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Multiscale information propagation in emergent functional networks

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Abstract: Complex biological systems consist of large numbers of interconnected units, characterized by emergent properties such as collective computation. In spite of all the progress in the last decade, we still lack a deep understanding of how these properties arise from the coupling between the structure and dynamics. Here, we introduce the multiscale emergent functional state, which can be represented as a network where links encode the flow exchange between the nodes, calculated using diffusion processes on top of the network. We analyze the emergent functional state to study the distribution of the flow among components of 92 fungal networks, identifying their functional modules at different scales and, more importantly, demonstrating the importance of functional modules for information content of networks, quantified in terms of network spectral entropy. Our results suggest that the topological complexity of fungal networks guarantees the existence of functional modules at different scales keeping the information entropy, and functional diversity, high.

Keywords: Information dynamics; Multiscale analysis; Networks entropy; Network density matrix; Fungal networks

1. Introduction

The underlying topology of complex systems is vital for their function, as it regulates the interactions between system’s units and constrain the pathways for information flow which are necessary for the system to operate. However, a direct analysis of the network – by using its adjacency matrix representation – only provides information about the local interactions, the direct ones encoded in the adjacency matrix, thus missing the dynamic and multiscale nature of node-node communications. To shed light on communication between system’s units at orders higher than one, i.e., beyond pairwise connections, a spectrum of possibilities are already available.

On the one hand, one can use random walk dynamics to gain insights about how information flow is locally trapped, allowing one, for instance, to uncover the functional mesoscale organization of classical [1], multilayer [2] and higher-order [3–5] networks, even across distinct scales [6–9]. Other exploration dynamics, such as walks, can be similarly used to capture the communicability between units and gain insights about the role of each node in exchanging information through the network [10–13]. On the other hand, one can employ higher-order models of the network structure [14,15] and study specific dynamics on the top of them, from synchronization [16–18] to social contagion [19–21], for instance.

Naturally, a rather holistic approach is required to capture the full information content of networks, their information dynamic and the multiscale communications between their units. In this work, we use diffusion processes to quantify the flow exchange between every pair of nodes and construct their emergent functional states, i.e., networks in which the links encode the information exchange between the nodes within a certain amount of time $\tau$. As a case study of biological relevance, we consider a broad set of networks of fungi and slime moulds, where structural links are obtained from cord conductances which, in turn, provide also a weight for pairwise connectivity (see [22] for details about the data set). This class of biological systems is peculiar in the way it stores and processes information: it has been shown that Physarum polycephalum (i.e., slime molds) grown on a rescaled...
map of Japan connects stations with a pattern resembling Tokyo’s rail system [23], and, even without a brain, they are able to externalize spatial memory to navigate complex environments [24], features which make them an exemplary model organism to understand the basic elements of computation in biological systems.

We show that such emergent functional states change with $\tau$ – acting as a tunable parameter encoding the propagation time of signals – while enabling one to analyze the multiscale nature of communications. Consequently, we show that analyzing the mesoscale organization of those emerging network states by means of an efficient algorithm, namely the Louvain method [25], it is possible to unravel the multiscale functional architecture of a system. In parallel, by using a recently developed statistical field theory of information dynamics [26], allowing for the definition of a suitable network density state [27,28] and multiscale analysis [29], we quantify the information content of networks in terms of their Von Neumann entropy, and show that the existence of functional modules are highly important for the macroscopic properties of systems.

Our results, when compared to randomized models of the fungal networks, suggest that their topological complexity plays a fundamental role in keeping their flow dynamics diverse, especially between middle- and long-range communication scales.

