

## Article

# The Akimotoite–Majorite–Bridgmanite Triple Point Determined in Large Volume Press and Laser-Heated Diamond Anvil Cell

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**Abstract:** The akimotoite–majorite–bridgmanite (Ak–Mj–Bm) triple point in  $\text{MgSiO}_3$  has been measured in large-volume press (LVP; COMPRES 8/3 assembly) and laser-heated diamond anvil cell (LHDAC). For the LVP data, we calculated pressures from the calibration by Leinenweber *et al.* [1]. For the LHDAC data, we conducted *in situ* determination of pressure at high temperature using the Pt scale by Dorogokupets and Dewaele [2] at synchrotron. The measured temperatures of the triple point are in good agreement between LVP and LHDAC at 1990–2000 K. However, the pressure for the triple point determined from the LVP is  $3.9 \pm 0.6$  GPa lower than that from the LHDAC dataset. The triple point determined through these experiments will provide an important reference point in the pressure–temperature space for future high-pressure experiments and allow mineral physicists to compare the pressure–temperature conditions measured in these two different experimental methods.

**Keywords:** triple point; bridgmanite; akimotoite; majorite; large-volume press; laser-heated diamond anvil cell

## 1. Introduction

Accurate determination of pressure ( $P$ ) and temperature ( $T$ ) is essential for laboratory experiments to contribute to the geophysical understandings of the deep interiors of Earth, other planets in the solar system, and exoplanets. As Orson Anderson demonstrated through his important work, thermal equations of state (EOS) of the standard materials can be used for estimating pressures from measurements of volume and temperature in high-pressure experiments [3].

The development of third generation synchrotron facilities in late 1990s and early 2000s enabled *in situ* measurements of phase boundaries and of physical properties, allowing mineral physicists to take advantage of the thermal EOS established by Anderson and others. However, early experiments found that results from different pressure scales and different experimental techniques differed by 2–3 GPa for the important phase boundaries near the 660-km discontinuity [4–8]. Intense efforts have been made since then [e.g., 9–11], but the discrepancy for some important phase boundaries still remain unresolved [12,13]. Other potential sources of the discrepancy have been investigated. For example, pressure effects on the thermocouple electromotive force (emf) calibration are important for large-volume press (LVP) experiments [14–16]. Spectroradiometry has been the standard method for temperature measurements in laser heating, but optical effects from diamond anvils and thermal gradients could potentially introduce artifacts in the measured temperature [17–20]. While the accuracy of the spectroradiometry through LHDAC has been studied [21], temperatures from spectroradiometry and thermocouples have never been cross examined at high pressure to our knowledge.

Comparing phase boundaries provides an opportunity to examine the pressure and temperature scales in high-pressure apparatuses. However, in such an effort, it is difficult to separate pressure effects and temperature effects [12]. It is desirable to have a reference point in the  $P$ – $T$  space for

such comparison. At pressures between 20 and 24 GPa, the triple point exists between akimotoite (Ak), majorite (Mj), and bridgmanite (Bm) in  $\text{MgSiO}_3$ . The understanding of the triple point not only provides a new opportunity to compare pressures and temperatures measured in different techniques but also the triple point itself is important in geophysics for understanding the origin of the seismic discontinuity structures near 660-km depth. Although the 660-km discontinuity has been related mainly to the post-spinel transition in  $\text{Mg}_2\text{SiO}_4$  for many decades [22], it has been well known that the phase boundaries in  $\text{MgSiO}_3$  can exist at depths very close to the 660-km discontinuity, and therefore affect the complex discontinuity structures at the bottom-most mantle transition zone [23,24].

In this paper, we report the triple point between Ak, Mj, and Bm in pure  $\text{MgSiO}_3$  measured in both LVP and LHDAC. We compare the pressure and temperature of the triple point from those two separate measurements and discuss possible sources of discrepancy between LVP and LHDAC.

## 2. Methods

### 2.1. Large-Volume Press (LVP)

High-pressure LVP experiments were conducted in the 1100 ton multi-anvil press at the Eyring Materials Center at Arizona State University. The LVP consists of eight second-stage anvils surrounding the cell assembly, six first-stage anvils surrounding the second-stage anvils, and a support ring surrounding the first-stage anvils. The second-stage anvils are tungsten carbide and have a 3.0 mm truncation on each corner and were used with injection-molded ceramic octahedra with 8.0 mm edges [1]. All LVP experiments reported here were conducted with the COMPRES 8/3 assembly [1,25].

