Serial and Parallel Iterative Splitting Methods: Algorithms and Applications

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Abstract

The properties of iterative splitting methods with serial versions have been analyzed since recent years, see [1] and [3]. We extend the iterative splitting methods to a class of parallel versions, which allow to reduce the computational time and keep the benefit of the higher accuracy with each iterative step. Parallel splitting methods are nowadays important to solve large problems, which can be splitted in subproblems and computed independently with the different processors. We present a novel parallel iterative splitting method, which is based on the multi-splitting methods, see [2], [10] and [15]. Such a flexibilisation with multisplitting methods allow to decompose large iterative splitting methods and recover the benefit of their underlying waveform-relaxation (WR) methods. We discuss the convergence results of the parallel iterative splitting methods, while we could reformulate such an error to a summation of the individual WR methods. We discuss the numerical convergence of the serial and parallel iterative splitting methods and present different numerical applications to validate the benefit of the parallel versions.

Key words: Multisplitting method; iterative splitting method; numerical analysis; operator-splitting method; initial value problem; iterative solver method; Waveform relaxation method.

AMS subject classifications: 35K45, 35K90, 47D60, 65M06, 65M55.

1. Introduction

Iterative splitting methods are nowadays important solver methods to solve large systems of ordinary, partial or stochastic differential equations, see [1], [3] and [7]. Iterative splitting methods are based on two solver ideas, while in the first part we separate the full operators into different sub-operators and reduce the computational time for such sub-computation, an additional benefit is the iterative part, which allows to solve a relaxation problem, like known in waveform-relaxation method or Picard's iterative method, see [6], [11] [12] and [13]. Both parts reduce the computational time and the complexity as if we solve all parts (full operator and direct method) together, see [5]. Such iterative splitting methods can be used to solve with less computational amount an approximate solution of the ordinary differential equations (ODEs) or semi-discretized partial differential equations (PDEs), see [5] and [13]. The bottleneck of all the iterative methods are given by the large sizes of the iterative matrices, see [10] and [15], therefore we

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considered parallel versions of the iterative splitting method, see [2] and [4]. We concentrate on solving linear evolution equations, such as the differential equation,

$$\partial_t c = Ac = \sum_{l=1}^{L} A_l c, \ c(0) = c_0,$$
 (1)

where $A \in \mathbb{R}^m \times \mathbb{R}^m$ is the full operator with m is the finite dimension of the operator and A_l are suboperators. Further $c \in C^1([0,T];\mathbb{R}^m)$ is the solution and $c_0 \in \mathbb{R}^m$ is the initial condition. Based on the decomposition of the large scale differential equation with operator A into different smaller sub-differential equations with operators A_l , where $l=1,\ldots,L$ and L is the number of processors, we distribute the computational time to many processors and reduce the computational time, see [15]. Further, we modify the synchronous parallel splitting method with chaotic (asynchronous) ideas, such that the computation and communication of the various processors can be done completely independently, see [14].

The outline of this paper is as follows. The serial iterative splitting method is explained in Section 2. The parallel iterative splitting method is introduced in Section 3. We discuss the theoretical results in Section 4. The numerical examples are presented in Section 5. In Section 6, we discuss the theoretical and practical results.

2. Serial iterative splitting method

We consider a two-level iterative splitting method, which is discussed for two operators in [3] and for L-operators in [8].

Based on the differential equation (1), we have the following decomposition of the operator A:

- $A = \sum_{l=1}^{L} A_l$, while A_l are the sub-operators for the iterative part (solver part),
- and $\overline{B_l} = A A_l$ are the sub-operators for the relaxation part (right hand side part). where we have $l = 1, \ldots, L$.

The serial iterative splitting method is given in the following algorithm. Here, we apply a fixed-splitting discretization step-size τ , namely, on the time-interval $[t^n, t^{n+1}]$. We solve the following sub-problems consecutively for $i = 0, L, \ldots, (m-1)L$:

$$\frac{\partial c_{i+1}(t)}{\partial t} = A_1 c_{i+1}(t) + B_1 c_i(t), \text{ with } c_{i+1}(t^n) = c^n,
\frac{\partial c_{i+2}(t)}{\partial t} = A_2 c_{i+2}(t) + B_2 c_{i+1}(t), \text{ with } c_{i+2}(t^n) = c^n,$$
(2)

$$\frac{\partial c_{i+2}(t)}{\partial t} = A_2 c_{i+2}(t) + B_2 c_{i+1}(t), \text{ with } c_{i+2}(t^n) = c^n,$$
(3)

$$\frac{\partial c_{i+L}(t)}{\partial t} = A_L c_{i+L}(t) + B_L c_{i+L-1}(t), \text{ with } c_{i+L}(t^n) = c^n,$$

$$(5)$$

where we assume for the first initialisation $c_0(t) = 0.0$, further c^n is the known split approximation at the time-level $t = t^n$. The split approximation at the time-level $t = t^{n+1}$ is defined as $c^{n+1} =$ $c_{(m-1)L}(t^{n+1})$. The stopping criterion is $||c_{(m-1)L} - c_{(m-2)L}|| \leq err$ and then we have the solution $c(t^{n+1}) = c_{(m-1)L}(t^{n+1}).$

The solutions are given as:

$$c_{i+1}(t) = \exp(A_1 (t - t^n))c^n + \int_{t^n}^t \exp(A_1 (t - s)) B_1 c_i(s) ds,$$
(6)

$$c_{i+2}(t) = \exp(A_2 (t - t^n))c^n + \int_{t^n}^t \exp(A_2 (t - s)) B_2 c_{i+1}(s) ds,$$
(7)

$$c_{i+L}(t) = \exp(A_L (t - t^n))c^n + \int_{t^n}^t \exp(A_L (t - s)) B_L c_{i+L-1}(s) ds,$$
(9)

where $t \in [t^n, t^{n+1}]$.

The integrals can be solved by Trapezoidal- or Simpsons-rule.

We define the error-function as $e_i(t) = c(t) - c_i(t)$ with $e_0(t) = c(t) - c_0$ and the maximum-norm $||e_i|| = \max_{t \in [0,T]} ||e_i(t)||_{\infty}$ and the maximum operator norm $||A_l|| = ||A_l||_{\infty}$.

Theorem 2.1 We have bounded operators $A_l \in \mathbb{R}^m \times \mathbb{R}^m$. Then the iterations (2)-(5), which are applied with i = 0, L, ..., (m-1)L, with L are the number of operators, for the Cauchy-problem (1) is of order $\mathcal{O}(\tau^{mL})$.

Proof 2.2 The proof is done for the 2-level method in the paper [8]. We apply a recursive argument of the iterative scheme with the L operators and obtain:

$$||e_{mL}|| \le \left(\prod_{l=1}^{L} C_l ||B_l||\right)^m ||e_0||,\tag{10}$$

where C_l is given with $C_l = \mathcal{O}(\tau)$, see also [8].

In the following, we concentrate extend the two-level iterative splitting method to a parallel iterative splitting method, while we could transform the two-level method with multisplitting method.

3. Parallel iterative splitting method

In the following, we parallelise the serial iterative splitting approach with the idea of multisplitting-approaches.

We present in the following approaches:

- Multi-splitting iterative approach,
- Two operator iterative splitting approach.

3.1. Multi-splitting iterative approach

The problem is given as $\frac{\partial c}{\partial t} = Ac(t) = \sum_{l=1}^{L} A_l c(t)$, $c = c_0$. The idea is a *multiple* decomposition of

$$A = A_l + B_l, \ l = 1, \dots, L,$$
 (11)

$$B_l = A - A_l, (12)$$

where A_l is a non-singular and B_l is the rest matrix.

Further, we have the decomposition of the parallel computable vectors:

$$c_i = \sum_{l=1}^{L} E_l c_{i,l}, \text{ and } E = \sum_{l=1}^{L} E_l,$$
 (13)

where c_i is the *i*-th iterative solution and $c_{i,l}$ are the parallel computable solutions in the *i*-th iterative step. E is the identity matrix and E_l are diagonal matrices with positive entries.

The multisplitting iterative approach is given as:

$$\frac{\partial c_{i,l}(t)}{\partial t} = A_l c_{i,l}(t) + B_l c_{i-1}(t), \text{ for } t \in [t^n, t^{n+1}]
\text{with } c_{i,l}(t^n) = c(t^n), l = 1, \dots L,$$
(14)

where the initalisation is $c_0(t) = c(t^n)$, we have i = 1, ..., I iterative steps. The stopping criterion is $||c_i - c_{i-1}|| \le err$ and then we have the solution $c(t^{n+1}) = c_i(t^{n+1})$.

The splitting error of the iterative splitting is of k+1 order, i.e. $O(\tau^{k+1})$, with

$$||err_{i+1}|| = K_i \tau_n^i ||err_0|| + O(\tau_n^{i+1}) , \qquad (15)$$

where $K_i = \sum_{l=1}^{L} \frac{1}{\omega_l} ||B_l||$ and $\sum_{l=1}^{L} \frac{1}{\omega_l} = 1$ are the weights $(\omega_l > 1)$ and $||err_0|| = ||c_0(t)||$, while $c_{-1}(t) = 0$.

