

What is the Real Shape of the Hubble Diagram, $z = H_0 \cdot D$ or $z + 1 = e^{H_0 \cdot t}$?

Analysis of the SN1a Supernovae and Gamma Ray Burst Redshift/Magnitude Data including the High Redshift Range up to $z = 8.1$

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Abstract

Analyses of the Hubble diagrams are presented for SN1a supernovae and gamma ray bursts in the redshift ranges $z = 0.01\text{--}1.3$ and $0.034\text{--}8.1$, respectively. Data are presented on the typical z/μ scale and also on the less common yet increasingly sensitive photon flight time $t/(z+1)$ scale. The primary conclusion is that on the basis of the presently accessible data in the redshift range $z = 0.01\text{--}8.1$, the slope of the Hubble diagram is, or is extremely close to, exponential.

1 Introduction

Hubble's constant (H_0) [1], in addition to the cosmic microwave background (CMB) and the big bang nucleosynthesis, is one of the key pillars of big bang cosmology. It is presumably the most significant cosmological parameter, since it determines the rate of expansion and the age and size of the universe, and influences numerous different parts of cosmology. As indicated by its great significance, the velocity interpretation of H_0 has achieved a dogmatic status and other hypotheses, such as the most discussed rival theory: an energy decrease of starlight with a constant rate, for example, were discarded. However, ongoing results have demonstrated

hitherto unrecognized vulnerabilities in the understanding of the redshift (RS, z) of starlight, which requires careful reevaluation.

H_0 is normally determined by direct measurement of the RS of atomic spectral lines emitted by distant galaxies, leading to $H_0 = 68.5$ (Planck reference set)– $73 \text{ km s}^{-1} \text{ Mpc}^{-1}$ (Millennium reference set) as the most probable range of H_0 .

A further method of determining H_0 is offered by the CMB power spectrum [2, 3]. Fitting the cosmological parameters of the lambda cold dark matter (Λ CDM) model with the cosmological constant to the baryonic acoustic oscillation (BAO) gives $H_0 = 68 \text{ km s}^{-1} \text{ Mpc}^{-1}$, which is close to the value derived from spectral lines.

This outcome is, however, not unambiguous. The BAO can be fitted with a similar, or even better, certainty using various parameters of the multi-parameter field, H_0 , Ω_M , Ω_Λ and w .

It was recently shown that the Einstein–de Sitter (EdeS) model with a zero cosmological constant can fit the CMB data as well as the concordance model. Calculating H_0 from different EdeS models without the cosmological constant gives a value of $H_0 = 30\text{--}40 \text{ km s}^{-1} \text{ Mpc}^{-1}$ [4, 5].

This result is fairly surprising. A H_0 of $30 \text{ km s}^{-1} \text{ Mpc}^{-1}$ is entirely inconsistent with values of $68.5\text{--}73 \text{ km s}^{-1} \text{ Mpc}^{-1}$. This logical inconsistency between the two extremely different H_0 values requires prudent clarification.

A further question emerges from the poor knowledge of the Hubble diagram (HD) in the high RS range of $1 < z < 8$. It is commonly acknowledged that the historical scenery to the expansion of the universe can be exactly represented by the concordance model, which makes explicit predictions regarding the shape of the HD.

In the RS range $z = 0.0104\text{--}1.3$, the RS/magnitude (z/μ) HD of SN1a supernovae appears to affirm the concordance model and the HD gives a reasonably good fit to the forecasts of the concordance model [6–8]. However, in addition, an exponential slope in the low RS range cannot be excluded on the basis of the presently available data [9–11]. Unfortunately, the fit

of the Λ CDM model to the SN1a z/μ data applies to only a limited range of distances. One explanation for this lies in the experimental difficulties involved, because at RSs $> \sim 1.3$, the optical light emitted by supernovae progressively dims with distance and exact measurements become difficult.

This constraint on the data has motivated several attempts to obtain cosmological parameters from gamma-ray burst (GRB) observations. GRBs are the most brilliant sources in the universe. They are acquired up to RSs of ~ 8 and higher, and there is hope that these objects may serve as dependable distance indicators.

