

2 **New Results in 5D Theory and Some Problems of** 3 **Astrophysics and Cosmology**

4 **Boris G. Aliyev** ¹

5 Architecture dep., V.I. Surickov Moscow State Art Institute
6 Tovarishcheskij per. 30, Moscow, 109004, Russia, bgaliyev@mail.ru

7 **Abstract:** It is shown that the 5D Ricci identities give us a way to create a new viewpoint on the
8 origin of the Maxwell equations, magnetic monopole problem, and also on some problems of the
9 Astrophysics and Cosmology. Specifically, the application of the identities together with the
10 monad and dyad methods makes it possible to introduce the new concept of the rest mass of the
11 elementary particles. The latter leads to the new connections between the General Relativity and
12 quantum field theories, as well as to a better understanding of the magnetic monopole problem and
13 the origins of the Maxwell equations. The obtained results also provide a new insight into the
14 mechanism of the accelerated expansion of the 4D Universe.

15 **Keywords:** monad and dyad methods; effective rest mass concept; scalar gravitational field; 5D
16 geodetic equation; cylindrical symmetry condition; 5D Ricci identities; Maxwell equations;
17 magnetic monopole; topological second-order transition in cosmology

18 **1. Introduction**

19 The article considers some old and new problems of modern physics, astrophysics, and cosmology
20 in the framework of the 5D theory, using the well-developed and well-known monad and dyad
21 methods [1] (pp. 184-207), [2,3] in General Relativity. Reviewing the papers of the different authors
22 who have worked in the framework of the 5D theories, one can see that almost all of them [4-7]
23 came to the 5D optics under the requirement that ratio e/m has to be constant. We demonstrate
24 that relaxing this requirement leads to new, fundamental, unexpected, and very surprising
25 consequences concerning the effective rest mass concept (ERM), the origins of the first and second
26 pairs of the Maxwell equations, and the role of the 5D Ricci identities in the understanding of the
27 some of these problems. Moreover, we hope that our findings can advance our current
28 understanding of the magnetic monopole problem, as well as the other important problems of
29 modern astrophysics and cosmology, including the ones of the dark matter and the dark energy.
30

31 **2. The basic elements of the monad and dyad methods in 5D theory**

32 Here I would like to represent briefly the basic information about the monad method in 5D theory
33 with the chronometrical gauging of the monad vector and also about the dyad method with the
34 mixed, chronometric and then, kinematic gauging of the dyad vectors. These methods allow us to
35 make the procedure of V_5 - reducing: $V_5 \mapsto V_4 \mapsto V_3$ more correctly or, in other words, the
36 orthogonal (4+1)-splitting and then (3+1+1)-splitting of the V_5 with the subsequent expressing of the
37 5D theory in terms of the 4D or 3D geometrical and physical values. In the framework of the monad
38 method with the chronometrical gauging of the monad vectors we can represent the 5D metric G_{AB}
39 as follows: $G_{AB} = g_{AB} - \lambda_A \cdot \lambda_B$ (here we have put the signature of the V_5 to be equal (+ - - -)).

¹ Communication address: Aliev, Boris; Altendorfer Str. 28, 09113, Chemnitz, Germany (Aliev, it is in a German manner and Aliyev, it is in an English manner).

40 The indexes $A, B, C, \dots = 0, 1, 2, 3, 5$. The 5D interval takes on the following form:
 41 $dI^2 = ds^2 - d\lambda^2$. Here $ds^2 = g_{AB} \cdot dx^A \cdot dx^B$ and $d\lambda = \lambda_A \cdot dx^A$. The space-like monad vectors
 42 can be written as follows:

$$43 \quad \lambda_A = \frac{G_{A5}}{\sqrt{-G_{55}}}, \lambda^A = \frac{G_5^A}{\sqrt{-G_{55}}}, \sqrt{-G_{55}} = \varphi.$$

