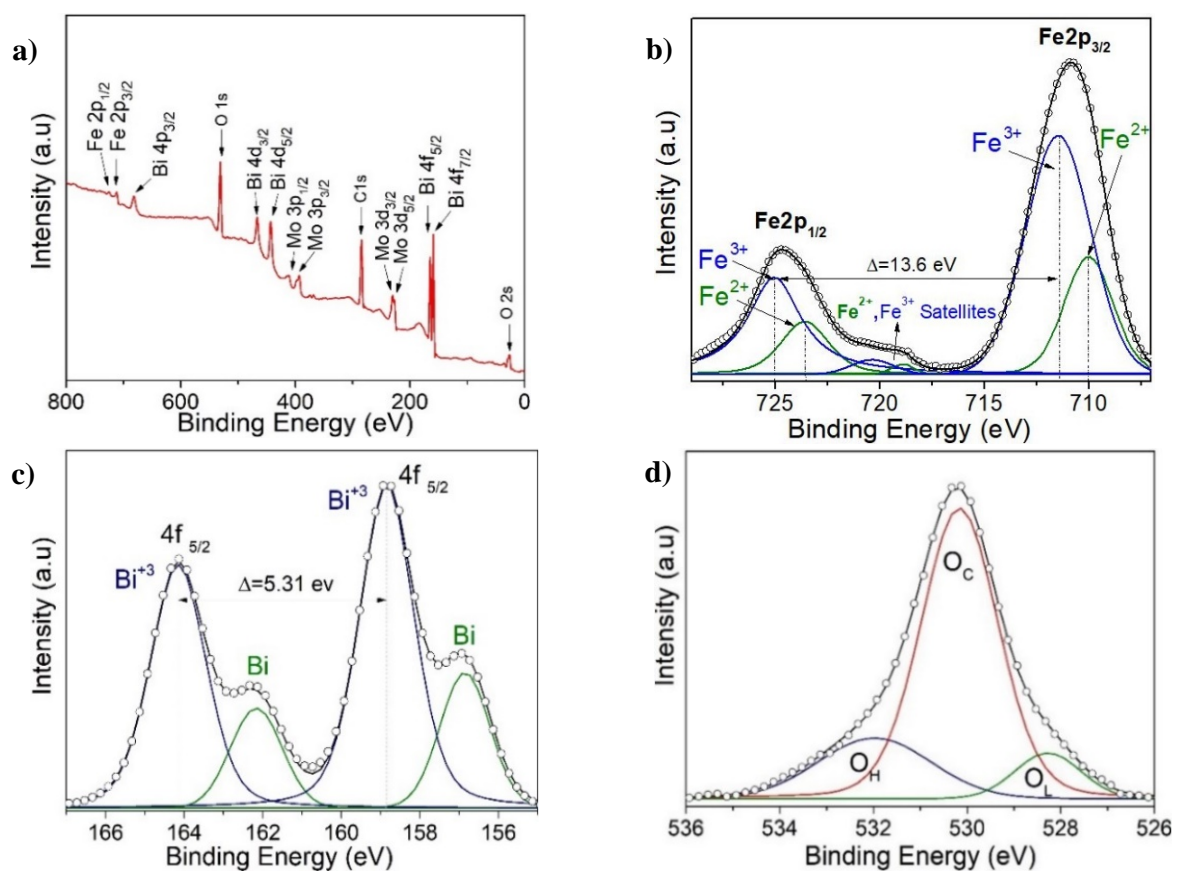
XPS measurements are presented for 4 samples, 2 monolayers ( $\text{VO}_2/\text{Al}_2\text{O}_3$  and  $\text{BFO}/\text{SrTiO}_3$ ) and 2 bilayers ( $\text{BFO}/\text{VO}_2/\text{Al}_2\text{O}_3$ , with BFO layer of 30 and 60 nm), respectively. These measurements were made in broad spectrum and high-resolution and are shown in **Figures 9** to **12**. All spectra were corrected by the binding energy of carbon at 284.8 eV and adjusted with mixed Lorentzian and Gaussian curves. The Mo signal is from the sample holder.



