

A new Lorentz violating model with particle's "maximum energy"

Jinwen Hu*, Huan Hu

School of physics and technology, Wuhan university, Wuhan 430060, China

Abstract

In Lorentz violating models, the rainbow model or theories of Quantum Gravity are usually discussed, and a common feature of these models is that they assume the particle's energy have a limited value rather than be infinite derived from the Lorentz model. The introduction of "maximum energy" is considered to be necessary in the combination of Quantum theory and Gravity. However, this paper shows that if we just follow the next three principles: (1)we can define the time in the whole space with a prescribed clock synchronization, (2)the time-space is uniform and the space is isotropic and (3)all the inertial systems are equivalent, then we can totally construct a general coordinate transformation to meet the symmetry of inertial systems, and with a special assumption on the speed of light, we can also construct a non-Lorentz transformation between inertial systems to make the particle's energy have a limited value.

Similar to the usual Lorentz violating models, the non-Lorentz transformation in this paper lead to a new modified disperse relation. We applied the obtained disperse relation to analyze the photon's arrival time lag effect and found that the "maximum energy" derived in our model is somewhat related to the "maximum energy" assumed in the rainbow model.

Keywords

Lorentz model; ultrahigh energy; Gamma ray burst; rainbow model; variable speed of light; time lag

1. Introduction

It is well known that the Special Relativity and General Relativity, which are proposed based on the Lorentz invariance, have already made great achievements, but at the same time the Lorentz violating models are also of some astrophysical interest. In the past few decades, the Lorentz violating models have been proposed in various frameworks and theories of Quantum Gravity, such as the string theory [1-3], the loop quantum theory [4,5], and the non-commutative time-space theory [6], which aim to combine the Quantum theory and the gravity theory. In many models called "Quantum Gravity", the "minimum length" or "maximum energy" independent of the

* Corresponding author at: School of physics and technology, Wuhan University, Wuhan 430060, China.
Email address: 2007.hujinwen@163.com (Jinwen Hu).

inertial systems, is assumed, which imply the Lorentz invariance is violated in the ultrahigh energy scale. For example, the rainbow model, which is a typical and often discussed Lorentz violating model and proposed in the doubly special relativity (DSR) theory, introduces the "minimum length", i.e., the Planck length, as another constant between inertial systems. In the framework of rainbow model, the particle's energy is finite rather than be infinite, and correspondingly, the particle's energy-momentum dispersion relation is modified as the following form (set $c=1$) [7,8]

$$\left[1 + \chi_1 \left(\frac{E}{E_{LV}}\right)^1 + \chi_2 \left(\frac{E}{E_{LV}}\right)^2 + O\left(\frac{E}{E_{LV}}\right)^3\right] E^2 - \mathbf{p}^2 = m_0^2 \quad (1)$$

where E denotes the total energy of the particle, m_0 denotes the rest mass, \mathbf{p} is the momentum, and E_{LV} denotes the energy scale at which Lorentz violating effects become strong, the couplings $\chi_{s=\pm 1, 2}$ are determined by the dynamical framework being studied.

The doubly special relativity (DSR) model provides a new perspective for the development of Quantum Gravity. For example, the application of Eq. (1) can make the black hole's temperature to be convergence [9,10] and alleviate the singularity problem arose in the Big Bang model [11,12], etc. The rainbow model have made progress in solving some problems with extreme conditions, but the specific form of rainbow model is still unknown. Because of this, a lot of work has been done to determine the parameters involved in the rainbow model. In the previous work, for example, Bolmont [13] applied the HETE-2 gamma ray burst data to constrain the energy scale E_{LV} to be at least 2×10^{15} GeV at the 95% Confidence Level, Nilsson [14] applied the Hubble+SN Ia+BAO (BOSS+Lyman α)+CMB data to constrain E_{LV} to be at least the order of 10^{16} GeV at 1σ , or even 10^{17} GeV at 3σ . And based on the first-order approximation of Eq. (1), Xu [15,16] analyzed the GRB 160509A event and claimed that there exists a linear relation between the variable speed of light and the photon's energy, and $E_{LV}=3.6 \times 10^{17}$ GeV was obtained.