44 The vector λ_A is tangential to the x^5 coordinate line and orthogonal to the space-time hypersurface
 45 V_4 . Also g_{AB} is the metric tensor of V_4 . Then, the orthonormal condition gives us: $\lambda^A \cdot \lambda_A = -1$;
 46 $\lambda^A \cdot g_{AB} = 0$; and, of course, we have the well-known relations: $G^{AB} \cdot G_{AB} = 5$; $G^{AB} \cdot G_{BC} = \delta_C^A$;
 47 $g^{AB} \cdot g_{AB} = 4$; $g^{AB} \cdot g_{BC} = \delta_C^A$. Now we construct the monad differentiation operators:
 48 $\lambda^A \cdot \partial_A = \partial_\Lambda^+$ and $g_A^B \cdot \nabla_B = \bar{\nabla}_A^+$. Then, we can build the x^5 and V_4 projected directional normal
 49 and covariant derivatives. In the case of the dyad method the 5D metric and 5D interval take on the
 50 following form: $G_{AB} = \tau_A \cdot \tau_B - \lambda_A \cdot \lambda_B - h_{AB}$; $dI^2 = d\tau^2 - d\lambda^2 - dl^2$; $d\tau = \tau_A \cdot dx^A$;
 51 $dl^2 = h_{AB} \cdot dx^A \cdot dx^B$. Now we can also write the orthonormal time-like (the case of the space-like
 52 vectors λ_A was commented above) vectors of the dyad in kinematic gauging as follows:

$$53 \quad \tau^A = \frac{G^{0A}}{\sqrt{G^{00}}}, \tau_A = \frac{G_A^0}{\sqrt{G^{00}}}, \tau^A \cdot \tau_A = 1, \lambda^A \cdot \tau_A = 0,$$

54 and $\tau^A \cdot h_{AB} = \lambda^A \cdot h_{AB} = 0$. Note, that 3D metric of the V_3 has the following properties:
 55 $h^{AB} \cdot h_{AB} = 3, h^{AB} \cdot h_{BC} = -\delta_C^A$. Vector τ^A is the tangential one to the x^0 coordinate line. We can also
 56 build the same type of the differentiation operators: $\partial_\tau^+, \partial_\Lambda^+$ are the τ - and λ - projected
 57 directional normal derivatives and $g_A^B \cdot \nabla_B = \bar{\nabla}_A^+ - V_3$ stand for the projected covariant
 58 derivative. We will mark these projected values with a tilde over the operators. See [1-3] for more
 59 details.

60 3. The new Rest Mass Concept and Some Cosmological Problems

61 Let us consider the geodesic equation in 5D theory, which has the standard form:

$$62 \quad \frac{d^2 x^A}{dI^2} = -P_{BC}^A \cdot \frac{dx^B}{dI} \cdot \frac{dx^C}{dI}. \quad (1)$$

63 Here 5D interval $dI^2 = G_{AB} \cdot dx^A \cdot dx^B$ and the indexes $A, B, C, \dots = 0, 1, 2, 3, 5$. G_{AB} is the 5D metric and
 64 P_{BC}^A stands for the 5D Christoffel symbols. Making here the (4+1)-splitting with the help of the monad method
 65 and eliminating the 5th coordinate, then imposing the cylindrical symmetry condition (CSC) along the 5th
 66 coordinate (it means that we have the Killing equations for the 5D metric G_{AB} along the tangent vector to the
 67 5th coordinate $\xi_A = \varphi \cdot \lambda_A$: Lie derivative $L_\xi G_{AB} = \xi_{A;B} + \xi_{B;A} = 0$) we may obtain from Eq. (1) the
 68 following system:

69
$$\frac{D^+(\varphi \cdot \hat{p})}{ds} = 0; \quad (2)$$

70
$$\frac{D^+ \bar{p}^\alpha}{ds} = \frac{Q_0}{c^2} \cdot \bar{u}^\beta \cdot F_{\beta \cdot}{}^\alpha + \partial^{\alpha+} \hat{m}_0. \quad (3)$$

71 Here ,

72
$$\frac{D^+}{ds} = \bar{u}^\alpha \cdot \bar{\nabla}_\alpha^+ - \hat{u} \cdot \partial_\Lambda^+.$$