**Figure 9.** Broad (a) and high-resolution (b) XPS spectrum of  $\text{VO}_2$  monolayer

**Figure 9** shows the spectra of  $\text{VO}_2/\text{Al}_2\text{O}_3$  monolayer. **Figure 9a)** can be seen the absence of impurity peaks except for molybdenum (Mo) coming from the sample holder used. However, a high-resolution spectrum, **Figure 9b)**, the respective curve adjustment was obtained for  $\text{V}^{+4}$  peaks (516.6 eV and 524.24 eV),  $\text{V}^{+5}$  peaks (517.94 eV and 525.58 eV),  $\text{O}_\text{L}$  peak (530.84 eV), and  $\text{O}_\text{H}$  peak (532.81 eV), respectively. Here,  $\text{O}_\text{L}$  and  $\text{O}_\text{H}$  peaks correspond to the bonds of oxygen with hydrogen and the crystal lattice of the compound studied and indicates that the films are very susceptible to oxidation in the environment. In addition, a distance between the vanadium doublet of 7.64 eV typical of vanadium oxides is observed, and it is evident that the predominant oxidation state is  $\text{V}^{+4}$ , which corresponds to  $\text{VO}_2$ . On the other hand, **Figure 10** shows the spectra of  $\text{BFO}/\text{SrTiO}_3$  monolayer. It is observed signals from bismuth, iron, oxygen, and molybdenum (from sample

holder), see **Figure 10a**). However, using high-resolution is possible to determine the composition in the thin film (**Figures 10 b**)-**d**). It is observed signals from Fe(2p) doublet (at 710.02 and 723.62 eV) and Bi(4f) doublet (at 156.84 and 162.15 eV), including metallic Bi. Their binding energies differences for each is 13.6 eV and 5.31 eV, respectively, which constitutes a response of iron oxides and bismuth oxides. On the other hand, the high-resolution spectra of oxygen (O1s, see **Figure 10d**) show the presence of the oxides of the lattice (metallic oxides, O<sub>L</sub>) and hydroxyl oxygen (O<sub>H</sub>) in the thin films. Therefore, it is clearly observed that the predominant oxidation states are Fe<sup>3+</sup> and Bi<sup>3+</sup> corresponding to the structure of BiFeO<sub>3</sub>.

