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# Electronic Surveillance and Security Applications of Magnetic Microwires

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#### **Abstract:**

Applications in security and electronic surveillance require combination of excellent magnetic softness with good mechanical and anti-corrosive properties and low dimensionality. We overviewed the feasibility of using glass-coated microwires for electronic article surveillance and security applications as well as, different routes of tuning the magnetic properties of individual microwires or microwires arrays making them quite attractive for electronic article surveillance and security applications.

We provide the routes for tuning the hysteresis loops non-linearity by the magnetostatic interaction between the microwires in the arrays of different types of amorphous microwires. The presence of neighbouring microwire (either Fe or Co-based) significantly affects the hysteresis loop of the whole microwires array. In a microwires array containing magnetically bistable microwires we observed splitting of the initially rectangular hysteresis loop with a number of Barkhausen jumps correlated with the number of magnetically bistable microwires. Essentially, non-linear and irregular hysteresis loops have been observed in mixed arrays containing Fe and Co-rich microwires. The obtained non-linearity in hysteresis loops allowed to increase the harmonics and tune their magnetic field dependencies. On the other hand, several routes allowing to tune the switching field by either post-processing or modifying the magnetoelastic anisotropy have been reviewed. Non-linear hysteresis loops have been also observed upon devitrification of amorphous microwires. Semi-hard magnetic microwires have been obtained by annealing of Fe-Pt-Si microwires.

Observed unique combination of magnetic properties together with thin dimensions and excellent mechanical and anti-corrosive properties provide excellent perspectives for the use of glass-coated microwires for security and electronic surveillance applications.

**Keywords:** magnetic microwires; magnetic bistability; magnetic tag; electronic surveillance; domain wall propagation; post-processing; magnetic anisotropy; magnetostatic interaction.

## 1. Introduction

Soft magnetic materials are essential part of magnetic sensors and devices demanded by several industries, including (but not limited to) microelectronics, electrical engineering, car, aerospace and aircraft industries, medicine, magnetic refrigerators, home entertainment, energy harvesting and conversion, informatics, magnetic recording or security and electronic surveillance [1-3]. In most of cases, like the case of security and electronic surveillance, in addition to excellent magnetic softness, a combination of mechanical and anti-corrosive properties and low dimensionality is required [4-6].

Almost all department stores, supermarkets, airports, libraries, museums, etc. are provided with different types of security and anti-thief systems. The principle of electronic article surveillance (EAS) systems operation is well established: articles are provided with tags that respond to electromagnetic fields generated by the gates at the store/supermarket/library exits [6]. The response is picked up by the antenna installed on the gate, switching on the alarm.

It is estimated that hundreds of thousands of such EAS systems have been installed and millions of tags are produced daily. Considering the great number of tags, they must be small, robust enough and inexpensive.

Usually, the magnetic field values generated by the gates are limited by electromagnetic regulations and therefore quite low, being typically below 100 A/m. Accordingly, the magnetic materials employed in tags must be magnetically soft enough. However, the magnetic softness of crystalline soft magnetic materials (Permalloy, Fe-Si) are affected by processing. Therefore, amorphous soft magnetic materials, prepared by rapid melt quenching, are considered as among the most suitable materials for tags containing soft magnetic materials [6,7].

Indeed, as a rule, amorphous materials present excellent magnetic softness together with superior mechanical properties [8-12]. Abrupt deterioration of the mechanical properties (such as tensile yield) upon the devitrification of amorphous precursor is reported [12]. Additionally, the fabrication process of amorphous materials involving rapid melt quenching is fast and inexpensive [4-8, 10]. Accordingly, amorphous soft magnetic materials are useful for the design of robust magnetic devices and magnetoelastic sensors [13-20].

Different rapid melt quenching methods allow the preparation of amorphous materials of planar (ribbons) or cylindrical (wires) shapes [4-7]. As discussed elsewhere, soft magnetic materials with squared hysteresis loops and relatively low coercivities are the preferred candidates for the EAS systems using magnetic tags [6]. The rectangular hysteresis loops can be easily implemented in different families of amorphous magnetic wires [21-25]. Therefore, considerable attention has been paid to applications of amorphous wires for magnetic tags for different kinds of EAS systems [7, 25-29].

The aforementioned squared hysteresis loops of magnetic wires are linked to the peculiar remagnetization process of magnetic wires running through a single and large Barkhausen jump [6, 7,21-25]. In such magnetic wires, a demagnetized state cannot be achieved [7, 21-23, 30-32]. Accordingly, such magnetic wires are also called magnetically bistable [30-32].

Magnetic wires can be prepared using different techniques involving rapid melt quenching [7, 21-23,33]. However, glass-coated magnetic microwires prepared by so called Taylor-Ulitovsky technique present the widest metallic nucleus diameters range (from 200 nm up to 100  $\mu$ m) [34-44]. In this way, the Taylor-Ulitovsky method is the unique technique allowing fabrication of nanowires by rapid melt quenching [34]. On the other hand, recently preparation of amorphous magnetic wires with diameter about 100  $\mu$ m coated by glass is recently reported [36]. The presence of a flexible, thin, biocompatible and insulating glass- coating allows to enhance the corrosive resistance and therefore makes these microwires suitable for novel applications including biomedicine, electronic article surveillance, non-destructive monitoring external stimuli (stresses, temperature) in smart composites or construction health monitoring through the microwire inclusions [37-42, 45-48]. It is worth noting, that in fact the Taylor-Ulitovsky technique is known since the 60-s [43] and has been used for the preparation of amorphous microwires since the 70-s [44]. The modern Taylor-Ulitovsky technique is suitable for preparation of continuous glass-coated microwires with up to 10 km long and roughly about 1 km of microwire can be prepared from 1 gram of metallic alloys [36,45].

Accordingly, considering dimensionality and combination of physical properties (magnetic, mechanical, corrosive), amorphous soft magnetic microwires are potentially suitable materials for electronic article surveillance and security applications. There are several original papers dealing with rather different (multi-bit or single-bit) security and EAS applications of magnetic microwires [25-27,47,48]. However, to our best knowledge, there are no any reviews summarizing published experimental results and analyzing trends in security and EAS applications of magnetic microwires.

Consequently, in this paper we will provide an overview of the trends related to EAS and security applications of glass-coated magnetic microwires.