73 The indexes $\alpha, \beta, \gamma, \dots = 0, 1, 2, 3$.

74 We have also used an integral of the Eq. (2) and inserted it to the Eq. (3). It should be also noted, that
 75 according to the monad method we have put here the ERM \hat{m}_0 of the Lorentz type in the following
 76 form (here m_0 is an initial rest mass of the 5D test particle and $\bar{p}^\alpha = \hat{m}_0 \cdot \bar{u}^\alpha$):

77
$$dI^2 = ds^2 - d\lambda^2, \hat{m}_0 = m_0 \cdot \hat{\beta}, \hat{\beta} = (1 - \hat{u}^2)^{-1/2}, \hat{p} = \hat{m}_0 \cdot \hat{u}, \hat{u} = d\lambda / ds.$$

78 A use of the above-mentioned integral of Eq. (2) directly leads to the following exact expression for
 79 the ERM:

80
$$\hat{m}_0 = \sqrt{m_0^2 + \frac{Q_0^2}{4 \cdot k_0 \cdot \varphi^2}} = \sqrt{m_0^2 + \frac{n^2 \cdot \hat{m}_{pl}^2}{\varphi^2}}, \quad (4)$$

81 Here k_0 is the Newtonian gravitational constant, $Q_0 = n \cdot e$ (e is an electric charge of the electron), and also we
 82 have put here the 5D Plank mass $\hat{m}_{pl} = e / (2 \cdot \sqrt{k_0})$. Expression (4) may be transformed to the more compact
 83 and convenient form by introducing the mass angle:

84
$$\chi_n = \operatorname{arcsinh} \frac{n \cdot \hat{m}_{pl}}{m_0 \cdot \varphi} : \hat{m}_0 = m_0 \cdot \cosh \chi_n.$$

85 On the other side, we may represent this expression as a modulus of a complex quantity: $\hat{m}_0 = \sqrt{m_{0z} \cdot \bar{m}_{0z}}$,

86 where $m_{0z} = m_0 + i \cdot n \cdot \hat{m}_{pl} / \varphi$. The same may be written in the exponential form $m_{0z} = \hat{m}_0 \cdot e^{i \cdot \psi_n}$, where

87 the phase $\psi_n = \arctan\{n \cdot \hat{m}_{pl} / (m_0 \cdot \varphi)\}$. Here \bar{m}_{0z} is a complex conjugate to the m_{0z} value. It motivates

88 us to assume a possible connection with the quantum properties of the particles. In these terms, one may

89 consider a photon as a complex null-particle $(0; 0)$, a neutrino as a real particle $(m_\nu; 0)$, and a hypothetical

90 tachyon as an imaginary particle $(0; m_\tau)$. It seems to be natural to extend out the equivalence principle of the

91 General Relativity on the scalar gravitational field (SGF) also, and because of it to put forward the hypothesis
 92 that every 5D particle must contain an electric charge. Then, the neutral 5D particles must be composite ones. In
 93 this sense, there are only two verily 4D particles: the photon and the neutrino. An additional argument to
 94 introduce the ERM is a fact that only in this case it is possible to do the further (3+1+1)-splitting in the 5D
 95 theory with an additional elimination of the time coordinate, using the dyad method [6]. Only in this case we
 96 obtain the well-known system of the geodetic equations, which is very familiar to the following system that we
 97 usually obtain in 4D theory, making the similar (3+1)-splitting of the 4D geodetic equation [8,9]:

$$98 \quad \frac{D^+ \hat{m}}{d\tau} = \hat{p}^i \cdot (\bar{F}_i - \bar{v}^k \cdot \bar{D}_{ik}) - \frac{Q_0}{c^2} \cdot \bar{v}^i \cdot \bar{E}_i + \frac{Q_0 \cdot \sqrt{1-\bar{v}^2}}{2 \cdot \sqrt{k_0} \cdot \varphi} \cdot \partial_\tau^+ \hat{m}_0, \quad (5a)$$

$$99 \quad \frac{D^+ \hat{p}^i}{d\tau} = \frac{Q_0}{c^2} \cdot (\bar{E}^i - \bar{v}^k \cdot \bar{H}_{.k}) - \hat{m} \cdot \bar{F}^i + 2 \cdot \hat{p}^k \cdot \bar{D}_k^i + \frac{Q_0 \cdot \sqrt{1-\bar{v}^2}}{2 \cdot \sqrt{k_0} \cdot \varphi} \cdot \partial^{+i} \hat{m}_0. \quad (5b)$$