Benefit:

- Parallel implementation (the method is designed for parallel contributions)

- Good error balance between the different operators Drawback:
- Balances in the decomposition of E_l important to damp large errors

3.2. Parallel splitting with two operators: Classical version

We have to apply the following algorithm, which is applied with synchronisation:

$$\frac{\partial c_{i,1}(t)}{\partial t} = Ac_{i,1}(t) + Bc_{i-1}(t), \text{ with } c_{i,1}(t^n) = c^n, c_0 = 0.0,
\frac{\partial c_{i,2}(t)}{\partial t} = Bc_{i,2}(t) + Ac_{i-1}(t), \text{ with } c_{i,2}(t^n) = c^n, c_0 = 0.0,
c_i(t) = \frac{1}{2} (c_{i,1}(t) + c_{i,2}(t)),$$
(16)

where c^n is the known split approximation at the time-level $t = t^n$. The split approximation at the time-level $t = t^{n+1}$ is defined as $c^{n+1} = c_i(t^{n+1})$. We have the stopping criterion $||c_i - c_{i-1}|| \le err$.

The solutions are given as:

$$c_{i,1}(t) = \exp(A(t-t^n))c^n + \int_{t^n}^t \exp(A(t-s)) B c_{i-1}(s) ds,$$
(17)

$$c_{i,2}(t) = \exp(B(t-t^n))c^n + \int_{t^n}^t \exp(B(t-s)) A c_{i-1}(s) ds,$$
(18)

$$c_i(t) = \frac{1}{2} \left(c_{i,1}(t) + c_{i,2}(t) \right), \tag{19}$$

where $t \in [t^n, t^{n+1}]$.

The integrals can be solved by higher order integration-rules, e.g., with the Trapezoidal- or Simpson's-rule.

3.3. Parallel splitting with two operators: Modern version

We have the following algorithm, which is applied without synchronization:

$$\begin{split} \frac{\partial c_{i,1}(t)}{\partial t} &= A c_{i,1}(t) \ + \ B c_{i-1,1}(t), \text{ with } \ c_{i,1}(t^n) = c^n \ , \ c_{0,1} = 0.0, \\ \frac{\partial c_{i,2}(t)}{\partial t} &= B c_{i,2}(t) \ + \ A c_{i-1,2}(t), \text{ with } \ c_{i,2}(t^n) = c^n, \ c_{0,2} = 0.0, \end{split}$$

where c^n is the known split approximation at the time-level $t = t^n$.

Processor l, l = 1, 2 runs the iterations independently until the stopping criterion $||c_{i,l} - c_{i-1,l}|| \le err$ is reached. Then, the approximations are synchronized

$$c_i(t) = \frac{1}{2} (c_{i_1,1}(t) + c_{i_2,2}(t)),$$

where processor l stops at iteration i_l .

The split approximation at the time-level $t = t^{n+1}$ is defined as $c^{n+1} = c_i(t^{n+1})$.

4. Theoretical Results

In the following, we deal with the m-dimensional initial value problem in the non-homogeneous form, see also the homogeneous form in Equation (1):

$$c'(t) = Ac(t) + f(t), \ x(0) = x_0, \tag{20}$$

where A = M + N and f is the right hand side.

Further, we deal in the following with the proof-ideas related to Waveform relaxation methods, see [9]

The initial value problem (20) is solved with the multisplitting Waveform-relaxation method, which is given as:

$$c'_{i+1}(t) = Mc_{i+1}(t) + Nc_i(t) + f(t), \ c(0) = c_0,$$
(21)

where A is given in Equation (1). Further $c_0(t) = c_0$ is the starting condition.

For the multisplitting approach, we have the following Definition 4.1.

Definition 4.1 We have $L \geq 1$ is the number of the splittings. Further, we have A, A_l, B_l, E_l , which are real-valued $m \times m$ matrices. Such that we obtain the multisplitting triple (A_l, B_l, E_l) for l = 1, ..., L: - $A = A_l + B_l$ and $B_l = \sum_{k=1, k \neq l}^{L} A_k$ with l = 1, ..., L,

- The matrices E_l are nonnegative diagonal matrices and satisfy: $\sum_{l=1}^{L} E_l = I$, where I is the identity matrix.
- $-s_l(i+1) \le i+1$ indicates the iteration, where the l-th component is computed prior to i+1.
- The multisplitting approach based on the Waveform-relaxation in the classical version is given in the following notation:

$$c'_{l,i+1}(t) = A_l c_{l,i+1}(t) + B_l c_i + f(t), \ c_{l,i+1}(0) = c_0,$$
(22)

$$c_{i+1}(t) = \sum_{l=1}^{L} E_l c_{l,i+1}(t). \tag{23}$$

- The multisplitting approach based on the Waveform-relaxation in the modern version is given in the following notation:

$$c'_{s_l(i+1)}(t) = A_l c_{s_l(i+1)}(t) + B_l c_i + f(t), \ c_{s_l(i+1)}(0) = c_0,$$
(24)

$$c_{i+1}(t) = \sum_{l=1}^{L} E_l c_{s_l(i+1)}(t).$$
(25)

4.1. Convergence Analysis

The solution of the multisplitting Waveform-relaxation method (22) and (23), is given as following: We solve the individual equations (22) as: can be written as:

$$c_{l,i+1}(t) = K_l c_i(t) + \phi_l(t),$$
 (26)

where we have

$$K_l c(t) = \int_0^t k_l(t-s) \ c(s) ds, \text{ for } l = 1, \dots, L,$$
 (27)

$$\phi_l(t) = \exp(tA_l)c_0 + \int_0^t \exp((t-s)A_l)f(s) \, ds, \text{ for } l = 1, \dots, L.$$
 (28)

where $k_l(t) = \exp(tA_l) B_l$ for l = 1, ..., L.

Further, we apply the multisplitting notation (23) and obtain the summations:

$$Kc(t) = \sum_{l=1}^{L} E_l K_l c(t), \tag{29}$$

$$\phi(t) = \sum_{l=1}^{L} E_l \phi_l(t), \tag{30}$$

where $k(t) = \sum_{l=1}^{L} E_l k_l(t)$ and we obtain the standard Waveform-relaxation method as:

$$c_{i+1}(t) = Kc_i(t) + \phi(t).$$
 (31)

We assume, that the Lemma 4.2 is fulfilled, see also [9].

Lemma 4.2 We assume that the following items are equivalent:

- We assume c(t) is a solution of the initial value problem (20).
- -c(t) is a solution of each multisplitted equation $c(t) = K_l c(t) + \phi_l(t), c(0) = c_0, l = 1, \ldots, L$.
- c(t) is the solution of the fixpoint equation $c(t) = Kc(t) + \phi_l(t)$.

We define $||c||_T = \max_{t \in [0,T]} |c(t)|$ as maximum norm and we also us $||\cdot||$ as a matrix norm induced by the vector norm $|\cdot|$.

Based on the assumptions, we derive in the following the errors and the convergence results.

In the Theorem 4.3 we derive the error of the i-th approximation, see also [9].

Theorem 4.3 There exists a constant $C := \sum_{l=1}^{L} C_l$, which is given to estimate the kernel k of the multisplitting waveform-relaxation operator, such that we obtain $||k||_T = C$. Then the error of the i-th approximation of the classical multisplitting WR method (22)-(23) is given by

$$||c^{i} - c||_{T} \le \frac{(CT)^{i}}{i!} (\exp(CT)||\phi||_{T} + ||c_{0}||_{T}).$$
(32)

Proof 4.4 We have given

$$c_i(t) = Kc_{i-1}(t) + \phi(t),$$
 (33)

We apply the Lemma 4.2 and follow with the iterative approach:

$$c_i(t) = K^i c_0(t) + \sum_{i=0}^{i-1} K^i \phi(t), \tag{34}$$

where is $K^{i}u(t)$ is the i-times convolution

$$K^{i}u(t) = \int_{0}^{t} k(t - s_{i}) \left(\int_{0}^{s_{i}} k(t - s_{i} - s_{i-1}) \left(\dots \int_{0}^{t - \sum_{j=0}^{i} s_{i-j}} k(t - \sum_{j=0}^{i} s_{i-j}) u(s_{1}) ds_{1} \dots \right) ds_{i-1} \right) ds_{i}.$$

Further we have $||u||_T = \max_{t \in [0,t]} |u(t)|$, where $|\cdot|$ is an appropriate Banach-Norm. It exists:

$$||k_l||_T \le C_l, \text{ for } l = 1, \dots, L,$$
 (35)

and we have

$$|k(t)| \le ||k||_T = ||\sum_{l=1}^L E_l k_l|| \le \sum_{l=1}^L C_l = C,$$
(36)

we apply the estimation of the Waveform-relaxation, see [11], and obtain:

$$||K^i||_T \le \frac{(CT)^i}{i!},\tag{37}$$

where we have $\lim_{i\to\infty} ||K^i||_T \to 0$.

The error estimate is than given as

$$||c - c^i||_T = ||(\lim_{j \to \infty} K^j c_0 - \sum_{j=0}^{\infty} K^j \phi(t)) - (K^i c_0 - \sum_{j=0}^{i-1} K^j \phi(t))||$$

$$\leq ||\sum_{j=i}^{\infty} K^{j} \phi(t) + K^{i} c_{0}|| \leq ||K^{i}||_{T} (||\sum_{j=0}^{\infty} K^{j}||\phi(t)||_{T} + ||c_{0}||_{T}). \tag{38}$$

We apply

$$\sum_{j=0}^{\infty} K^j = K^0 + K^1 + \dots + K^{\infty} \le \exp(CT).$$
(39)

Then we obtain the estimation

$$||c - c^{i}||_{T} = \leq \frac{(CT)^{i}}{i!} (\exp(CT)||\phi||_{T} + ||c_{0}||_{T}).$$
(40)

We have the following new convergence Theorem 4.5 based on the extension of the classical convergence Theorem version 4.3.