This review is organized as follows. In Section 2 the experimental methods as well as the microwires characteristics analyzed in this review are provided. Section 3 deals with results on feasibility of using magnetic microwires for magnetic tags followed by overview of tuning of hysteresis loop non-linearity by the magnetostatic interaction between microwires and then by multi-bit magnetic tags applications of magnetic microwires.

#### 2. Materials and Methods

Fe-, Co- Ni- rich glass-coated amorphous microwires have been prepared by Taylor-Ulitovsky preparation method described in details elsewhere (the compositions and diameters of microwires are provided in Table 1) [36,45].

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Composition	Metallic Nucleus Diameter, d (µm)	Total Diameter, D (µm)	Ratio ρ= d/D	Magnetostriction Coefficient, $\lambda_s \times 10^{-6}$
Fe <sub>74</sub> B <sub>13</sub> Si <sub>11</sub> C <sub>2</sub>	10	20	0.5	38
$Fe_{74}B_{13}Si_{11}C_2$	12.3	15	0.82	38
$Fe_{74}B_{13}Si_{11}C_2$	17.3	28.2	0.61	38
$Fe_{74}B_{13}Si_{11}C_2$	19.4	26.6	0.73	38
Fe75B9Si12C4	15.2	17.2	0.88	38
Fe65Si15B15C5	12.6	20	0.63	38
$C_{069.2}Fe_{3.6}Ni_1B_{12.5}Si_{11}C_{1.2}Mo_{1.5}$	22.8	23.2	0.98	-1
$Co_{67}Fe_{3.9}Ni_{1.5}B_{11.5}Si_{14.5}Mo_{0.6}$	29.2	31	0.94	-0.5
$Fe_{71.7}B_{13.4}Si_{11}Nb_3Ni_{0.9}$	103	158	0.65	35
$Co_{69.2}Fe_{4.1}B_{11.8}Si_{13.8}C_{1.1}$	25.6	30.2	0.85	-0.03
$Co_{64,04}Fe_{5,71}B_{15,88}Si_{10,94}Cr_{3.4}Ni_{0,3}$	94	126	0.75	2
$Fe_{16}Co_{60}Si_{13}B_{11}$	12	29	0.41	15
$Fe_{47.4}Ni_{26.6}Si_{11}B_{13}C_{2}$	29	32.2	0.9	20
$Fe_{49.6}Ni_{27.9}Si_{7.5}B_{15}$	14.2	33.85	0.42	20
Fe71,8Cu1Nb3,1Si15B9,1	7.0	24.8	0.282	30
Fe71,8Cu1Nb3,1Si15B9,1	18.2	39	0.467	30
Fe38.5Co38.5B18Mo4Cu1	10	16.6	0.6	
$Fe_{50}Pt_{40}Si_{10}$	8	21	0.38	

Table 1. Compositions and geometry of studied glass-coated microwires.

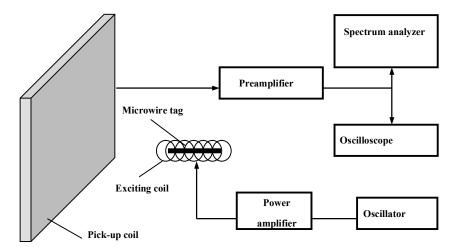
Generally, we analyzed three different types of magnetic microwires: i) amorphous microwires with high positive magnetostriction coefficients,  $\lambda_s$ , (Fe-Si-B-C, Fe-Ni-Si-B-C or Fe-Ni-Si-B), ii) amorphous microwires with vanishing  $\lambda_s$  (Co-Fe-Ni-B-Si-Mo, Co-Fe-Ni-B-Si-Mo, Co-Fe-B-Si-Cr-Ni or Co-Fe-B-Si-C) or iii) partially crystalline (nanocrystalline) (Fe-Cu-Nb-Si-B, Fe-Co-B-Mo-Cu or Fe-Pt-Si) microwires.

The amorphous structure of all the microwires has been proved by the X-ray Diffraction (XRD) method. All studied microwires present a broad halo in the XRD patterns. The XRD patterns have been obtained by the Bruker (D8 Advance) X-ray diffractometer with Cu K $_{\alpha}$  ( $\lambda$  = 1.54 Å) radiation. Several samples have been annealed in a conventional furnace at temperatures below the crystallization temperature. Typically, the crystallization of amorphous microwires was observed at  $T_{ann} \ge 500$  °C [49].

The induction method previously is used for the hysteresis loops measurements. The details of the experimental set-up are described in details elsewhere [46]. The hysteresis loops were represented as the magnetic field, H, dependence of the normalized magnetization,  $M/M_0$ , being M the magnetic moment at a given magnetic field, and  $M_0$  the magnetic moment at the maximum magnetic field amplitude  $H_m$ . Such hysteresis loops are useful for comparison of the samples with different chemical compositions (and, hence, different saturation magnetization).

In several cases, the hysteresis loops were measured with a conventional superconducting quantum interference device, SQUID.

The experimental set-up allowing to measure the electromagnetic response of a magnetic tag consisting of an exciting coil, a pick-up coil, a preamplifier and registration facilities is described elsewhere (see Fig.1) [47]. The magnetic tag (up to 4 cm long) is magnetized by 5 cm long exciting coil producing nearly uniform magnetic field up to 640 A/m with frequency f=332 Hz. We used a squared pick-up coil containing 20 turns with a side of 20 cm. An exciting coil with a magnetic tag inside was located perpendicular to the pick-up coil plane. The electromagnetic signal from the tag was amplified by a preamplifier with a voltage gain ~ 100 and detected by the registration facilities, i.e., spectrum analyzer CF 5210 and a digital oscilloscope. The noise voltage level at a frequency higher than 1 kHz is given by ~ 10  $\mu$ V/Hz<sup>1/2</sup>. The tag placed inside the exiting coil produces a periodic signal of negative and positive pulses detected in the digital oscilloscope.



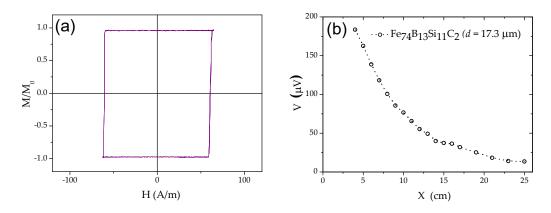
**Figure 1.** Scheme of experimental set-up designed to detect the electromagnetic signals of the magnetic labels. Reprinted with permission from ref. (47).

We also measured the amplitudes of the first seven harmonics of the voltage induced in the pick-up coil using a lock-in amplifier.