A number of endeavors have been made to use GRB data to calculate the HD [12–16], with varying degrees of success. Gupta [16] has shown that by weighting a number of different cosmological models with supernovae and GRB z/μ data, the results slightly favor the Λ CDM model but do not support firm conclusions. All models can fit the data well with a fairly high χ^2 probability. Lerner [17] also achieved a similar result and has shown that no single model provides a statistically better fit to the same z/μ dataset than another.

Ongoing work in the high RS range of up to $z = 8.1$ using mixed SN1a and GRB data [18–20], or solely GRB data [21], has demonstrated that the slope of the HD is, or is extremely close to, exponential, and this contradicts the expectations of the Λ CDM model.

The aim of the present study is to perform a comparative HD test using SN1a supernova and calibrated, cosmology independent GRB RS/ μ data points. In the high RS range, it should be possible to definitively verify whether the HD shows the linear

$$z = H_0 D \quad (1)$$

or the exponential

$$z + 1 = e^{H_0 * t} \quad (2)$$

relationship, an issue that is not easily detectable in the $z < 1$ region. Explaining this important question could set significant limits on the real rate of expansion and on the matter and energy content of the universe.

In Equations (1)–(3), t represents the flight time of the photons (s) from the co-moving radial distance D (cm) to the observer, which is proportional to D (Mpc) as used in the Hubble law.

2 Preparation of Hubble Diagram

The HD is a linear plot of the measured distance (usually Mpc) versus the measured RS, which is often represented on the less sensitive logarithmic RS/ μ scale.

Since the difference between the measured and the calculated data becomes more pronounced on the more sensitive linear scale, in addition to the z/μ HD, a plot of the photon flight time (t) versus RS was used for representation of the HDs. The photon flight time was calculated from

$$t = \frac{D}{c} = \frac{10^{(\mu+5)/5}}{(z+1) \cdot 3 \cdot 10^{10}} \cdot 3,085 \cdot 10^{18} \quad (3)$$

For visualization of the data the $t \cdot 10^{-14}/z$ data presentation was chosen.

In addition to the increased sensitivity, an obvious advantage of the $t/(z+1)$ representation is the direct illustration of the shape of the HD, which can be directly compared with the predictions following from Equations (1) and (2).

3 Data Collection and Processing

3.1 Dataset for SN1a Supernovae

The joint light-curve analysis (JLA) data index, the most recent list of 740 spectroscopically affirmed SN1a supernovae with excellent light curves, is believed to currently be the most exact supernovae z/μ database. From these data, 32 binned RS/ μ data points were calculated by Betoule et al. [22], which give an exact representation of the 740 data points.

3.2 Processing of SN1a Data

The high precision of the JLA data allows for the direct conversion of the observed z/μ data points into the corresponding $t/(z+1)$ data set without previous smoothing of the data.

3.3 Dataset for Gamma Ray Bursts

A total of 138 calibrated, cosmology independent GRB z/μ data points collected by Liu and Wei [23] from 557 Union2 compilations were utilized as the starting dataset. Compared with the 2010 data release containing 109 z/μ data points [24], the 2015 dataset contains 30 more data points and the data are more precise, particularly in the low z range.

3.4 Processing Gamma Ray Burst Data

In view of the experimental difficulties in determining the z/μ data, it is likely that in addition to the observed considerable scatter, large datasets taken from different observations and from different sources will contain outliers. If these are not removed in the refinement procedure, they will overwhelm the fit. Any single distant outlier can impose considerable bias in the best fit of the HD and ruin the result. It is therefore important to clean the data by eliminating the outliers and to use these cleaned data for statistical analysis. For this purpose, the data smoothing and for the following identification of outliers, best fit curves were calculated using the empirical potential function

$$\mu = a * z^n, \quad (4)$$

which was found in earlier studies to be the best mathematical approximation for describing the slope of the observed z/μ diagrams [18, 19, 25].

3.5 Identification of Outliers: The Grubbs Test

The well-known Grubbs test [26] was used for the identification of erroneous luminosity indicators. The Grubbs test is used to detect outliers in a data set of N values that are nearly normally distributed. Assuming a normal distribution of the sample, as confirmed by the very low skew in Table 1, the test is performed by computing x_0 , defined as:

$$X_o \geq \frac{G_G * STABW + x_{Mean}}{\sqrt{\frac{N}{N-1}}} \quad (5)$$

where:

x_0 is the suspected outlier;

x_{Mean} is the absolute value of the mean of N data points;

N is the number of data points;

STABW is the standard deviation of N values;

G_G is the Grubbs number.