2. Materials and Methods

2.1. Propagator of information dynamics

The nodes of a complex network can be represented as canonical vectors $|x_i\rangle, i = 1, 2, \ldots, N$, shaping a discrete and finite Euclidean space [30]. In this space, the pairwise connections between nodes can be encoded into an adjacency operator $\hat{W}$, where $\langle x_j | \hat{W} | x_i \rangle = W_{ij}$ provides the intensity of the link between node $i$ and $j$. To capture flow dynamics within the network, one can assume the presence of an information field $|\varphi(\tau)\rangle$ on the top of its topology, free to evolve according to the diffusion equation

$$\partial_\tau |\varphi(\tau)\rangle = -\hat{L} |\varphi(\tau)\rangle,$$

$\hat{L}$ being the Laplacian operator $\langle x_j | \hat{L} | x_i \rangle = K_i \delta_{ij} - W_{ij}$, where $K_i = \sum_{j=1}^{N} W_{ij}$ is the strength defined as summation of weights of links emanating from node $i$, and $\delta_{ij}$ is the Kronecker delta. The solution of the master equation in terms of the propagator $\hat{U}(\tau) = e^{-\tau \hat{L}}$ is given by

$$|\varphi(\tau)\rangle = \hat{U}(\tau) |\varphi(0)\rangle.$$

To quantify the propagation of the field with an initial value $\varphi_0$ originating from node $i$, one can set the initial condition as $|\varphi(0)\rangle = \varphi_0 |x_i\rangle$ and calculate the flow of the field from node $i$ to any node $j$ as $\varphi_0 \langle x_j | \hat{U}(\tau) | x_i \rangle$. 


2.2. Emergent functional states

Figure 1. Emergent functional states associated to four different network types, at four different temporal scales, $\tau$. On the left, the Structure column represents the considered synthetic networks exhibiting random geometric, Barabasi Albert, Watts Strogatz and Stochastic Blocks topology. In front of each network type, its emergent functional state is visualized at four temporal scales: $\tau = 0.1, 1, 10, 100$. Subplots show the emergent functional state of a (A) random geometric network with radius 0.18, (B) Barabasi Albert network with $m = 1$, (C) Watts Strogatz network with mean degree of 4 and rewiring probability of 0.05 and (D) a Stochastic Blocks network having 4 communities with intra-community connectivity probability of $10^{-3}$ and inter-community connectivity probability of 0.5. The intensity of each link in the emergent functional state, characterized by the flow exchange between the nodes at the specific temporal scale, is colored from low (dark blue) to high (yellow).

Instead of focusing on the adjacency matrix, that provides a static and localized picture of interactions in the system, one can consider the propagator encoding pairwise flow exchange between the nodes at multiple scales, characterized by a tunable temporal parameter $\tau$ to span short-, middle- and long-range functional interactions between the
nodes. To this aim, we define the emergent functional state as \( \hat{W}(\tau) \) such that the interaction between node \( i \) and \( j \) follows

\[
\langle x_i | \hat{W}(\tau) | x_j \rangle = \phi_0 U_{ij}(\tau)(1 - \delta_{ij}),
\]

(3)

where \( \hat{W}(\tau) \) is, in practice, equivalent to the propagator rescaled by the constant \( \phi_0 \), excluding the diagonal entries \( \langle x_i | \hat{U}(\tau) | x_i \rangle, i = 1, 2, ... N \).

It is worth noting that, in the limit of very small temporal scale, the propagator can be linearized as \( \hat{U}(\tau) = e^{-\tau L} \approx I - \tau \hat{L} \), where \( I \) is the identity matrix, and the functional network adjacency reduces to the adjacency matrix \( \hat{W}_{ij}(\tau) = \phi_0 (1 - \delta_{ij})U_{ij}(\tau) \approx \phi_0 (1 - \delta_{ij})(\delta_{ij} - \tau K_i \delta_{ij} + \tau W_{ij}) = \phi_0 \tau W_{ij}(1 - \delta_{ij}) \), encoding local interactions between the nodes, while excluding self-loops. For clarity, a visualization of the emergent functional state corresponding to four synthetic network is visualized in Fig. 1. Moreover, for the emergent functional state \( \hat{W}(\tau) \), we represent the strength of node \( i \) using small letter \( k_i = \sum_{j=1}^{N} W_{ij} \). The strength distribution depends on the temporal scale of interactions \( \tau \) (See Fig. 2).