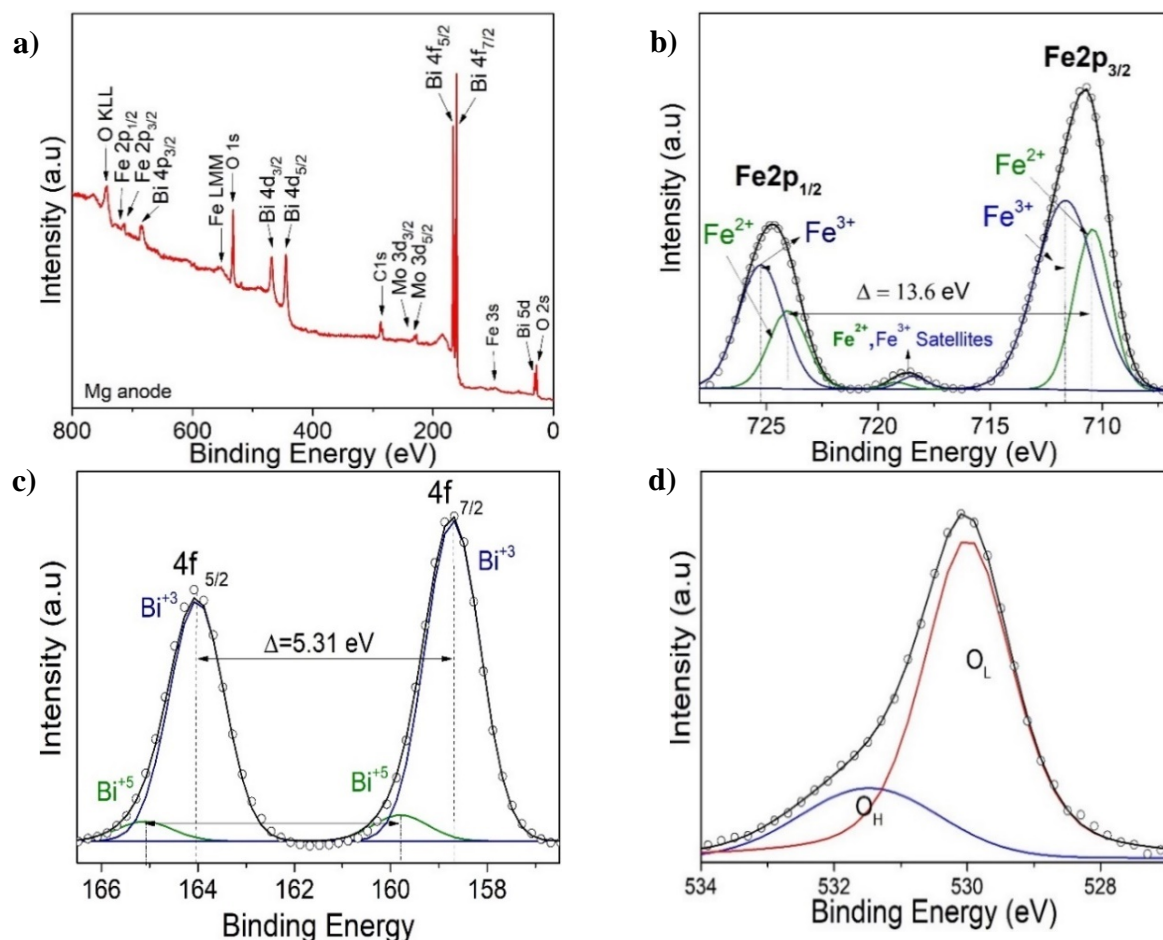


**Figure 10.** Broad (a) and high-resolution (b, c, d) XPS spectra of BFO/SrTiO<sub>3</sub> monolayer

For the bilayer thin films, **Figure 11** and **Figure 12**, we can observe that the Fe(2p) doublet and Bi(4f) doublet binding energies difference of 13.6 eV and 5.31 eV are maintained. Therefore, the BFO thin films with layer of 30 nm exhibits more oxidation states as well as a shift towards higher binding energies. However, an ion bombardment revealed that the amount of Fe<sup>3+</sup> increases in the BFO thin films, unlike the surface where there was a higher presence of the Fe<sup>2+</sup> oxidation state. In the bilayers, there was a small displacement of the maximums towards higher values of binding energy, however, this did not increase the number of oxidation