Theorem 4.5 There exists a constant $C := \sum_{l=1}^{L} C_l$, which is given to estimate the kernel k of the multisplitting waveform-relaxation operator, such that we obtain $||k||_T = C$. Then the convergence of the modern multisplitting WR method (24)-(25) is given by

$$||c^{i_{\min}} - c||_T \le \frac{(C T)^{i_{\min}}}{i_{\min}!} ||c^0 - c||_T,$$
 (41)

where $i_{\min} = \min_{l=1}^{L} s_l(i)$, where $s_l(i) \leq i$ are the retarded iterations of the l-th processor.

Proof 4.6 We start with the estimation of the i-th iteration:

$$||c^{i_{min}} - c||_T \le ||\sum_{l=1} E_l K_l(c^{s_l(i)} - c)||_T \le$$
(42)

$$\leq ||\sum_{l=1}^{\infty} E_l K_l(c^{i_{min}-1} - c)||_T \leq ||K||_T ||(c^{i_{min}-1} - c)||_T.$$

$$(43)$$

Then, we have the recursion:

$$||c^{i_{min}} - c||_T \le ||K^{i_{min}}||_T ||(c^0 - c)||_T. \tag{44}$$

where we apply $||K^{i_{min}}||_T \leq \frac{(C\ T)^{i_{min}}}{i_{min}!}$, based on the idea in [11], and we obtain:

$$||c^{i_{min}} - c||_T \le \frac{(C T)^{i_{min}}}{i_{min}!} ||(c^0 - c)||_T, \tag{45}$$

where $i_{\min} = \min_{l=1}^{L} s_l(i)$.

Remark 4.7 For the parallel error, we have the order $\mathcal{O}(\tau^m)$, if we assume that all processors have at least m iterative cycles, while for the serial error, we have the order $\mathcal{O}(\tau^{mL})$. Means in the serial version, we have to apply mL iterative steps in sum to obtain the result, while in the parallel version, we only apply m iterative steps, while L processors share the computation to solve the L sub-equations. Further, we can assume, that the sub-equations are faster to solve, while the sub-operators are much more smaller and simpler to handle. Such that we have $t_{sub} \leq \frac{t_{full}}{L}$, while t_{sub} is the time to solve a sub-problem and t_{full} the time to solve the full problem. Therefore, we have a benefit in the parallel distribution and we obtain faster the higher order $\mathcal{O}(\tau^{mL})$ as with the serial version.

Remark 4.8 We can modify the parallel method with inner and outer iterative cycles. Means we deal with L inner cycles for each processor and apply m outer cycles for all the processors. For such methods, we also obtain the higher order $\mathcal{O}(\tau^{mL})$ but more faster than in the serial version.

5. Numerical examples

In the following, we deal with different numerical example to verify and test the theoretical results. We deal with:

- Only time-dependent problem: We apply ordinary differential equation to verify the theoretical results.
- Linear time and spatial dependent problem: We apply a diffusion equation with different spatial dependent operators and test the application to partial differential equations.
- Nonlinear time and spatial dependent problem: We apply a mixed diffusion-convection with Burgers
 equation to test and verify the application to nonlinear problems.

5.1. First Example: Matrix problem

For a first test example, consider the matrix equation,

$$u'(t) = \begin{bmatrix} 1 & 2 \\ 0 & 1 \end{bmatrix} u, \quad u(0) = u_0 = \begin{pmatrix} 1 \\ 1 \end{pmatrix},$$
 (46)

the exact solution is

$$\exp\left(\begin{bmatrix} 1 & 2 \\ 0 & 1 \end{bmatrix} t\right) = \begin{pmatrix} \exp(t) & 2t & \exp(t) \\ 0 & \exp(t) \end{pmatrix}. \tag{47}$$

We split the matrix as,

- Two operator approach

$$A + B = \begin{bmatrix} 1 & 2 \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} 0.3 & 1 \\ 0 & 0.3 \end{bmatrix} + \begin{bmatrix} 0.7 & 1 \\ 0 & 0.7 \end{bmatrix}$$
(48)

- Multiple operator approach

$$A_1 + B_1 = \begin{bmatrix} 1 & 2 \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} 0.1 & 1.0 \\ 0 & 0.1 \end{bmatrix} + \begin{bmatrix} 0.9 & 1.0 \\ 0 & 0.9 \end{bmatrix}$$

$$(49)$$

$$A_2 + B_2 = \begin{bmatrix} 1 & 2 \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} 0.5 & 0.1 \\ 0 & 0.5 \end{bmatrix} + \begin{bmatrix} 0.5 & 1.9 \\ 0 & 0.5 \end{bmatrix}$$
 (50)

where the E_1 and E_2 are given as:

$$E_1 = \begin{bmatrix} 0.9 & 0 \\ 0 & 0.1 \end{bmatrix}, E_2 = \begin{bmatrix} 0.1 & 0 \\ 0 & 0.9 \end{bmatrix}.$$
 (51)

We include Tables 1 and 2 corresponding to Multi-splitting iterative approach classical and modern version with the above partitions and using different discretizations in [0, 1] of step h allowing a maximum of 10 iterations and a tolerance of 10^{-3} . We can see in the results the relative and absolute errors for each component of the solution and the mean of iterations performed in order to reach the tolerance.

h	$\max e_1$	$\max e_2$	$\operatorname{rel}e_1$	rel e_2	it 1	it 2
0.1	0.008308	0.00040332	0.0010188	0.00014837	4	4
0.05	0.00043184	2.906e-05	5.2955e- 05	1.069 e - 05	3	3
0.025	0.00010623	7.2291e-06	1.3027e-05	2.6595e-06	3	3
0.0125	2.6335e-05	1.8026e-06	3.2294e-06	6.6316e-07	3	3
0.00625	6.5553e-06	4.5007e-07	8.0385e-07	1.6557e-07	3	3

Table 1 Multisplitting classic version

Remark 5.1 We applied the multisplitting method with the classical (synchronous) and modern (chaotic) approach. We receive the same accuracy of the numerical results, means the methods are as same accurate. We obtain some more benefits of the modern method, if we apply large time-steps, such that the solution of one sub-problem can be achieved faster and benefit the solution of the second sub-problem. For such small unbalances computations the modern approach is more efficient.

h	$\max e_1$	$\max e_2$	rel e_1	rel e_2	it 1	it 2
0.1	0.0081586	0.0003996	0.0010005	0.000147	4	4
0.05	0.00030027	3.1292 e- 05	0.00012447	1.1512 e-05	3.35	3
0.025	0.00027348	1.3123 e-05	3.3536e-05	4.8277 e - 06	3	3
0.0125	6.8274 e-05	3.281e-06	8.3723e-06	1.207 e-06	3	3
0.00625	1.7056e-05	8.2026e-07	2.0915e-06	3.0176e-07	3	3

Table 2 Multisplitting Modern version

5.2. Second Example: Diffusion problem

We deal with the following diffusion problem:

$$u'(\mathbf{x},t) = \Delta u(\mathbf{x},t), \ (\mathbf{x},t) \in \partial \Omega \times [0,T], \tag{52}$$

$$u(\mathbf{x},0) = \sin x \, \sin y \, \sin z, \, \mathbf{x} \in \Omega, \tag{53}$$

$$u(\mathbf{x},t) = 0, \ (\mathbf{x},t) \in \partial\Omega \times [0,T],$$
 (54)

where we have the analytical solution $u_{an}(\mathbf{x},t) = \exp(-3t) \sin x \sin y \sin z$, with $\mathbf{x} = (x,y,z)^t$ and $\Omega = [-\pi, \pi] \times [-\pi, \pi] \times [-\pi, \pi].$

In operator notation, we write as following:

$$A = A_1 + A_2 + A_3, (55)$$

where $A_1 = \frac{\partial^2}{\partial x^2}$, $A_2 = \frac{\partial^2}{\partial y^2}$, $A_3 u = \frac{\partial^2}{\partial z^2}$ and we assume, that the zero-boundary conditions (Dirichlet boundary conditions) are embedded.

The problem is discretized by using a 4-D mesh in $\Omega \times [0,T]$. Denote by $u_{i,j,k,t}$ the approximated value of the solution at node (x_i, y_i, z_k, t) for a given t. For the time-integration, we apply the integral formulation, see Equations (6)-(9).

For the spatial discretization, we test a second and fourth order scheme:

$$\frac{\partial^2}{\partial x^2} u_{i,j,k,t} = \frac{u_{i+1,j,k,t} - 2u_{i,j,k,t} + u_{i-1,j,k,t}}{\Delta x^2}, \text{ second order approach},$$

$$\frac{\partial^2}{\partial x^2} u_{i,j,k,t} = \frac{-u_{i+2,j,k,t} + 16u_{i+1,j,k,t} - 30u_{i,j,k,t} + 16u_{i-1,j,k,t} - u_{i-2,j,k,t}}{12\Delta x^2}, \text{ 4th-order approach},$$
(56)

$$\frac{\partial^2}{\partial x^2} u_{i,j,k,t} = \frac{-u_{i+2,j,k,t} + 16u_{i+1,j,k,t} - 30u_{i,j,k,t} + 16u_{i-1,j,k,t} - u_{i-2,j,k,t}}{12\Delta x^2}, \text{ 4th-order approach}, (57)$$

where we have the analogous operators for the y and z derivations.