However, the rainbow model is somewhat complicated, especially the correction terms of Eq. (1) are still unclear in physics. But the rainbow model is a special and typical Lorentz violating model, it inspired us that there may be other possible relationship between the speed of light and the symmetry of inertial systems. So this paper re-investigated the relationship between the symmetry of inertial systems and the Lorentz model. We found that by imposing some new special rules on the speed of light, we can construct a general function on the speed of light to make the inertial systems to be equivalent, which is shown in the content of Sect. 2. In Sect. 3, by taking a specific assumption on the speed of light, we obtained a non-Lorentz transformation between inertial systems to make the particle's energy have a limited value, which is similar to the rainbow model. In Sect. 4, we applied the modified energy-momentum dispersion relation derived from the non-Lorentz transformation to analyze the photon's arrival time lag effect, and found that the particle's "maximum energy" derived in our paper is somewhat related to the "maximum energy" assumed in the rainbow model. In Sect. 5 we summarized the paper.

2. Variable speed of light

It is well known that the rainbow model presents that the speed of light maybe associated with the photon's energy [15,16], then here we proposed a general hypothesis that: For a light source in vacuum, when it moves at a velocity \mathbf{v} relative to

an observer in vacuum, then the observed (by the observer) speed of light emitted by the light source is nc , where n is a dimensionless quantity, c is the speed of light in vacuum. Obviously, in order not to violate some fundamental principles and experiments, we should impose some rules on the parameter n as follows

1. Firstly, as stated in the special relativity, we should define the (proper) time (note that the only time that makes sense in special relativity is the proper time measured by the clocks carried by physical observers) in the whole space with a prescribed clock synchronization, that is, for a specific inertial system, based on the speed of light emitted by a light source that fixed in the specific inertial system, we can calibrate the clock fixed in the inertial system to synchronize. So it requires that

$$n(v = 0, c) = 1 \quad (2)$$

2. Secondly, according to the general concept of time and space that the time-space is uniform and the space is isotropic, it requires that $n(v, c)$ is independent of the direction of vector v and c , namely,

$$n(v, c) = n(-v, c) = n(v, -c) = n(-v, -c) \quad (3)$$

3. In addition, we also follow the principle that all the inertial systems are equivalent.

Based on the above assumptions on the speed of light, next we will derive the coordinate transformation between the two inertial systems $S(x, y, z, t)$ and $S'(x', y', z', t')$. And here we assume that S' is moving at a velocity v relative to S .

Firstly, for simplicity, we assume that the three spatial coordinates of the two coordinate systems are parallel to each other, and the direction of v is along the x -axis or x' -axis, then we obtain $y=y', z=z'$.

Secondly, since the time-space is uniform, the coordinate transformation between S and S' should be in a linear form (note that in the rainbow model, the coordinate transformation between S and S' is assumed to be in a nonlinear form, here we follow that the time-space is uniform and thus the coordinate transformation between S and S' should be in a linear form), and we assume that

$$x = \gamma(x' + vt') \quad (4)$$

Where $\gamma=\gamma(v, c)$ is a proportionality factor.

The inverse transformation of Eq. (4) is

$$x' = \gamma'[x + (-v)t] \quad (5)$$

Where γ' is a proportionality factor.

Because of the symmetry of space, we will obtain

$$\gamma(v, c) = \gamma(-v, -c) \quad (6)$$

Note that the above we didn't distinguish the direction of vector v and c , i.e., the direction of vector v and c maybe along the positive x -axis (or x' -axis) or along the negative x -axis (or x' -axis). And if we distinguish the direction of vector v and c by the positive and negative signs, we will obtain four different combinations, i.e., (v, c) , $(v, -c)$, $(-v, c)$, $(-v, -c)$. Based on Eq. (6) we can obtain $\gamma(v, c) = \gamma(-v, -c)$, $\gamma(-v, c) = \gamma(v, -c)$.

-c).