100 Here

$$101 \quad \frac{D^+}{d\tau} = \partial_\tau^+ + \hat{v} \cdot \partial_\Lambda^+ + \bar{v}^i \cdot \bar{\nabla}_i^+.$$

102 The indexes $i, j, k, \dots = 1, 2, 3$.

103 In another case, it occurred that it is simply impossible to perform the (3+1)-splitting of the equation (3) because
 104 the equations (5a) and (5b) are interlaced so strongly that we have no chances to disconnect them. The last terms
 105 in equations (5a) and (5b) are the scalar 3-forces which are caused by the dependence of the ERM on the SGF.

106 Here we have put $\hat{m} = \hat{m}_0 / \sqrt{1-\bar{v}^2}$, $\bar{v}^i = d\bar{x}^i / d\tau$, $\hat{v} = d\lambda / d\tau$. Then, $\hat{p}^i = \hat{m} \cdot \bar{v}^i$,
 107 $\hat{p} = \hat{m} \cdot \hat{v} = \hat{m}_0 \cdot \hat{u}$, and $d\tau = c \cdot dt$. It should be noted also that the scalar 3-forces in equations (5a) and
 108 (5b) vanish for the zero rest mass particles, as it occurs in all of the scalar-tensor theories due to the factor
 109 $\sqrt{1-\bar{v}^2}$. The ERM \hat{m}_0 in the 5D theory appears just as the relativistic mass m in 4D theory:

$$110 \quad dl \rightarrow ds \Rightarrow ds^2 = d\tau^2 - dl^2; m = m_0 \cdot (1-v^2)^{-\frac{1}{2}} = m_0 \cdot \beta; v = \frac{dl}{d\tau}; d\tau = c \cdot dt$$

111 and it is easy to prove that the same result we obtain whenever we increase the dimension: $V_n \mapsto V_{n+1}$.

112 Thus, following the idea of P. Ehrenfest [10], one can see here that the increase in the dimension adds some new
 113 features to the physical nature of the particles.

114 The geodetic equation (3) we can easily bring to the following form:

$$115 \quad \hat{m}_0 \cdot \frac{D^+ \bar{u}^\alpha}{ds} = f_L^\alpha + f_{BD}^\alpha. \quad (6)$$

116 Here

117 $f_L^\alpha = \frac{Q_0}{c^2} \cdot \bar{u}^\beta \cdot F_{\beta}^\alpha$ is the 4D Lorenz force and $f_{BD}^\alpha = -\frac{Q_0^2 \cdot P^{\alpha\beta} \cdot \Phi_\beta}{4 \cdot k_0 \cdot \hat{m}_0 \cdot \varphi^2}$ is the 4D scalar one or, as we

118 have proposed [9], the Brance-Dicke force. Here $P^{\alpha\beta} = g^{\alpha\beta} - \bar{u}^\alpha \cdot \bar{u}^\beta$ is a tensor of the orthogonal

119 projection upon the 4D-velocity \bar{u}^α direction and $\Phi_\beta = \bar{\nabla}_\beta^+ \ln \varphi$. It should be marked that the Lorentz force

120 for big masses usually is equal to zero because it depends on the Q_0 linearly. But the Brance-Dicke force

121 depends on the Q_0^2 and although it is very weak it can be very significant, being accumulated for the big

122 masses (BME – the big masses effect). In addition, the scalar force is negative and has the repellent property.