### 2.3. Emergent functional modules

A variety of algorithms have been developed to capture the communities in complex networks, including the Louvain method [25], a widely used and efficient one. Here, instead of analyzing the community structure of the underlying network, we focus on emergent functional states \( \hat{W}(\tau) \). Each functional module consists of nodes that exchange higher amount of flow between themselves than with the rest of the nodes. Naturally, the content of functional modules depend on the underlying network and the propagation time scale. When short range communications are considered, the number of communities is expected to be large, due to the large number of possible local neighborhoods. Conversely, at long range interactions, a smaller number of modules containing large numbers of nodes are expected to emerge.

It is worth mentioning that dynamical processes have been previously used to devise methods, like the Markov Stability framework, to identify communities. In the Markov Stability framework, a dynamical process like random walk at its equilibrium state is assumed on top of the network. In this case, meaningful communities are identified based on the probability that the random walker originating from these communities stays inside them over time.

Note that despite the similarity in using dynamical processes, our approach is fundamentally different. Firstly, the emergent functional state can be calculated for any general propagation process, regardless of the initial conditions. For instance, assume the fraction of the initial amount of the field \( \phi_0 \) on top of node \( i \) is given by \( \Phi_0 \). In this case, this initial condition can be built into the propagator as follows:

\[
\hat{U}(\tau) = \Pi e^{-\tau L},
\]

(4)

where \( \Pi_{ij} = \phi_0 \delta_{ij} \), generalizing Eq. 3 to initial distributions other than uniform \( \hat{I} = \hat{I} \).

Therefore, given that there is no constraint on the initial distribution, the modules appearing in the emergent functional state can be fundamentally different from the ones identified by the Markov Stability framework, that assumes the equilibrium distribution. Secondly, even if the initial distributions are taken to be the same, it is worth noting that by applying the Louvain method to the emergent functional states, one optimizes the modularity [31]. Of course, the result of such algorithm can be coherent with the Markov Stability framework. However, the quality functions these two algorithms optimize are different and further studies are required to provide a comprehensive picture of their similarities and differences.
2.4. Statistical physics of complex information dynamics

While approaches based on classical information theory and statistical physics have been widely used to study complex networks [32,33], they often rely on network descriptors such as degree and miss the full information content of interconnected systems. Conversely, inspired by developments in quantum thermodynamics, network density matrices have been proposed to study complex systems through the lens of Von Neumann entropy and without the information loss due to reducing the network to its descriptors [34]. Early definitions of the density matrix have been successfully applied to a range of problems, including the dimensionality reduction of multilayer networks [35], while they often were not physically interpretable and did not satisfy requirements such as sub-additivity.

Resolving the discussed issues related to the previous definitions, the network Gibbs state has been introduced [26,36,37] to quantify and compare the information content of complex networks. This framework has been successfully applied to a wide variety of problems in network science, from multilayer reducibility and its effect on the transport phenomena [38], to the human microbiome [36], the human brain [39,40], network robustness [41] and pan-viral interactomes [42].

In the following, we review the basics of this framework. The network Gibbs state is defined as

\[ \hat{\rho}(\tau) = \frac{\hat{U}(\tau)}{\text{Tr}[\hat{U}(\tau)]}. \]  (5)

Remarkably, it has been shown that the density matrix provides the smallest element to describe the flow dynamics in the system [26]. Furthermore, the density matrix is a superposition of a ensemble of operators, acting like streams directing the flow in the system. The mixedness of streams can be quantified using the Von Neumann entropy as

\[ S(\tau) = -\text{Tr}[\hat{\rho}(\tau) \log \hat{\rho}(\tau)], \]  (6)

which is a measure of diversity of flow in the system. Later in the text, we will show that the number of functional modules in fungal networks is proportional to the Von Neumann entropy.