G_G can be found in statistical tables for different levels of confidence and numbers of data points. For 138 data points, the value of G_G is 1.958 at the 95% confidence level. **If the value of x_0 calculated from ($\mu_{\text{measured}} - \mu_{\text{calculated on basis of the best fit}}$) is found to be greater than the numerical value of the right-hand side of Equation (5), the data point in question is discarded. Then, on the basis of the reduced dataset and new a and b coefficients, the mean and the new STABW are calculated and so on.**

The goodness of fit was calculated using the likelihood estimator

$$\chi^2 = \sum_{i=1}^N \frac{(\mu(t)_{\text{obs.}} - \mu(t)_{\text{calc.}})^2}{\mu(t)_{\text{calc.}}} \quad (6)$$

3.6 Luminosity Distances

Luminosity distances were calculated using the ICRAR cosmological calculator [27].

3.7 Data Presentation

Excel and Excel solver were used for data fitting, refinement and data presentation.

4 Results

4.1 Hubble Diagram for SN1a Supernovae in the RS Range 0.0104–1.3 on Logarithmic z/μ and Linear $t/(z+1)$ Scales

4.1.1 Logarithmic Data Presentation

Figure 1 shows the logarithmic z/μ HD for the observed z/μ values (squares) together with the z/μ data points calculated with ICRAR, with $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_M = 0.295$ and $w = -1.018$ (triangles) [20, 27].

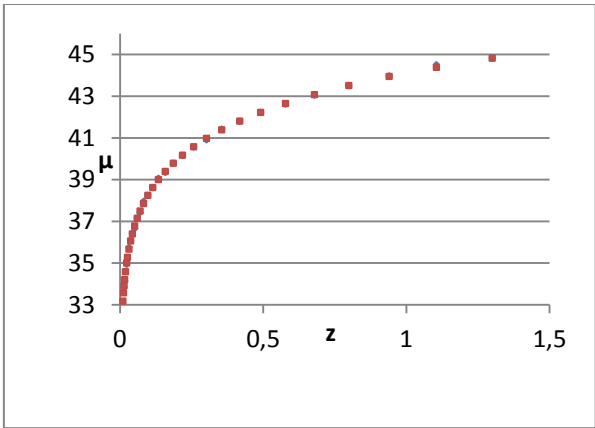


Figure 1

Figure 1: Logarithmic z/μ HD for the observed z/μ data (squares) together with the z/μ data points calculated with ICRAR, with $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_M = 0.295$ and $w = -1.018$ (triangles) [20, 27].

The fit results are summarized in Table 1 below.

Table 1

SN1a supernovae

Descriptive statistics for the observed z/μ data related to the calculated z/μ .

Valid cases	R^2	F-test $\mu_{\text{obs}}/\mu_{\text{calc}}$	$\sum \chi^2$ $(\mu_{\text{obs}} - \mu_{\text{calc}})$
31	0.9994	0.9932	0.0062
Std. deviation $(\mu_{\text{obs}} - \mu_{\text{calc}})$	Variance $(\mu_{\text{obs}} - \mu_{\text{calc}})$	Skew $(\mu_{\text{obs}} - \mu_{\text{calc}})$	Mean $(\mu_{\text{obs}} - \mu_{\text{calc}})$
0.084	0.0072	-1.09	-0.0061

As can be seen from Figure 1 and Table 1 (which represents the difference between the observed and calculated data quantitatively), the two data sets are practically congruent.

4.1.2 Data Presentation on Linear $t/(z+1)$ Scale

In Figures 2 and 3, the corresponding HDs are plotted on the linear $t \cdot 10^{-14}/(z+1)$ scale for the observed and calculated μ values.