Each assembly used a cylindrical Re-capsule made from a  $3 \times 6 \text{ mm}^2$  sheet of rhenium foil that was filled with pure  $\text{MgSiO}_3$  glass synthesized using the containerless laser levitation method [26,27]. This capsule was placed inside of an  $\text{MgO}$  sleeve within the Re-furnace and  $\text{LaCrO}_3$  sleeve. The Re-furnace and  $\text{LaCrO}_3$  sleeve were placed inside the octahedron. The completed assembly was placed on top of four second stage anvils with four other second stage anvils placed on top once inside of the LVP.

Pressures for each run were raised at a rate of 400 psi/hour, and temperatures were raised to target temperatures at a constant rate of 100 K/min. Heating of the assemblies lasted 30–60 minutes to avoid kinetic effects. Temperature was monitored just above the capsule in the assembly by a W5%Re-W%26Re thermocouple. The run was held at target temperature for the heating duration until quenched. Pressure began being pumped down immediately after quench at a rate of 400 psi/hour. Once at ambient conditions, the assembly was recovered for analysis (Tab.1).

**Table 1.** Large-volume press runs performed in this study. Estimated uncertainty for pressure is 0.5 GPa. Estimated uncertainty without the pressure effects on thermocouple is 5 K. We assigned an uncertainty of  $\pm 50$  K for the runs where we estimated temperature through applied voltage calibrated for thermocouple.  $T$ : temperature,  $P$ : pressure, O.P.: oil pressure,  $F$ : force,  $t$ : heating duration.

Run ID	$T$ (K)	O.P. (psi)	$P$ (GPa)	$F$ (ton)	$t$ (min)	Product
BB1434JD	$2173 \pm 5$	5000	20.5	515	30	Bm
BB1438JD	$2023 \pm 5$	4500	20.0	463	30	Bm
BB1469BK	$1873 \pm 5$	4937	22.0	611	60	Bm
BB1470BK	$1873 \pm 5$	4365	20.0	449	60	Ak
BB1472BK	$2073 \pm 50$	4103	19.0	421	30	Mj
BB1475BK	$1873 \pm 5$	4892	21.0	504	60	Bm
BB1476BK	$1875 \pm 50$	4609	20.5	475	60	Bm
BB1477BK	$1973 \pm 5$	4025	19.0	414	30	Mj
BB1478BK	$1923 \pm 50$	4192	19.5	432	30	Ak+Mj
BB1479BK	$1948 \pm 5$	4344	20.0	447	30	Ak

For the pressure calculation, we used the calibration for COMPRESS 8/3 assembly reported by Leinenweber *et al.* [1]. In that study, *in situ* X-ray diffraction acquired in LVP at the APS synchrotron,

Argonne National Lab. The oil pressure was calibrated for the capsule pressure using tungsten thermal EOS [28] (Tab. 1).

Recovered LVP samples were sliced axially and were carefully polished using alumina sandpaper. For phase identification, Raman spectroscopy was used for all LVP products in 1D point and 2D scanning modes at ASU. Laser wavelength was 532.22 nm and the beam size was 2  $\mu\text{m}$ . Measurements were taken for 1000 accumulations of 1 second exposure each (16 minutes total) on each spot at 30–45 mW of laser power.

## 2.2. Laser-Heated Diamond Anvil Cell (LHDAC)

The same  $\text{MgSiO}_3$  glass used in LVP was loaded in DAC #1 in LHDAC experiments (Tab. 2). The sample powder was mixed with 10 wt% Pt powder which served as the pressure calibrant and laser coupler for heating. Pressure was calculated by combining the measured unit-cell volume and temperature with the thermal EOS by Dorogokupets and Dewaele [2]. The powder mixture was pressed into a foil which was 100–200  $\mu\text{m}$  in size and 15–20  $\mu\text{m}$  thick. The foil was loaded into a laser-drilled hole in a pre-indented rhenium gasket using a micro-manipulator (Microsupport Axis Pro SS). The foil was supported by 5 spacers that consisted of <15  $\mu\text{m}$  pure sample grains on each side of the diamond culet. Argon gas (Ar) was loaded as a pressure medium and thermal insulator in a gas loading system (Sanchez Technologies GLS1500) at 1450 bar pressure at ASU.