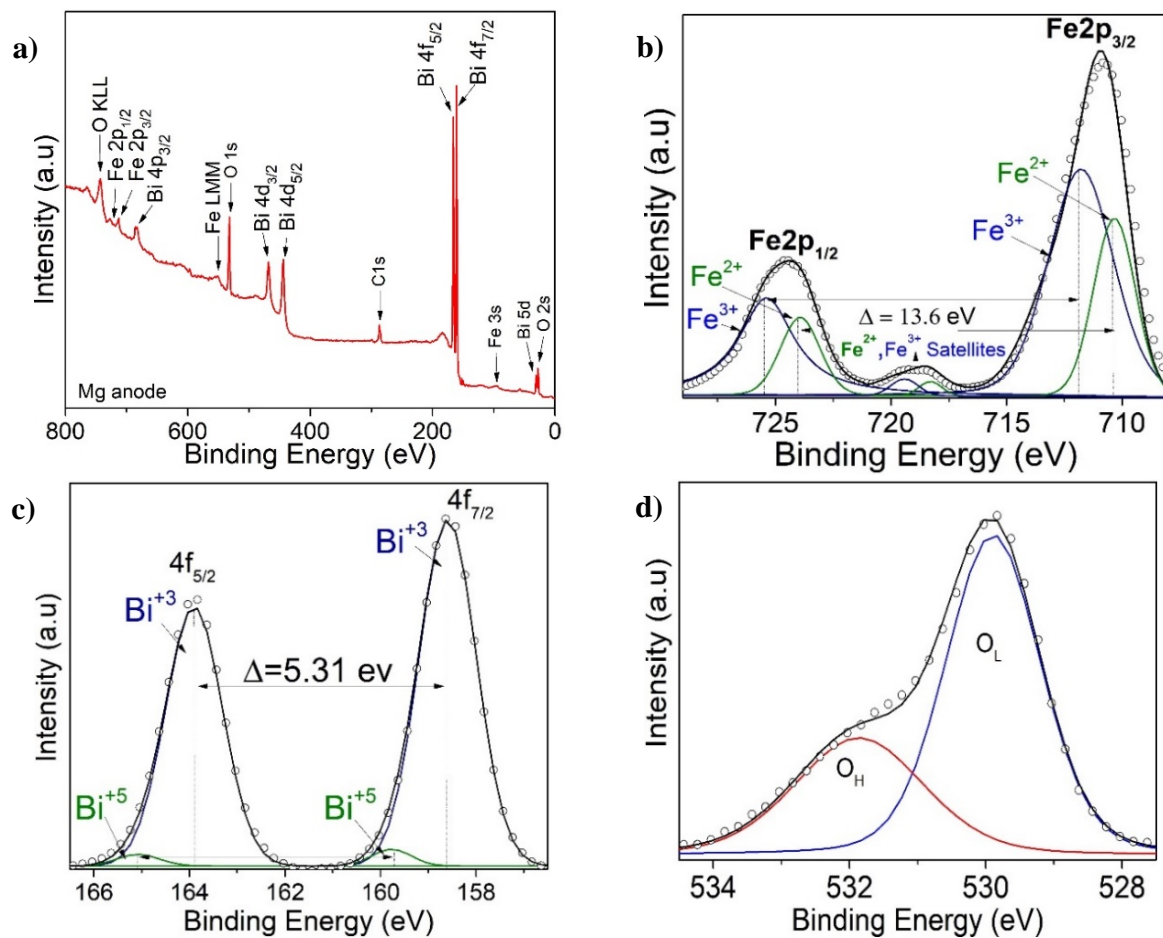
states and it can be noted that at less layer,  $\text{Fe}^{3+}$  increases in proportion. On the other hand, it is evident that  $\text{Bi}^{3+}$  is very predominant in both monolayers and bilayers. Instead, there are an increase in the contribution of the oxides of the lattice (metallic oxides,  $\text{O}_\text{L}$ ) for the bilayers, which is a good indication that the surface area is more homogeneous and that there is less contamination. In addition, hydroxyl oxygen ( $\text{O}_\text{H}$ ) increases in the thin films with lower growth power and in the eroded sample, which may suggest that these thin films could be more photocatalytic.

**Figure 13** shows a summary of the atomic percentages of the thin films analyzed in this work. It is observed that the  $\text{VO}_2$  films most indicated to be coupled with the BFO films are those with 60 min of growth, since they have a higher percentage of vanadium on the surface and presented a higher proportion of  $\text{V}^{4+}$ . On the other hand, the most suitable BFO films to couple in the bilayer with  $\text{VO}_2$  are those with a BFO layer of 30 nm, because they show iron with a higher atomic percentage in the surface area and a higher  $\text{Fe}^{3+}$  ratio.