2.5. Functional diversity

Units of complex systems exhibit undeniable similarities in appearance, while they can take very different functional roles according to how they are positioned in the network. For instance, neurons belonging to different functional areas of brain are involved in different functions of the system [43]. The functional diversity of networks can be determined, from the perspective of information dynamics, in terms of the differences among the nodes as senders or receivers of information. It has been previously shown that the Von Neumann entropy is a proxy of functional diversity in the system [26], and the presence of topological complexity boosts the functional diversity of networks at middle to large propagation time scales \( \tau \). The functional diversity, in that study, is quantified in terms of the average cosine distance of the flow distribution vectors emanating from pairs of nodes.

Comparing the Von Neumann entropy of empirical networks with their configuration models—i.e., a model that generates a random network with the same degree distribution as the real network under study [44]—, or types of randomized versions, has been used to investigate the benefits of topological complexity often observed in real systems for their functional diversity [26,40].

2.6. Normalizing the temporal scales across networks

For a connected undirected network, let the eigenvalues of the Laplacian operator be \( \lambda_\ell, \ell = 1, 2, ..., N \), ordered as \( \lambda_\ell \leq \lambda_{\ell+1} \) and \( \lambda_1 = 0 \). Therefore, the eigenvalues of the propagator follows \( e^{-\tau \lambda_\ell} \), \( \ell = 1, 2, ..., N \). The diffusion process reaches equilibrium at \( \tau \to \infty \), where all the eigenvalues of the propagator decay except the one corresponding to the first eigenvalue of the Laplacian \( e^{-\tau \lambda_1} = 1 \). The second eigenvalue of the Laplacian determines
the diffusion time $\tau_d = 1/\lambda_2$, that is the temporal scales at which the dynamics is close to equilibrium and the second eigenvalue of the propagator has decayed to $e^{-\tau_d \lambda_2} = \frac{1}{2}$.

The diffusion time is different from network to network, depending on the connectivity, number of nodes, topology, etc. To allow for comparison across networks, one can divide the propagation time scale $\tau$ by the diffusion time to obtain the rescaled temporal parameter $\tau/\tau_d$.

2.7. Fungal networks

Here, we consider three species of fungi and slime molds: *Physarum polycephalum* (Pp), *Phanerochaete velutina* (Pv) and *Resinicium bicolor* (Rb). For each species we analyze a certain number of networks to capture the average topological properties of the species. In these networks the weight of the structural links are proportional to the cord conductance for pairwise interactions (See [22] for details about the data set). While the full data set consists of 270 networks, we focus on a subset of 92 networks with the number of nodes equal or smaller than 500, for computational convenience.

2.8. Randomized networks

Often, properties of networks are understood in comparison with null models. The most famous null models for networks are Erdos-Renyi random models, networks of the same size and the same number of links as the original network while the links are randomly distributed between the nodes, and configuration models, where the degree distribution is fixed and the connections are randomized.

In case of very sparse networks, like trees, the randomization of links leads a connected network to a network that is no longer connected, consisting of a number of disconnected sub networks. Undoubtedly, the resulting disconnected network exhibits profoundly different information dynamics and, therefore, can only work as a null model that trivially and significantly differs from the original network. For instance, it has been shown that at large temporal scales, the Von Neumann entropy of networks can be approximated as the number of isolated components in disconnected networks [41], as an example of trivial differences between these null models and sparse networks.

The fungal networks under investigation here are both sparse and weighted. We exploit weights to generate suitable null models, defined as randomized-weight configuration model (RWCM), as follows: One fixes the position of links in the original network, while assigning random weights, from a uniform distribution, to the links. Finally, all the links are normalized by the summation of weights in the random network and multiplied by the summation of links in the original network. This normalization guarantees the the overall weight strength of the null model equals the original network. Overall, RWCM keeps the network connected while randomizes the distribution of weights, providing null models for sparse weighted networks.

