HTS Accelerator Magnets and Conductor development in Europe

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Abstract: In view of the preparation for a post-LHC collider, the high-energy physics (HEP) community started from 2010 to discuss various options, including the use of HTS for very high field dipoles. Therefore, a small program was set in Europe aiming at exploring the possibility of using HTS for accelerator quality magnets. Based on various EU funded programs, though at modest levels, has enabled the European community of accelerator magnets to start getting experience in HTS and addressing a few issues. The program was based on use of REBCO tapes to form 10 kA Roebel cables, to be used to wind small dipoles of 30-40 mm aperture in the 5 T range. The dipoles are designed to be later inserted in a background dipole field (in Nb:Sn), to reach eventually a field level in the 16-20 T range, beyond the reach of LTS. The program is currently underway: more than 1 km tape of high performance (Jc > 500 A/mm² at 20 T, 4.2 K has been manufactured and characterized, various 30 m long Roebel cables have been assembled and validated up to 13 kA, a few dipoles have been wound and tested, reaching at present 4.5 T in stand-alone (while a dipole made from race track coils with no-bore exceeded 5 T using stacked tape cable) and a test in a background field is being organized.

Keywords: High current density HTS; Roebel cables; Accelerator magnets; Collider Magnets; Superconducting HTS Magnets.

1. Introduction

Accelerators have been among the most powerful instruments for scientific discoveries. The identification of the Higgs boson [1-2] at CERN on July 4, 2012 has been the first big discovery of the LHC, the largest instrument for High Energy physics, put in operation in 2008 and then operating in collider mode since march 2010. The discovery of the Higgs boson and its characterization has opened new perspectives for the Physics beyond the Standard Model. While the LHC is striving to reach its design beam energy by pushing the Nb-Ti superconducting magnets [3] to their design value of 8.3 tesla in the next LHC Run 3 of 2021-2024, the HEP (High Energy Physics) accelerator community is actively pursuing the next step: the High Luminosity LHC (HL-LHC or HiLumi) machine [4-5] for which a new generation of superconducting accelerator magnets is now under construction [6-7], meanwhile preparing the further step to go even further with magnet technology

2. Accelerator and magnets

The discovery potential of a hadron collider depends critically on strength of the magnetic field in various way. In a circular accelerator or collider, the energy of the particle beam depends directly on the size of the accelerator and on the main dipole field according to the simple relation (valid for relativistic particle): E ≅ 0.3 B R, where E is the beam energy in TeV (Teraelectronvolt), B is field of the dipoles in tesla and R is the effective radius of the accelerator in km. In LHC the 8.3 T dipole field with the 2.8 km effective radius (physical radius is 4.2 km, however the bending field covers only 2/3 of the 26.7 km tunnel length) yield the 7 TeV beam energy, allowing 14 TeV in the center of mass during collisions.
There is another important figure of merit to qualify a collider as a discovery instrument: the luminosity, i.e. the collision rate per unit of reaction cross-section. It is not enough to produce one only Higgs boson, indeed. Plenty of it must be generated in order to allow its identification and for measuring its subtle property. In a field, like HEP, where the interesting phenomena are more and more rare, surrounded by billions of “useless” interaction (noise), luminosity, i.e. the quantity of useful particles produced during collision and that can be detected, is becoming of paramount importance. That is why CERN is investing more than 1 B$ (material budget only), that is 25-30% of the LHC cost, just to improve the luminosity with the above-mentioned HL-LHC project. It turns out that luminosity depends on the peak field of special magnets in the low-beta insertion. In particular, the so-called Inner Triplet (IT) quadrupoles, controlling the size of the beam (beta) at the collision, are the most critical magnets in term of peak field since they feature both high field gradient and large aperture. For the HL-LHC, peak fields in the range 11-12 T are then required for the large aperture IT quadrupoles for the previously mentioned reasons. In addition, also a few dipoles necessary for increasing the beam intensity (another important factor to improve luminosity), needs to be designed at 11 T, in order to accommodate beam cleaning collimators in the cold region. Therefore, the High Luminosity LHC, with its magnets 30-40% more powerful than the LHC ones, is a novel step in improving the LHC as discovery instrument. These magnets constitute a breakthrough in accelerator technology, indeed, as can be seen in Fig. 1 reporting the progress in magnetic fields for past and future accelerators.

![Magnetic field evolution for Hadron Collider](image)

Figure 1. Progress of the operating field in circular hadron colliders (ill-fated SSC reported for comparison). LHC dipoles are actually working at 7.7 T and are supposed to operate at 8.3 T design field in 2022-2023. HL-LHC has already shown its field level in a few prototypes. HE-LHC (Malta workshop 2010) is a pure paper design. FCC-hh has started dedicated R&D in 2014.

As a further step beyond LHC and its High Luminosity upgrade, CERN has investigated in 2010 an energy upgrade of LHC, called High Energy LHC (HE-LHC) [8], based on filling the present 27 km long LHC tunnel with 20 T dipole field [9], see Fig.1. Based on use of HTS for the inner coil layer, and LTS for the outer coil layers, this project, also in view of the long R&D for the 12 T field level of HL-LHC, soon appeared to be much too ambitious as regarding the technology and too little ambitious as regarding the energy reach (33 TeV center-of mass energy). In 2013, based on the outcome of the new update of the European Strategy of Particle Physics, as a possible next step post-
LHC and post-HiLumi, CERN has proposed an 80-100 km circular collider called FCC (Future Circular Collider) to reach 100 TeV center-of-mass energy. This would require the dipole field to reach in operation 20 T (for 80 km) or 16 T (for 100 km) [10]. Soon the baseline of the dipole field value was fixed to be 16 T, to stay possibly in the reach of Nb3Sn, see Fig.1. HTS was clearly perceived as a technology not mature enough for accelerators, for good reasons. It appears evident on the plot of Fig.1 that going 15-16 T would have been a huge jump. However, the needs of HTS for use in a next higher energy collider either for main dipoles, in the far future, or in special regions of FCC (of high-radiation environment and/or high thermal load) was clear. In addition it was also evident that only a dedicated R&D could assessed if HTS could meet the very demanding specifications for collider magnets. For this reason, a continuous, though modest, effort started in Europe in 2012-2013, devoted to accelerator magnets using HTS.

2. The FP7-EuCARD HTS racetrack magnet

Well before the above-mentioned program for FCC, in 2007 an integrated research activity for accelerator R&D was launched, applying for EU funding in the program FP7 with the name FP7-EuCARD (European Coordination of Accelerator R&D). The EuCARD magnet work package was mainly devoted to develop conductor and magnet technology for the HiLumi LHC magnets, featuring the design and construction of the Fresca2 large bore dipole, with 100 mm free coil aperture and 13 T nominal field at 1.9 K, wound in Nb3Sn. Fresca2 did reach 14.6 T [11] which is at present the record field for a dipole magnet. See Fig. 2 for a schematic view of Fresca 2 and for its training curve.

![Figure2. 3-D sketch and training curve of Fresca2 dipole. The letters a,b and c in the quench plot refer to various assemblies and different prestress configuration. (from ref. 11).](image)

The program had also a task dedicated to the design and construction of an HTS race track with no bore. The EuCARD HTS racetrack was conceived as a basic R&D exercise on HTS for the European accelerator magnet community, and its main features are:

1. The magnet is a stack of flat coils in racetrack shape, with no accessible bore. The coils are wound as flat pancakes. Three pancakes, or coil layers, are above the midplane and three exactly symmetrical are in the bottom coil.
2. The conductor is composed of two REBCO tapes, 12 mm wide, soldered face to face to form a pair that sandwich a pure Cu ribbon in the center (for a total of 70 μm of copper in the sandwich). The stack is about 200 μm thick and then a 130 μm Cu-Be ribbon is soldered on each side, so the total conductor thickness is about 460 μm. Conductor is then insulated with polyester film. Finally, two conductor units are co-wound to form a cable of large current. A schematic of the cable is shown in Fig. 3.
3. The cable above described is not transposed, and the inner conductor of the cable has smaller inductance than the outer one. To compensate for this effect, each pancake on the top part (wrt to the midplane) is connected to its bottom symmetrical companion in such a way that the current in the inner conductor is in series with the outer conductor of the same cable of the bottom pancake. The coil layout and magnet mechanics are depicted in Fig. 4. The magnet is first assembled for the stand-alone test with a demountable structure, where force
and prestress can be adjusted. The structure is also easily demountable to allow for various assembly trials. For the final configuration, as high field insert inside Fresca 2, a more compact structure is necessary to keep the e.m. forces, see Fig. 4.