For measurements in DACs #2 and #3, we used Ak as a starting material. The Ak sample was synthesized in the LVP at ASU. For sufficient amount of akimotoite sample, we conducted the synthesis in a 10/5 COMPRES assembly with Re foil capsule filled with  $\text{MgSiO}_3$  glass. The synthesis conditions were 22 GPa and 1473 K for 1 hour. After quenching, the capsule was cut in half with a diamond wire saw and analyzed with Raman to ensure pure Ak product. The Ak sample was removed from the capsule and ground in an agate mortar under acetone for 20 min. The sample powder was mixed with 10 wt% Pt powder which once again served as the pressure calibrant [2] and laser coupler. The mixture was again pressed into a foil of the same size as the  $\text{MgSiO}_3$  sample. For these samples, we used NaCl as a pressure medium which was dried at 100 °C for 24 hours. Thin foils of NaCl were loaded into the sample chamber above and below the sample foil. All samples were compressed in symmetric-type DAC utilizing type Ia standard design 400  $\mu\text{m}$  diamond anvils.

LHDAC data was acquired using monochromatic XRD in DAC at sector 13-IDD of the GeoSoilEnviroCARS (GSECARS) at the Advanced Photon Source (APS) utilizing the double-sided laser heating system. The laser beams were co-axially aligned with the X-ray beam to measure diffraction patterns on the heating spot. The typical beam diameters for the X-ray beam and laser heating spot are 5  $\mu\text{m}$  and 20  $\mu\text{m}$ , respectively.

The 2D diffraction images were acquired using a Dectris Pilatus detector at GSECARS 13-IDD and then integrated to 1D diffraction patterns using calibration parameters obtained from the  $\text{LaB}_6$  standard in DIOPTAS software [29]. Diffraction images were exposed for 5–10 seconds and were collected before, during, and after heating. The diffraction images were analyzed and phases identified using PeakPo software [30]. Pt peaks were fitted with pseudo-Voigt profile function to obtain the peak positions in PeakPo. Pressure was calculated with pytheos [31]. The unit-cell volume of Pt was calculated with the 111 and 200 peaks and some with the 220 peak as well.

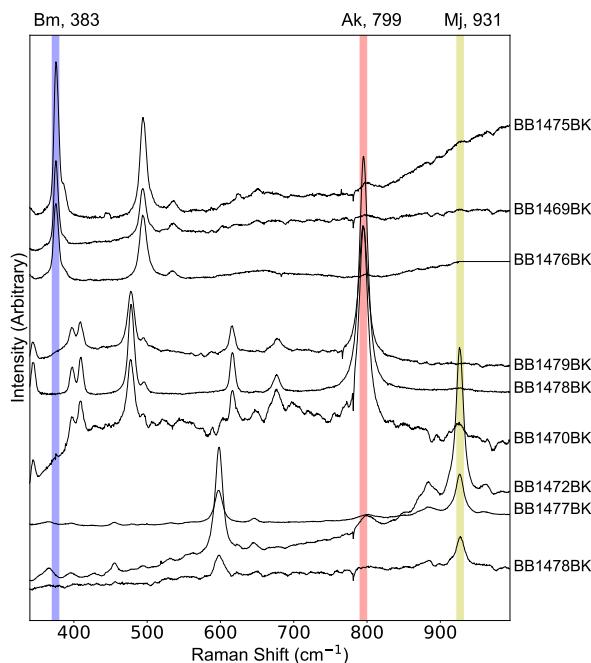
## 3. Results

### 3.1. Large-Volume Press

Raman spectroscopy was used to identify the recovered samples. Bm was identified by the strong modes at 383 and 501  $\text{cm}^{-1}$  (Fig. 1), and a weaker mode at 542  $\text{cm}^{-1}$  also appeared clearly in spectra. Ak was identified by the strong modes at 480 and 799  $\text{cm}^{-1}$  along with a few other weaker modes. Mj was identified by the strong modes at 602 and 931  $\text{cm}^{-1}$ . The peak positions for all three phases are in good agreement with those reported in literature [32–35].