In order to establish the convergence of the algorithms, we compute the solution $u(\Delta x, h)$ obtained using spatial and temporal steps $\Delta x = \Delta y = \Delta z$ and h, respectively. We use different measures to estimate the convergence. On one hand, we can compare the outcome of the method $u(\Delta x, h)$ with the exact solution u_{ana} for every point of the mesh, which shows the convergence of the method. On the other hand, we can compare $u(\Delta x, h)$ with the result obtained halving the time or space steps, h/2 or $\Delta x/2$, at the points shared by the corresponding meshes. This allows to analyze how the results depend on these steps. Denote by $e_{i,j,k}(\Delta x, h)$ the difference between the results at a mesh point (x_i, y_i, z_k, T) , obtained using two different time or space steps, and by $\delta_{i,j,k}(\Delta x,h)$ the difference with the analytical solution at the same point. In the tables, we will denote the maximum errors by

$$e_{\max} = \max_{i,j,k} |e_{i,j,k}(\Delta x, h)|,$$
 (58)

and

$$\delta_{\max} = \max_{i,j,k} |\delta_{i,j,k}(\Delta x, h)|, \tag{59}$$

and the mean errors by

$$e_{\text{mean}} = \frac{1}{N} \sum_{i,j,k} |e_{i,j,k}(\Delta x, h)|,$$
 (60)

and

$$\delta_{\text{mean}} = \frac{1}{N} \sum_{i,j,k} |\delta_{i,j,k}(\Delta x, h)|, \tag{61}$$

where N is the number of spatial nodes at time T.

In the following, we discuss different decompositions of the multi-operator splitting approach:

- Directional decomposition: We decompose into the different directions:

$$A_1 = \frac{\partial^2}{\partial x^2}, \ A_2 = \frac{\partial^2}{\partial y^2}, \ A_3 = \frac{\partial^2}{\partial z^2}. \tag{62}$$

Here, we have the benefit of decomposing the different directions.

The drawback is related to the unbalanced decomposition, while the matrices have different sparse entries. Therefore the exponential matrices of the operators are different in their sparse behaviour and the error can not be optimal reduced.

We can reduce the unbalanced problem, if we deal with the idea to use $\Delta t \approx \Delta x$, see [8]. Then, we obtain at least a second order scheme (related to the spatial discretization).

At least, we test also the strong relation of the time- and spatial scales with respect to the explicit CFL condition, means, we have:

$$\Delta t \approx \frac{\Delta x^2}{6},\tag{63}$$

where $\Delta x = \Delta y = \Delta z$.

We apply the splitting algorithms with the directional decomposition in [0, 1]. The splitting is iterated until a tolerance of 10^{-8} or a maximum of 10 iterations are reached. The values shown in the tables correspond to maximum or mean values at T=1. Tables 3 and 4 present the results obtained using different number of temporal steps and the second and fourth order schemes for the spatial discretization, respectively. In tables 5 and 6, we fix the number of temporal steps to 1280 and vary the number of spatial intervals.

Splitting algorithm	Time steps	$e_{ m mean}$	$e_{ m max}$	$\delta_{ m mean}$	$\delta_{ m max}$	Average iterations
	10	3.0021e-05	8.8220e-03	1.3021e-04	2.0943e-02	4.0000
Serial Iterative	20	1.6139e-05	4.7546e-03	1.1744e-04	2.0304 e-02	3.0000
Serial Relative	40	7.8770e-06	2.2831e-03	1.1352e-04	2.0135e-02	3.0000
	80	3.8391e-06	1.1013e-03	1.1245e-04	2.0059e-02	2.0688
	10	7.5244e-06	1.1242e-03	1.1498e-04	2.0319e-02	9.9500
Classical Parallel	20	4.1720e-06	6.1563e- 04	1.2189e-04	2.0942e-02	7.0000
Classical I aranei	40	2.1715 e-06	3.1759 e- 04	1.2579 e-04	2.1295 e-02	5.0750
	80	1.1042e-06	1.6070 e-04	1.2784 e-04	2.1479e-02	4.8688
	10	6.9978e-06	8.6818e-04	1.1445e-04	1.9636e-02	9.9000
Modern Parallel	20	4.2214 e - 06	5.4241e-04	1.2128e-04	2.0487e-02	7.9750
Modern Faranei	40	2.3180e-06	3.0696e-04	1.2536e-04	2.1029e-02	6.0000
	80	1.2147e-06	1.6316e-04	1.2758e-04	2.1336e-02	4.9938

Table 3

Results using the second order scheme for the spatial discretization, with 10 spatial subintervals and different number of temporal steps

Table 3 shows that, using the second order scheme to approximate the spatial derivatives, the order of convergence of all the considered methods with respect to the temporal step h is 1, because the estimated errors are approximately proportional to h. The differences with respect to the analytical solution are almost constant, due to the spatial discretization error error. The serial iterative splitting method is unstable for h = 0.1, whereas the parallel splitting methods converge for all the cases shown in the table. When the methods converge, there is no difference in the final approximation among the considered methods, but the parallel methods require about double number of iterations to converge.

Splitting algorithm	Time steps	$e_{ m mean}$	e_{max}	$\delta_{ m mean}$	$\delta_{ m max}$	Average iterations
	10	8.7714e-02	1.1920e+03	8.7714e-02	1.1920e+03	3.0500
Serial Iterative	20	8.1984e-09	5.3739e-05	4.6038e-08	3.1645 e-04	3.0000
	40	2.0483e-09	1.3425 e-05	3.7888e-08	2.7126e-04	3.0000
	80	5.1200e-10	3.3557e-06	3.5951 e- 08	2.6299e-04	2.0063
	10	3.2878e-08	2.1529e-04	7.8916e-08	5.3159e-04	7.1000
Classical Parallel	20	8.1984e-09	5.3681 e-05	4.6038e-08	3.1644e-04	6.0250
Ciassicai i arailei	40	2.0483e-09	1.3412e-05	3.7888e-08	2.7125e-04	5.0125
	80	5.1204e-10	3.3528e-06	3.5951 e- 08	2.6299e-04	4.0062
	10	3.2881e-08	2.1529e-04	7.8918e-08	5.3157e-04	7.5500
Modern Parallel	20	8.1984e-09	5.3684 e-05	4.6038e-08	3.1644e-04	6.0250
	40	2.0483e-09	1.3412e-05	3.7888e-08	2.7125e-04	5.0125
	80	5.1204e-10	3.3528e-06	3.5951 e- 08	2.6299e-04	4.0062

Table 4
Results using the fourth order scheme for the spatial discretization, with 10 subintervals in each dimension and different number of temporal steps

Table 4 shows that, using the fourth order scheme, the order of convergence of with respect to the temporal step h is 2, being the estimated errors proportional to h^2 .

Splitting algorithm	Space intervals	$e_{ m mean}$	e_{\max}	$\delta_{ m mean}$	$\delta_{ m max}$	Average iterations
	10	6.0092 e- 05	1.0185 e-02	1.1203e-04	2.0006e-02	2.0031
Serial Iterative	20	2.7133e-05	4.9543e- 03	5.3066e-05	9.8264 e-03	2.0219
	40	_	_	2.6073e- 05	4.8721 e- 03	3.0000
	10	5.9473e-05	9.5024e-03	1.2940e-04	2.1621e-02	4.0031
Classical parallel	20	3.4999e-05	5.9655e-03	7.0607e-05	1.2119e-02	4.2938
	40	_	_	3.5660e-05	6.1534 e-03	6.0219
	10	5.9809e-05	9.5336e-03	1.2933e-04	2.1583e-02	4.0031
Modern Parallel	20	3.5212e-05	6.0031e-03	7.0192e-05	1.2050e-02	5.9781
	40	_	_	3.5029e-05	6.0466e-03	9.9812

Table 5
Results using the second order scheme for the spatial discretization, with 320 temporal steps and different number of spatial subintervals

In Table 5, the spatial step is reduced simultaneously in the three dimensions, in order to keep the increments equal, because the performance of the method is better in this case. The number of time steps is big enough to ensure the convergence in the case of smaller spatial subinterval and is used in all the cases to allow the comparison depending only on the spatial step. The error is halved when the step in each spatial dimension is halved, which increases the computational cost by a factor of 8, thus the behavior of the method is poor. The parallel methods obtain slightly less approximation to the analytical result and require more iterations than the serial iterative method. The modern parallel method needs more iterations per step to converge than the classical parallel method.

Table 6 compares the results obtained varying the number of spatial subintervals. The error is divided by about 20 when the step in each spatial dimension is halved, which indicates a convergence of order higher than linear with respect to the cube of the spatial step, improving the behavior of the method with respect to the order two scheme. There is no noticeable difference in the results among the different methods, although the parallel ones require double number of iterations than the sequential method to converge.