Figure 3. Left: picture of the cross section of the basic HTS unit (two HTS tapes soldered to a central Cu ribbon); center: scheme of the conductor, with the double tape in between two CuBe ribbons with insulation; right: the cable composed by two independent conductors (courtesy of A. Ballarino, CERN and M Durante, CEA).

Figure 4. The layout of the EuCARD HTS magnet by CEA. Left: view of the three coil layers. Center. First demountable assembly for stand alone test. Right: compact assembly for test in Fresca2. Courtesy of M. Durante, CEA.

The project has been relatively long. Launched initially in 2008 as part of FP7-EuCARD, it suffered from lack of resources and was completed in the years 2013-2017 in the frame of the CERN-CEA collaboration. CERN procured the conductor, manufactured by Superpower, Inc., and CEA was in charge of magnet design, fabrication and test, with some financial support by CERN. The magnet after a few tests successfully reached 5.4 T in stand-alone, which is well consistent with the design (originally was 6 T but then the number of turns was reduced).

Now it is prepared for testing in the Fresca2 facility, and, despite delays due to the Covid-19 emergency, a test is foreseen by 2021.

3. The FP7-EuCARD2 and H2020-ARIES programs: overview

Despite that in 2012-2013 the EuCARD HTS magnet was far from completion, in view of the potential use of HTS for any post-LHC collider, a program called EuCARD-2 (European Collaborative Accelerator R&D-2) was proposed to European Commission-framework programme7, EC-FP7, in 2011, see basic ideas in [12]. The program was then approved with 80% financing in 2012 and finally it started in 2013. The effort on HTS conductor and magnets was named WP10-Future Magnets and was part (about 15%) of the much larger EuCARD-2 programme. Led by CERN and including most Institutions active in accelerator magnets in Europe, as well as a few other laboratories, the EuCARD-2- WP10-Future Magnets program [13] had the main goals of
manufacturing and qualifying in real demonstrator coils and magnets an HTS cable with characteristics useful for accelerator magnets.

The EU funded program lasted four years, ending in April 2017; however since the beginning all work was managed as a long term collaboration, extending for most partners two or three years more, to properly complete the first R&D base. The main tasks were: Conductor development (tapes, cable, characterization), Magnet development (design, technology, construction) and Test of the coil/magnets as stand-alone in a suitable facility capable of variable temperature and high current. The test in a background field (Fresca2) was left out of the EU program, to be pursued later by CERN and CEA. However, the test in high background field was actually dictating the whole design of the magnet.

EC-FP7 provided about 1.3 M€, i.e. about 40% of the initial 3.3 M€ direct cost of the program. The actual direct cost (without overheads) was in total 5.8 M€: the 1.2 M€ additional funding for more conductor and coils were added by the Institutes (mainly CERN). Eleven Institutes collaborated in EuCARD2-Future Magnets:

1. CERN, Geneva, CH, (general coordinator and participating to all tasks and, in particular, responsible for design and construction of one type of magnet and of magnet testing)
2. CEA-Saclay, FR, (in charge of coordinating magnet design and responsible for design and construction of one type of magnet)
3. Bruker HTS (BHTS), Alzenau, DE, the Industry in charge of developing and manufacturing the REBCO tapes (including re-coating after tape punching).
4. KIT (Karlsruhe Institute of Technology), DE, in charge of producing the Roebel cable (punching of the tapes and assembly in cable of meandered tapes)
5. University of Geneva (CH), University of Twente (NL) and University of Southampton (UK), in charge of the various characterizations of conductor, both in the form of tape and cable.
6. INFN-LASA (Milan branch of the Italian Institute for Nuclear Physics), in charge of preparing a test station and carrying out one test.
7. Tampere University in Finland, in charge of quench protection simulation.
8. INP of Grenoble (FR), in charge of a special magnet design
9. DTI, the Danish Institute of Technology, giving support to magnet design and construction.

One of the most successful initiatives of the collaboration was the organization of a series of workshops, WAMHTS, i.e., workshop on accelerator magnet in HTS. The workshops are open to additional laboratories or universities, both in Europe and all over the world, especially the US and Japan. In many of these workshops we had participation of Industry. The first workshop was held in Hamburg in May 2014, the second year of EuCARD-2 and the series is now under the coverage of ARIES (see subsection 3.2). The need to share experience in detail and of mutualizing the resources makes this workshop very useful. The workshops are attended by 50-80 people and the last one, the fifth of the series, was held in Budapest on 11-12 April 2019.

3.1 Objectives of EuCARD-2.

A high Jc material is the necessary, however not sufficient, condition to enable the performance of magnets. Accelerator magnets are demanding, indeed:

- High current density in the coil package (typically 400 A/mm²), at the relevant field, in order to make magnets of reasonable size and affordable cost. For our case, this translates into the following requirements:
  - High current density over the whole tape or wire cross section, called \( J_{\text{engineering}} \). Our goal was \( J_c > 400 \text{ A/mm}^2 \) at 20 T and 4.2 K (taken as a reference operating temperature.), with magnetic field perpendicular to the broad face of the tape (worst direction).
  - A compact cable with high filling factor, to avoid excessive dilution of the current density in the coil package.
  - Thin and robust insulation.
• Operating current in the conductor in the 5-20 kA range, i.e., a 10 kA-class cable.
• Multi-strand conductor with strand transposition and contact resistance among strands low enough to enable current transfer, but high enough to avoid field quality disruption during ramp. Strands are flat REBCO taper rather than round wire as in the LTS cables.
• Control of field quality within a few units (1 unit being 100 ppm of the main field). In this initial R&D phase, a few tens of units are considered sufficient, both for magnetization and for winding geometry.

The scope of the program could be condensed into two main objectives:

1. Produce a 10 kA-class conductor of at least 20-30 m unit length (requiring in total approximately 1 km of 12 mm wide tape); For this we set goals for tape of $J_c = 400-600$ A/mm$^2$ at 20 T, 4.2K.
2. Build various small dipole magnets, of some accelerator characteristics, to qualify the conductor in near-to-operating conditions.

Further objectives, besides getting experience with HTS technology, were learning how to protect a HTS magnet with high current density (we want to operate at 4.2 K), qualifying new quench detection technologies and dealing with screening effects for field quality and quench. Finally, an important objective was learning how to test high current density HTS magnets.

The conductor is fully described in section 4. Here it is important to underline that the critical choice was selecting the type of conductor to use, see ref. [12] where the possibility of using Bi-2212 was initially envisaged. The selection was made at the very beginning of the project, based on various considerations, including magnet design and construction. The first choice was to select the type of superconductor: REBCO tape was preferred to Bi-2212 round wire and the reasons are explained in an official EuCARD-2 document [14], mainly driven by magnet construction considerations and diversification of sources. Roebel cable was chosen for its high filling factor (like the classical Rutherford cable) and because transposition was considered a key feature.

The magnet was fixed with a free bore of 30-40 mm and 200-300 mm of straight part. Three types of magnet were pursued:

1. A coil block design, called Aligned Blocks or AB dipole, i.e., with Roebel cable positioned such as the tape broad face is parallel to field lines for maximizing the current density) and flared ends. This design was pursued mainly by CERN.
2. A classical cosφ design, to explore synergy with a design vastly used in the accelerator domain. This design was taken up by CEA.
3. A coil block design based on stacked tape (rather than Roebel cable like the previous two layouts). The design was proposed and developed by INPG.