**Table 2.** Experimental runs in LHDAC in this study. PM: pressure medium, SM: starting material, XE: X-ray energy. Other abbreviations are the same as Table 1.

DAC ID	Setting	Spot ID	P (GPa)	T (K)	Product
#1	PM: Ar	1	20.2	1506	Ak
	SM: MgSiO <sub>3</sub>	2	21.3	1897	Mj
	XE: 37 KeV	3	24.1	1903	Ak
		4	21.0	1610	Ak
		5	21.4	2056	Mj
		6	22.0	1531	Ak
		7	23.2	1870	Ak
		8	30.3	2110	Bm
		9	26.4	2097	Bm
		10	24.5	1717	Bm
		11	27.2	1884	Bm
		12	23.5	1717	Ak
		13	29.2	1706	Bm
		14	24.3	2123	Bm
		15	27.1	2096	Bm
		16	26.3	1992	Bm
		17	27.0	2021	Bm
		18	25.5	1934	Bm
		19	21.4	1541	Ak
		20	22.5	1897	Mj
		21	23.6	2050	Mj
		22	22.0	1637	Ak
#2	PM: NaCl	1	16.8	1616	Mj
	SM: Ak	2	17.5	1761	Mj
	XE: 30 KeV	3	22.4	2019	Mj
		4	20.1	1964	Mj
		5	24.0	2349	Mj
		6	20.6	1931	Mj
		7	24.0	2136	Mj
		8	20.2	1922	Mj
		9	21.2	1905	Mj
		10	26.1	1809	Bm
		11	21.7	1798	Mj
		12	22.8	1804	Ak
		13	22.8	2004	Mj
		14	27.5	1835	Bm
		15	29.3	2003	Bm
#3	PM: NaCl	1	26.1	1625	Bm
	SM: Ak	2	27.0	1600	Bm
	XE: 30 KeV	3	24.6	1900	Bm
		4	25.3	1831	Bm
		5	25.1	1908	Bm
		6	24.6	1600	Bm
		7	26.9	2126	Bm



**Figure 1.** Raman spectra of the samples recovered from the LVP runs. The vertical lines highlight the key peaks for the identification of Bm (blue) at  $383\text{ cm}^{-1}$ , Ak (red)  $799\text{ cm}^{-1}$ , and Mj (yellow)  $931\text{ cm}^{-1}$ . A LVP run ID is provided to each spectrum (Tab.1).

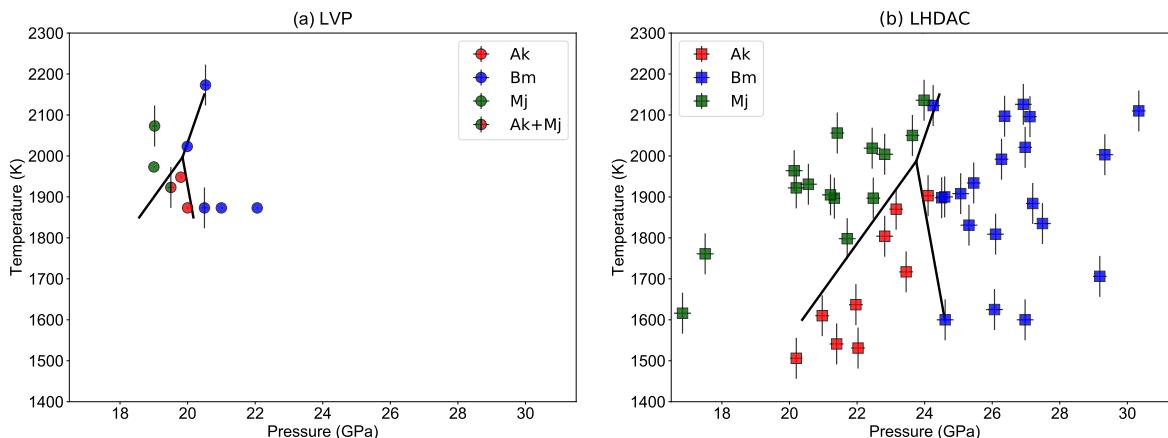
For sample BB1478BK, we identified both Ak and Mj phases located at different areas within the sample: Ak was found within the cooler region of the sample (outer area close to capsule) and Mj was found within the warmer region of the sample (center of capsule). This sample with two phases allows us to tightly constrain the Ak–Mj phase boundary (and therefore the triple point), and will later be discussed.