#### 3. Results and Discussion

#### 3.1. Feasibility of using magnetic microwires for magnetic tags.

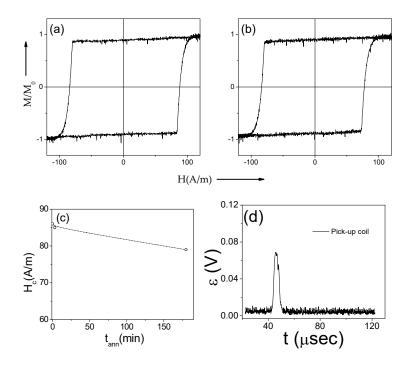
For magnetic tag applications the magnetic response must be as high as possible. Fe-rich microwires present a higher saturation magnetization. Additionally, as-prepared Fe rich microwires have perfectly rectangular hysteresis loops (see Figure 2a). To assess the feasibility of using Fe-rich microwires for magnetic tags, we measured the 5th harmonic as a function of the distance between the tag and the pick-up coil. As can be seen from Figure 2b, the 5-th harmonics of 3 cm long Fe<sub>74</sub>B<sub>15</sub>Si<sub>11</sub>C<sub>2</sub> microwire (metallic nucleus diameter, d=17.3  $\mu$ m) can be detected at a distance up to 25 cm. Similar studies of Fe –rich with d≈100  $\mu$ m show that in this case, the signal can be detected at a distance up to 50 cm [49].



**Figure 2.** Hysteresis loop (a) and the amplitudes of the 5-th harmonics generated by a 3 cm long  $Fe_{74}B_{13}Si_{11}C_2$  microwire with diameter  $d = 17.3 \,\mu\text{m}$  as a function of the distance from the pick-up coil (b).

Perfectly rectangular hysteresis loops of Fe-rich microwires are quite stable: the character of hysteresis loops remains the same even after long annealing (180 min) at elevated annealing temperature,  $T_{ann}$ =400 °C (see Figure 3a,b). Only slight coercivity,  $H_c$ , decrease is observed upon annealing at  $T_{ann}$ =400 °C (Figure 3c). On the other hand, quite sharp voltage peaks (about 10 µsec) in the pick-up coils are produced upon magnetization switching in such Fe-rich microwires (Figure 3c).

## 3.2. Tuning of hysteresis loop non-linearity by the magnetostatic interaction between microwires.



**Figure 3.** Hysteresis loops of as- prepared (a), and annealed at  $T_{amm}$ = 400 °C for 180 min (b) Fe<sub>75</sub>B<sub>9</sub>Si<sub>12</sub>C<sub>4</sub> microwires, dependence of coercivity on annealing time (c) and the time dependence of the EMF signal in the pick-up coil (d).

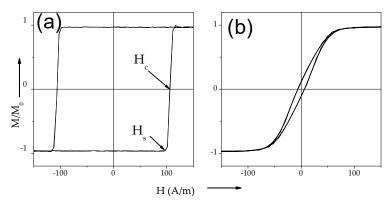
Magnetic tags applications require a non-linear hysteresis loop that contains the characteristic distribution of harmonic frequencies. It is believed that the steeper the magnetization reversal, the

higher the harmonic content of the signal. Accordingly, perfectly rectangular hysteresis loops with low coercivity observed in Fe-rich microwires (Figures 2, 3) are attractive for use as magnetic tags.

On the other hand, the non-linearity of the hysteresis loop of the magnetic microwires can be further improved using the magnetostatic interaction of microwires. Below, we will present several experimental results on magnetic response of two kinds of individual microwires (Co<sub>67</sub>Fe<sub>3.9</sub>Ni<sub>1.5</sub>Si<sub>14.5</sub>Mo<sub>6</sub> and Fe<sub>74</sub>Bi<sub>3</sub>Si<sub>11</sub>C<sub>2</sub>) as well as the arrays containing either microwires of the same type or arrays containing two different kinds of microwires. Microwires in each array were located close to each other, that is, the magnetic nucleus were separated only by the glass coatings.

The hysteresis loops of such microwires are rather different: microwires with high and positive magnetostriction coefficient,  $\lambda_s$ , exhibit perfectly rectangular hysteresis loops with  $H_c \approx 100$  A/m (Figure 4a), however, inclined hysteresis loop with quite low  $H_c$  ( $H_c \approx 5$  A/m) is observed in Co<sub>67</sub>Fe<sub>3.9</sub>Ni<sub>1.5</sub>Bi<sub>14.5</sub>M<sub>0.6</sub> microwires (see Figure 4b).

As discussed elsewhere, the hysteresis loop of even individual Fe-rich magnetically bistable microwires is remarkably affected by the magnetic field amplitude. The most relevant hysteresis



**Figure 4.** Hysteresis loops of Fe<sub>75</sub>B<sub>9</sub>Si<sub>12</sub>C<sub>4</sub> microwires with positive (a) and Co<sub>67</sub>Fe<sub>3.9</sub>Ni<sub>1.5</sub>Si<sub>14.5</sub>M<sub>0.6</sub> with vanishing (b) magnetostriction coefficients.

loop change in a single Fe-rich magnetically bistable microwire takes place when the magnetic field amplitude,  $H_0$ , exceeds the switching field,  $H_s$ , value [50,51]: below certain "critical" magnetic field amplitude,  $H_0$ , value (for studied Fe<sub>75</sub>B<sub>9</sub>Si<sub>12</sub>C<sub>4</sub> microwire at  $H_0 \approx 100$  A/m) the hysteresis loop abruptly disappears (see Figure 5a). Above such a critical field, the magnetization switching by single and large Barkhausen jump occurs. Accordingly, such critical magnetic field is commonly referred as the aforementioned switching field,  $H_s$ , at which the irreversible magnetization switching begins. It is worth noting, that in AC hysteresis loops at low  $H_0$  and magnetic field frequency, f,  $H_s \approx H_c$  (see figures 4a, 5a). However, with increasing  $H_0$ , one can observe a deviation from the perfectly rectangular hysteresis loop typical of magnetically bistable Fe-based microwires (Figure 5a). This modification of the hysteresis loop shape (more noticeable for high  $H_0$  -values) was explained taking into account the counterbalance between the sweep rate, dH/dt, and the magnetization switching time required for single domain wall (DW) propagation over the sample [50,52]. In the case of triangular input signal dH/dt is given as [52]:

$$dH/dt = 4fH_0 \tag{1}$$

Accordingly, increasing of  $H_0$  or f results in faster sweep rate, dH/dt.