Two full dipoles have been engineered, built and tested in AB layout, called FeatherM2 dipoles, and a third one is still being manufactured at present. One dipole in Cosφ layout has been manufactured and is in the final assembly stage for testing. The third design, with stacked table conductor, did not go beyond the conceptual design level.

3.2 Objectives of ARIES

At the end of 2015, well before EuCARD-2 was finished in 2017, the successive European program, ARIES was determined. Since the magnet community was very busy with HiLumi LHC, and scarce R&D resources were still engaged with the new high field Nb:Sn program for FCC and with completion of EuCARD-2, there was little room for a new initiative on HTS for accelerators in Europe. Therefore, we could get only a small amount of resources, about 500 k€, that were devoted to one scope: improve the engineering current density of the REBCO tape produced for EuCARD-2 by Bruker. ARIES, being a pure conductor development program only, its description is entirely included in the next section 4.

4. The conductor R&D programs in EuCARD-2 and ARIES

Even if HTS conductors have been used to generate high field in solenoids for many years [15-18], their validation as technical conductors suitable for accelerator magnets in the 20 T range is in its
infancy. As previously mentioned in the section 3., one of the main requirements for accelerator magnets is a high engineering current density in the basic conductor element (tape in our case), indicated as $J_e$ and defined as the ratio of the critical current to the total conductor cross section. The performance target for $J_e$ at the operating field and temperature is of the order of 500 A/mm$^2$. Only two industrial HTS conductors have achieved the desired $J_e$ performance in the 20 T range: YBCO (or REBCO, used to indicate the use of a Rare Earth, RE, other than Yttrium, Y) and BSCCO-2212. REBCO comes in the form of a thin tape ready for use, where the thin layer of Rare Earth superconductor (few µm thickness) is deposited on a thin (tens of nm to few µm) film which acts as a buffer layer between REBCO and a strong substrate (Hastelloy, stainless steel or NiW) 30 to 100 µm thick. The tape is then coated with a few µm of silver to protect the superconductor and finally some copper (5-100 µm) is added for stabilization either via electrodeposition or via soft soldering. Industrial BSCCO-2212 is produced as a round wire, using procedures and tooling that resemble closely those used for the production of LTS wires. BSCCO-2212 has recently demonstrated to have record $J_c$ values close to 1000 A/mm$^2$ at 4.2 K and 30 T [19]. However, the necessity to perform a high temperature heat treatment at high pressure (up to 100 bar) after winding in order to reach the best performance represents a complication to its penetration in the applications. A strategic decision on conductor taken in the three above mentioned European programs for accelerator magnets was to focus almost exclusively on one of the two materials, i.e. REBCO, and thus avoid excessive diversification. This choice, mainly driven by technical considerations on better mechanical properties and simpler magnet technology, see [14], also provided complementarity among the EU and US-based programs, where BSCCO-2212 was supported by the US-DOE Conductor Development Program [20] and the Bismuth Strand and Cable Collaboration (BSCCO) [21]. After the high $J_e$, another important requirement for accelerator magnets is of small inductance, for use in long magnet strings; hence, the conductor in the winding must have large current carrying capacity, in the range of 10 kA. As the current carrying capability of BSCCO-2212 wires and REBCO tapes at the envisaged operating conditions is in the 100 – 1000 A range, it is an imperative to use multi-strand cable, of compact type. While the choice of the Rutherford cable configuration is straightforward in the case of BSCCO-2212 round wire, it is not compatible with REBCO thin tapes. However, alternative geometries appear to be suitable for the high-current density required by accelerator magnets: stacked tapes in various configurations [22-24], Conductor on Round Core (CORC®) [25] and Roebel Cable [26]. Special stacked tape cables were used for the EuCARD race track insert [27], while in EuCARD-2 the Roebel configuration was preferred after a review of the various cable options [28].

Differently from the EuCARD activity on HTS, which contained only a task on the design and construction of an insert, EuCARD-2 focused on a strong HTS conductor program linked to a magnet R&D part, the goal being the creation of a bridge between materials and large scale magnets. REBCO tapes were mainly provided by Bruker HTS (BHTS), a member of the EuCARD-2 consortium. While the development of other tape manufacturers has been initially driven by the perspectives of applications in the electrical utility sector, the R&D program of BHTS was dedicated from its beginning to develop and produce long-length tapes for high resolution NMR spectrometers [29], whose targets in terms of high current density at high field and low temperature have a considerable synergy with the requirements for accelerator magnets. The targets and the minimum required performance for the REBCO tapes considered within the scope of EuCARD-2 and ARIES are reported in Table 1.

<table>
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<th>Parameter</th>
<th>EuCARD-2</th>
<th>ARIES</th>
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<tr>
<td>$J_e(4.2$ K, 20 T) (A/mm$^2$)</td>
<td>600</td>
<td>1000</td>
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<td>400</td>
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EuCARD-2 set a \( J_c \) ultimate target of 600 A/mm\(^2\) at 4.2 K and 20 T in a perpendicular field orientation, with a minimum \( J_c \) of 400 A/mm\(^2\) to be considered for magnet design. The goal of ARIES was to increase further \( J_c \), aiming at a minimum of 600 A/mm\(^2\) and targeting finally to 1000 A/mm\(^2\). Being these the first R&D activities on HTS conductors for accelerator magnets, it was decided to reduce demands for the homogeneity of the critical current in a unit length and for the magnetization amplitude. Both values reported in Table 1 are more than ten times larger than what was achieved in production Nb-Sn wires of HiLumi. As for the mechanical characteristics, a maximum transverse compressive stress of 100 MPa and longitudinal applied strain of ±0.3 % were specified, under which the superconductor shall retain mechanical integrity and a minimum of 90 % of its virgin \( J_c \) (before loading). The levels above are minimum requests for a high field accelerator magnet, and are not particularly challenging for REBCO tapes, whose metallic substrate is typically a hard metal with high elastic modulus and large elastic strain limit. The last target concerns the unit length (UL) of superconductor. A minimum length of the order of 50 m was required, which was largely in excess of the length of a pole in both the aligned-blocks and the cos9 designs. At the end piece lengths of 30 m were also accepted (minimum length to wind a pole was 26 m).

4.1. Progress in the performance of REBCO tapes: fabrication

The development of practical superconductors based on REBCO relies on two common features: a bi-axially textured template, consisting of a flexible metallic tape (Hastelloy, stainless steel or NiW) coated with a multifunctional oxide barrier, and an epitaxial REBCO layer. The textured template, which is needed to eliminate all but low-angle grain boundaries in REBCO, is created by either deforming the metal substrate with rolling assisted bi-axially textured substrate technology (RABiTS) [30] or by texturing the buffer layers deposited on the metal substrate by so-called ion beam assisted deposition (IBAD) [31] and by its variant, alternating beam assisted deposition (ABAD) [32]. The epitaxial REBCO layer is grown either by physical routes, such as pulsed laser deposition (PLD) [33–35] and reactive co-evaporation (RCE) [36, 37], or by chemical routes, such as metal organic deposition (MOD) [38] and metal organic chemical vapor deposition (MOCVD) [39].