Now we will solve the expression for γ . If the light signal is emitted by the light source at the moment that the origin of S and S' are coincides, then based on the above assumption on the speed of light, we will obtain

$$\begin{cases} x^2 + y^2 + z^2 = (ct)^2 \\ x'^2 + y'^2 + z'^2 = (nct')^2 \\ y = y' = 0 \\ z = z' = 0 \end{cases} \quad \text{and} \quad \begin{cases} x'^2 + y'^2 + z'^2 = (ct')^2 \\ x^2 + y^2 + z^2 = (nct)^2 \\ y = y' = 0 \\ z = z' = 0 \end{cases} \quad (7)$$

The first formula of Eq. (7) represents that when the light source is fixed in S , then for the observer in S , the observed speed of light is c (based on Eq. (2)), while for another observer in S' , the observed speed of light is nc .

Similarly, since S and S' are symmetric, when the light source is fixed in S' , then for the observer in S' , the observed speed of light is c , while for another observer in S , the observed speed of light is nc , which case corresponds to the second formula of Eq. (7).

Based on Eqs. (4)~(7), and similar to the derivation method shown in the special relativity theory, we can solve the expression for γ and obtain the coordinate transformation between S and S'

$$\begin{cases} x' = \gamma(x - vt) \\ t' = \gamma(t - \frac{v}{k^2(v, c)}x) \end{cases} \quad (8)$$

where $\gamma(v, c) = 1 / \sqrt{1 - v^2 / k^2}$, $k(v, c) = \sqrt{nv c^2 / (nc - c + v)}$.

From Eq. (8) it can be seen that $k(v, c) = k(-v, -c)$, $k(-v, c) = k(v, -c)$, and $\gamma(v, c) = \gamma(-v, -c)$, $\gamma(-v, c) = \gamma(v, -c)$, which is agree with Eq. (6).

Here it should be noted that, from Eqs. (4)~(7), it doesn't require that $k(v, c)$ is equal to $k(-v, c)$. In fact $k(v, c)$ and $k(-v, c)$ represent the two different cases or two independent events, they cannot be linked together by the symmetry in one event.

Based on Eq. (8) we can obtain

$$\frac{dx'}{dt'} = \frac{dx - v dt}{dt - \frac{v}{k^2} dx} = \frac{dx / dt - v}{1 - \frac{v}{k^2(v, c)} \frac{dx}{dt}} = f(v, c) \quad (9)$$

As stated above that the direction of vector v and c maybe along the positive x -axis (or x' -axis) or along the negative x -axis (or x' -axis), and if we distinguish the direction of vector v and c by the positive and negative signs, we will obtain four different combinations, i.e., (v, c) , $(v, -c)$, $(-v, c)$, $(-v, -c)$. Then based on Eq. (9), we can obtain that

Case 1: Note that the above assumed that S' is moving at a velocity v relative to S . When the light source is fixed in S , then for the observer in S , the observed speed of light is c , while for the observer in S' , it has

$$\left\{ \begin{array}{l} \text{when } \frac{dx}{dt} = c, \frac{dx'}{dt'} = f(v, c) = \frac{c - v}{1 - \frac{v}{k^2(v, c)}c} = f(-v, c) = \frac{c + v}{1 - \frac{-v}{k^2(-v, c)}c} = nc \\ \text{when } \frac{dx}{dt} = -c, \frac{dx'}{dt'} = f(v, -c) = \frac{-c - v}{1 - \frac{v}{k^2(v, -c)}(-c)} = f(-v, -c) = \frac{-c + v}{1 - \frac{-v}{k^2(-v, -c)}(-c)} = -nc \end{array} \right. \quad (10)$$

Case 2: Similarly, since S and S' are equivalent, when the light source is fixed in S' , then for the observer in S' , the observed speed of light is c , while for the observer in S , it has (note that in this case the velocity of S relative to the light source is $-v$)

$$\left\{ \begin{array}{l} \frac{dx}{dt} = \frac{dx' / dt' - (-v)}{1 - \frac{-v}{k^2(-v, c)} \frac{dx'}{dt'}} = f'(v, c) \\ \text{when } \frac{dx'}{dt'} = c, \frac{dx}{dt} = f'(v, c) = \frac{c + v}{1 + \frac{v}{k^2(-v, c)}c} = f'(-v, c) = \frac{c - v}{1 + \frac{-v}{k^2(v, c)}c} = nc \\ \text{when } \frac{dx'}{dt'} = -c, \frac{dx}{dt} = f'(v, -c) = \frac{-c + v}{1 + \frac{v}{k^2(-v, -c)}(-c)} = f'(-v, -c) = \frac{-c - v}{1 + \frac{-v}{k^2(v, -c)}(-c)} = -nc \end{array} \right. \quad (11)$$