We inspected the samples under a stereomicroscope. The Bm samples are white, light-gray semi-translucent with  $5\text{--}10\text{ }\mu\text{m}$  crystals. The Ak samples are dark-gray with light-gray speckled throughout with  $2\text{--}5\text{ }\mu\text{m}$  crystals. The Mj samples are gray to light-gray in color with the middle of the sample being lighter in color with  $<5\text{ }\mu\text{m}$  crystals.

We chose to conduct LVP experiments at a narrow  $P\text{--}T$  field to make high-resolution determination of the Ak–Mj–Bm triple point. The observations of both Mj and Ak in one sample (BB1478BK) provide a tight constraint on the location of the Ak–Mj boundary (Fig. 2a). The observations of Ak at 19.5 GPa and 1923 K and Mj at 19 GPa and 1973 K indicate that the Ak–Mj boundary temperature should not exceed  $\sim 2000\text{ K}$  at 19–20 GPa. The observations of Ak at 20 GPa and 1873–1948 K and Bm at 20.5–22 GPa and 1873–2023 K suggest that the Ak–Bm boundary should be located between the pressures in the LVP experiments. In this determination we adapted the Clapeyron slopes of the three boundaries from Ishii *et al.* [36]. Combining these constraints with other data points, the  $P\text{--}T$  condition for the triple point is determined to be  $19.9 \pm 0.4\text{ GPa}$  and  $2000 \pm 50\text{ K}$  in our LVP experiments on  $\text{MgSiO}_3$ . We attempted to fit the data with different Clapeyron slopes [e.g., 6]. Because of the narrow  $P\text{--}T$  region of our study and the better agreement in the Clapeyron slopes of these boundaries in the literature, our determined  $P\text{--}T$  conditions for the triple point is not sensitive to the choice of the Clapeyron slopes of the boundaries.

### 3.2. Laser-Heated Diamond Anvil Cell

Three LHDAC were used to acquire data for the triple point. Cell #1 loaded with the  $\text{MgSiO}_3$  glass starting material was used to determine the Mj–Bm phase boundary through rapid heating into the Mj field and then slowly increasing temperature to gain pressure through thermal pressure effects until



**Figure 2.** Data points from our (a) LVP and (b) LHDAC experiments. The red, green, and blue symbols represent Ak, Mj, and Bm, respectively. The estimated phase boundaries and triple point are shown. We fixed the Clapeyron slopes of the involved phase boundaries to those reported in Ishii *et al.* [36] and then adjusted the location of the phase boundaries to fit our data points. Our LVP triple point is located at  $2000 \pm 50$  K and  $19.9 \pm 0.4$  GPa. Our LHDAC triple point is located at  $1990 \pm 100$  K and  $23.8 \pm 0.6$  GPa.