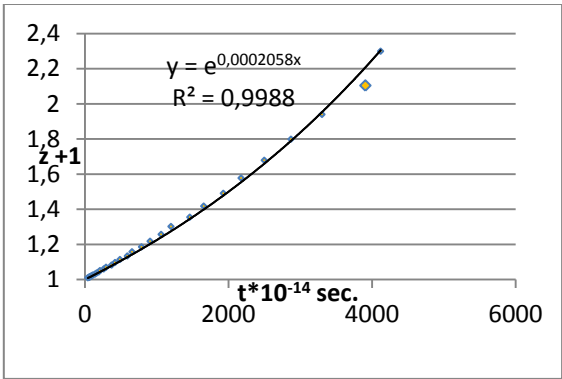


Figure 2

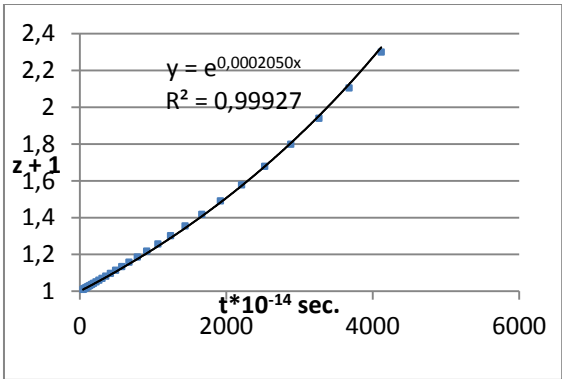


Figure 3

Figure 2: $t/(z+1)$ HD: trendline (solid line), observed data points (dotted line).

Figure 3: $t/(z+1)$ HD: trendline (solid line), calculated data points (dotted line).

(The orange point at $z+1 = 2.105$ in Figure 2 is obviously an ‘outlier’ and was therefore excluded from the calculation of the goodness of fit indicators.)

Descriptive statistics are summarized in Tables 2 and 3.

Table 2

SN1a supernovae

Data fitting with exponential function, observed data

Descriptive statistics $t/(z+1)$

Valid cases	H_0	R^2	F-test z_{obs}/z_{calc}	$\sum \chi^2$ $(z_{obs}-z_{calc})$
30	0.0002058 $H_0 = 2.058 \cdot 10^{-18}$ s^{-1}	0.9988	0.9088	0.00256
Std. deviation $(z_{obs}-z_{calc})$	Variance $(z_{obs}-z_{calc})$	Skew $(z_{obs}-z_{calc})$	Mean $(z_{obs}-z_{calc})$	
0.01199	0.00014	2.21	-0.0033	

Table 3

SN1a supernovae

Data fitting with exponential function, calculated data

Descriptive statistics $t/(z+1)$

Valid cases	H_0	R^2	F-test z_{obs}/z_{calc}	$\sum \chi^2$ $(z_{obs}-z_{calc})$
31	0.0002050 $H_0 = 2.05 \cdot 10^{-18} \text{ s}^{-1}$	0.99927	0.9236	0.00176
Std. deviation $(z_{obs}-z_{calc})$	Variance $(z_{obs}-z_{calc})$	Skew $(z_{obs}-z_{calc})$	Mean $(z_{obs}-z_{calc})$	
0.0089	$7.59 \cdot 10^{-5}$	1.82	0.0032	

In the RS range 0–1.3, the HD for the observed data and the HD based on the calculated z/μ data can be described with high accuracy by both the Λ CDM model and the exponential functions $z+1 = e^{0.0002058 \cdot t \cdot 10^{-14}}$ and $z+1 = e^{0.0002050 \cdot t \cdot 10^{-14}}$, respectively. For both the linear and logarithmic scales, the two cases are, within statistical uncertainties, similar.

This result is rather unexpected and probably of an accidental nature that leaves two options open: either, in the low RS range at certain values of the parameters Ω_M , Ω_Λ , H_0 and w , the shape of the HD is very close to exponential, or, the shape of the observed HD is really exponential not only in the low but also in the high RS range. Measurements including the high RS region are necessary for clarifying this important question.

4.2 Hubble Diagram for Gamma Ray Bursts on Logarithmic z/μ and Linear $t/(z+1)$ Scales including the High Redshift Range up to $z = 8.1$

4.2.1 Logarithmic Data Presentation

From Figure 4, we can see that the GRB data are affected by considerable scatter and can therefore not be used for a precise evaluation. Careful data cleaning is necessary in order to obtain a consistent outlier free dataset. Starting from 138 data points, four refinement steps were necessary to produce 96 outlier-free data points (Figure 5), which considerably

improved the goodness of fit indicators (Tables 5 and 6). The best fit line $\mu = 44,035 \cdot z^{0.0588}$

allows a scatter free conversion of the z/μ data into the $t/(z+1)$ scale.