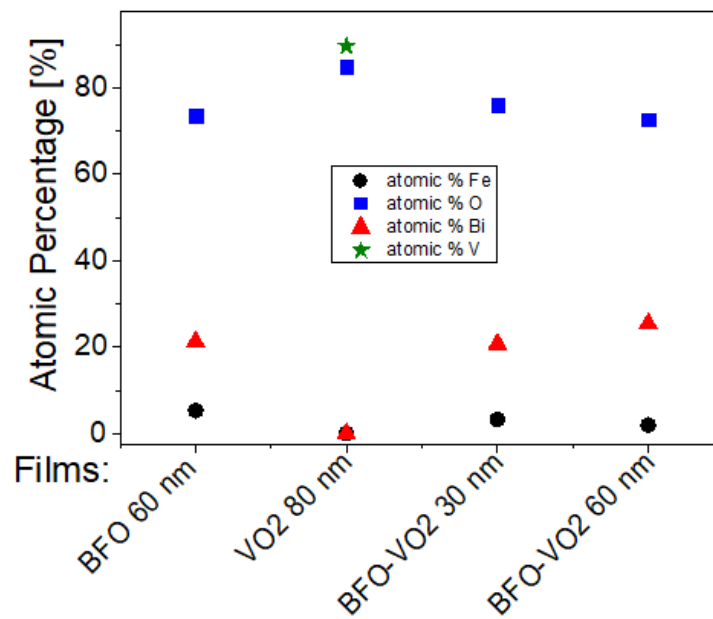
**Figure 11.** Broad (a) and high-resolution (b, c, d) XPS spectra of BFO/ $\text{VO}_2$  bilayer with BFO layer of 30

nm



**Figure 12.** Broad (a) and high-resolution (b, c, d) XPS spectra of BFO/VO<sub>2</sub> bilayer with BFO layer of 60 nm

Finally, it should be noted that growing this type of film in bilayer without previous antecedents constituted a challenge, not only because of the coupling of their crystalline structures before and after the VO<sub>2</sub> transition temperature, but also because the upper layer (BFO) is grown at a higher temperature than the lower VO<sub>2</sub> layer and even its annealing temperature at 470 °C. This has therefore deoxidation of the VO<sub>2</sub> film turning metallic and losing its metal insulator transitions (MIT), which is not desirable, since coupling with a material that presents a structural phase transition is essential to generating a new line of applications and studies of the physical properties of these bilayers. Therefore, the VO<sub>2</sub> films were carefully over-oxidized, anticipating these losses when the upper BFO film grew and in this way a growth methodology of these bilayers was obtained.



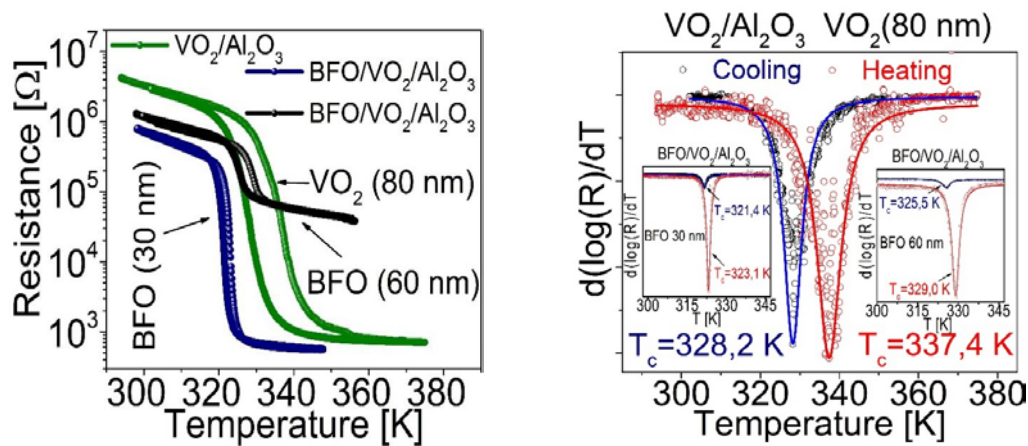
**Figure 13.** Atomic percentages on the surface of the selected films.