Such change of the hysteresis loops is linked with  $H_c$  increase. Previously, the frequency and magnetic field amplitude dependence of coercivity in various magnetic materials has been described as [52,53]:

$$H_c = H_{co} + B(fH_o)^{1/n}$$
 (2)

being  $H_{co}$  - the static coercivity,  $H_{o}$  - the magnetic field amplitude, and n is a coefficient ranging from 1 to 4, which depends on the sample geometry and the type of the hysteresis loop of the studied materials, and B - a coefficient depending on the intrinsic material parameters [52,53].

Additionally, even the switching field,  $H_s$ , increases with increasing  $H_o$  and f (see Figure 5a). However,  $H_s$  increases slower than  $H_c$  with increasing  $H_o$  and f (see Figure 5a for  $H_o$ ). The origin of such  $H_c$  (f),  $H_s$ (f),  $H_c$  ( $H_o$ ) and  $H_s$ ( $H_o$ ) dependencies has been discussed considering a reversible magnetization process associated with reversible DW movement at low magnetic field (below magnetization switching) and the irreversible DW movement associated to large and single Barkhausen jump [52,53].

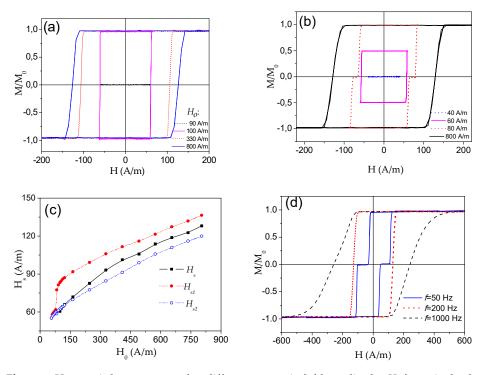


Figure 5. Hysteresis loops measured at different magnetic field amplitudes  $H_0$  for a single glass-coated Fe<sub>74</sub>B<sub>13</sub>Si<sub>11</sub>C<sub>2</sub> microwire (a), array with two microwires (b); dependence of switching field,  $H_s$  on  $H_0$  for a single microwire (solid line) and for an array containing 2 microwires  $H_{s1}$  and  $H_{s2}$  (dot-line) (c); hysteresis loops of the two microwires measured at different magnetic field frequencies f (d). Reprinted with permission from ref. [50].

The hysteresis loop of an array containing two  $Fe_{74}B_{13}Si_{11}C_2$  microwires is rather different from that of a single  $Fe_{74}B_{13}Si_{11}C_2$  microwire. Two Barkhausen jumps can be observed at Ho>80 A/m (see Figure 5b). Such peculiar hysteresis loop shape has been explained considering the magnetostatic

interaction in the two-microwire array [50,51]. Such magnetostatic interaction is a consequence of stray fields created by magnetically bistable microwires: the superposition of external and stray fields causes magnetization reversal in one of the samples, when the external field is below the switching field of a single microwire. Single rectangular hysteresis loop (similar to the case of single microwire shown in Figure 5a) is observed for 60 A/m < 80 A/m (see Figure 5b).

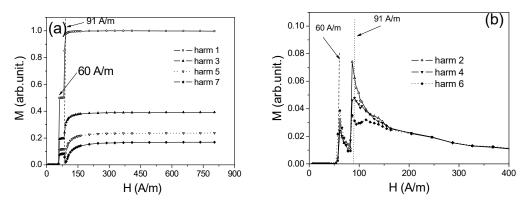
In the array consisting of two microwires, the lower switching field of the first Barkhausen jump,  $H_{s1}$ , decreases, while the switching field of the second Barkhausen jump,  $H_{s2}$ , increases (Figure 5c). Such difference must be attributed to the stray field created by neighboring microwire [50,51]. The origin of the different  $H_s$  values of individual microwires can be related with metallic nucleus diameters or glass-coating thickness fluctuations, stresses induced by cutting, etc.

Increasing magnetic field amplitude (approximately at  $H_0 > 250$  A/m), this splitting of the hysteresis loop disappears (Figure 5d). Such dependence of the hysteresis loop of two microwires array can be understood from the counterbalance between the dH/dt and the switching time determined by the velocity of the DW propagation along the whole wire.

As can be appreciated from eq. (2),  $H_c$  is also affected by the frequency, f. Accordingly,  $H_c$ , as well as overall hysteresis loops of two microwires array, are affected by f in a similar way as by  $H_0$  (see Figures 5c,d). For a two microwires array, two-steps hysteresis loops are observed for f<150 Hz. At f > 150 Hz, the hysteresis loop splitting disappears, and at 150 < f<1000 Hz, a single smooth magnetization jump is observed.

Accordingly, the odd and even harmonics of the signal of two Fe-rich microwires array are affected by  $H_0$  and f (see Figure 6a,b).

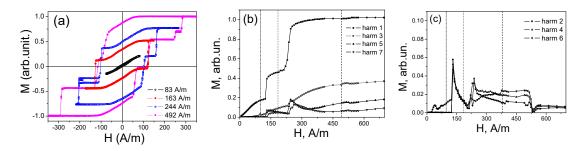
A sharp increase in the harmonics amplitudes is observed when  $H_0$  exceeds  $H_{s1}$  and  $H_{s2}$  (see Figures 6a,b). The even harmonics amplitudes are significantly inferior to the odd harmonics amplitudes. The field dependences of odd harmonics have a "plateau" between 60 and 90 A/m,



**Figure 6.** Dependences of odd harmonics (a) and even harmonics (b) on magnetic field amplitude in Fe<sub>74</sub>B<sub>13</sub>Si<sub>11</sub>C<sub>2</sub> microwires. Reprinted with permission from ref. (51).

which reflects the hysteresis loops splitting (see Figure 6a).

Another example of tuning the non-linearity of hysteresis loops and harmonics is the magnetostatic interaction of microwires with different character of hysteresis loops. Rather non-linear hysteresis loops can be obtained in an array consisting of one  $Fe_{74}B_{13}Si_{11}C_2$  and one  $Co_{67}Fe_{3.9}Ni_{1.5}Si_{14.5}M_{0.6}$  microwires (see Figure 7a). In such array, at  $H_0$  <90 A/m (which



**Figure 8** Hysteresis loops of 4-wires system consisting of three Fe<sub>74</sub>B<sub>13</sub>Si<sub>11</sub>C<sub>2</sub> microwires and a Co<sub>67</sub>Fe<sub>3.9</sub>Ni<sub>1.5</sub>B<sub>11.5</sub>Si<sub>14.5</sub>M<sub>0.6</sub> microwire (a); odd and even harmonics vs. H<sub>0</sub> are given in (b) and (c), respectively. Reprinted with permission from ref. (51).