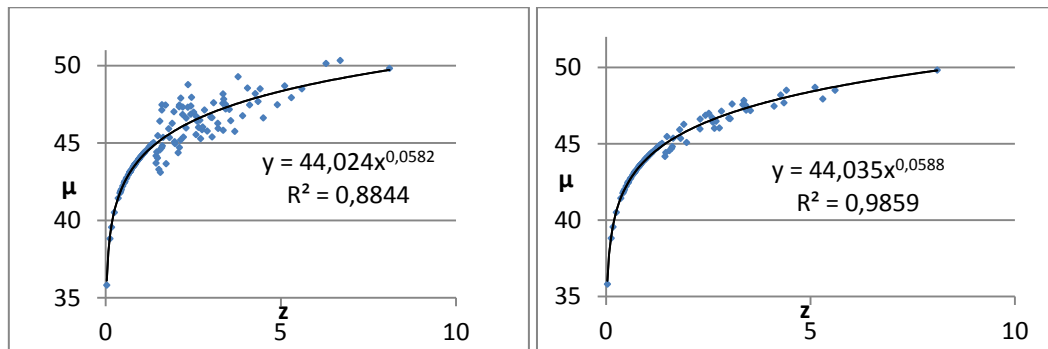


Figure 4

Figure 5

Figure 4: z/μ diagram for the original 138 data points.

Figure 5: z/μ diagram for the 96 cleaned data points.

Table 4

Results of regression with $\mu = a \cdot z^b$

Descriptive statistics z/μ for 138 data points

Valid cases	a	b	R^2	$\sum \chi^2$ ($\mu_{\text{obs}} - \mu_{\text{calc}}$)
138	44.024	0.582	0.8844	1.925
F-test	Std. deviation	Variance	Skew	Mean
$\mu_{\text{obs}} / \mu_{\text{calc}}$	($\mu_{\text{obs}} - \mu_{\text{calc}}$)		($\mu_{\text{obs}} - \mu_{\text{calc}}$)	($\mu_{\text{obs}} - \mu_{\text{calc}}$)
0.4376	0.8057	0.65	0.027	0.073

Table 5

Results of regression with $\mu = a \cdot z^b$

Cleaned dataset

Descriptive statistics z/μ for 96 cleaned data points

Valid cases	a	b	R^2	$\sum \chi^2$ ($\mu_{\text{obs}} - \mu_{\text{calc}}$)
96	44.035	0.0588	0.9859	0.1607
F-test	Std. deviation	Variance	Skew	Mean

$\mu_{\text{obs}}/\mu_{\text{calc}}$	$(\mu_{\text{obs}}-\mu_{\text{calc}})$	$(\mu_{\text{obs}}-\mu_{\text{calc}})$	$(\mu_{\text{obs}}-\mu_{\text{calc}})$	$(\mu_{\text{obs}}-\mu_{\text{calc}})$
0.9668	0.2793	0.078	-0.834	-0.0128

4.2.2 Data Presentation on Linear $t/(z+1)$ Scale

Figure 6a shows a plot of the linear $t*10^{-14}/(z+1)$ HD from the best fit parameter (Table 5) of the observed data (line a) together with the curve inferred from the calculated Λ CDM data with $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_M = 0.295$ and $w = -1.018$ (line b) [22].