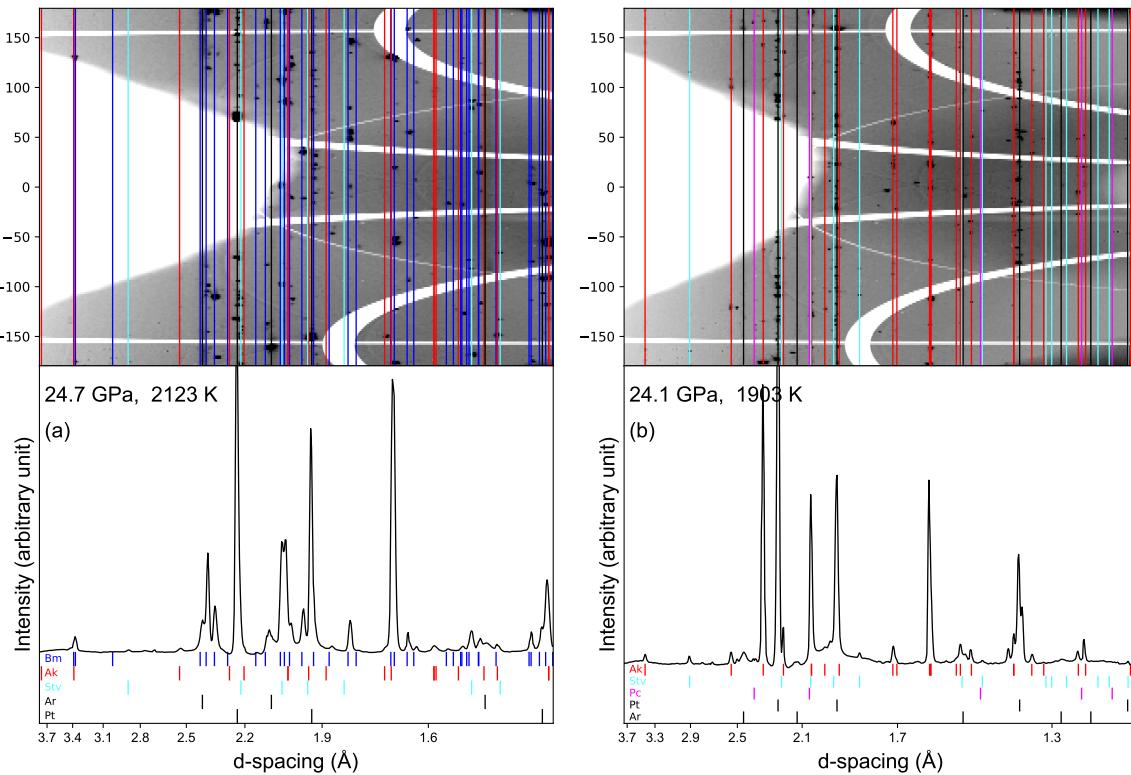
the first Bm peaks are observed (Fig. 3a). In order to avoid forming Ak, rapid heating was conducted because Ak could remain as a metastable phase and overlap with the Bm and Mj diffraction peaks. Rapid heating was accomplished by aiming the heating laser beam at a previously heated spot. Once the desired temperature was reached, the sample was rapidly translated by moving the LHDAC such that the heating laser was aimed at an adjacent unheated spot, and thus reached target temperature quickly without a slow laser power ramp up. Once the Mj peaks are observed, the temperature was increased slowly until the first Bm peaks form. The location in  $P-T$  space where this happens was used to determine the Mj–Bm phase boundary (Fig. 2b).

The experiments with cells #2 and #3 loaded with the Ak starting material were designed to constrain the boundaries between Ak and Mj and between Ak and Bm. Temperature was raised slowly until the first observation of the Bm or Mj diffraction peaks.

Mj was mainly identified through the 040, 323, and 431 peaks. The 444 and 046 peaks were observable in some patterns and these helped to further confirm the presence of Mj. Ak was identified with the 10 $\bar{2}$ , 104, 110, 20 $\bar{4}$ , and 11 $\bar{6}$  peaks. The 113 peak was sometimes used, but it overlaps enough with the Mj 125 peak that it was less useful. Finally, Bm was identified with the 111, 200, 120, 210, 022, 122, 121, 023, and 221. Higher 2 $\theta$  peaks were used to confirm the presence of Bm, if available. Similar to our LVP analysis, the triple point was constrained by moving the three phase boundaries in the  $P-T$  space while we fix the Clapeyron slopes of the boundary to those reported in Ishii *et al.* [36]. The  $P-T$  condition for the triple point from our LHDAC dataset is  $23.8 \pm 0.6$  GPa and  $1990 \pm 100$  K.