 $H\left(A/m\right)$   $H\left(A/m\right)$ 

**Figure 7.** (a) Hysteresis loops of the Fe<sub>74</sub>B<sub>13</sub>Si<sub>11</sub>C<sub>2</sub> + Co<sub>67</sub>Fe<sub>3.9</sub>Ni<sub>1.5</sub>B<sub>11.5</sub>Si<sub>14.5</sub>M<sub>0.6</sub> array; (b) dependences of odd harmonics on magnetic field amplitude and (c) dependences of even harmonics on magnetic field amplitude. Reprinted with permission from ref. (50).

corresponds to  $H_8$  of Fe<sub>74</sub>B<sub>13</sub>Si<sub>11</sub>C<sub>2</sub> microwire) the hysteresis loops character is typical of those for a single Co<sub>67</sub>Fe<sub>3.9</sub>Ni<sub>1.5</sub>Si<sub>14.5</sub>M<sub>0.6</sub> microwire.

Essentially, non-linear hysteresis loop have been observed at  $H_0$  >110 A/m (Figure 7a). Such peculiar hysteresis loops can be interpreted as the superposition of two hysteresis loops: one from magnetically bistable Fe<sub>74</sub>B<sub>13</sub>Si<sub>11</sub>C<sub>2</sub> microwire (shown in Figure 4a) and the other one from Co<sub>67</sub>Fe<sub>3.9</sub>Ni<sub>1.5</sub>B<sub>11.5</sub>Si<sub>14.5</sub>M<sub>0.6</sub> microwire with linear hysteresis loop (shown in Figure 4b). At intermediate  $H_0$  -values the shape of the hysteresis loop depends on  $H_0$ .

The peculiar hysteresis loop character at  $H_0 \le 120$  A/m can be explained by the partial magnetization reversal of the magnetically bistable wire under the influence of the stray field from the Co-based wire. The stray field is affected by the sample demagnetizing factor and the sample magnetization [54,55]. In the case of Co-rich microwire the magnetization and hence, the stray field are affected by the applied magnetic field (as can be appreciated from the hysteresis loops shown in Figure 4b). In contrast, the magnetization of Fe-rich sample change by abrupt jump and below and above  $H_s$  is almost independent on magnetic field (see Figure 4a).

Accordingly, such microwire array consisting of two microwires (Fe-rich and Co-rich) with different hysteresis loops presents odd and even harmonics quite different from the case of the array with two Fe-rich microwires (see Figures 7 b,c). A single sharp jump of odd and even harmonics is observed at  $H_0 \approx H_s$ . There is also a change in the odd and even harmonics in the weak ( $H_0 < H_s$ ) field region (see Figures 7 b, c).

Further tuning of harmonic spectra is observed in the array consisting of three Fe<sub>74</sub>B<sub>13</sub>Si<sub>11</sub>C<sub>2</sub> and one Co<sub>67</sub>Fe<sub>3.9</sub>Ni<sub>1.5</sub>B<sub>11.5</sub>Si<sub>14.5</sub>M<sub>0.6</sub> microwires (see Figure 8). The hysteresis loop of such array, consisting of three magnetically bistable microwires and one microwire with inclined hysteresis loop, is essentially non-linear and has a complex shape (see Figure 8a).

Basically, the hysteresis loop observed in low  $H_0$  region is similar to those of a single Co-rich microwire (see Figure 8a). The superposition of a linear hysteresis loop and three rectangular hysteresis loops with three Barkhausen jumps is observed with increasing  $H_0$  (see Figure 8a).

The harmonic spectra also reflects the multistep magnetization process of the array consisting of four microwires with different hysteresis loops with regions of gradual changes and abrupt jumps (see Figures 8 b,c).

Thus, the use of arrays consisting of magnetic microwires allows us to create a complex and unique spectra of magnetic harmonics in magnetic microwires.

One of the main features of magnetically bistable amorphous microwires is that such microwires behave similarly to single domain magnets. Such behavior is linked to perfectly rectangular hysteresis loops of Fe-rich microwires (see Figures 4a, Figure 9a) and the magnetostatitc interaction described above. Accordingly, the hysteresis loop of a single magnetically bistable microwire and of the array of magnetically bistable microwires are substantially different. In the case of a single Fe<sub>65</sub>Si<sub>15</sub>B<sub>15</sub>C<sub>5</sub> microwire and arrays consisting of 2, 5 and 10 Fe<sub>65</sub>Si<sub>15</sub>B<sub>15</sub>C<sub>5</sub> microwires the hysteresis loops are rather different (see Figure 9).

The hysteresis loops shown in Figure 9 have been obtained when the microwires in the array were placed touching each other: the distance between the magnetic nucleus was equal to the double glass-coating thickness (7.4  $\mu$ m). For a single microwire, a single Barkhausen jump is observed (Figure 9a). An increase in the number of microwires causes an increase in the number of jumps (see Figures 9b-d) that correlates with the number of microwires.

Observed splitting of the hysteresis loops of the microwires array depends on the distance between the microwires (see Figure 9e). At a distance of about 2 mm such splitting becomes negligible (Figure 9e). It is worth mentioning, that such magnetostatic interaction in Co-rich

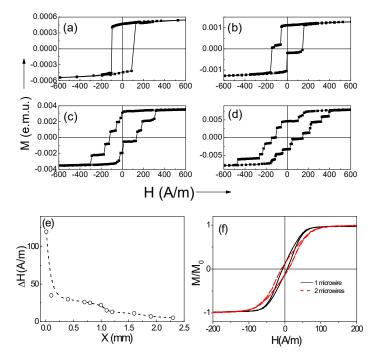


Figure 9. Hysteresis loops of a single glass-coated Fe<sub>65</sub>Si<sub>15</sub>B<sub>15</sub>C<sub>5</sub> amorphous microwire (a) and of arrays consisting of 2 (b), 5 (c) and 10 (d) Fe<sub>65</sub>Si<sub>15</sub>B<sub>15</sub>C<sub>5</sub> microwires, dependence of the hysteresis loop splitting on the distance between Fe<sub>65</sub>Si<sub>15</sub>B<sub>15</sub>C<sub>5</sub> microwires in the two-microwires array (e) and hysteresis loop of the array consisting of 2 glass-coated microwires (f). The line in Figure e is just a guide to the eyes. Figures (a)-(d) and (f) are reprinted with permission from ref. (55), figure (e) is reprinted with permission