123 Thus, it may be one of the reasons of the Universe expansion. I believe also that the expression (4) for the ERM

124 permits us to think that an electric charge e superposes simultaneously the role of the scalar charge, so it can

125 explain to us why we could not find this scalar charge so far. Also, one can hope that an interplay between the

126 SGF and an electric charge can explain how the dark matter may contribute to the total mass of the 4D Universe,

127 which is immersed in some sense into the scalar ocean. It should be added also that the SGF contributes to the

128 braking radiation force (BRF) [11] and permits us to generalize the concept of the one:

129 $g^\alpha = g_E^\alpha + g_S^\alpha + g_{ES}^\alpha$. Here we should tell the total BRF may be represented as a sum of the electromagnetic,

130 scalar and mixed parts, where the first one has the following form:

$$131 \quad g_E^\alpha = \frac{2 \cdot e^3 \bar{u}^\gamma}{3 \cdot \hat{m}_0 \cdot c^3} \cdot \left(\bar{u}^\beta \cdot \bar{\nabla}_\gamma^+ F_{\beta}^\alpha + \frac{e \cdot P^{\alpha\delta}}{\hat{m}_0 \cdot c^2} \cdot F_{\beta\delta} \cdot F_{\gamma}^\beta \right);$$

132 and the second one has the form as below:

$$133 \quad g_S^\alpha = -\frac{e^4 \cdot P^{\alpha\beta} \cdot \bar{u}^\gamma}{6 \cdot c \cdot k_0 \cdot \varphi^2 \cdot \hat{m}_0^2} \cdot \left(\bar{\nabla}_\gamma^+ \Phi_\beta - 2\Phi_\beta \cdot \Phi_\gamma + \frac{3 \cdot e^2 \cdot \Phi_\beta \cdot \Phi_\gamma}{4 \cdot k_0 \cdot \varphi^2 \cdot \hat{m}_0^2} \right);$$

134 And at last the third one has the next form:

$$135 \quad g_{ES}^\alpha = -\frac{e^5}{6 \cdot c^3 \cdot k_0 \cdot \varphi^2 \cdot \hat{m}_0^3} \cdot \left(g^{\beta\delta} \cdot P^{\alpha\gamma} - 3 \cdot g^{\alpha\gamma} \cdot \bar{u}^\beta \cdot \bar{u}^\delta \right) \cdot F_{\beta\gamma} \cdot \Phi_\delta.$$

136 See [9] for more details.

137 4. The 5D Ricci identities and some problems of the Astrophysics and Cosmology

138 We should remember some interesting and perspicacious ideas of the scientific classics about the

139 connections between world geometry and physical interactions. Following one of the greatest

140 mathematicians of the XIXth century - W.K. Clifford [12] – let us consider the well-known in

141 Riemannian geometry Ricci identities in the case of the 5D theory:

$$142 \quad R_{BCD}^A + R_{DBC}^A + R_{CDB}^A = 0 \quad . \quad (7)$$

143 The (4+1)-splitting of the V_5 gives us the following relations between the 4D physical and
144 geometrical values [9] starting with the 4D Ricci identities:

$$145 \quad R_{\beta\gamma\delta}^{\alpha} + R_{\delta\beta\gamma}^{\alpha} + R_{\gamma\delta\beta}^{\alpha} = 0 . \quad (8)$$

146 Then, we can obtain the following very important connection between the electromagnetic tensor and the curls
147 of the SGF gradients. It gives us a very specific trigger of the changing electromagnetic process:

$$148 \quad \sqrt{\bar{\kappa}_0} \cdot \partial_{\Lambda}^+ \varphi \cdot F_{\alpha\beta} = \Phi_{\alpha;\beta} - \Phi_{\beta;\alpha} = -m_{\alpha\beta} . \quad (9)$$

149 Here $\bar{\kappa}_0 = 4 \cdot k_0 / c^4$ and the SGF curls $m_{\alpha\beta}$ we may consider as the quasi-particles of the magnetic
150 monopole type. In some sense, these quasi-particles are similar to solitons [13,16]. Importantly, that imposition
151 of the CSC gives rise to the vanishing of these scalar curls, since only in this case the necessary and sufficient
152 condition for SGF to be laminar ($\Phi_{\alpha;\beta} = \Phi_{\beta;\alpha}$) may be satisfied.