#### 4. Discussion

The  $P-T$  conditions of the triple point we obtain are  $19.9 \pm 0.4$  GPa and  $2000 \pm 50$  K for LVP and  $23.8 \pm 0.6$  GPa and  $1990 \pm 100$  K for LHDAC (Fig. 4). Despite the fact that they can be biased by different systematic error sources (such as pressure effects on thermocouple emf calibration in LVP and optical effects on spectroradiometry through diamond anvils in LHDAC), a remarkable agreement was found in temperature for the triple point from both techniques. To our knowledge, this is the first direct cross examination of the two temperature measurement techniques. Such an agreement may not necessarily be applicable for pressures much lower or much higher than the range we conducted measurements, i.e., 20–24 GPa, because some of the perceived systematic error sources could be pressure dependent, such as pressure effects on thermocouple emf calibration. However, for the pressure range of the mantle transition zone, our result provides important experimental confirmation for comparing temperature measurements from LVP and LHDAC experiments.



**Figure 3.** 2D diffraction images of LHDAC data (top) with integrated 1D diffraction patterns (bottom): (a) the formation of Bm, shown in blue, across the Mj/Bm phase boundary, and (b) the stability of Ak and absence of Bm. We provide  $P-T$  conditions of the diffraction measurements. The colored vertical bars in the 2D image and the colored ticks in the 1D pattern show the diffraction peak positions of different phases. We found some weak diffraction features possibly from Ak and Stv (stishovite) in (a) and Stv and P<sub>c</sub> (pericline). They are likely formed as transient phases due to kinetics and differential elemental diffusion during phase transition.

Our results suggest that pressure calibration is the most important issue to resolve in comparing LVP and LHDAC datasets, 3.9 GPa difference. So far LHDAC studies have reported systematically higher pressures for the phase boundaries in the mantle transition zone by 2–3 GPa compared with LVP studies [4–6,8,12]. Our new results reported here also confirm the trend but with a greater magnitude. The difference is particularly important to resolve because COMPRES 8/3 assembly and its pressure calibration [1] have been widely used in high-pressure studies.

The calibration for 8/3 assembly was conducted through *in situ* measurements using the tungsten EOS [28] in Leinenweber *et al.* [1]. For our LHDAC experiments, we chose to use the thermal EOS of Pt by Dorogokupets and Dewaele [2]. This pressure scale is known to yield better agreements with the Au and MgO scales according to Ye *et al.* [13] for a wide pressure range. In order to ensure the consistency between LHDAC and LVP results, it would be useful to examine the agreements between the W scale and the Pt scale at *in situ* high  $P-T$ . However, the high shear strength of tungsten can be a potential issue for the accurate determination on its EOS [37,38]. Therefore, the COMPRES 8/3 assembly can be calibrated using other pressure standards. Pt, Au, and MgO could be a good candidate as they have been used widely in LHDAC. However, despite the agreements over a larger pressure range, the MgO, Au, and Pt scale by Dorogokupets and Dewaele [2] are different by ~2 GPa at 20–40 GPa and high temperatures according to Ye *et al.* [13]. Therefore, an important challenge still remains for the thermal EOS of important materials at the pressure range for the mantle transition zone. It is notable that some of the standard materials (particularly Au) have low melting temperature compared with the mantle geotherm at pressures of  $\leq 30$  GPa, potentially causing significant anharmonic effects in their EOS or even abnormal pre-melting behaviors.