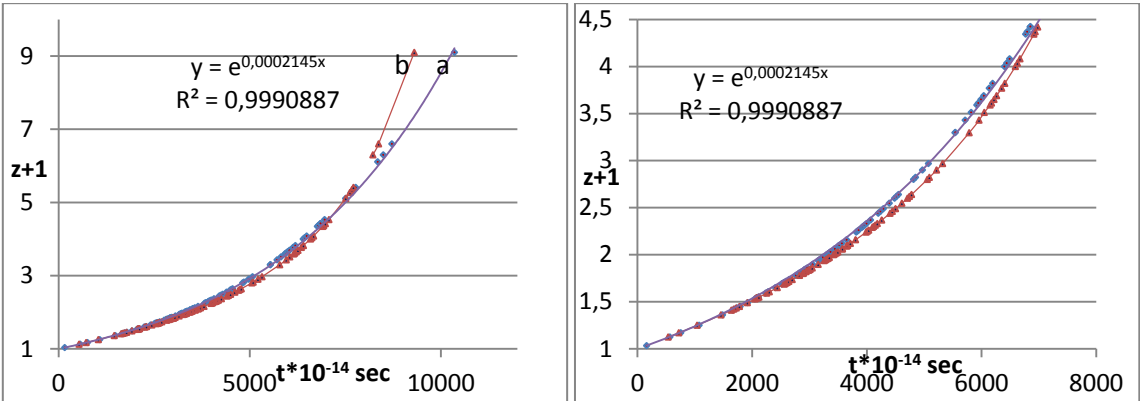


Figure 6a

Figure 6b

Figure 6a: Representative $t^{*10^{-14}}/(z+1)$ HD based on 97 statistically verified GRB data points (line a) and the calculated Λ CDM data (line b).

Figure 6b: Section of Figure 6a in the RS range from 1–4.5.

The shape of the HD for the observed data (line a in Figure 6a) is very close to exponential, whilst the Λ CDM model shows systematic deviations from the exponential best fit curve (line b) in Figure 6a):

$$\sum \chi^2 t \text{ (exponential best fit – observed data)} = 108,$$
$$\sum \chi^2 t \text{ (exponential best fit – } \Lambda\text{CDM model)} = 791.$$

The HD diagram on basis of the Λ CDM model with $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ deviates below the trendline of the best-fit curve for $z + 1 < \sim 4.5$ to the bottom (Figure 6b) and above it for $z + 1 > \sim 4.5$ (Figure 6a). These deviations are of a non-statistical nature and thus the Λ CDM model does not reflect the observed exponential slope.

Table 6

Descriptive statistics for the exponential line (a) in Figure 6a based on z

Valid cases	H_0	R^2	F-test z_{obs}/z_{calc}	$\sum \chi^2$ $(z_{obs}-z_{calc})$
96	0.0002145 $H_0 = 2.14 \cdot 10^{-18} s^{-1}$	0.9991	0.9584	0.0484
Std. deviation $(z_{obs}-z_{calc})$	Variance $(z_{obs}-z_{calc})$	Skew $(z_{obs}-z_{calc})$	Mean $(z_{obs}-z_{calc})$	
0.0434	0.0001	1.098	-0.022	

5 Conclusions

A remarkable result of our analysis of the SN1a supernovae data in the RS range 0.0104–1.3 is that on the basis of the z/μ HD, the good fit of the cosmological parameters to the observed data cannot be seen as a definitive proof in favor of the redshift interpretation according to Equation (1). Figures 2 and 3 show that even in the more sensitive $t/(z + 1)$ representation and even on basis of the most accurate JLA data, the slope of the Λ CDM-HD cannot be distinguished from an exponential slope according to Equation (2), which is characteristic for an energy decay with a constant rate. In the RS range $< \sim 1.3$, the two curves are practically congruent.

A significant deviation of the Λ CDM model from the observed data emerges at higher RSs of $z > \sim 3$. Based on this work, the conclusion can be drawn that in the whole redshift range up to $z \sim 8$, the slope of the HD on the linear scale may be (or is very close to) exponential. Since Hubble’s constant represents probably the most important cosmological parameter, this important issue requires careful clarification and explanation. An exponential slope of the HD would require the most significant proof for the universal expansion, the velocity interpretation of the Hubble law, to be discarded, which most astronomers are not prepared to acknowledge. The presented results are a strong indication for an exponential slope of the HD

in the whole RS range from $z = 0.031$ – 8.1 ; nevertheless, much more precise data, especially in the high RS range, are required to confirm our results. Due to the work of Swift (The Neil Gehrels Swift Observatory, NASA), the fundamental highlights of GRBs have become better known as of late. A large number of blasts have been seen in the range of $0.034 < z \leq 8.1$, which opens up the possibility of estimating the expansion history back to the formation of the first stars. The hope is justified that this can be achieved with future surveys.

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Additional Information

The author declares no conflict of interest

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