Eq. (10) and Eq. (11) can be also expressed in a vector form

$$\left\{ \begin{array}{l} \text{when } \frac{dx}{dt} = c, \frac{dx'}{dt'} = nc \\ \text{when } \frac{dx'}{dt'} = c, \frac{dx}{dt} = nc \end{array} \right. \quad (12)$$

Eq. (12) also shows that Eq. (8) is just the solution of Eq. (7).

Importantly, it can be seen that the forms of Eq. (8) are similar to the Lorentz transformation, i.e., replacing c in the Lorentz transformation with k we can obtain Eq. (8), and similar to the proof method shown in the special relativity, it is easy to prove that the Maxwell's Equations are also covariant based on Eq. (8).

Based on Eq. (8) we can obtain the time-space metric

$$ds^2 = -k^2 dt^2 + dx^2 \quad (13)$$

Correspondingly, the particle's energy-momentum dispersion relation is

$$E^2 = p^2 k^2 + E_0^2 \quad (14)$$

Where $E_0 = m_0 k^2$ denotes the particle's energy, $E = \gamma m_0 k^2$ is the total energy of the particle, $p = \gamma m_0 v$ denotes the particle's momentum.

Here it should be noted that, similar to the other Lorentz violating models, at above we just try to allow the speed of light to be possible to vary between inertial systems, and subsequently we try to impose some constraints on the parameter n to

make the time-space meet the above three principles, thus we obtain a self-consistent coordinate transformation between inertial systems. The idea is different from the DSR model, which introduce a second constant besides the speed of light, and also different from the Standard Model Extension (SME) model, which assumed there exists extra fields resulting in the observer Lorentz transformation is reserved while the particle Lorentz transformation is violated. In this paper we didn't introduce any constants or fields, and the speed of light observed by the different observers is not constant, but we obtained a special non-Lorentz model. And the forms of the coordinates transformation, the particle's energy-momentum dispersion relation and the time-space metric obtained in this paper are similar to which obtained from the Lorentz model.

3. Particle's "maximum energy"

As we know, in Lorentz model, the particle's energy tends to be infinite when the particle's velocity is close to the speed of light, however, the DSR model introduces a new constant as the energy limit of the particles, which is considered to be necessary in the theory of Quantum Gravity. In this paper, as Eq. (8) shown that if $n \equiv 1$ then Eq. (8) returns to the Lorentz model. But here we would like to discuss another interesting case where n is not always equal to 1, as the general Lorentz violating models suggested [1-6].

Inspired by the idea of rainbow model, we found that Eq. (8) imply that it is possible the particle's energy have a limit.

Based on Eq. (8), the time-space scaling factor is

$$\gamma = \frac{1}{\sqrt{1 - v^2 / k^2}} = \frac{1}{\sqrt{\frac{1 - v/c}{n} (n + \frac{v}{c})}} \quad (15)$$

Eq. (15) inspires us that when $v=c$, it is possible that γ does not tend to be infinite if we assume $n=0$ when $v=c$. So here we can construct an expression for n . Since n has been constrained in Eq. (2) and Eq. (3), that is

$$\begin{cases} n(v=0, c) = 1 \\ n(v, c) = n(v, -c) = n(-v, c) = n(-v, -c) \\ \lim_{v \rightarrow c} \frac{1 - v/c}{n} = \text{const} \end{cases} \quad (16)$$

Based on Eq. (16), we can only take a few but finite kinds of expressions for n , and the following expression for n is one of them

$$n = \frac{1}{1 - Q} (1 - Q^{1 - v^2/c^2}) \quad (17)$$

where Q is a constant.