- Balanced decomposition: We decompose into the different directions:

Splitting algorithm	Space intervals	$e_{ m mean}$	e_{\max}	$\delta_{ m mean}$	$\delta_{ m max}$	Average iterations
1	10	3.3808e-08	2.4779e-04	3.5363e-08	2.6041e-04	2.0031
1	20	1.5518e-09	1.2258 e-05	1.5936e-09	1.2625 e - 05	2.0031
1	40	_	_	4.5839e-11	3.6643e-07	2.0031
2	10	3.3808e-08	2.4779e-04	3.5363e-08	2.6041e-04	4.0031
2	20	1.5518e-09	1.2258 e-05	1.5936e-09	1.2625 e - 05	4.0031
2	40	_	_	4.5838e-11	3.6642 e-07	4.0031
3	10	3.3808e-08	2.4779e-04	3.5363e-08	2.6041e-04	4.0031
3	20	1.5518e-09	1.2258e-05	1.5936e-09	1.2625 e - 05	4.0031
3	40	_	_	4.5838e-11	3.6642e-07	4.0031

Table 6

Results using the fourth order scheme for the spatial discretization, with 320 temporal steps and different number of subintervals in each dimension

$$A_1 = \frac{1}{3}A, \ A_2 = \frac{1}{3}A, \ A_3 = \frac{1}{3}A.$$
 (64)

Here, we have the benefit of equal load balances of the matrices, such that the exp-matrices have the same sparse structure.

Splitting algorithm	Time steps	$e_{ m mean}$	e_{max}	$\delta_{ m mean}$	$\delta_{ m max}$	Average iterations
	10	3.2825e-08	2.1449e-04	6.2681 e-07	4.0959 e-03	3.0500
Serial Iterative	20	8.1862e-09	5.3493e- 05	6.5964 e-07	4.3104e-03	3.0000
	40	2.0453e-09	1.3365e-05	6.6783e- 07	4.3639e-03	3.0000
	80	5.1125e-10	3.3408e-06	6.6987e-07	4.3773e- 03	2.0063
	10	3.2825e-08	2.1449e-04	6.2681e-07	4.0959e-03	7.0500
Classical Parallel	20	8.1862e-09	5.3493e- 05	6.5964 e-07	4.3104e-03	6.0000
Classical Farallel	40	2.0453e-09	1.3365e-05	6.6783e- 07	4.3639e-03	5.0125
	80	5.1128e-10	3.3410e-06	6.6987e-07	4.3773e-03	4.0062
	10	3.2825e-08	2.1449e-04	6.2681e-07	4.0959e-03	7.0500
Modern Parallel	20	8.1862e-09	5.3493e- 05	6.5964 e-07	4.3104e-03	6.0000
Modern Parallel	40	2.0453e-09	1.3365e-05	6.6783e-07	4.3639e-03	5.0125
	80	5.1128e-10	3.3410e-06	6.6987e-07	4.3773e-03	4.0062

Table 7

Results using the second order scheme for the spatial discretization, with 10 spatial subintervals and different number of temporal steps

- Mixed decomposition: We decompose into the different directions:

$$A_1 = (1 - \epsilon) \frac{\partial^2}{\partial x^2} + \frac{\epsilon}{2} (\frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}), \ A_2 = (1 - \epsilon) \frac{\partial^2}{\partial y^2} + \frac{\epsilon}{2} (\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2}), \tag{65}$$

$$A_3 = (1 - \epsilon) \frac{\partial^2}{\partial z^2} + \frac{\epsilon}{2} (\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}), \tag{66}$$

where $\epsilon = [0, 2/3]$, means for $\epsilon = 0$, we have the directional decomposition, while for $\epsilon = 2/3$ we have the balanced decomposition.

Here, we have the benefit, e.g., for small $\epsilon = 0.1$, we have nearly a directional decomposition scheme. Based on the $\epsilon \neq 0$, we stabilise the scheme.

Remark 5.2 Based on the balanced decomposition with $\epsilon = 2/3$, we do not have problems with the splitting approaches and obtain optimal results. For the pure unbalanced decomposition, means $\epsilon = 0$, we decompose into different directions. Here, we have to restrict us to the exact second order approach, which is $\Delta t \approx \Delta x$.

Remark 5.3 We obtain the benefit of the classical and modern parallel iterative splitting method based on larger time-steps and more iterative steps. In such an optimal version, we are much faster than the serial version and also the result is more accurate. For very fine time-steps, we do not see an improvement in the accuracy, but we see a benefit in the computational time, means the parallel versions are more faster.

5.3. Third Example: Mixed convection-diffusion and Burgers equation

We deal with a partial differential equation, which is a 2D example of a mixed convection-diffusion and Burgers equation. For such equations, we can find analytical solution. The model problem is given as:

$$\partial_t u = -1/2u\partial_x u - 1/2u\partial_y u - 1/2\partial_x u - 1/2\partial_y u$$

+ $\mu(\partial_{xx} u + \partial_{yy} u) + f(x, y, t), (x, y, t) \in \Omega \times [0, T],$ (67)

$$u(x, y, 0) = u_{\text{ana}}(x, y, 0), (x, y) \in \Omega,$$
 (68)

$$u(x, y, t) = u_{\text{ana}}(x, y, t) \text{ on } \partial\Omega \times [0, T],$$

$$(69)$$

where $\Omega = [0, 1] \times [0, 1]$, T = 1.25, and μ is the viscosity.

The analytical solution is given as

$$u_{\rm ana}(x,y,t) = (1 + \exp(\frac{x+y-t}{2\mu}))^{-1} + \exp(\frac{x+y-t}{2\mu}),\tag{70}$$

where we compute f(x, y, t) accordingly.

As in the previous example, denote by $u(\Delta, h)$ the numerical solution obtained using spatial subintervals of amplitude $\Delta = \Delta x = \Delta y$, time steps h and a tolerance of $tol = 10^{-6}$, allowing a maximum of 40 iterations. On one hand, we will compare the numerical solution with the exact one u_{ana} for every point of the mesh, which shows the convergence of the method. On the other hand, we will compare $u(\Delta, h)$ with the result obtained halving the time steps, h/2, at the points shared by the corresponding meshes. Denote by $e_{i,j,k}(\Delta, h)$ the difference between the results obtained using two different time steps, h and h/2, at a common mesh point (x_i, y_j, t_k) , and by $\delta_{i,j,k}(\Delta, h)$ the difference with the analytical solution at the same point. In the tables, we will use the error estimates given by

$$e_{\max} = \max_{i,j,k} |e_{i,j,k}(\Delta, h)|, \tag{71}$$

and

$$\delta_{\max} = \max_{i,j,k} |\delta_{i,j,k}(\Delta, h)|, \tag{72}$$

and the mean errors by

$$e_{\text{mean}} = \frac{1}{N} \sum_{i,j,k} |e_{i,j,k}(\Delta, h)|, \tag{73}$$

and

$$\delta_{\text{mean}} = \frac{1}{N} \sum_{i,j,k} |\delta_{i,j,k}(\Delta, h)|, \tag{74}$$

where N is the total number of nodes (x_i, y_j, t_k) .

In order to obtain the temporal cost of the parallel schemes, we measure the time consumed by processor l in each iteration m, in $[t_k, t_{k+1}]$, $tp_{k,l,m}$. In the classical parallel scheme the processors synchronize at each iteration, so the cost for this time interval is $tp_k = \sum_{m} \max_{l=1,2} tp_{k,l,m}$, whereas, in the modern parallel scheme,

the processors iterate independently in $[t_k, t_{k+1}]$ performing m_l , l=1, 2 iterations until convergence, thus, the cost is $tp_k = \max_{l=1,2} \sum_m tp_{k,l,m}$. The final cost is obtained adding the results of all time subintervals.

The values have been obtained using etime, running Matlab 2015a on a desktop computer with an Intel Core i7 processor in a 64 bits operating system with Windows 7 Professional.

In the following, we discuss different decompositions of the multi-operator splitting approach:

- Directional decomposition: We decompose into the different directions (x and y direction):

$$A(u)u = -1/2u\partial_x u - 1/2\partial_x u + \mu\partial_{xx}u, \tag{75}$$

$$B(u)u = -1/2u\partial_y u - 1/2\partial_y u + \mu\partial_{yy} u + f(x, y, t).$$
(76)

Before studying the convergence of the method, we analyze the influence of parameter μ . Tables 9 and 8 compare the results obtained for different values of μ using the second and the fourth order discretization scheme, respectively, taking 10 subintervals in each spatial dimension and 640 time steps. As we can observe, the error generally increases with μ . Along this example, we will take $\mu = 0.5$, because for this value, the methods require the minimum number of iterations, except in one case.

Splitting algorithm	μ	e_{max}	e_{mean}	δ_{max}	δ_{mean}	Average iterations	Time
	0.25	4.3464e-02	5.3355e-04	2.2106e-01	8.4732e-03	2.9719	6.1689
Generalized Iterative	0.50	5.1944e-03	5.5729 e-04	6.0042 e-03	9.4742e-04	2.5703	5.4561
	1.00	1.3508e-02	2.0793e- 03	1.7504 e-02	2.7070e-03	2.7891	5.9197
	2.00	4.9564 e - 02	8.9366e-03	6.2889 e-02	1.1465 e-02	3.0187	6.5253
	0.25	1.6915e-02	3.7593e-04	2.2038e-01	8.4779e-03	4.1391	4.4844
Classical Parallel	0.50	4.0117e-03	4.0302e-04	4.6768 e - 03	7.1258e-04	4.0047	4.8457
Classical I arallel	1.00	8.2553e-03	1.2533e-03	1.1192e-02	1.7195e-03	4.0094	4.3426
	2.00	2.1603e-02	3.9868e-03	2.9915e-02	5.5900e-03	5.0422	5.5451
	0.25	1.7583e-02	3.7876e-04	2.1864e-01	8.4754e-03	4.1578	4.4456
Modern Parallel	0.50	4.2148e-03	4.3094e-04	4.9121e-03	7.4444e-04	4.0062	4.3670
Modern Parallel	1.00	9.0609e-03	1.4340e-03	1.2127e-02	1.9278e-03	8.0031	8.6506
	2.00	2.5669e-02	5.1563e-03	3.4678e-02	6.9560e-03	30.0156	33.3511

Table 8 Results of the directional decomposition using the second order scheme for the spatial discretization, with 10 spatial subintervals, 640 temporal intervals and different values of μ .