Figures (a)-(d) and (f) are reprinted with permission from ref. (55), figure (e) is reprinted with permission from ref. (50).

microwires with inclined hysteresis loops is not quite pronounced. Thus, the presence of the second Co<sub>67</sub>Fe<sub>3.9</sub>Ni<sub>1.5</sub>Bi<sub>1.5</sub>Si<sub>14.5</sub>M<sub>0.6</sub> microwire in two Co<sub>67</sub>Fe<sub>3.9</sub>Ni<sub>1.5</sub>Bi<sub>1.5</sub>Si<sub>14.5</sub>M<sub>0.6</sub> microwires array causes a slight increase in the effective anisotropy field (see Figure 9f). The hysteresis loop shape remains almost the same.

## 3.3. Multi-bit magnetic tags applications of magnetic microwires.

Hysteresis loops with several sharp jumps, observed in magnetically bistable microwires arrays, seem to be suitable for multi-bit magnetic tags. Such magnetic tags, consisting of several magnetically bistable microwires and presenting overall hysteresis loop with several Barkhausen jumps, have been proposed for the magnetic codification method [25,27,56]. In such a magnetic tag, exposed to an *AC* magnetic field, each particular microwire is remagnetized in a different magnetic field, giving rise to an electrical signal on the detection system (see figure 10).

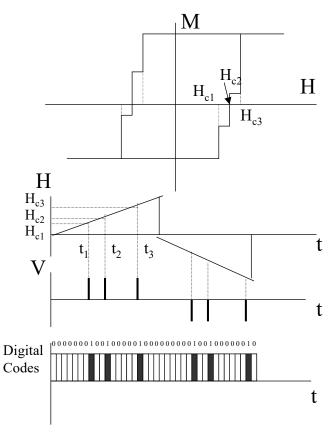
In a magnetic tag consisting of magnetically bistable microwires of the same composition, the hysteresis loop splitting, among other factors, is substantially affected by the distance between microwires. Therefore, multi-bit magnetic tags consisting of magnetically bistable microwires with rather different switching fields are considered more suitable for such application [25,56]. The extended range of switching fields provides a possibility to use a large number of combinations for magnetic codification.

A variety of  $H_s$ -values can be achieved either by compositional  $H_s$  dependence, by the effect of internal stress or thermal treatments on  $H_s$ .

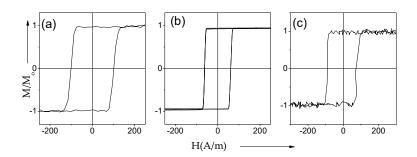
The influence of chemical composition on the  $H_s$ -values of amorphous microwires is originated by the compositional dependence of  $\lambda_s$ : a decrease in  $\lambda_s$  is observed in Fe<sub>x</sub>Co<sub>1-x</sub> amorphous alloys upon doping Fe-rich microwires with  $\lambda_s \approx 40 \times 10^{-6}$  for x=1, by Co up to  $\lambda_s \sim -(5-3) \times 10^{-6}$  for x=0 [57-59]. Similarly, a decrease in  $\lambda_s$  is reported in Fe<sub>x</sub>Ni<sub>1-x</sub> alloys with an increase in Ni content [59,60].

A tendency of decrease in coercivity in Fe<sub>x</sub>Co<sub>1-x</sub> amorphous based microwires can be appreciated from Figure 4, where the coercivity,  $H_c$ , drops from  $H_c \approx 100$  A/m for Fe<sub>75</sub>B<sub>9</sub>Si<sub>12</sub>C<sub>4</sub> ( $\lambda_s \approx 40 \times \cdot 10^{-6}$ ) up to  $H_c \approx 5$  A/m for Co<sub>67</sub>Fe<sub>3.9</sub>Ni<sub>1.5</sub>B<sub>11.5</sub>Si<sub>14.5</sub>Mo<sub>0.6</sub> microwire ( $\lambda_s \approx -0.5 \times \cdot 10^{-6}$ ). The other example illustrating  $H_s$  compositional dependence is shown in Figure 11, where hysteresis loops of Fe<sub>77.5</sub>Si<sub>7.5</sub>B<sub>15</sub> ( $\lambda_s \approx 38 \times \cdot 10^{-6}$ ), Fe<sub>47.4</sub>Ni<sub>26.6</sub>Si<sub>11</sub>B<sub>13</sub>C<sub>2</sub> ( $\lambda_s \approx 20 \times \cdot 10^{-6}$ ) and Fe<sub>16</sub>Co<sub>60</sub>Si<sub>13</sub>B<sub>11</sub> ( $\lambda_s \approx 15 \times \cdot 10^{-6}$ ) microwires are shown.

However, even for microwires with fixed chemical composition the  $H_s$ -values can be tuned by the internal stresses,  $\sigma_i$ , values. The main (though, not the unique) origin of the internal stresses in glass-coated microwires is the different thermal expansion coefficients of the metallic nucleus and



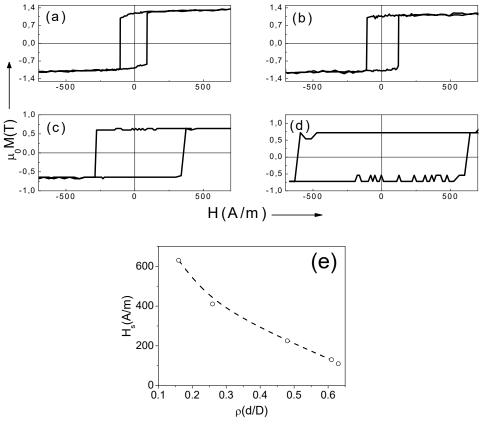
**Figure 10**. Schematic representation of the encoding system based on magnetic bistability of the microwires. Reprinted with permission from ref. (56).



**Figure 11.** Hysteresis loops of as-prepared Fe77.5Si7.5B15 (a), Fe47.4Ni26.6Si11B13C2 (b), and Fe16C060Si13B11 (c) microwires. Reprinted with permission from ref. (60).

the glass coating [36, 61-64]. Accordingly, it is assumed (and experimentally confirmed) that  $\sigma_i$  -magnitude inside the metallic nucleus is affected by the  $\rho$ -ratio between the metallic nucleus diameter, d, and the total microwire diameter, D ( $\rho$ =d/D) [36, 62-64].