BHTS was the partner responsible for manufacturing the tape in both EU funded programs EuCARD-2 and ARIES. CERN procured further 500 m tapes with its own funding, of the same specifications of EuCARD-2, doubling the quantity manufactured by BHTS to give some continuity to the production line and making possible to manufacture more coils. The technological chain of BHTS was established for the deposition of YBCO on Cr-Ni stainless steel tapes with widths of 4 or 12 mm. ABAD is the proprietary process of BHTS employed for the deposition of the yttria-stabilized zirconia (YSZ) buffer layer, typically 2–3 μm thick, which serves as a template for the bi-axially textured growth of YBCO. The YSZ layer has thus an equivalent function to the MgO layer resulting from the IBAD process used by other manufacturers, but it is approximately 5 to 10 times thicker. The YBCO layer, whose typical thickness lies in the 1.5–2 μm range, is deposited on the buffered stainless steel tape by PLD using a drum-based tape transport. The tape is wound on a tubular drum and transported through a deposition zone without any rewinding. For the EuCARD-2 and ARIES productions, a new machine, called PLD-300 and jointly owned by BHTS and CERN, was set up for this process: its drum is capable of producing unit lengths of about 300 m of 4 mm-wide tape or, alternatively, 100 m of 12 mm-wide tape. While an obvious limitation of length is present, this technique provides the benefit of avoiding contamination of the tape by friction, which may be generated in a reel-to-reel deposition process.
Figure 5. Evolution of the engineering critical current performance of BHTS YBCO tapes and EuCARD-2 and ARIES $J_e$-targets. $J_e(B)$ is measured at 4.2 K in perpendicular orientation for three tapes: T002, produced at the beginning of EuCARD-2, T2290, produced during EuCARD-2, and Q065-18, produced during ARIES.

Figure 6. Schematic view of the layered structure of a REBCO tape.

The production process of BHTS was fine-tuned for deposition on a 100 μm-thick tape. This relatively high thickness enables favorable mechanical properties that are important for ultra-high field solenoidal magnets, but is deleterious for the $J_e$ performance. The value of $J_e$ at 4.2 K, 20 T for the material produced at the beginning of the EuCARD-2 collaboration, in 2013, was below 300 A/mm$^2$ (see Fig. 5). Given the layered structure of REBCO tapes, shown in Fig. 6, there were three possible strategies to enhance $J_e$:

i) Increasing the thickness of the superconducting layer;

ii) Increasing the critical current density, $J_c$, of the superconducting layer, introducing artificial pinning centers;

iii) Reduction of the substrate thickness.

The growth of thick REBCO films appears to be a straightforward opportunity in improving $J_e$. However, common to most REBCO film growth techniques, the superconducting critical current density $J_c$ was typically found to strongly deteriorate with the thickness of the REBCO layer, $t_{REBCO}$, when $t_{REBCO} \geq 3 \mu m$ [40-42]. As $J_e = J_c \cdot t_{REBCO}/t_{tape}$, where $t_{tape}$ is the total thickness of the tape, the increase of $J_e$ is jeopardized by the competition between the increase of $t_{REBCO}$ and the resulting degradation of $J_c$. A significant research effort has been devoted worldwide towards improving the in-field $J_c$ by introducing nanoscale defects, e.g. BaZrO$_3$ (BZO) [43,44] and BaHfO$_3$ (BHO) [45], in the REBCO matrix. In-field $J_c$ depends importantly on the angle between the field direction and the REBCO crystallographic orientation: the anisotropy makes $J_c$ maximum when the field is parallel to
the tape surface and minimum when it is perpendicular to the tape surface. Moreover, the ratio of the current densities in the two orientations tends to increase with increasing magnetic field. By itself, the anisotropy is not detrimental but it introduces an extra layer of complexity in the magnet design process. In a magnet, the conductor is exposed to a magnetic field whose angle depends on its position in the winding. This implies that the load line at each point of a REBCO-based coil virtually intercepts a different Jc(B) curve. A careful adaptation of the magnet design process to the specific anisotropic tape characteristic is necessary in order to determine safety margins of each design. Depending on the technique adopted for the REBCO film growth, the morphology and size of the artificial nanoscale defects and thus their effects on the pinning properties are different. In particular, PLD and MOCVD lead to the formation of anisotropic nanorods (~5 nm in diameter, >100 nm in length) of BZO or BHO that grow along the c-axis [46,47]. The intrinsic weak pinning in the magnetic field parallel to the c-axis, i.e. perpendicular to the tape surface, is thus reinforced by the interaction between nanorods and vortices.

Substantial progress was done in this direction by BHTS. A significant increase of Jc in high magnetic field was recorded in tapes manufactured employing the so-called double-disorder (DD) route [29,48,49], which creates intrinsic and extrinsic defects in the YBCO layer. The intrinsic disorder is caused by local stoichiometry deviations with respect to the nominal YBa2Cu3O7-x composition and is achieved through a tailored modulation of the oxygen pressure during the PLD process. The O2 pressure variation causes a change of mean-free path for the different atoms in the plasma plume generated by laser ablation of the YBCO target. The result is the formation of different defects of crystallinity (e.g., stacking faults, nanoislands, nano-precipitates, etc.) in the YBCO film that serve as pinning centers. The extrinsic disorder is obtained by 5 wt% content of BZO in the YBCO target, resulting in the appearance of local and organized BZO-nanorods mostly oriented along the c-axis of the YBCO film. A total length of more than 600 m of YBCO tape was produced by BHTS in the frame of EuCARD-2, to which 500 m of CERN-procured post-EuCARD-2 length should be added, in lengths from short samples (order of 1 m), to widths and lengths that were relevant to the conductor needs for the dipole demonstrators (longest 12 mm-tape with homogeneous properties is about 90 m). As shown in Fig. 5, the target Jc of 600 A/mm² at 4.2 K, 20 T set at the beginning of the EuCARD-2 activities was achieved in perpendicular field orientation, i.e. when Jc is minimum. Thanks to the optimization of the DD-route, the latest tapes systematically exceeded the target value, with the highest Jc extrapolating towards 900 A/mm² at 4.2 K, 20 T [48].
Based on the technical success of EuCARD-2, the primary goals of ARIES were set to further increase the target $J_c$ by a factor of 1.5–2 and to industrialize the process for long lengths (> 25 m). The most straightforward way to reach the desired performance appeared to be the reduction of the stainless steel substrate thickness from 100 to 50 μm. Nonetheless, the implementation of this strategy proved to be harder than expected. Major core steps in the ABAD and PLD processing routes had to be re-tuned because of a substantial variation of mechanical and thermal parameters of the thinner substrate [49]. In particular, the deposition of the relatively thick YSZ buffer layer on the 50 μm substrate determined a pronounced transverse bending of the tapes, a bow-shape that was traced back to the thermal contraction mismatch between the ceramic layer and the stainless steel (see Fig. 7). Even though this effect does not seem to have a great influence on the critical current, it can be detrimental for tape manipulation during cabling and winding. In particular, the mechanical positioning for tape cutting, i.e. the meandering formation during Roebel cable formation described in section 4.3, certainly becomes more difficult. To mitigate the bow-shape, BHTS also started depositing the YSZ film on the back of the tape. As a result, the tape has a significantly lower bowing height, as shown in the tape on the left-side of Fig. 7. Despite all these challenges, BHTS has been able to produce tapes with a very high critical current density. The Cu stabilizer has been tailored to obtain a good $J_c$ overall, with a mix of standard 2 x 20 μm Cu thickness and 2 x 7 μm Cu thickness. Several units between 15 m and 75 m reached or exceeded the minimum $J_c$(4.2 K, 20 T) of 600 A/mm². The best unit, a 12 mm-wide tape originally 25 m-long and prepared with 2 x 7 μm Cu stabilizer, achieved a $J_c$ of 1153 A/mm² at 4.2 K, 19 T, in perpendicular field orientation (see Fig. 5). Figure 8 reports the magnetic field dependence of $J_c$, in perpendicular orientation, at various temperatures between 4.2 K and 40 K for an equivalent tape prepared with a thicker 2 x 20 μm Cu stabilizer: the value of $J_c$(20 K, 19 T) is 430 A/mm², i.e. it exceeds at 20 K the minimum value set for $J_c$ at 4.2 K in EuCARD-2. This gives evidence of the substantial impact of the developments driven by the accelerator magnet R&D to the recent progress in the technology of long REBCO tapes.