Splitting algorithm	μ	e_{max}	e_{mean}	δ_{max}	δ_{mean}	Average iterations	Time
	0.25	4.9976e-02	5.1652e-04	6.8988e-02	8.8945e-04	2.9437	6.5982
Serial Iterative	0.50	6.4457 e-03	6.0827 e-04	8.4526e-03	8.0368e-04	2.3594	5.2817
Seliai Itelative	1.00	1.3637e-02	2.1752e-03	1.7452 e- 02	2.8197e-03	2.4156	5.5056
	2.00	5.1088e-02	9.2044e-03	6.3508 e- 02	1.1818e-02	3.0922	6.5146
	0.25	4.9306e-02	4.3517e-04	6.8056e-02	7.9501e-04	4.2078	4.7986
Classical Parallel	0.50	4.8864e-03	4.0313e-04	6.6107e-03	5.4922 e-04	4.0047	4.5927
Classical Faranei	1.00	7.4319e-03	1.1897e-03	1.0073e- 02	1.6334e-03	4.0109	4.6604
	2.00	1.8283e-02	3.6989e-03	2.5167e-02	5.1564 e-03	6.1125	6.5804
	0.25	4.8756e-02	4.3735e-04	6.7317e-02	7.9550e-04	4.1141	4.7061
Modern Parallel	0.50	5.1847e-03	4.3964e-04	6.9543e- 03	5.9040e-04	5.0047	5.7085
Modern Paranei	1.00	8.1250e-03	1.4075e-03	1.0876e-02	1.8863e-03	10.0078	11.5772
	2.00	2.5335e-02	5.5574 e-03	3.3627e-02	7.3088e-03	40.0000	42.8804

Table 9 Results of the directional decomposition using the fourth order scheme for the spatial discretization, with 10 spatial subintervals, 640 temporal intervals and different values of μ .

As we can observe, the mean error is approximately proportional to the inverse of the viscosity, μ . Instead, the number of iterations and the execution time is not affected by the viscosity changes, except for the last case of the modern parallel algorithm, where the maximum number of iterations is reached in each step, consuming also more execution time. In what follows, we will take the viscosity parameter $\mu = 0.5$.

We will first analyze the influence on the convergence of the number of time steps, using second or fourth order approximations for the spatial derivatives.

Table 10 show that the estimated errors are roughly proportional to the square of the time step, although, in the end, the differences with the analytical solution δ decrease more slowly, due to the discretization error. The classical and modern parallel methods converge more slowly in the case nt=160, but, in the other cases, they have better performance in terms of error estimates and temporal cost. The utilization of fourth order approximations instead of second order approximations for the discretization spatial derivatives also gives better results for more than 160 time intervals, at a slightly higher cost.

Splitting algorithm	Time steps	e_{max}	e_{mean}	δ_{max}	δ_{mean}	Average iterations	Time
	160	1.5166e-01	1.2191e-02	1.8951e-01	1.5624e-02	5.9875	3.1916
Serial iterative	320	2.9397e-02	2.6622 e-03	3.7849e-02	3.4590 e - 03	3.0094	3.5150
	640	6.4457 e - 03	6.0827 e-04	8.4526e-03	8.0368e-04	2.3594	5.4186
	1280	1.5092 e-03	1.4657e-04	2.0076e-03	1.9688e-04	2.0008	9.1334
	160	6.1519e-02	4.8300e-03	8.6168e-02	6.7805e-03	39.8937	10.8616
Classical Parallel	320	1.8052 e-02	1.4475 e-03	2.4663e-02	1.9953e-03	5.0188	2.9691
Classical Laranei	640	4.8864 e - 03	4.0313e-04	6.6107e-03	5.4922 e-04	4.0047	4.6626
	1280	1.2736e-03	1.0740e-04	1.7251 e-03	1.4737e-04	3.0023	6.9907
	160	1.1044e-01	7.5050e-03	1.3589e-01	9.7029e-03	40.0000	10.8955
Modern Parallel	320	2.0231e-02	1.7146e-03	2.7185e-02	2.3017e-03	10.0156	5.9235
Modern Parallel	640	5.1847e-03	4.3964e-04	6.9543e- 03	5.9040e-04	5.0047	5.8144
	1280	1.3121e-03	1.1194e-04	1.7697e-03	1.5233e-04	3.0023	6.9861

Table $\overline{10}$

Results of the directional decomposition using the fourth order scheme for the spatial discretization, with 10 spatial subintervals and different number of temporal steps

Now we analyze the dependence on the number of spatial subintervals. Tables 11 and 12 display the errors obtained varying the number of space subintervals, considering 1280 time steps, $\mu = 0.5$, tolerance 10^{-6} and maximum number of iterations 40.

Splitting algorithm	Space steps	e_{max}	e_{mean}	δ_{max}	δ_{mean}	Average iterations	Time(s)
Serial Iterative	5	1.5415e-04	1.5346e-05	6.7133e-03	1.0434e-03	2.0008	3.9673
	10	1.2967e-03	1.3550 e- 04	2.2505 e-03	4.2699 e-04	2.0008	8.2839
	20	3.0299e-02	1.6372 e-03	3.9211e-02	2.1837e-03	2.4992	24.8612
	5	1.5192e-04	9.4243e-06	6.7124 e-03	1.0434e-03	3.0008	2.9936
Classical Parallel	10	1.0957e-03	1.0800e-04	2.2581 e-03	3.5285 e-04	3.0016	6.3319
	20	1.8765 e-02	1.0091e-03	2.5443e-02	1.4653 e-03	4.0047	20.5138
	5	1.5292e-04	9.4575e-06	6.7112e-03	1.0434e-03	3.0008	2.9776
Modern Parallel	10	1.1220e-03	1.1151e-04	2.2857e-03	3.5492 e- 04	3.0023	6.3142
	20	2.0590e-02	1.1583e-03	2.7563e-02	1.6374 e-03	7.0039	35.8080

Table $\overline{11}$

Results of the directional decomposition using the second order scheme for the spatial discretization, with 1280 time subintervals and different number of spatial steps

The best approximations are obtained for 10 subintervals in each spatial dimension. The error increment for more space intervals can be due to the fact that the condition number of the exponential matrix utilized to solve the differential system increases fast with the number of spatial intervals, making the solution less reliable.

The parallel splittings obtain better results in error and temporal cost than the serial method in almost all cases. The utilization of fourth order approximations for the spatial derivatives gives less error, mainly for small number of spatial intervals.

The temporal cost increases between linearly and quadratically with the number of spatial steps, being higher for the fourth order scheme.

Splitting algorithm	Space steps	e_{max}	e_{mean}	δ_{max}	δ_{mean}	Average iterations	Time(s)
Serial Iterative	5	1.0914e-04	1.4811e-05	1.9332e-04	2.9784e-05	2.0008	4.5106
	10	1.5092 e-03	1.4657e-04	2.0076e-03	1.9688e-04	2.0008	9.0779
	20	3.4036e-02	1.7974 e-03	4.3702e-02	2.3171e-03	2.1875	22.9909
	5	9.1209e-05	1.0079e-05	1.6718e-04	2.3934e-05	3.0008	3.3897
Classical Parallel	10	1.2736e-03	1.0740e-04	1.7251e-03	1.4737e-04	3.0023	6.9723
	20	1.9388e-02	1.0103e-03	2.6359e-02	1.3954e-03	4.0055	21.6915
Modern Parallel	5	9.1244e-05	1.0081e-05	1.6743e-04	2.3956e-05	3.0016	3.3446
	10	1.3121e-03	1.1194e-04	1.7697e-03	1.5233e-04	3.0023	7.0148
	20	2.1437e-02	1.2132e-03	2.8739e-02	1.6326e-03	9.0047	48.6330

Table $\overline{12}$

Results of the directional decomposition using the fourth order scheme for the spatial discretization, with 1280 time subintervals and different number of spatial steps

- Convection and diffusion decomposition:

Here, we decompose to an explicit part, which is the convection, and into an implicit part, which is the diffusion.

$$A(u)u = -1/2u(\partial_x u + \partial_y u) - 1/2(\partial_x u + \partial_y u), \tag{77}$$

$$Bu = \frac{1}{2}\mu(\partial_{xx}u + \partial_{yy}u) + f(x, y, t). \tag{78}$$

Tables 13 and 14 analyze the convergence of the different splitting methods for the convection and diffusion decomposition varying the time step.