3. Results

3.1. Functional Networks

We numerically calculate the emergent functional networks corresponding to all considered fungal networks, at different scales. When the temporal parameter is small, the flow would not penetrate into the network and the information exchange is limited to the nodes and their first neighbors. While, at large temporal scales, the functional network becomes a fully connected network where each node is connected to all others (See Fig.2 for an example). Note that the average degree (strength) of functional networks can be defined as

$$\bar{k} = \frac{1}{N} \sum_{i,j=1}^{N} \tilde{W}_{ij},$$

which can be divided by $\phi_0$ to make it independent of the initial value of the field.
The emergent functional states corresponding to one realization of *Resinicium bicolor* networks, at 3 different scales (Markov time: $\tau = 0.0001, 0.01, 1.5$) are shown. Below each functional state, the corresponding strength distribution is reported. Tuning the temporal parameter from small to large values, one reaches a state where the functional state is fully entangled.

We limit the temporal range to $\tau/\tau_d \in [10^{-10}, 10]$, to capture the full range of flow dynamics within all the fungal networks. Consequently, it is straightforward to show that the connectivity of functional networks increases with the rescaled temporal parameter (See Fig.3).
The average strength of nodes $\bar{k}$ in the emergent functional state (See Methods) is plotted against the rescaled temporal parameter $\tau / \tau_d$ (See Methods). The relevant temporal scales for analysis of information dynamics depends on a variety of parameters, such as the number of nodes and their topology, reflected in the spectrum of the Laplacian matrix. During the transition from the extremely small rescaled temporal scale to the large rescaled temporal scale, it is observed that the average strength of the emergent functional state (See Methods) goes from zero, where no exchange flow is possible between the nodes, to $\phi_0$, where all the field originated from each node is completely distributed to the others. Here, we numerically observe that the relevant temporal scale for analysis of information dynamics in all three fungal species lays in $\tau / \tau_d \in [10^{-10}, 10]$.

Using the Louvain algorithm, we find the number of functional modules and their members in each functional fungal network ($\hat{W}$) at each temporal scale (See Fig.4). Also, we calculate the entropy of the fungal networks ($\hat{W}$), to show that the number of functional modules scales is logarithmically proportional to the Von Neumann entropy of the network at each temporal scale.
Discussion

Structural metrics are generally unable to capture different aspects of information flow within complex systems such as the fungal networks. Instead, a variety of methods have been developed, mostly based on coupling stochastic processes to the structural data and taking the higher order interactions into account, to study the functional interactions, transport phenomena, mesoscopic organization and information content of networks [1–21].

In this article, we proposed a novel perspective in terms of emergent functional states, represented as networks where links encode the flow exchange between the nodes, according to diffusion processes. The emergent functional states can be calculated at different scales.
signal propagation time scales, describing short- to long-range functional interactions between the components of the system and providing a multiscale lens for systemic analysis. Consequently, we analyzed classical metrics of the emergent functional networks, such as degree distribution, encoding the total flow exchange of nodes within the networks, and the number and content of functional modules, obtained from feeding the emergent functional state as a network into the Louvain algorithm, which is based on modularity maximization. Moreover, we find numerical relations between the knowledge gained from the emergent functional networks and our information theoretic measures, to provide a broader picture. We show that the multiscale modular organization of the emergent functional states is highly important for systems’ Von Neumann entropy, a proxy of the information content of networks [26]. Finally, we introduce the method of randomized-weight configuration models (RWCM), to obtain null models enabling us to compare the properties of empirical systems having highly sparse and weighted networks. We compare the fungal networks with their RWCM and show that they exhibit significantly larger Von Neumann entropy, indicating the efficiency of natural self-organization in providing structures that keep the flow dynamics diverse, even when long range communications are considered.

Our work provides a unifying picture of multiscale information dynamics and information content of fungal networks, further relating information dynamics to information-theoretic measures, a step towards understanding of how collective computation occurs and emerges among interconnected biological units.

Funding: This research received no external funding.

Data Availability Statement: The data are publicly available from the cited references.

Conflicts of Interest: The authors declare no conflict of interest.