As can be appreciated from Figure 12, even for the same microwire composition,  $H_s$  can be modified by almost an order of magnitude (from 85 to 630 A/m) by changing the  $\rho$ -ratio. The correlation of  $H_s$  and  $\rho$ -ratio is evidenced by the  $H_s(\rho)$  shown in Figure 12e.



**Figure 12.** Hysteresis loops of Fe<sub>70</sub>B<sub>15</sub>Si<sub>10</sub>C<sub>5</sub> amorphous microwires with different metallic nucleus diameters d and total diameters D: with  $\rho = 0.63$ ; d=15 µm (a);  $\rho = 0.48$ ; d=10.8 µm (b);  $\rho = 0.26$ ; d=6 µm (c);  $\rho = 0.16$ ; d=3 µm (d) and  $H_s(\rho)$  dependence of the same microwires. The line in Figure (e) is just a guide to the eyes. Adapted from ref. (63).

The main problem with magnetic tags consisting of microwires with different d-values is that the magnetic moments of microwires with different d are rather different. Accordingly, the alternative approach lies in the use of heat treatment allowing internal stresses relaxation keeping the magnetization values the same.

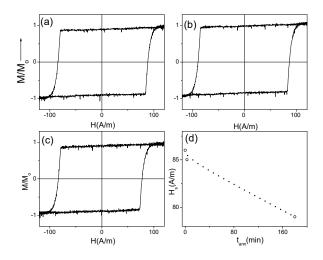
For the Fe<sub>75</sub>B<sub>9</sub>Si<sub>12</sub>C<sub>4</sub> microwires, annealing is not very effective: annealing allows only a slight  $H_s$  decrease (see Figures 13a-d).

Annealing is the more effective route for  $H_s$  tuning in Fe<sub>49.6</sub>Ni<sub>27.9</sub>Si<sub>7.5</sub>B<sub>15</sub> microwires with positive magnetostriction ( $\lambda_s \approx 20 \text{ x}10^{-6}$ ) [61,65,66].

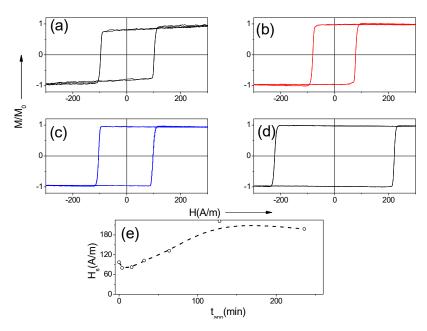
As-prepared Fe<sub>49.6</sub>Ni<sub>27.9</sub>Si<sub>7.5</sub>B<sub>15</sub> microwires present rectangular hysteresis loops (see Figure 14a), as expected for microwires with positive  $\lambda_s$  –values (about 20 x10<sup>-6</sup>). All annealed Fe<sub>49.6</sub>Ni<sub>27.9</sub>Si<sub>7.5</sub>B<sub>15</sub> microwires (at  $T_{ann}$ =410 °C) present perfectly rectangular hysteresis loops. However, a remarkable increase in  $H_s$  is observed in Fe<sub>49.6</sub>Ni<sub>27.9</sub>Si<sub>7.5</sub>B<sub>15</sub> microwires upon annealing (see Figure 14 a-d).

Magnetic hardening, observed in amorphous Fe<sub>49.6</sub>Ni<sub>27.9</sub>Si<sub>7.5</sub>B<sub>15</sub> microwires upon annealing has been explained considering the effect of DW stabilization [60,67,68]. The mechanism of such DW stabilization is linked to directional atomic pair ordering along a preferred magnetization direction during the annealing and generally reported for amorphous alloys with two or more ferromagnetic elements [60].

Observed possibility to tune coercivity of Fe-Ni based microwires by annealing, make them suitable for multi-bit tags applications.



**Figure 13.** Hysteresis loops of as- prepared (a), and annealed at  $T_{ann}$ = 400 °C for 3 min (b) and 180 min (c) Fe<sub>75</sub>B<sub>9</sub>Si<sub>12</sub>C<sub>4</sub> microwires and dependence of coercivity on annealing time (d). The line in Figure (d) is just a guide for the eyes.



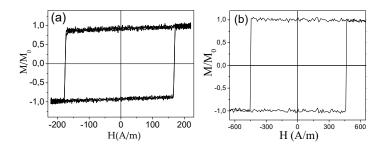
**Figure 14.** Hysteresis loops of as-prepared (a) and annealed at  $T_{ann}$ =410 °C for  $t_{ann}$ = 4 min (b) 32 min (c) 128 min (d) and  $H_c(t_{ann})$  dependence (e) for Fe<sub>49.6</sub>Ni<sub>27.9</sub>Si<sub>7.5</sub>B<sub>15</sub> microwires. The lines in Figure (e) is just a guide to the eyes. (Reproduced with permission from ref. (63) Fig. 4 Copyright © 2020 AIP Publishing LLC, Open Access).

Even wider range of coercivities can be achieved by partial or complete devitrification of amorphous microwires. The main attractive feature of nanocrystalline materials is their magnetic softening upon nanocrystallization [17, 60]. Such magnetic softening is commonly attributed to mixed amorphous nanocrystalline (average grain size of 10–15 nm) structure of properly annealed amorphous Fe– based alloys doped by Cu and Nb [17,60,69]. Such nanocrystalline FeSiBCuNb alloys

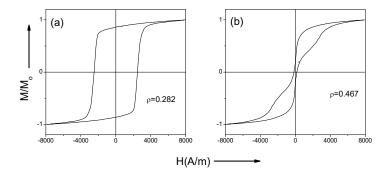
are commonly known as Finemet [60,69]. More recently, another family of nanocrystalline FeCoB-M-Cu (Hitperm) alloys has been proposed [69].

From the viewpoint of tags applications the main advantage of nanocrystalline alloys is the high saturation magnetization [60,69-74].