153 At last, making use of Eq. (9), we obtain the first pair of the Maxwell equations with a nonzero r.h.s. of the
154 magnetic monopole type, namely:

$$155 \quad F_{\alpha\beta;\gamma} + F_{\gamma\alpha;\beta} + F_{\beta\gamma;\alpha} = \frac{2}{\sqrt{\bar{\kappa}_0} \cdot \partial_{\Lambda}^+ \varphi} \cdot \begin{vmatrix} \Phi_{\alpha} & \nabla_{\alpha}^+ & \Phi_{\alpha} \\ \Phi_{\beta} & \nabla_{\beta}^+ & \Phi_{\beta} \\ \Phi_{\gamma} & \nabla_{\gamma}^+ & \Phi_{\gamma} \end{vmatrix} . \quad (10)$$

156 Imposition of the CSC here also gives rise to the vanishing of the r.h.s of Eq. (10), which reduces this equation
157 to the conventional form:

$$158 \quad F_{\alpha\beta;\gamma} + F_{\gamma\alpha;\beta} + F_{\beta\gamma;\alpha} = 0 . \quad (11)$$

159 Thus, we can conclude that the first pair of the Maxwell equations is obliged to the Riemannian structure of the
160 V_5 and has another origin than the second one. The r.h.s. R in Eq. (10) we can rewrite in a more convenient
161 form:

$$162 \quad R = \frac{2}{\sqrt{\bar{\kappa}_0} \cdot \partial_{\Lambda}^+ \varphi} \cdot (\Phi_{\alpha} \cdot m_{\beta\gamma} - \Phi_{\beta} \cdot m_{\alpha\gamma} + \Phi_{\gamma} \cdot m_{\alpha\beta}) . \quad (12)$$

163 This gives us the possibility to represent the second pair of the Maxwell equations, which we have obtained
164 from the 5D variation principle [15], with the same monopole type of the r.h.s., if we believe that R is the linear
165 combination of the magnetic monopoles currents. With this purpose we write down the mixed, x^5 - and V_4 -
166 projected, 5D field equations with the energy-momentum tensor of the 5D dust as a r.h.s.:

$$167 \quad \bar{\nabla}_{\nu}^+ (\varphi^3 \cdot F^{\mu\nu}) = -\frac{8 \cdot \pi}{c^2} \cdot \sqrt{k_0} \cdot \varphi \cdot Q_5^{\mu} . \quad (12)$$

168 Here

$$169 \quad Q^{AB} = \mu_0 \cdot c \cdot \frac{dx^A}{dI} \cdot \frac{dx^B}{d\tau} ,$$

170 where μ_0 is a matter density [15]. But before we have not imposed the CSC, we may easily bring the Eq. (12) to
 171 the following form (further we will avoid to mark the projected values with a tilde), casting temporarily aside
 172 the r.h.s. in Eq. (12) :

$$173 \quad \nabla_{\nu}^{+} F^{\mu\nu} = 3 \cdot \Phi_{\nu} \cdot m^{\nu\mu} = -\frac{4 \cdot \pi}{c} \cdot j_m^{\mu} . \quad (13)$$

174 Here we can interpret j_m^{μ} as a magnetic monopole current, where we can believe that curl $m_{\mu\nu}$ corresponds,
 175 for example, to the «north»-particle of the magnetic monopole type and $m_{\nu\mu}$ - to the «south» one, or vice versa.

176 Let us call them «n-monopole» and «s-monopole». Thus, we believe that there exist the two types of the
 177 magnetic monopole «charges»: n and s ones.