Besides Eq. (17) one may consider another expression for n , but as we know there are many experiments restricting the violation of Lorentz model, so we should

choose an expression for n to not violate the previous experiments' results. Figure 1 shows the $n \sim v$ curve when taking $Q = (1/e)^{10^6}$ randomly as an example.

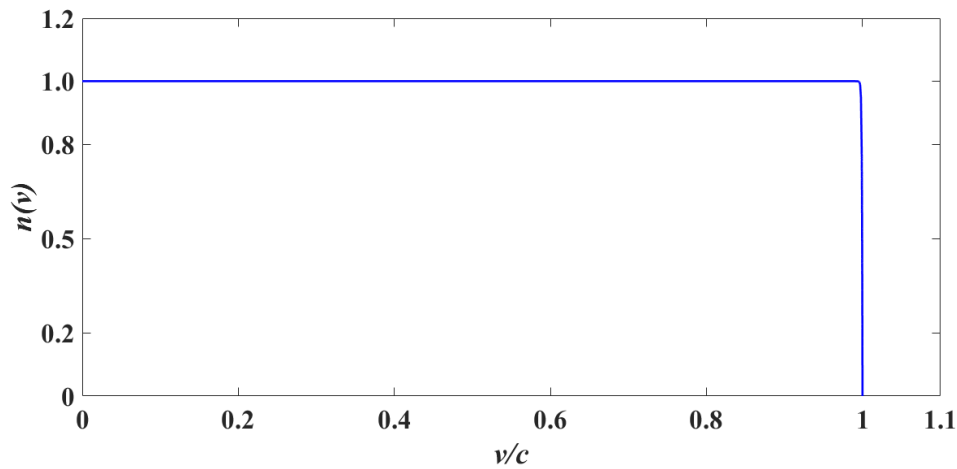


Fig. 1 $n(v) \sim v$ curve when taking $Q = (1/e)^{10^6}$

Thus, the time-space scaling factor limit and the particle's total energy limit are respectively

$$\lim_{v \rightarrow c} \gamma = \lim_{v \rightarrow c} \frac{1}{\sqrt{1 - v^2/c^2}} = \lim_{v \rightarrow c} \frac{1}{\sqrt{\left(\frac{1 - v/c}{n}\right)\left(n + \frac{v}{c}\right)}} = \sqrt{\frac{2 \ln Q}{Q - 1}}$$

$$\lim_{v \rightarrow c} E = \lim_{v \rightarrow c} \gamma m_0 c^2 = E_{QG} = \frac{m_0 c^2}{[1 - 0.5(Q - 1) / \ln Q]} \sqrt{\frac{2 \ln Q}{Q - 1}} \quad (18)$$

4. Modified arrival time lag effect

It can be seen from Fig.1 that the modified particle's energy-momentum dispersion relation will return to the Lorentz case at low or medium energy. So next we will discuss the behavior of particles with ultrahigh energy.

When $v \sim c$ for an ultra-relativistic particle, it can be obtained from Eq. (17) that (set $c=1$)

$$n = \frac{1}{1 - Q} (1 - Q^{1-v^2}) = \frac{1}{1 - Q} [1 - Q^{(1+v)(1-v)}] \approx \frac{1}{1 - Q} [1 - Q^{2(1-v)}] \approx \frac{2 \ln Q}{Q - 1} (1 - v) \quad (19)$$

Then

$$\begin{aligned}
\frac{E}{E_{QG}} &= \frac{m_0 k^2}{\sqrt{1 - v^2/k^2}} / \left[\frac{m_0 c^2}{[1 - 0.5(Q-1)/\ln Q]} \sqrt{\frac{2 \ln Q}{Q-1}} \right] \\
&= \frac{[1 - 0.5(Q-1)/\ln Q]}{\sqrt{2 \ln Q/(Q-1)}} \frac{nv/(n-1+v)}{\sqrt{(1-v)(1+v/n)}} \\
&\approx \frac{[1 - 0.5(Q-1)/\ln Q]}{\sqrt{2 \ln Q/(Q-1)}} \frac{2 \ln Q/(Q-1)v}{2 \ln Q/(Q-1) - 1} \frac{1}{\sqrt{1-v+v(Q-1)/(2 \ln Q)}} \\
&\approx \frac{[1 - 0.5(Q-1)/\ln Q]}{\sqrt{2 \ln Q/(Q-1)}} \frac{2 \ln Q/(Q-1)}{2 \ln Q/(Q-1) - 1} \frac{1}{\sqrt{(Q-1)/(2 \ln Q) + (1-v)}} \\
&\approx \frac{1}{\sqrt{2 \ln Q/(Q-1)}} \left[\sqrt{\frac{2 \ln Q}{Q-1}} - \frac{1}{2} \left(\frac{2 \ln Q}{Q-1} \right)^{3/2} (1-v) \right] \\
&= 1 - \frac{\ln Q}{Q-1} (1-v)
\end{aligned} \tag{20}$$