Splitting algorithm	Time steps	e_{max}	e_{mean}	δ_{max}	δ_{mean}	Average iterations	Time(s)
	160	1.1984e-02	1.3870e-03	2.2076e-02	3.0691e-03	3.7250	5.2347
Serial Iterative	320	4.8721e-03	6.7652 e- 04	1.0137e-02	1.6865 e-03	3.0062	8.3601
Serial Iterative	640	2.1637e-03	3.3484e-04	5.2782 e- 03	1.0112e-03	2.0031	10.6301
	1280	1.0151e-03	1.6667e-04	3.1179e-03	6.7675 e-04	2.0008	21.1887
	160	9.7517e-03	1.4278e-03	1.9648e-02	3.1946e-03	10.1000	7.0953
Classical Parallel	320	4.5454 e - 03	7.1452e-04	9.9089e-03	1.7712e-03	5.0125	7.0217
Classical I arallel	640	2.1597e-03	3.5723e- 04	5.3681e-03	1.0580 e-03	4.0031	10.9344
	1280	1.0506 e - 03	1.7865e-04	3.2100e-03	7.0114e-04	3.0016	16.4434
Modern Parallel	320	4.5794e-03	7.1442e-04	9.9479e-03	1.7711e-03	9.6687	13.3257
	640	2.1612e-03	3.5716e-04	5.3734e-03	1.0579 e-03	4.8344	15.4879
	1280	1.0537e-03	1.7866e-04	3.2137e-03	7.0116e-04	3.0047	17.1489

Table $\overline{13}$

Results for the convection diffusion decomposition using the second order scheme for the spatial discretization with 10 spatial subintervals and different number of temporal steps

Tables 13 and 14 shows that, for a fixed spatial step, the estimated errors are proportional to the temporal step, in contrast with the results for the directional decomposition where the errors are of second order with respect time. This behavior can be due to the fact that in the directional decomposition, the solution can be obtained decomposing the matrices in blocks corresponding to the variables in a row or column of the spatial mesh because the spatial derivatives are separated in the split, whereas in the case

of the convection and diffusion decomposition, in each part of the split, they appear derivatives in both directions, preventing the solution in blocks. In the fourth order case, the modern parallel algorithm with 160 temporal steps diverges.

Splitting algorithm	Time steps	e_{max}	e_{mean}	δ_{max}	δ_{mean}	Average iterations	Time(s)
	160	1.1045e-02	1.3653e-03	1.9508e-02	2.6934e-03	4.3000	6.4748
Serial Iterative	320	4.4569 e-03	6.6767e-04	8.4817e-03	1.3315e-03	3.0062	9.1641
Serial Iterative	640	2.0588e-03	3.3259e-04	4.0263e-03	6.6459 e-04	2.5734	15.3656
	1280	9.9555e-04	1.6628e-04	1.9676e-03	3.3219e-04	2.0008	24.9010
	160	8.4411e-03	1.3887e-03	1.6781e-02	2.7980e-03	39.8875	30.5468
Classical Parallel	320	4.1997e-03	7.0339e-04	8.3408e-03	1.4133e-03	5.0156	7.6316
Classical I arallel	640	2.0774 e-03	3.5459 e-04	4.1414e-03	7.1074e-04	4.0047	11.9168
	1280	1.0361e-03	1.7814e-04	2.0640e-03	3.5636e-04	3.0016	17.6236
	320	4.1843e-03	7.0237e-04	8.3219e-03	1.4119e-03	25.3500	39.5635
Modern Parallel	640	2.0741e-03	3.5440e-04	4.1377e-03	7.1047e-04	4.7797	14.6983
	1280	1.0348e-03	1.7807e-04	2.0636e-03	3.5628e-04	3.0031	17.9165

Table $\overline{14}$

Results for the convection diffusion decomposition using the fourth order scheme for the spatial discretization with 10 spatial subintervals and different number of temporal steps

The dependence on the spatial step is analyzed in tables 15 and 16. The behavior is similar to the case of the directional decomposition, presenting an increment of the estimated error with the number of spatial subintervals for a fixed time step. In the second order scheme, the δ -errors decrease with the space step. In the fourth order scheme, the δ -errors are lower, but they slightly increase when the space step decreases. For this scheme, the modern parallel algorithm with 20 spatial steps diverges. The temporal cost is relatively high in the case of 20 subintervals, due tho the computational overhead for dealing with big matrices.

Splitting algorithm	Spatial intervals	e_{max}	$e_{ m mean}$	$\delta_{ m max}$	$\delta_{ m mean}$	Average iterations	Time(s)
Serial Iterative	5	5.6621e-04	1.0099e-04	8.0418e-03	1.2428e-03	2.0008	12.6049
	10	1.0151e-03	1.6667e-04	3.1179e-03	6.7675 e-04	2.0008	25.1408
	20	1.5741e-03	2.1460e-04	3.1037e-03	5.2448e-04	2.0016	244.2272
	5	6.4980e-04	1.1055e-04	8.2248e-03	1.2617e-03	3.0008	9.1133
Classical Parallel	10	1.0506 e - 03	1.7865e-04	3.2100e-03	7.0114e-04	3.0016	19.2526
	20	1.4562 e-03	2.2646e-04	2.9778e-03	5.4990e-04	4.0039	244.5342
Modern Parallel	5	6.4782e-04	1.1048e-04	8.2242e-03	1.2617e-03	3.0016	9.0470
	10	1.0537e-03	1.7866e-04	3.2137e-03	7.0116e-04	3.0047	19.5866
	20	1.4709 e-03	2.2651e-04	2.9890e-03	5.4993e- 04	6.0141	339.3991

Table 15

Results for the convection diffusion decomposition using the second order scheme for the spatial discretization with 10 temporal steps and different number of spatial subintervals

The use of order approximations to the space derivatives has little influence in the results of the algorithm.

- Balanced decomposition: We decompose into:

$$A(u)u = (1 - \epsilon) \left(-1/2u\partial_x u - 1/2\partial_x u + \mu\partial_{xx} u \right) + \epsilon \left(-1/2u\partial_y u - 1/2\partial_y u + \mu\partial_{yy} u \right) + \epsilon f(x, y, t), \tag{79}$$

$$B(u)u = \epsilon \left(-1/2u\partial_x u - 1/2\partial_x u + \mu\partial_{xx} u \right) + (1 - \epsilon) \left(-1/2u\partial_y u - 1/2\partial_y u + \mu\partial_{yy} u \right) + (1 - \epsilon) f(x, y, t). \tag{80}$$

where ϵ is an arbitrary parameter that can be tuned in order the achieve the maximum efficiency.

We first examine the influence of the parameter ϵ on the convergence of the different splitting schemes. Tables 17 and 18 show the influence of the parameter on the convergence of the methods. The method

Splitting algorithm	Time steps	$e_{ m max}$	$e_{ m mean}$	$\delta_{ m max}$	$\delta_{ m mean}$	Average iterations	Time(s)
	5	6.0051e-04	1.0212e-04	1.1690e-03	2.0887e-04	2.0008	14.2460
Serial Iterative	10	9.9555e-04	1.6628e-04	1.9676e-03	3.3219e-04	2.0008	25.2884
	20	1.4979e-03	2.1332e-04	2.8794 e-03	4.2570 e-04	2.3289	290.9178
	5	6.7972e-04	1.1140e-04	1.3249e-03	2.2717e-04	3.0008	10.6702
Classical Parallel	10	1.0361e-03	1.7814e-04	2.0640 e-03	3.5636e-04	3.0016	18.2544
	20	1.3859 e-03	2.2483e-04	2.7619e-03	4.5066e-04	4.0047	268.6941
	5	6.7894e-04	1.1130e-04	1.3246e-03	2.2718e-04	3.0016	10.2447
Modern Parallel	10	1.0348e-03	1.7807e-04	2.0636e-03	3.5628 e-04	3.0031	18.4238

Table 16

Results for the convection diffusion decomposition using the fourth order scheme for the spatial discretization with 1280 temporal subintervals and different number of spatial steps

has the same behavior for values of ϵ which are symmetric with respect to 0.5. The results are quite uniform for ϵ in the range [-0.1, 1.1] whereas for other parameter values, the method diverges. Only the classical parallel algorithm obtains the result for $\epsilon = 2$, whereas the serial iterative method even fails for $\epsilon = 1.1$ in the second order as well as in the fourth order schemes.

Splitting Algorithm	ϵ	e_{max}	e_{mean}	δ_{max}	δ_{mean}	Average iterations	Time(s)
	0.50	2.0353e-03	3.7750e-04	5.2893e-03	1.1017e-03	2.8812	18.8881
	0.70	2.0305e-03	3.7711e-04	5.3163e- 03	1.1028e-03	2.7359	17.5882
Serial Iterative	0.90	1.9874 e-03	3.7510e-04	5.3979e-03	1.1078e-03	2.0031	12.7734
Seriai Iterative	1.00	1.8680e-03	3.7128e-04	5.4439e-03	1.1173e-03	2.0031	13.3356
	0.50	2.0337e-03	3.7747e-04	5.2887e-03	1.1017e-03	3.9859	13.1540
	0.70	2.0241e-03	3.7712e-04	5.2762 e-03	1.1013e-03	3.9859	13.0457
Classical Parallel	0.90	1.9944e-03	3.7604 e-04	5.2378e-03	1.0999e-03	3.9859	13.2031
Classical Faranei	1.00	1.9713e-03	3.7521e-04	5.2081e-03	1.0988e-03	3.9859	13.5399
	0.50	2.0337e-03	3.7747e-04	5.2887e-03	1.1017e-03	3.9859	13.1767
	0.70	2.0246e-03	3.7713e-04	5.2776e-03	1.1013e-03	4.0047	13.3876
Modern Parallel	0.90	1.9957e-03	3.7606e-04	5.2423e-03	1.1000e-03	4.0047	13.3008
Modern Parallel	1.00	1.9721e-03	3.7521e-04	5.2139e-03	1.0990e-03	4.0062	13.3750

Table 17

Results for the balanced decomposition using the second order scheme for the spatial discretization with 640 temporal subintervals and different values of parameter ϵ

We fix the parameter $\epsilon=0.9$ for the analysis of the method. Table 19 compare the results of the considered splitting algorithms for different number of temporal steps using the fourth order schemes, respectively. The fourth order scheme provides better approximations than the second order one. The parallel methods behave better than the serial method for more than 160 temporal steps.