### 3.3 Electrical characterization

One of the best indications of the good coupling in the bilayers based on BFO and VO<sub>2</sub>, is that these also present hysteresis curves with a decrease of several orders of magnitude of electrical resistance in a fairly short time, see **Figure 14 left**), even so, there is a leftward shift in the transition temperature with respect to the 340 K of VO<sub>2</sub>, this lower value (around ~325 K) can be caused by the increase in interfacial tension in VO<sub>2</sub>, since, there are studies such as the reported in Refs [20] and [25], which show us that there is a relationship between the increase in the tension of the crystal lattice in VO<sub>2</sub> and the decrease in critical temperature. This result corroborates the XRD analysis that show tensions in its crystal lattice, this would be, not only with the substrate of Al<sub>2</sub>O<sub>3</sub> but also with the film on it (BFO). Additionally, this possibly contributes to increase the voltage of the lattice in the VO<sub>2</sub> and therefore to decrease the transition temperature.

The electrical properties of the BFO/VO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> heterostructures are shown in **Figure 14**. A well-defined MIT with a resistance change of two and three orders of magnitude ( $\Delta R \sim 100$  and  $\sim 1000$ ) can be seen in a) and b) with an  $T_{MIT}$  between  $\sim 322$  and  $327$  K (hysteresis width between  $\Delta H \approx 4$  and  $3$  K) for the 30 and 60 nm bilayers respectively.  $T_{MIT}$  is estimated by the well-known expression  $T_{MIT} = \frac{1}{2}(T_c - T_h)$ . Here,  $(T_c - T_h)$  represent the transition temperatures at the center of the derived curve ( $d[\log(R)]/dT$ ) during the cooling and heating procedure, respectively (inset in **Figure 14 right**)). The observation of an MIT, which improves the conductivity of the studied system, confirms the presence of the VO<sub>2</sub>(M) phase in the lower film. Here, it should be noted that the VO<sub>2</sub>(B) phase exhibits neither MIT nor thermal hysteresis phenomena [26]. The resistance of VO<sub>2</sub>(B) decreases exponentially with increasing temperature [27]. It is striking that the measured

$T_{MIT}$  of the polycrystalline film  $VO_2$  in the heterostructure is relatively low compared to that reported for individual crystals or epitaxial thin films [28]. This finding, along with the observed resistance change of several orders of magnitude, are in agree with the results reported for polycrystalline  $VO_2$  thin films [29], [30].



**Figure 14.** left) Resistance curves as a function of temperature for  $VO_2/Al_2O_3$  ( $VO_2$  80 nm), BFO/ $VO_2/Al_2O_3$  (BFO layer of 30 nm) and BFO/ $VO_2/Al_2O_3$  (BFO layer of 60 nm), right) derived from electrical resistance as a function of temperature, for the films obtained.