Based on Eq. (20) we can obtain

$$\frac{v}{c} = 1 - \frac{Q-1}{\ln Q} + \frac{Q-1}{\ln Q} \frac{E}{E_{QG}} \tag{21}$$

Eq. (21) shows that for an ultra-relativistic particle governed by Eqs. (14) and (17), its velocity is proportional to its energy.

Multiplying mc^2 on both sides of Eq. (21), we can obtain the equation for the photons

$$pc = E \left(1 - \frac{Q-1}{\ln Q} + \frac{Q-1}{\ln Q} \frac{E}{E_{QG}} \right) \tag{22}$$

where p is the photon's momentum, and E is the photon's energy.

If $E/E_{QG} \sim 0$, then based on Eq. (22), we can obtain that

$$\frac{\partial E}{\partial p} = \frac{c}{1 - \frac{Q-1}{\ln Q} + \frac{Q-1}{\ln Q} \frac{E}{E_{QG}}} \approx \left(1 - \frac{Q-1}{\ln Q} \frac{E}{E_{QG}} \right) c \tag{23}$$

Note that $Q \sim 0$ is used in Eq. (23). Now we apply Eq. (23) to study the photon's arrival time lag effect in astronomy. Based on Ref. [17], we can obtain the co-moving distance traversed by a photon, emitted at red-shift z and traveling up to red shift 0

$$x(z, E) = \frac{c}{H_0} \int_0^z \left[1 - \frac{Q-1}{\ln Q} \left(\frac{E^0}{E_{QG}} \right) (1+z') \right] \frac{dz'}{\sqrt{\Omega_m (1+z')^3 + \Omega_\Lambda}} \tag{24}$$

where E^0 is the photon's energy measured at present, Ω_m , Ω_Λ and H_0 are the cosmological parameters evaluated today.

Then based on Eq. (24) we can obtain the photon's arrival time lag due to the modified energy-momentum dispersion relation in Eq. (14) as

$$\Delta t = \frac{Q-1}{\ln Q} \frac{1}{H_0} \frac{E^0}{E_{QG}} \int_0^z \frac{(1+z') dz'}{\sqrt{\Omega_m (1+z')^3 + \Omega_\Lambda}} \propto E^0 \tag{25}$$

In Ref. [16], the first-order approximation of Eq. (1) (that is, $\chi_1=1$, $\chi_2=0$) is taken

to analyze the photon, and the corresponding arrival time lag was obtained as

$$\Delta t = \frac{1}{H_0} \frac{E^0}{E_{LV}} \int_0^z \frac{(1+z') dz'}{\sqrt{\Omega_m (1+z')^3 + \Omega_\Lambda}} \propto E^0 \quad (26)$$

Comparing Eq. (25) with Eq. (26), we can obtain that

$$E_{QG} = \frac{Q-1}{\ln Q} E_{LV} \quad (27)$$

Eq. (27) imply that the modified energy-momentum dispersion relation (i.e., Eq. (14)) is somewhat related to the rainbow model.