4.2. Progress in the performance of REBCO tapes: characterization

The high $J_c$ in high magnetic fields of REBCO tapes is a necessary but not sufficient condition in the perspective of their use in accelerator magnets. To qualify REBCO tapes for operation in very high-field coils, magnet design relies on the control of electrical, mechanical and thermal properties of the superconductor. The performance in these respects is strongly influenced by fabrication route,
Figure 9. Radar charts summarizing the main properties of REBCO tapes from six industrial manufacturers: American Superconductor (a), BHTS (b), Fujikura (c), SuNAM (d), SuperOx (e) and SuperPower (f). Compilation of data from Refs [36-38].

tape layout and materials, and these parameters vary from one manufacturer another. It was already pointed out how a detailed knowledge of the anisotropy of the critical current, $I_c$, at the operating temperature and field is important for magnet design. Moreover, the characterization of the temperature dependence of $I_c$ between the operating temperature and the current sharing temperature, $T_{cs}$, is necessary to consolidate the electromagnetic models simulating the coil behavior in case of quenches. Concerning the quench protection, the operating temperature margin is significantly larger compared to LTS-based magnets. This results in a much slower normal zone propagation velocity that makes quench detection difficult. Management of the mechanical loads imposed by the electromagnetic forces during operation is another critical issue.

In the frame of EuCARD-2, the Group of Applied Superconductivity at the University of Geneva performed an extensive measurement campaign to compare the electromagnetic, electromechanical and thermophysical properties of tapes from various industrial manufacturers and a detailed analysis is reported in Refs. [50-54]. The results obtained on tapes from American Superconductor, BHTS, Fujikura, SuNAM, SuperOx and SuperPower are summarized as radar charts in Figure 9, which provides a snapshot of the performance of industrial tapes in 2015, at the time of the project. The axes
of the radar charts report \( J_e(4.2 \text{ K}, 19 \text{ T}) \) and \( J_c(77 \text{ K}, \text{ self-field}) \), the residual resistivity ratio (RRR) and the area fraction of the Cu stabilizer, the cross section area of the tape, the irreversible strain and irreversible stress limits. Some of these properties are discussed in more detail in the following. However, Figure 9 gives an instant picture of the large spread of performance among manufacturers.

Figure 10. The stress–strain dependence of six industrial REBCO tapes is shown at 4.2 K, with the indication of the Young modulus and yield strength values. Adapted from Ref. [51].

Figure 11. Normalized critical current \( I_c/I_{c,\text{max}} \) versus longitudinal stress \( \sigma \) and irreversible stress limits \( \sigma_{\text{irr}} \) at 4.2 K, 19 T for five industrial REBCO tapes. Adapted from Ref. [51].

Figure 10 shows the relation between longitudinal strain and mechanical stress at 4.2 K for the various tapes. Young’s moduli are in the 155–187 GPa range, highest for BHTS and lowest for SuperPower [51]. All curves are nonlinear even at low strains due to the presence of the Cu stabilizer which quickly reaches its yield point. The BHTS stress-strain curve is more rounded compared to the other samples because of its stainless steel substrate, which makes it the sample with the lowest yield strength, 736 MPa and 499 MPa for the tapes with the substrate of 100 \( \mu \text{m} \) and 50 \( \mu \text{m} \), respectively. The highest yield strength is observed for the REBCO tape from SuperOx, about 1000 MPa, and this is mainly due to its strong Hastelloy substrate and to the low Cu cross section area. From the diverging stress-strain behavior one would expect these tapes to also differ significantly in their electromechanical properties. However, there are no differences in the longitudinal stress dependencies of the critical current below 600 MPa, as shown in Fig. 11 at 4.2 K, 19 T. The irreversible stress limits of all examined samples are in the 740–840 MPa range [51]. Due to its rounded stress-strain relation combined with low yield strength, the BHTS tape with 100 \( \mu \text{m} \)-thick substrate exhibits...
the lowest irreversible stress limit among those reported in Fig. 11. The BHTS tape with 50 µm-thick substrate, produced in the frame of ARIES, has not yet been tested, but it is reasonable to expect a further reduction of the irreversible stress limit, due to the lowered yield strength.

REBCO tapes are thermally and electrically stabilized with Cu, added to the tape either by colamination or by electrodeposition, the total thickness of the stabilizer ranging between 10 and 100 µm. If a local heating occurs during operation, heat and current transfers along the tape through the stabilizer are the main channels for preventing a quench in insulated coils. In the case of non-insulated (NI) windings, conduction in adjacent turns is supposed to become the dominating mechanism. In any case, an adequate amount of Cu is needed, with sufficiently high RRR. RRR > 100 is routinely achieved in LTS wires. The Cu stabilizer in REBCO tapes exhibits significantly lower RRR values, ranging between 14 and 61 for the samples reported in Fig. 9. In particular, the BHTS tape from one of the EuCARD-2 batches had a value of RRR = 23. Through a progressive tailoring of the electrodeposition process, the tapes produced for ARIES regularly reached RRR > 50.

Figure 8. Temperature dependence of the thermal conductivity of REBCO tapes from six industrial manufacturers at B = 0 T (a) and B = 19 T (b). Adapted from Ref. [52].

Focusing on the thermophysical properties, the thermal conductivity of the tape, $\kappa$, which is one of the most relevant parameters governing the thermal stability, can be estimated with a reasonably good accuracy in zero field from the RRR and the cross sectional area of the Cu stabilizer using an analytical formula. The $\kappa(T)$ curves measured at $B = 0$ T and $B = 19$ T on the six industrial REBCO tapes are reported in Fig. 12 (a) and (b), respectively. Comparing the results shown in Fig. 12 (a) and (b), it follows that the differences among tapes from various manufacturers observed at $B = 0$ T are strongly reduced after the application of an intense field [52]. This is the consequence of two competing effects: (i) different relative reduction of $\kappa$ on applying the magnetic field; (ii) different amount of Cu present in the tapes. The field-induced reduction of $\kappa$ in Cu depends on the RRR and samples with higher purity, i.e. Fujikura and SuNAM, show larger variations. On the other hand, since the Cu contribution to the longitudinal thermal conductivity of the tape is proportional to the
Figure 13. Schematic illustration of a Roebel cable. Adapted from Ref. [28].

A Roebel cable, shown in Fig. 13, is made of meandered-shape punched tapes, assembled in a compact configuration. The compaction, defined as the ratio of tape to cable engineering current density, reaches values above 80%, which is similar to Rutherford cables made out of round wires. In EuCARD-2 Roebel cable was also preferred to other conductor geometries for its full transposition, which is needed to enforce equal current distribution among the tapes. This is a basic paradigm that holds for LTS magnet conductors and its verification for HTS was one of the tasks of the magnet development activity. The main disadvantage of the Roebel cable is that it requires punching of the tape, with loss of costly material and potential degradation associated with the fact that one edge of the tape is cut open. This last issue was solved in EuCARD-2 by changing the fabrication sequence, i.e. punching the tapes right after the Ag layer deposition and then coating the resulting meandered tape with Cu. The so-called punch-and-coat approach, opposed to the usual coat-and-punch sequence, offered the advantage of sealing the slit sides of the tape and thus avoiding exposing the superconducting REBCO layer to air. In a later stage of EuCARD-2 this process was further refined by a special Ag coating of the punched surface before coating with copper.

A cable width of 12 mm was chosen for the winding of both magnet variants developed in EuCARD-2, i.e. aligned-blocks and cosθ (see section 5). The width of the cable corresponds to the initial width of the tape, which is then cut in meanders with a residual tape width of 5.5 mm. The initial design proposed by KIT, which was the EuCARD-2 partner in charge for the manufacture of Roebel cables, accommodated 15 tapes with a 226 mm transposition pitch. However, it turned out to be delicate to use when winding over small radii, as tapes tend to slide differentially in the cable, as well as for the layer jump because the tapes buckle when the cable is bent in the non-easy direction. The experience from winding tests thus triggered an iteration on the cable geometry. This was an important step to include in the cabling and winding validation as well as cables made of tapes that were procured from manufacturers other than BHTS. In practice, two cable geometries were defined.
to achieve the required critical current target: (i) a cable made of 15 tapes, based on 100 μm-thick tapes (as obtained with a 50 to 60 μm substrate and 40 to 50 μm of Cu stabilizer); (ii) a cable made of tapes of 140 μm thickness (as obtained starting with a 100 μm-substrate). These are typically the BHTS tapes produced for EuCARD-2 and described in section 4.1. Because of the higher $J_c$ compared to other manufacturers, the number of tapes in the cable was reduced to 13, still achieving a sufficiently high current density. For both geometries the cable transposition pitch was increased to 300 mm, to leave some additional space for slippage of single tapes during winding.