178 Finally, we can establish that just after we have imposed the CSC, as a result of it, the SGF curls disappear and
 179 the Eq. (2) gives us an integral of the movement along the 5th coordinate connected with an electric charge of
 180 the 5D test particle, so thus the energy-momentum tensor of the 5D dust (the r.h.s. in Eq. (12)) easily transforms
 181 in the second pair of the Maxwell equations and gives the new r.h.s. of the conventional type as follows below
 182 [15]:

$$183 \quad \nabla_{\nu}^{+} F^{\mu\nu} = -\frac{4 \cdot \pi}{c \cdot \varphi^2} \cdot j_e^{\mu} . \quad (14)$$

184 Here $j_e^{\mu} = \rho_0 \cdot \bar{v}^{\mu}$ is the 4D electric current (ρ_0 is an electric charge density).

185 The analysis of these results leads us to the fundamental idea about the evolution of the 4D Universe. The
 186 process of the transition to the CSC in the 4D Universe seems to be very familiar with the second-order
 187 transition in liquid helium. Then, we may hypothesize that here we have some kind of a second-order phase
 188 transition in cosmology. Let us call it the topological one [16]. It is quite possible that this transition is
 189 connected with the cooling of the expanding Universe. Also, we can suppose that, as a result of this transition,
 190 we have some kind of the superfluidity state of the scalar matter and it may accelerate the expansion of the 4D
 191 Universe [16]. This acceleration, as well known, was discovered more than twenty years ago with the help of
 192 the observations after the type Ia supernovas explosions. Also, we may add that this transition and, then, further
 193 compactification of the 5th coordinate is possibly caused by the Casimir effect [17]. Besides it, one can assume
 194 that in modern times we cannot find the magnetic monopoles in the 4D Universe, maybe only a few of the relict
 195 ones [9,18]. Also one can hypothesize, basing on the expression (4), that the ERM of the particles possibly
 196 varies with the cosmological time, depending on the scalar φ , which is connected with the 5D metric G_{55} . In
 197 the work [19] the authors have considered the 5D metric of the Kasner type, which depends on the cosmological
 198 time.

199 5. Discussion

200 This article was written because for a long time the author has accumulated a lot of very interesting
 201 results in the 5D theory of the Kaluza type, which has been almost forgotten during the last fifty
 202 years. It has occurred that if one goes beyond the 5D optics, certain new and very non-trivial

203 properties of the matter and 4D Universe may appear. Following this approach, the author has
 204 succeeded to generalize the rest mass concept and understanding more deeply the quantum nature
 205 of the matter. The implementation of the monad method the (4+1)-splitting of the 5D Ricci identities
 206 makes it possible to understand how the Riemannian structure of the World affects its physical
 207 properties. It permits one to approach closer to the understanding of the magnetic monopole
 208 problem and the origins of the Maxwell equations. The obtained results also provide new insight
 209 into the mechanism of the accelerated expansion of the 4D Universe. The author believes that the
 210 application of the developed approach extends far beyond the specific problem discussed in the
 211 present paper. Its generalization beyond the 5D case may be also worth-while.

212 6. Conclusions

213 Finally, the author would like to express the hope that these results may open a new page in the
 214 investigations of the 5D theory and will attract more attention of the physical community.

215 **Funding:** This research received no external funding.

216 **Acknowledgments:** The author would like to thank Yu.S. Vladimirov, A.V. Solovyov, M.I. Tribelsky, and V.A.
 217 Berezin for useful discussions. The author is also grateful to A.V. Solovyov for valuable help in the manuscript
 218 preparation.

219 **Conflicts of Interests:** The author declares no conflicts of interest.