Importantly, it is well known that there are many experiments restricting the violation of Lorentz model, such as the experiments in Ref. [18-26], and these experiments' result imply the value of Q is close to 0 (when $Q=0$ then $n=1$). In our previous work [27], we have restricted that $Q < (1/e)^{10^6}$, and the corresponding time-space scaling factor γ limit is bigger than 1414.2, while the above conclusion, i.e., for photons, it has $E_{QG}=(Q-1)/(\ln Q)E_{LV}$, indicates that the previous work on the restriction of E_{LV} can equally apply to constrain E_{QG} once we obtain the value of Q . But unfortunately, so far we still can't obtain the specific value of Q , for that the value of Q can only be obtained in the experiment of particle with rest mass and ultrahigh energy. And we expect the specific value of Q can be obtained in the future's neutrino experiment.

5. Conclusion

In this paper, we re-investigated the relationship between the speed of light and the symmetry of inertial systems based on the following three principles: (1)we can define the (proper) time in the whole space with a prescribed clock synchronization, (2)the time-space is uniform and the space is isotropic and (3)all the inertial systems are equivalent. We found that it was not necessary to restrict the speed of light to be a constant to meet the symmetry of inertial systems. And by imposing some special rules on the speed of light, we can construct a general coordinate transformation between inertial systems to meet the symmetry of inertial systems. Similar to the rainbow model, with a special assumption on the speed of light we also construct a non-Lorentz coordinate transformation between inertial systems to make the particle's energy have a limited value.

By analyzing the photon's arrival time lag effect, we found that the particle's energy limit derived in our model is somewhat related to the particle's energy limit assumed in the rainbow model. That is, the first-order approximation of the rainbow model is somewhat related to the dispersion relations derived in our paper.

Compared with the rainbow model, Eq. (14) providing a new specific form for the generic modified dispersion relations (MDP) [28] has clear correction terms in physics, and the equation has just one undetermined parameter (the value of Q). Moreover, as Eq. (18) shown, the particle's total energy limit depends on the particle's rest mass, which is also different from the rainbow model or the usual theories of Quantum Gravity that the assumed or introduced "maximum energy or Planck energy" is independent of the particle's mass.

It should be noted that, although we construct a non-Lorentz transformation in above to make the particle's energy have a limited value, particles still cannot travel faster than the speed of light emitted by a stationary light source. And here we didn't

claim certainly that n deviates from 1 at the ultrahigh energy scale, but as the above shown, in theory and in no violation of the present experiments, it is possible that n is not always equal to 1. And similar to the rainbow model, if the value of Q is not equal to 0 (if $Q=0$ then $n=1$ and Eq. (8) returns to the Lorentz case), then we will have to modify the current black hole model obtained from the general relativity theory, and therefore it may hopefully solve some problems faced by the current black hole model, which we will discuss in the next paper.