![Figure 14](image_url)

**Figure 14.** Critical current versus applied transverse pressure for three 15-tape REBCO Roebel cables: Cable #1 and Cable #2 are assembled from SuperPower tape and impregnated with Araldite CY5538 and CDT-101K, respectively. Cable #3 is made with BHTS tape and impregnated with CDT-101K. Adapted from Ref. [61].

The use of cables with an odd number of tapes in both designs emerged from a study of the effective section that experiences stress under transverse loads [57]. Once wound in a dipole, electromagnetic forces are mainly acting normal to the broad face of the cable. REBCO tapes have, in principle, an excellent transverse compression strength when the load is applied on the broad face [58]. However, Roebel cables have an inhomogeneous thickness along their length and across their width, which results from the punching procedure, non-uniformities in the tape thickness as well as from the presence of crossovers. This induces an uneven distribution of the transverse compressive stress applied to the cable and may lead to local stresses in excess of the mechanical limits of REBCO tapes. Moreover, a small transverse effective section may result in large inter-tape contact resistance. Therefore, management of the transverse effective section is essential for both mechanical and electrical reasons. The 2D geometrical model of Fleiter et al. [57], which determines the relation between crossing angle of the meander tape, transposition length and number of tapes, provided the guidelines to optimize the cable design and obtain effective sections above 50% of the total transverse section.

4.3.2. Transverse stress tolerance of REBCO Roebel cables

REBCO Roebel cables are being developed in view of 20 T class accelerator magnets. The electromagnetic forces generated in such magnets will lead to sizeable transverse stress on cables, which in the present designs is of the order of 150 MPa [59]. Due to the reduced effective surface, the critical current of bare Roebel cables already starts to degrade at stress levels as low as 40 MPa [58,60]. However, proper impregnation with epoxy resin can dramatically improve the transverse pressure tolerance helping the stress redistribution by increasing the effective section supporting the pressure.
At the University of Twente, Roebel cables with an architecture directly relevant for the EuCARD-2 magnet program were investigated with a variable transverse mechanical load at 4.2 K in a 10.5 T perpendicular magnetic field [61]. Two alternative impregnation methods were tested: (i) CTD-101K, which is used at CERN for the impregnation of large-scale magnets, and (ii) Araldite CY5538 mixed with FW600 EST fused silica powder, following a recipe from KIT [61]. Figure 14 reports the results of the experiments performed on three cables. All cables were manufactured at KIT and comprise 15 tapes with a transposition length of 226 mm. Cable #1 and Cable #2 were assembled from SuperPower tape and impregnated with Araldite CY5538 and CDT-101K, respectively. Cable #3 was made with BHTS tape and impregnated with CDT-101K. The critical current at 4.2 K, 10.5 T of the Roebel cable made with BHTS tape is 7.75 kA, i.e. 2.5 to 3 times higher than that of the cables made with SuperPower tapes with similar architecture. This is a consequence of the tailored developments of BHTS to enhance the performance of the tapes at low temperature, high field. All impregnated cables exhibit a remarkable tolerance to transverse stress and satisfy by far the design requirements of the presently envisaged HTS accelerator magnet demonstrators: no critical current degradation is observed up to 440 MPa in Cable #1 and Cable #2, and up to 370 MPa in Cable #3. Only at higher stress levels, the cables show a gradual but irreversible reduction in critical current. The cross-sections of the cables were examined by optical microscopy to check for visible damage of the tapes due to the transverse loading. Interestingly, Cable #3, which was made with the punch-and-coat approach, is the only one that does not exhibit signs of tape delamination [61].

4.3.3. AC loss and inter-tape resistance of REBCO Roebel cables

EuCARD-2 also promoted investigations on the magnetization, AC loss and inter-tape resistance of REBCO Roebel cables [62-64]. Two measurement campaigns were carried out at the University of Twente and at the Southampton University on 15-tape SuperPower-type Roebel cable samples prepared at KIT and impregnated with CDT-101K at CERN. The experimental results were used as input for an advanced electrical cable network simulation model proposed by Van Nugteren et al. [62], which was needed to analyze critical issues in the design of the magnets developed in EuCARD-2, such as thermal stability, quench propagation and dynamic magnetic field quality.

At Southampton University, AC loss inductive measurements were performed with an alternating magnetic field applied perpendicular to the broad face of the cable in a temperature range between 5 K and 90 K and at a frequency of 5 Hz [63]. The main result was that the 15-tape Roebel sample behaves as two in-line magnetically coupled stacks, each stack of 7–8 tapes. Moreover, the AC loss is clearly hysteresis-dominated. This conclusion was confirmed by the group at the University of Twente: measurements were performed calorimetrically and inductively at 4.2 K and low frequency of 0.01–0.1 Hz, both in perpendicular and parallel field orientation. The coupling loss was found to be lower than the hysteresis loss in both orientations, within the range of the experiments [64]. However, these observations substantially differ from earlier results extracted from a similar cable but impregnated with the alumina-filled epoxy resin CTD-101G, which showed considerable coupling loss when exposed to magnetic field parallel to the broad face of the cable [62].

The experiments from Gao et al. [64] were also complemented by direct measurements of the inter-tape resistance. A tape in a Roebel cable changes position along the longitudinal direction, thus within one transposition length every tape is in contact with two neighboring tapes. The inter-tape resistance is defined as the contact resistance over a single transposition length. It was found that its value at 4.2 K is in the range of 1–10 μΩ, leading to a surface resistance of 0.5–10 mΩm², which is not significantly different from the values measured for the LTS wires in a Rutherford cable. These are encouraging results as they show that a balance is possible between coupling loss and current sharing to obtain thermal stability and adequate magnetic field quality.

5. Magnet(s) design and technology

The strong Jc anisotropy vs. field direction, the large tape with shielding of perpendicular field, the intrinsic variation of the field inside the coil and the ramping regime, require a big effort for proper modeling. A quite detailed multi-physics e.m. model has been set up and is described
elsewhere [65,66]. It determines the peak field for each direction and the critical surface in each point and the current distribution. This last is also computed during a quench, which is a major breakthrough in the e.m. modelling, given the difficulty in protecting HTS magnets. This code has allowed to design and optimize our EuCARD-2 magnets.

5.1 Reference magnet design: AB Feather Magnets

The main objective of the magnet demonstrator was the validation of the HTS conductor by generating a 5 T dipole field in a free cold bore of 40 mm, as envisaged for High-Energy LHC dipoles [9], the reference project at the start of EuCARD-2. The design is based on a rectangular coil block dipole lay-out where the conductor (i.e the wide face of the tapes composing the Roebel cable) in each block is aligned to the main field, the Aligned Block (AB) dipole FeatherM2 [65, 66], see Fig. 15. The field direction considered for the alignment is the one when the AB Feather magnet is excited inside the 13 T Fresca2 background field. The AB Feather magnet generates a field, at nominal current density of $J_e = 400 \text{ A/mm}^2$, of 5 T. Inside Fresca2 the total field should sum up to 18 T.

In stand-alone mode, this dipole should generate 5 T, too, despite that the field lines are not perfectly parallel to the conductor wide face, due to the different field configuration. The extra margin due to lower field (5 T instead of 18 T) should compensate for the reduction due to non-negligible transverse field.