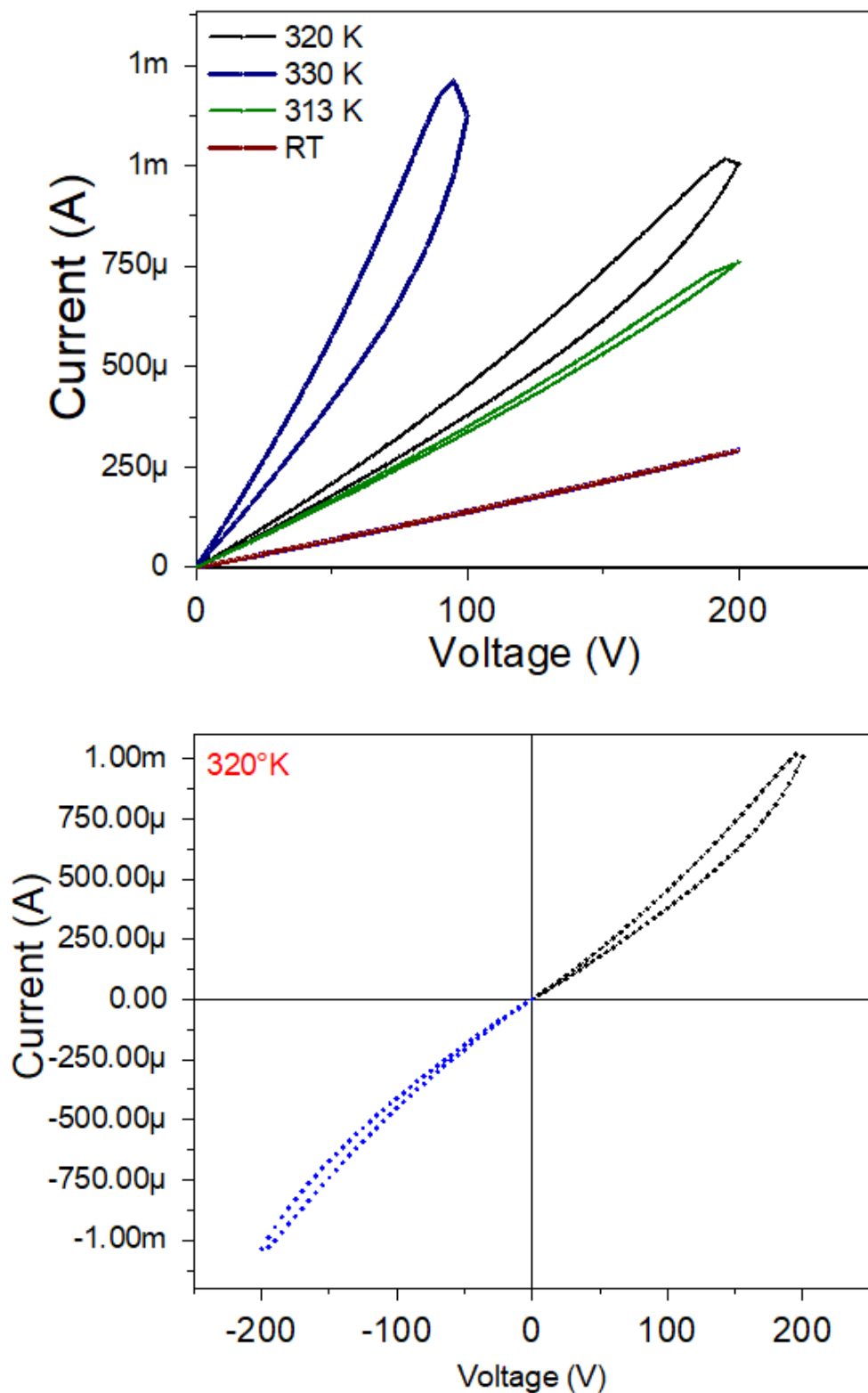
It is worth mentioning that the growth of polycrystalline thin films of  $VO_2$  with reduced  $T_{MIT}$  is technologically interesting and constitutes a current research topic in solid state physics [31]. Although the  $T_{MIT}$  of the polycrystalline thin film  $VO_2$  can be effectively adjusted and various approaches have been adopted to explain the origin of the reduced  $T_{MIT}$ , there is no consensus on this matter. Therefore, the true physical mechanism behind the reduction in the  $T_{MIT}$  value of polycrystalline  $VO_2$  thin films remains an open question.