References

- [1]G. Amelino-Camelia, J.R. Ellis, N.E. Mavromatos, D.V. Nanopoulos, Distance measurement and wave dispersion in a Liouville string approach to quantum gravity, *Internat. J. Modern Phys. A* 12, 607-624 (1997).
- [2]G. Amelino-Camelia, Testable scenario for relativity with minimum length, *Phys. Lett. B* 510, 255-263 (2001) .
- [3]G.P. de Brito, P.I.C. Caneda, Y.M.P. Gomes, J.T.G. Junior, V. Nikoofard, Effective models of quantum gravity induced by Planck scale modifications in the covariant quantum algebra, 2016. arXiv:1610.01480.
- [4]N. Abbasvandi, M.J. Soleimani, S. Radiman, W.A.T. Wan Abdullah, Quantum gravity effects on charged microblack holes thermodynamics, *Internat. J. Modern Phys. A* 31 (23), 1650129 (2016) .
- [5]A. Tawfik, H. Magdy, A.F. Ali, Lorentz invariance violation and generalized uncertainty principle, *Phys. Part. Nucl. Lett.* 13, 59-68 (2016) .
- [6]S. Doplicher, K. Fredenhagen, J. E. Roberts, Spacetime quantization induced by classical gravity, *Phys. Lett. B* 331, 39 (1994).
- [7]G. Amelino-Camelia, J.R. Ellis, N.E. Mavromatos, D.V. Nanopoulos, S. Sarkar, Tests of quantum gravity from observations of gamma-ray bursts, *Nature* 393, 763-765 (1998).
- [8]J. Ellis, K. Farakos, N.E. Mavromatos, V.A. Mitsou, D.V. Nanopoulos, A search in gamma-ray burst data for nonconstancy of the velocity of light, *Astrophys. J.* 535 (1) , 139 (2000) .
- [9]C. Liu, H. Liu, Spectroscopy of a charged black hole in gravities rainbow via an action invariance, *Astrophys. Space Sci.* 357 (2), 114 (2015) .
- [10]S. Gangopadhyay, A. Dutta, Constraints on rainbow gravity functions from black hole thermodynamics, *Europhys. Lett.* 115 (5), 50005 (2016) .
- [11]Yi Ling, Rainbow universe, *JCAP* 0708, 17 (2007) .
- [12]A. Awad, A.F. Ali, B. Majumder, Nonsingular rainbow universes, *J. Cosmol. Astropart. Phys.* 1310, 052 (2013).
- [13]J. Bolmont, A. Jacholkowska, J.-L. Atteia, F. Piron, G. Pizzichini, Study of time lags in HETE-2 Gamma-Ray Bursts with redshift: search for astrophysical effects and Quantum Gravity signature, 2007. arXiv:astro-ph/0606725v6.
- [14]Nils A. Nilsson, Mariusz P. Dabrowski, Energy scale of Lorentz violation in Rainbow Gravity, *Physics of the Dark Universe* 18, 115 (2017) .
- [15]H. Xu, B.-Q. Ma, Light speed variation from gamma-ray bursts, *Astropart. Phys.* 82, 72 (2016).
- [16]H. Xu, B.-Q. Ma, Light speed variation from gamma ray burst GRB 160509A, *Phys. Lett. B* 760, 602 (2016) .
- [17]U. Jacob, T. Piran, Lorentz-violation-induced arrival delays of cosmological particles, *J. Cosmol. Astropart. Phys.* 0801, 031 (2008) .
- [18]M. E. Tobar, P. Wolf, S. Bize, G. Santarelli, and V. Flambaum, Testing local Lorentz and position invariance and variation of fundamental constants by searching the derivative of the comparison frequency between a cryogenic sapphire oscillator and hydrogen maser, *Phys. Rev. D* 81, 022003 (2010).

- [19]M. Nagel, S. Parker, E. Kovalchuk, P. Stanwix, J. Hartnett, E. Ivanov, A. Peters, and M. Tobar, Direct terrestrial test of Lorentz symmetry in electrodynamics to 10^{-18} , Nat. Comm. 6, 8174 (2015).
- [20]H. P. Robertson, Postulate versus Observation in the Special Theory of Relativity, Rev. Mod. Phys. 21, 378 (1949).
- [21]R. Mansouri and R. U. Sexl, A Test Theory of Special Relativity: I. Simultaneity and Clock Synchronization, Gen. Rel. Grav. 8, 497 (1977).
- [22]R. Mansouri and R. U. Sexl, A Test Theory of Special Relativity: II. First Order Tests, Gen. Rel. Grav. 8, 515 (1977) .
- [23]R. Mansouri and R. U. Sexl, A Test Theory of Special Relativity: III. Second-Order Tests, Gen. Rel. Grav. 8, 809 (1977).
- [24]C. Lämmerzahl, C. Braxmaier, H. Dittus, H. Müller, A. Peters, and S. Schiller, Kinematical Test Theories for Special Relativity, Int. Jour. Mod. Phys. D 11, 1109 (2002).
- [25]S. Liberati, Tests of Lorentz invariance: a 2013 update, Class. Quant. Grav. 30, (2013) .
- [26]F. W. Stecker, S. T. Scully, S. Liberati, D. Mattingly, Searching for traces of planck-scale physics with high energy neutrinos, Phys. Rev. D91 (4), 045009 (2015).
- [27]Jinwen Hu, The non-Lorentz transformation corresponding to the symmetry of inertial systems and a possible way to the quantization of time-space, Physics Essays 30, 3 (2017).
- [28]G. Amelino-Camelia, J. R. Ellis, N. E. Mavromatos, D. V. Nanopoulos and S. Sarkar, Tests of quantum gravity from observations of γ -ray bursts, Nature 393, 763 (1998).