![Figure 15. Crosse section of the coils of the Aligned Block dipole. The alignment is optimized when the dipole is inserted in the Fresca2 background field (see plot at the right), so is not as good as when the dipole is powered in stand-alone mode (left plot), [65].](image)

The structure is based on an external support which is a pre-compressed stainless steel shell [67,68], see Fig. 16. A key feature is that the pre-compression does not need to completely counteract the e.m. forces during magnet excitation: because of the large temperature margin of HTS, small movements of the conductor should not lead to a quench. Therefore, the coil package can be inserted in the structure with some tolerance, making the assembly quite easy. The conductor during magnet excitation leans against the structure that has to be rigid enough to minimize the deformation. Interference with the Fresca2 structure is indeed forbidden by design. To make the structure rigid enough, avoiding excessive thickness of the outer restraining shell, the adopted solution was to link the outer shell to the inner structure with a midplane plate, see Fig.10. This reduces the free bore to 35 mm, but it is a temporary solution related to the test condition of Fresca2. The structure has an easy job in standalone at 5 T, however at 18-20 T the forces are quite high. The e.m. model predicts that current is strongly non-uniform, because of the shielding properties of high-$J_c$ REBCO tape and – to lesser extent - the partial coupling between tapes in the cable. A detailed analysis of the stress concentration due to high $J$ and B, as well as to non-uniform current distribution is reported in [68], where the mechanical structure of FeatherM2 is also described.
Figure 16. One-quarter cross section of the AB FeatherM2 dipole (courtesy of J. van Nugteren and J. Murtomaki, CERN). The midplane steel connecting the inner and outer shell to keep stresses at 18 T is indicated. At the top: the Roebel cable used to wind the coils.

The coils are insulated with glass fiber braid that is sleeved on the cable before winding (the short length of the cable, less than 30 m, makes this operation quite easy). Then the coils are vacuum impregnated with epoxy resin (CTD 101K), carefully avoiding resin-rich zones to avoid delamination due to the high thermal contraction.

One further remarkable characteristic of the AB dipole is the use of copper rings in between coil layers and the outer shell. The ring imparts to the coil the radial force it receives from the restraining cylinder; however, most important, these rings are a well coupled inductor, to quickly extract the energy out of the coil following a quench and avoid an excessive hot spot temperature due to low quench propagation velocity. The University of Tampere (Fin) has developed a model to evaluate the quench propagation for our HTS coils from basic properties [69,70], which greatly complemented the CERN e.m. numerical model.

To get experience with Roebel cable, the AB FeatherM2 dipoles were preceded by a smaller single coil winding, a six-turn simple flat racetrack, called FeatherM0 (with no cable alignment), see [65,66 and 71]. A sketch of FeatherM0 and Feather M2 is reported in Fig. 17.

Two FeatherM0 coils and two AB FeatherM2 dipoles were manufactured and tested, see Fig. 18. The cable for winding two further coils in order to assemble a third AB FeatherM2, has just now been completed [72].
Figure 17. Sketch of the AB FeatherM2 dipole (left) and, in scale, of the coil test FeatherM0 (right).

Figure 18. FeatherM0.5 coil after construction (top left). AB dipole FeatherM2.3-4 during winding (top right) and once assembled before yoking (bottom).

5.2 The EuCARD-2 Cos⁹ dipole

The EuCARD-2 collaboration has also pursued the use of Roebel cable in the classical cos⁹ dipole configuration, the reference lay-out for accelerators [73]. The Cos⁹ configuration inherently has the largest field perpendicular to the wide face of the cable, a penalty for REBCO Roebel cable. However, this demonstrator is a good direct comparison between LTS and HTS technology. Since recent REBCO tapes have reduced the anisotropic behavior, thanks to various additional pinning mechanisms, though the “isotropization” is more marked at higher temperature, we think this is a very useful exercise. The effort is carried out by CEA, Saclay [74-76], see Fig. 19. The design is complex, both because of mechanical reasons and because Roebel does not behave as good as a Rutherford cable made of round strands. Also, the mechanical design for high field impeded the use of two layers since forces were too high, so the coil is a single layer cos⁹. The coil ends turned to be difficult and various tests and a full coil prototype with dummy cable were needed to work out this and other technological issues, see Fig. 20. However, CEA has now overcome all issues, and the three coils (two needed and one spare) have just been wound and impregnated, see Fig. 20. The insulation
and impregnation technology is quite similar to the one of the AB Feather dipoles. The Cos9 magnet assembly in its iron yoke is foreseen to end in autumn 2020 and a stand-alone test is scheduled at the beginning of 2021.

Figure 19. Cross section showing the conductor layout (left) and the whole magnetic circuit of the CEA Cos9 dipole wound with Roebel cable for EuCARD2.

Figure 20. Top: practice winding of the CEA dipole with Roebel dummy cable (test with different end spacers concepts, the blue ones on the left have been finally chosen). Bottom left: detail of the cable turning during winding; center and right: pictures of the one of the three coils just after the impregnation.

5.3 The EuCARD-2 stacked tape dipole

Another design was explored, based on a cable composed of a simple stack of parallel tapes. This design has been pursued by INP Grenoble (Fr), based on a simple stack of 4 mm wide tapes, such as to have a square cross-section [77,78]. To avoid difference in length among tapes, and to provide a minimum of transposition, the cable is twisted along its longitudinal axis right before and after the turning of the coil ends, see Fig.21. A full design has not yet been completed. In particular a more detailed analysis would be needed for mechanics, to limit the shear stress in the cable at the points where it is twisted, and for field quality. However, its simple structure, the high compaction factor and the efficiency in use of HTS make this layout attractive as a possible alternative, once the basic issues are resolved.
5.4 Magnet technology

Three FeatherM0 coils were made from dummy cable composed of stainless steel tapes, punched like the HTS tape for Roebel, to learn and test various technologies. An attempt to use glass ball-charged resin for impregnation failed, since the braid acts as a filter for the micro-balls (nano-size balls would be needed, which is quite complicated and expensive).

The first coil wound with superconducting cable, FeatherM0.4 [79], was used for benchmarking various diagnostic systems, especially for quench detection, which with current as high as 10 kA must happen in a few tens of ms. We developed a classical voltage comparison method with good sensitivity. Sensitivity was extremely high with the first dipole, FeatherM2.1-2, that had an insulated metal wire inside the cable, used as a reference loop perfectly coupled with the coils to cancel inductive signals, and good enough for the second dipole FeatherM2.3-4, that did not have such a coupled wire. Three additional types of sensors were placed in FeatherM0.4 to detect the quench, a) a temperature array placed as near as possible to the coils; b) an array of small flat pick-up coils capable of detecting the flux variation associated with current redistribution in tapes and cable; c) a couple of optical glass fibers in Bragg grating configuration to detect temperature and/or strain status variation [80]. Optical fibers showed good potential to detect temperature rise well below the current sharing temperature. The other two systems were less successful, and better position was identified for pick-up coils in the FeatherM2 magnet. In any case, classical voltage detection, with refined data acquisition, showed sufficiently high sensitivity to detect the quench with sufficient anticipation to safely shut down the magnet even with large current (>10 kA).

An interesting feature of the FeatherM2 dipole design is the use of a copper ring around the coil as a structural element that acts as an inductive coupler capable of extracting a part of the energy in a very fast and safe mode from the coil. More details can be found in [81].

The first FeatherM0.4 test in 2016 [79] was extremely useful in improving connections with large transport current (10-15 kA) in conductors dominated by stainless steel, with little copper. Following the test on FeatherM0.4 a new type of connection was developed, call a “Fin-Joint”, to capture each single tape with an appropriate transfer length, as discussed in [82]. However, the test was successfully concluded by injecting 13 kA in FeatherM0.4 at 12 K, after having first explored the critical surface from 80 K down to low temperatures. The results were quite satisfactory, showing good current sharing among tapes.
6. Magnet test and results

All test so far have been carried out at CERN, where an adaptation of an existing set up allows testing from 80 K down to 5 K and then in liquid at 4.2 K, with current up to 20 kA. A standard and fast quench detection system (QDS) is available for these tests as well as a fast switch based on IGBT technology to open the circuit in a few ms. From 2021 the INFN-LASA facility with similar characteristics to the one of CERN, will also be available.