220 References

- 221 1. Vladimirov, Yu.S. *Classical gravity theory*; Krasand: Moscow, Russia, 2018; pp. 184-207.
- 222 2. Aliyev, B.G. Motion equations in the 5D unified field theory. In *Abstracts of the IX-th Int. Conf. on Gen. Relat.*
 223 *Grav.*, Vol. 3, Jena, Germany (GDR), 1980, p. 679.
- 224 3. Aliyev, B.G. 5D theory of the scalar-tensor gravitation and electromagnetism in dyad form. In *Problems of*
 225 *the theory of gravitation and elementary particles*, Stanyuckovich K.P., Ed.; Atomizdat: Moscow, Russia, 1979;
 226 issue 10, pp. 141-149, (in Russian).
- 227 4. Fock, V.A. Über die invariante Form der Wellen- und der Bewegungsgleichungen für einen geladenen
 228 Massenpunkt. *Zeitschr. für Physik*, **1926**, Bd 39, S. 226-232.
- 229 5. Klein, O. Quantentheorie und fünfdimensionale Relativitätstheorie. *Zeitschr. für Physik*, **1926**, Bd 37, S. 895.
- 230 6. De Broglie, L. L'Univers a cinq dimensions et la mecanique ondulatoire. *J. de Physique*, Serie VI **1927**, Vol. 8,
 231 № 2, p. 65.
- 232 7. Rumer, Yu.B. *Investigations on 5D Optics*; State Publishing House of the Technical and Theoretical
 233 Literature: Moscow, Russia, 1956, (in Russian).
- 234 8. Aliyev, B.G. The behavior of the charged particles in 5D gravity theory. In *The modern problems of the*
 235 *general relativity theory*, Physical Institute of the Byelorussian Academy of Sciences: Minsk, Byelorussian
 236 republic, 1979; pp. 154-160, (in Russian).
- 237 9. Aliyev, B.G. The effective rest mass concept and magnetic monopole problem in 5D Theory. In *Gravitation,*
 238 *Astrophysics, and Cosmology*, Proceedings of the ICGAC-12, Moscow, Russia, June 28-July 4, 2015;
 239 Melnikov, V. and Jong-Ping Hsu, Eds; World Scientific: Singapore, 2016, pp.321-326.
- 240 10. Ehrenfest, P. In what way does it become manifest in the fundamental laws of Physics that space has three
 241 dimensions. In *Proc. Amsterdam Ac.* **1917**, v. 20, p. 200.
- 242 11. Landau, L.D. and Lifshitz, E.M. *The classical theory of fields, v.2, 3th Ed.*; Pergamon Press, Maxwell House:
 243 New York-London, USA-GB, 1971.
- 244 12. Clifford, W.K. *Mathematical Papers*; MacMillan: New York-London, 1968.
- 245 13. Gross, D., Perry M. Magnetic monopoles in Kaluza-Klein theories. *Nucl. Phys. B*, **1983**, v. 226, p.29.
- 246 14. Aliyev, B.G. The rest mass concept and some problems of Cosmology in 5D Theory. In *Abstracts of the*
 247 *RUSGRAV-16 (16th Rus. Int. Conf. on Gen. Relat. Grav.)*; BFU named I. Kant: Kaliningrad, Russia, 2017, p.
 248 91.

- 249 15. Aliyev, B.G. On the energy-momentum tensor of the 5D dust. In *Abstracts of the reports of the Int.*
250 *school-seminar "Multidimensional Gravity and Cosmology"*; RGA (Rus. Grav. Assoc.): Moscow, Russia, 1994,
251 p. 1.
- 252 16. Aliyev, B.G. The solitons and the topological second-order transition in 5D Theory. In *Abstracts of the*
253 *RUSGRAV-16 (16th Rus. Int. Conf. on Gen. Relat. Grav.)*, BFU named I. Kant: Kaliningrad, Russia, 2017, p. 91.
- 254 17. Mostepanenko, V.M. and Trunov, N.N. *Casimir Effect and its Applications*; Oxford University Press: Oxford,
255 GB, 1997; pp. 191-193.
- 256 18. Chodos, A., Detweiler, S. Where has the fifth dimension gone? *Phys.Rev. D* **1980**, v. 21, p.2167.
- 257 19. Aliyev, B.G. Where has the magnetic monopole gone? In *Abstracts of the ICGAC-12, Moscow, Russia, June*
258 *28-July 4, 2015*; PFUR: Moscow, Russia, 2015, p. 110.
- 259



© 2020 by the authors. Submitted for possible open access publication under the terms and conditions of the Creative Commons Attribution (CC BY) license (<http://creativecommons.org/licenses/by/4.0/>).