Finally, we compare the temporal cost of the different combinations of algorithms for different number of time steps and 10 space subintervals. The results are shown in table 20. But for the case of 160 time steps, the classical parallel algorithm obtains the best results for the different analyzed decomposition methods. For 1280 time steps, the modern parallel algorithm also outperforms the serial iterative algorithm.

Here, we have the benefit, e.g., for small $\epsilon = 0.1$, we have nearly a directional decomposition scheme. Based on the $\epsilon \neq 0$, we stabilize the scheme.

Remark 5.4 We also compared the computational time of the serial iterative, classical parallel iterative and modern parallel iterative method. We obtain a speedup for the parallel versions of a factor nearly $L_{class} \leq L$, if we apply L-processors. For the modern parallel version, we obtain with the same accuracy of the results a small reduction of the speedup with $L_{modern} < L_{class}$.

Splitting Algorithm	ϵ	e_{max}	e_{mean}	δ_{max}	δ_{mean}	Average iterations	Time(s)
	0.50	2.0372 e-03	3.7851e-04	4.1693e-03	7.5936e-04	2.8953	18.0265
	0.70	2.0322e-03	3.7806e-04	4.1578 e-03	7.5847e-04	2.7594	17.4235
Serial Iterative	0.90	1.9852 e-03	3.7580 e- 04	4.0627 e-03	7.5404 e-04	2.0031	14.1016
	1.00	1.8577e-03	3.7189e-04	3.8137e-03	7.4641e-04	2.0031	12.6249
	0.50	2.0351e-03	3.7850e-04	4.1684e-03	7.5935e-04	4.0000	13.0915
	0.70	2.0238e-03	3.7818e-04	4.1532 e-03	7.5894 e-04	4.0000	12.7921
Classical Parallel	0.90	1.9891 e- 03	3.7724 e-04	4.1063e-03	7.5767e-04	4.0000	14.7150
Classical Laranei	1.00	1.9620e-03	3.7651e-04	4.0699e-03	7.5671e-04	4.0000	12.7658
	0.50	2.0351e-03	3.7850 e-04	4.1684 e-03	7.5935e-04	4.0000	12.9446
	0.70	2.0237e-03	3.7817e-04	4.1532 e-03	7.5892e-04	4.0047	13.3572
Modern Parallel	0.90	1.9872 e- 03	3.7717e-04	4.1051e-03	7.5759e-04	4.0062	12.9882
Modern Faranei	1.00	1.9573e-03	3.7639e-04	4.0660e-03	7.5656e-04	4.0062	12.7415

Table 18 Results for the balanced decomposition using the fourth order scheme for the spatial discretization with 640 temporal subintervals and different values of parameter ϵ

Splitting Algorithm	Time steps	e_{max}	e_{mean}	δ_{max}	δ_{mean}	Average iterations	Time(s)
	160	7.7146e-03	1.4517e-03	1.4819e-02	2.9435e-03	5.3063	9.1072
Serial Iterative	320	3.7147e-03	7.4258e-04	7.7759e-03	1.4957e-03	3.0094	10.5303
Seliai Iterative	640	1.9852 e- 03	3.7580 e- 04	4.0627 e-03	7.5404 e-04	2.0031	14.1477
	1280	1.0283e-03	1.8906e-04	2.0776e-03	3.7847e-04	2.0008	27.5724
	160	6.2090e-03	1.4422e-03	1.3622e-02	2.9381e-03	39.4500	35.0620
Classical Parallel	320	3.6203 e-03	7.4307e-04	7.7242e-03	1.4998e-03	5.0156	8.4539
Classical Laranei	640	1.9891e-03	3.7724 e-04	4.1063e-03	7.5767e-04	4.0000	13.1319
	1280	1.0447e-03	1.9009e-04	2.1175e-03	3.8066e-04	3.0016	19.8095
	160	6.1419e-03	1.4425e-03	1.3720e-02	2.9384e-03	40.0000	35.7971
Modern Parallel	320	3.6291e-03	7.4300e-04	7.7320e-03	1.4997e-03	6.5125	12.2648
	640	1.9872 e- 03	3.7717e-04	4.1051e-03	7.5759e-04	4.0062	17.2202
	1280	1.0454 e-03	1.9008e-04	2.1182e-03	3.8065e-04	3.0023	24.0136

Table $\overline{19}$

Results for the balanced decomposition with $\epsilon=0.9$ using the fourth order scheme for the spatial discretization with 10 spatial subintervals and different number of temporal steps

Remark 5.5 Based on balancing the decomposition of the operators with ϵ , we also obtain the benefit of the classical and modern parallel iterative splitting method with larger time-steps and more iterative steps to improve the accuracy. We can accelerate the speed of computations with the parallel versions and obtain speed ups with the modern parallel version.

6. Conclusion

The extensions of iterative splitting methods to parallel versions allow to reduce the computational time. We obtain the same accuracy as in the serial version. The improvements are obtained with larger time-steps and additional iterative steps, while we could reduce the computational time with the parallel methods. The benefit is of course the balance to multiple processors with additional memories. Further, we could apply the resources to improve with additional steps the accuracy of the approximations. We circumvent the problem of the *memory of the algorithm*, see [8], which we have if we obtain only a serial method. Based on the parallel distribution, we have additional iterative steps for each processor and we distribute such a memory to all processors. For large scale numerical experiments, we could present the

Derivative approximation	T:	2nd order	4th order	2nd order	4th order	2nd order	4th order
Splitting method	Time steps	Directional	decomposition	Convection	Diffusion dec.	Balanced d	ecomposition
	160	2.7961	3.1916	5.2347	6.4748	10.7499	9.1072
Serial Iterative	320	3.2505	3.5150	8.3601	9.1641	10.2080	10.5303
Serial Iterative	640	5.5462	5.4186	10.6301	15.3656	14.0321	14.1477
	1280	8.4610	9.1334	21.1887	24.9010	29.1536	27.5724
	160	3.2553	10.8616	7.0953	30.5468	9.4198	35.0620
Classical Parallel	320	2.7616	2.9691	7.0217	7.6316	9.1697	8.4539
Classical I arallel	640	4.4090	4.6626	10.9344	11.9168	14.8338	13.1319
	1280	6.4774	6.9907	16.4434	17.6236	21.0049	19.8095
	320	4.3984	5.9235	13.3257	39.5635	9.5343	12.2648
Modern Parallel	640	4.4004	5.8144	15.4879	14.6983	13.1658	17.2202
Modern Faranei	1280	6.4701	6.9861	17.1489	17.9165	21.9695	24.0136

Table 20

Cost in time of the different algorithms.

benefit of the parallel resources.

In future, we consider more real-life problems and extend the parallel iterative methods to stochastic differential equations.

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References

- [1] I. Farago and J. Geiser. *Iterative Operator-Splitting methods for Linear Problems*. International Journal of Computational Science and Engineering, 3(4), 255-263, 2007.
- [2] A. Frommer and D.B. Szyld. On asynchronous iterations. J. Comput. Appl. Math., 123:201–216, 2000.
- [3] J. Geiser. Iterative Splitting Methods for Differential Equations. Numerical Analysis and Scientific Computing Series, Taylor & Francis Group, Boca Raton, London, New York, 2011.
- [4] J. Geiser. Multi-stage waveform Relaxation and Multisplitting Methods for Differential Algebraic Systems. arxiv:1601.00495 (http://arxiv.org/abs/1601.00495), January 2016.
- [5] J. Geiser. Multicomponent and Multiscale Systems: Theory, Methods, and Applications in Engineering. Springer, Cham, Heidelberg, New York, Dordrecht, London, 2016.
- [6] J. Geiser. Picards iterative method for nonlinear multicomponent transport equations. Cogent Mathematics, 3(1), 1158510, 2016.
- [7] J. Geiser. Iterative Splitting Methods for Coulomb Collisions in Plasma Simulations. arxiv:1706.06744 (http://arxiv.org/abs/1706.06744), June 2017.
- [8] T. Ladics and I. Farago. Generalizations and error analysis of the iterative operator splitting method. Cent. Eur. J. Math., 11(8):1416-1428, 2013.
- [9] R. Jeltsch and B. Pohl. Waveform Relaxation with Overlapping Splittings. SIAM J.Sci.Comput., 16(1):40-49, 1995.
- [10] D.P. O'Leary and R.E. White. Multi-splittings of matrices and parallel solution of linear systems. SIAM J. Algebraic Discrete Methods, 6:630–640, 1985.
- [11] U. Miekkala and O. Nevanlinna. Convergence of dynamic iteration methods for initial value problems. SIAM J. Sci. Stat. Comput., 8: 459-482, 1987.
- [12] U. Miekkala and O. Nevanlinna. Iterative solution of systems of linear differential equations. Acta Numerica, 5: 259-307, 1996.
- [13] S. Vandewalle. Parallel Multigrid Waveform Relaxation for Parabolic Problems. Teubner Skripten zur Numerik, B.G. Teubner Stuttgart, 1993.
- [14] D. Yuan and K. Burrage. Convergence of the parallel chaotic waveform relaxation method for sti systems. Journal of Computational and Applied Mathematics, 151: 201-213, 2003.
- [15] R.E. White. Parallel algorithms for nonlinear problems. SIAM J. Algebraic Discrete Methods, 7:137–149, 1986.