Tests were carried out at CERN on the FeatherM0.4 coil at variable temperature that was then subsequently tested, also at various temperatures, in the high field split coil Sultan facility in the frame of a collaboration between CERN and EPFL-CRPP. A further test coil FeatherM0.5 was tested directly in the Sultan facility. The two AB dipoles, FeatherM2.1-2 and FeatherM2.3-4 were tested in the CERN variable temperature facility.

6.1 Results on FeatherM0 coils

The results of the first test at CERN on FeatherM0 are reported in Fig. 22. The maximum current was limited by heating in the joints (from the coil terminals to current leads), rather than conductor in the coil. At each warm-up the configuration of the joints was modified and the cooling of the joints improved (by pouring LHe directly on the limiting joints). Finally, we could reduce the joint resistance down to the ten nΩ range and, at the fourth cool-down, we could see the coil quenching, see red solid square in Fig. 22. In this way we could validate the cable in terms of current transport since we exceeded the 10 kA limit. The cable itself was a very preliminary prototype out of the EuCARD-2 collaboration. The coil was limited to 12 kA at 23 K, so we can extrapolate some 15-20 kA to be the quench limit at 4.2 K (our system was not suitable to sustain such high current in a safe mode). We did not observe degradation in the four cooldown-warmup cycles.

![Figure 22. Results (maximum current vs operating temperature) of the power test of FeatherM0 at CERN.](image)

Then the coil was tested in the Sultan facility [83]. After cooldown the coil reached the same performance as at CERN, however with higher coil resistance (35 nΩ). Then it reached 10 kA in steady state under an 11 T background field. Given the modest self-field of the coil (about 0.2T/kA) the conductor was exposed to a field of 13 T total. The internal resistance of the coils generated substantial
heating. Then, in the subsequent warmup-cooldown cycle the resistance increased dramatically in a way that it was no longer possible to energize the coil. The same behavior, extremely high resistance, happened straight away at the first cooldown of the FeatherM0.5 coil, manufactured with tape not from EuCARD-2. Therefore, this test coil never saw important transport current.

The reason for this behavior is not clear: the most likely hypothesis is that it may be a sign of degradation under field with delamination between the REBCO layer and substrate. In case of FeatherM0.5 it seems that delamination happened after cooldown. However, investigation is still under way at CERN to obtain a definitive answer.

6.2 Results of AB dipoles: first two FeatherM2 magnets

Because of moving the production plant from one site to a new one, Bruker could not deliver the EuCARD-2 tape on time for the first FeatherM2 dipole. Therefore, it was decided to wind the first EuCARD-2 dipole with a cable quickly procured by CERN from SuperOx. The tape was manufactured by SUNAM and the Roebe cabling was taken care by SuperOx as final responsible of the product [84].

The FeatherM0.4 test was important also to debug the test facility and, as mentioned above, to improve some critical issues, like the HTS joints. This allowed a fast test of FeatherM2.1-2. The first test happened in April 2017, see Fig. 23, just in time to meet the deadline of the EuCARD-2, reaching a field of 3.35 T when powered by 6.5 kA at 5 K, above expectations based on Ic evaluation with scaling law (we do not have a direct Ic cable measurement). Actually, the maximum steady state field was 3.1 T. The extra 0.25 T are gained by allowing the coil to go up smoothly in temperature, up to a certain point when the total voltage increase starts to accelerate and the power supply needs to shut-off. This behavior points to a very soft transition, smearing the voltage increase, dominated by current sharing mode. This is extremely useful because the quench starts smoothly and it is easy to detect. Actually, since the heating and voltage increase are very reproducible, one may think to run the magnet in a voltage mode rather than in current mode (to reach maximum current).

![Figure 23. First test of AB dipole FeatherM2.1-2 (quench current vs operating temperature).](image)

A full analysis of the test results is reported in [85, 86]. The main points are:

- There is evidence of current sharing among tapes, as anticipated by the measurement of losses on cables. However, it seems that current sharing is larger than expected, with
contact resistance in the range of 10 \( \mu \Omega \). This helps short-circuiting local defects, but also makes the quench onset easy to detect.

- The transition is so smooth that we could increase performance by 10%, and reach 3.35 T by increasing the ramp rate.
- The effect of energy dissipation in the copper ring, a feature that may be useful for future high field magnets [81] is visible in accelerating the current decay.
- The absence of training is remarkable: even considering that it is a low field and small energy magnet, this confirms the benefit of a large temperature margin.
- First cool down did not degrade the quench performance. However a subsequent thermal cycle saw a degradation of about 10% but the cool down was not properly controlled. More verifications are under way.
- The first campaign of magnetic measurements showed a considerable amount of coupling current decaying on an 80 s time constant and a relatively small contribution of persistent current [87], of about 15-20 units. The current decay might be triggered by a - low but not negligible - internal resistance. While favorable from a field quality point of view, this last effect needs to be understood through more investigation.

Finally, the second EuCARD-2 AB dipole, with the full performance tape of Bruker and the Roebel cabling by KIT, as from of EuCARD-2 collaboration, was manufactured in 2018 and tested in 2019 and 2020. We performed three cooldown-warmup cycles. In the first two we were limited by bad connections and probably by a damaged conductor in the leads just at the exit of the coils (probably due to bad manipulation). After some fix was put in place, we were able, at the third test cycle, to reach 9 kA, i.e 4.5 T in steady state, with a quench starting in the coil, see Fig.24. While very near to design (5 T) this is visibly below the 6-7 T we could expect from the extrapolated short sample value on tape. We do not have, like for FeatherM2.1-2 any direct measurement of \( I_c \) on the cable. One good point is the absence of noticeable degradation after three cooldowns. The 9 kA limitation may actually not be due to the conductor in the coil: we suspect that the outer lead is heating (in the damaged point near the exit of the coil) and the coil is driven normal by thermal conduction.

![Figure 24. First powering campaign of the second AB dipole, FeatherM2.3-4, always in terms of maximum current vs operating temperature. Unpublished data, courtesy of G. de Rijk, G. Kirby, J. van Nugteren and F. Mangiarotti, CERN.](image)

Cold tests, also to assess the field quality, will be resumed as soon as the Covid-19 emergency is over. This HTS program is important for the long term future of CERN, however it has at present very low priority with respect to LHC operation and High Luminosity LHC construction. Therefore, the
reduced activity due to restarting after the Covid-19 lock-down, makes this complementary program to be more severely delayed than other programs.

7. Conclusions

The European program EuCARD-2, and to some extent also EuCARD and ARIES, have constituted a small, however continuous, program of HTS accelerator magnets and conductor development. It has been instrumental in bringing together various laboratories and companies.

So far we have demonstrated capability of a high filling factor 10 kA class cable, Roebel type (up to a length of a few tens of meters) and we developed a new dipole layout, based on Aligned Block and flared ends, to allow an aperture of 40 mm in the coils (reduced to 35 mm to allow insertion in the Fresca2 high field facility without mechanical coupling). The highest field reached so far is 4.5 T in standalone with the FeatherM2.3-4. But we hope with some fix to increase beyond this level.

In any case we have two further magnets to test: a cos9-dipole by CEA (not meant for record field but to learn technology in a well-known configuration), and a third AB dipole: FeatherM2.5-6 by CERN, for which the cable has been procured according to the EuCARD-2 specification. It should be manufactured and tested by Spring 2021 and should go beyond the design goal of 5 T.

In general, we should say that HTS, despite the intrinsic difficulties for use in accelerator magnets, keep their main promise of training-free magnets and therefore offer solutions that should be exploited with novel designs and different approaches than LTS.

Our program now also includes the test of various inserts inside a high field facility. The first one, FeatherM2.1-2 inside Fresca2, may happen before the end of 2020. If it will reach a field above the 16 T limit of LTS for dipoles, this would constitute a tremendous boost for HTS technology for particle accelerators and for the design of even higher field dipoles [88].

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