Although understanding the phase transition in thin-film  $VO_2$  is challenging, it is probably related to a network distortion process. In a simplified image, it has been argued that a charge density wave along the c-axis of rutile [32], with the wave vector  $2cR$ , forms during transition. As a result, the unit cell along the c-axis is duplicated, generating a periodic distortion of the network in this direction. A slight rotation of the dimers with respect to the c-axis has also been verified [32]. As for the phase transition of  $VO_2$  polycrystalline films, it is evident that the phase transition in this case is easier than that of  $VO_2$  monocrystals [21]. The main feature of  $VO_2$  polycrystalline films is the deformation compatibility between differently oriented grains. Therefore, it is to be expected that grains oriented along specific directions may have irregular distortions of the network.

To finish with the characterizations carried out, it was decided to study I-V curves to analyze the electrical behavior of the bilayer thin film with BFO layer of 30 nm, as a first approximation to future studies of these bilayers in functional applications for storage devices, such as memresistors. The bilayer thin film with BFO

layer of 30 nm was chosen because it was the one that, in percentage, varied its electrical and magnetic responses more compared to the MIT of  $\text{VO}_2$ , therefore, the film that has the greatest potential to work with nanotechnological applications or prototypes.

**Figure 15** shows the I-V curves of BFO/ $\text{Al}_2\text{O}_3$  (BFO layer of 30 nm) varying voltage from 0 V to 200 V at different temperatures (**Figure 15 (up)**) and from -200V to 200V at a temperature of 320 K (**Figure 15 (down)**). Here, it is evident how the film progressively goes from having an ohmic behavior at room temperature to having a hysteresis cycle typical of resistive memories, which has a maximum width of  $\sim 20$  V and as the temperature decreases this hysteresis cycle closes.



**Figure 15.** I-V curves for a BFO/VO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> bilayer, with BFO layer of 30 nm, with different temperatures (up), and at 320 K with positive and negative voltage values (down).

Studying more thoroughly the I-V of the bilayer at 320 K, we can observe a difference in the symmetry of the resistive cycle with positive and negative voltages, however, the negative voltage cycle is closed more, an

effect that agrees with what is reported in Refs. [33] and [34]. Here, it is denoted that it is a typical behavior of these memresistor materials, that they change their response with each measurement.

On the other hand, effects such as magnetoelectric coupling and its behavior at different fields and with temperatures before and after the transition would be a next step for the study of these heterostructures with high mismatch and FE/ AF/MIT coupling.

#### 4. CONCLUSIONS

It was possible to obtain the growth parameters for films based on BFO and VO<sub>2</sub>, together with a methodology that allows obtaining a good coupling between films and the desired oxidation states of vanadium, bismuth, and iron. BFO and VO<sub>2</sub> monolayers were obtained with ordered growth, good adherence, roughness (between 0.2 and 2.3 nm), and grain size (between 20 and 32 nm). In addition, an ordered growth without cracks is evidenced. AFM measurements show differences in the topography of the individual layers and bilayers, a change in the average grain diameter from 20 nm to 67 nm and a change in roughness from 0.2 nm to almost 16 nm in BFO/VO<sub>2</sub> bilayers.

AFM measurements show differences in the topography of individual monolayers and bilayers, a change in average grain diameter from 20 nm to 67 nm was observed, and a change in roughness from 0.2 nm to nearly 16 nm was observed in BFO/VO<sub>2</sub> bilayers. In addition, an orderly growth without cracks (AFM) is evident. The homogeneity observed in the films with the best results will possibly contribute to obtaining good results in characterizations of current, voltage, resistance and behavior against electric and magnetic fields. It was observed by means of XPS that the predominant oxidation state in the individual and bilayer thin films is VO<sub>2</sub>. However, there is also a presence in the films of the V<sub>2</sub>O<sub>5</sub> phase. The BFO and BFO/VO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> films show the presented BiFeO<sub>3</sub> as the predominant phase.

Furthermore, thin films show a structural phase change with resistance change of several orders of magnitude, both for monolayers and for bilayers. Also, hysteresis loops are presented here, which is an excellent indication to work these films in memresistor and solar cells.

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## Conflicts of Interest

There are no conflicts to declare

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