As shown in Figure 15, as-prepared Finemet-like and Hitperm-like glass-coated microwires also present perfectly rectangular hysteresis loops. In the present case, the Fe<sub>38.5</sub>Co<sub>38.5</sub>B<sub>18</sub>Mo<sub>4</sub>Cu<sub>1</sub> microwire present nanocrystalline structure in as-prepared state [70, 71]. The advantage of as-prepared nanocrystalline materials is that they can present better mechanical properties [70,72]. One the other



**Figure 15.** Hysteresis loops of as-prepared Fe<sub>70.8</sub>Cu<sub>1</sub>Nb<sub>3.1</sub>Si<sub>14.5</sub>B<sub>10.6</sub> ( $\rho$  = 0.38) (a) and Fe<sub>38.5</sub>Co<sub>38.5</sub>B<sub>18</sub>Mo<sub>4</sub>Cu<sub>1</sub> ( $\rho$  = 0.6) (b) microwires. Adapted from refs. (70,71), respectively.



**Figure 16.** Hysteresis loops of Fe<sub>71.8</sub>Cu<sub>1</sub>Nb<sub>3.1</sub>Si<sub>15</sub>B<sub>9.1</sub> microwires with  $\rho$ =0,282 (a) and  $\rho$ = 0,467 (b) annealed at 700 °C. Adapted from ref. (74).

hand, rectangular hysteresis loop can be observed in nanocrystalline microwires devitrified by annealing of an amorphous precursor [70,74]. Thus, rectangular hysteresis loop with  $H_c\approx 2000$  A/m is observed in Fe<sub>71,8</sub>Cu<sub>1</sub>Nb<sub>3,1</sub>Si<sub>15</sub>B<sub>9,1</sub> microwire ( $\rho$  =0.282) annealed at  $T_{ann}$  = 700 °C (see Figure 16a).

In several cases, partial devitrification also allows to obtain peculiar step-wise hysteresis loops [74,75]. Such two-jump-like hysteresis loop observed in Fe<sub>71,8</sub>Cu<sub>1</sub>Nb<sub>3,1</sub>Si<sub>15</sub>B<sub>9,1</sub> microwire ( $\rho$  =0.467) (see Figure 16b) has been explained by mixed amorphous- crystalline (bi-phase) structure [74].

Accordingly, an alternative route allowing to avoid the problems with high precision tag design is the development of partially devitrified microwires presenting multi-step hysteresis loops [74-76].

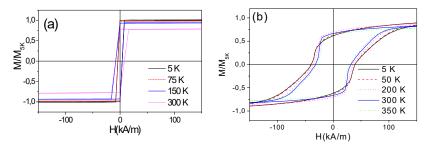
The second magnetic phase can also be created on the glass shell. Accordingly, bimagnetic glass-coated microwires consisting of glass-coated microwire surrounded by an external magnetic microtube have been reported [77]. However, such technology requires one more technological process related to precise sputtering or electroplating of the magnetic microtube [77]. Considering

thousands of security systems and millions of tags required for such applications on a daily basis, such technological scheme can be challenging.

The above described magnetostatic interaction between various microwires requires certain precision and special attention to magnetic multi-bit tag design. Accordingly, appropriate digital algorithms have been developed for the multi-bit tag recognition [26].

On the other hand, magnetically hard and semi-hard microwires are required for the development of smart markers for the electronic article surveillance [78]. A semi-solid magnetic material is proposed as a "deactivating element". When the deactivating element is magnetized, it creates a stray magnetic field that saturates the neighboring soft magnetic element, making the soft magnetic element undetectable by the interrogator used in the interrogation zone.

One of the routes allowing magnetic hardening is the use Fe-Pt base microwires and proper annealing allowing formation of L10-type superstructure [79]. Elevated coercivity,  $H_c \approx 40 \text{ kA/m}$ , has been achieved in properly annealed Fe<sub>50</sub>Pt<sub>40</sub>Si<sub>10</sub> microwires upon devitrification of the amorphous precursor (see Figure 17).



**Figure 17.** Hysteresis loops of as-prepared (a) and annealed at 500 °C for 1 h (b) Fe<sub>50</sub>Pt<sub>40</sub>Si<sub>10</sub> microwires measured at different temperatures. Adapted from ref. (78).

Several alternative routes allowing magnetic hardening include controllable crystallization of Co- or Fe-rich microwires either by Joule heating [80,81] or by conventional furnace annealing [74].

Provided routes for design of non-linear hysteresis loops allow us to consider amorphous and devitrified Fe-, Co-Fe, Fe-Ni and Fe-Pt rich microwires as quite promising candidates for the use in security and electronic surveillance applications. We were able to tune the switching field of magnetically bistable microwires, either by chemical composition of the metallic nucleus, by the internal stresses value (through the glass-coating thickness) or by heat treatment as well as the magnetic field dependencies of even and odd harmonics by magnetostatic interaction between magnetic microwires. A predictable design of non-linear hysteresis loops can serve as a good basis for magnetic tags application using glass-coated microwires.

### 4. Conclusions

We overviewed the properties of soft magnetic glass-coated microwires and the routes allowing to obtain non-linear hysteresis loops either by different post-processing or using magnetostatic interaction between the microwires, making them quite attractive for electronic article surveillance and security applications.

The feasibility studies show that the 5-th harmonics of 3 cm long typical Fe–rich microwire can be detected at a distance up to 25 cm.

We showed, that the presence of neighbouring microwire (either Fe- or Co-based) significantly affects the hysteresis loop of the whole microwires array. In a microwires array containing magnetically bistable microwires we observed splitting of the initially rectangular hysteresis loop with a number of Barkhausen jumps correlated with the number of magnetically bistable

microwires. Essentially, non-linear and irregular hysteresis loops have been observed in mixed arrays containing Fe and Co-rich microwires. The observed non-linear hysteresis loops allowed to increase the harmonics and to tune their magnetic field dependencies.

Non-linear hysteresis loops have been also observed upon devitrification of amorphous microwires.

On the other hand, several routes allowing to tune the switching field by either post-processing or modifying the magnetoelastic anisotropy have been reviewed.

Observed unique combination of magnetic properties, together with thin dimensions and excellent mechanical and anti-corrosive properties, provide excellent perspectives for the use of glass-coated microwires for security and electronic surveillance applications

**Author Contributions:** Conceptualization, A.Z. and V.Z.; methodology, M.I., J.M.B., and V.Z.; validation, A.Z.; formal analysis, A.Z. and V.Z.; investigation, A.Z., P.C.-L, and V.Z.; resources, A.Z., and J.G.; data curation, V.Z., P.C.-L., J.M.B., and M.I., writing—original draft preparation, A.Z., P.C.-L, and V.Z.; writing—review and editing, A.Z., P.C.-L., and V.Z.; supervision, A.Z.; funding acquisition, A.Z., J.G., and V.Z. All authors have read and agreed to the published version of the manuscript.

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