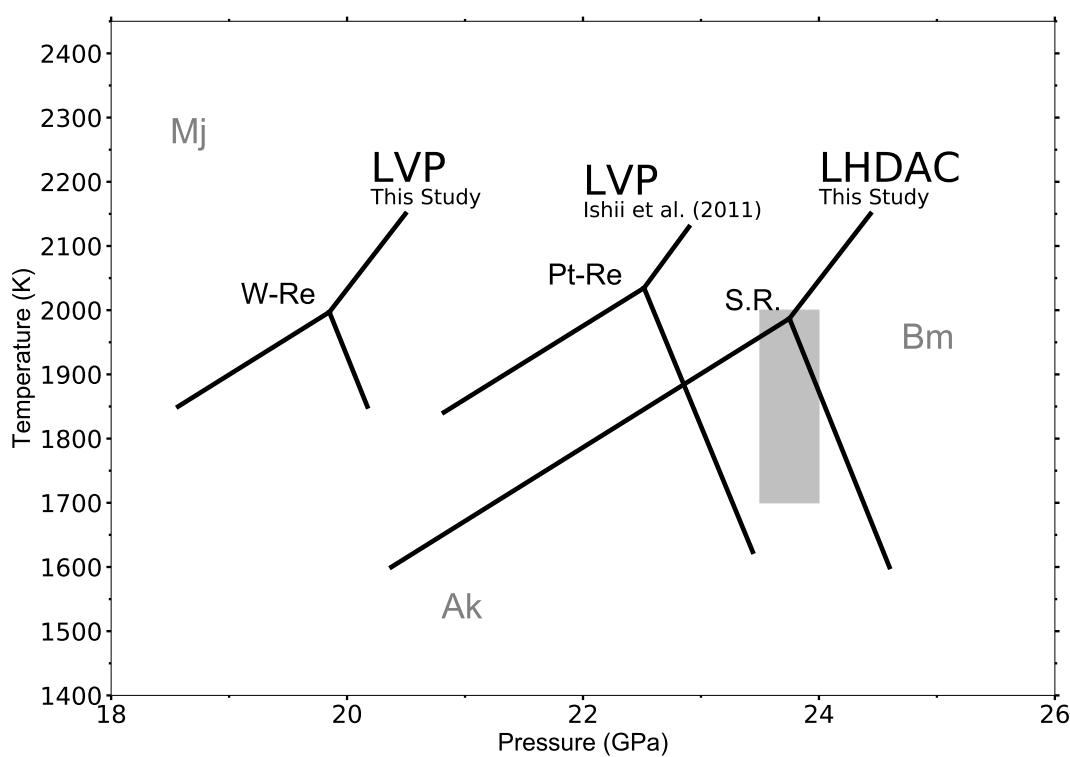
The seismic properties of the 660-km discontinuity are in general agreement with those of the post-spinel transition [39]. Ishii *et al.* [36] showed that the Ak–Bm boundary should be close to the post-spinel boundary within 1 GPa [36]. Fig. 4 shows that the Ak–Bm and Mj–Bm boundaries in  $\text{MgSiO}_3$  measured by LHDAC are closer to the  $P-T$  conditions expected for the 660-km discontinuity. Does it mean that LHDAC yields more reliable results for the location of the phase boundary? However, this approach is not desirable as the experimental methods should be able to address the question of whether the mantle phase boundaries are indeed the source of the 660-km discontinuity rather than the other way around.

Some former LVP studies have measured the Ak–Mj, Mj–Bm, and Ak–Bm boundaries in pure  $\text{MgSiO}_3$  and inferred triple point [6,36,44]. However, the reported  $P-T$  conditions for the boundaries and the triple point do not agree with each other: the discrepancy can be as large as 2 GPa in pressure and 300 K in temperature among the LVP measurements. In terms of temperature, our results are in best agreement with the most recent report by Ishii *et al.* [36]. They reported stability of Ak up to 1973 K at 22.3 GPa and the triple point in their phase diagram can be inferred to be  $2035 \pm 60$  K, which is in agreement with our LVP and LHDAC results on the temperature of the triple point within 100 K. This agreement is encouraging in that at least the recent studies converge on the temperature of the Ak–Mj–Bm triple point within 100 K even between different high-pressure techniques (LVP and LHDAC) and between different temperature measurement techniques (thermocouple, W5%Re–W26%Re in our LVP study and Pt/Pt-13%Re in Ishii *et al.* [36]’s LVP study, and spectroradiometry in our LHDAC study).

The pressure inferred for the triple point in Ishii *et al.* [36] is between our LVP and LHDAC results, located approximately in the middle (Fig. 4). They calibrated pressures at high temperatures based on previously reported boundaries in  $\text{Mg}_2\text{SiO}_4$ ,  $\text{MgSiO}_3$ , and  $\text{MgAl}_2\text{O}_4$  (see Ishii *et al.* [36] for references) which are all different in pressure calculation methods. The difference re-iterate the fact that the challenge in comparing LVP and LHDAC remains in the pressure measurements.

## 5. Conclusions

Both our large-volume press (LVP) and laser-heated diamond anvil cell (LHDAC) measurements show that the temperature of the akimotoite–majorite–bridgmanite triple points is 1990–2035 K which



**Figure 4.** The phase boundaries between akimotoite, majorite, and bridgmanite and the triple point between them. The gray rectangular area represents pressure–temperature conditions expected for the 660-km discontinuity [40–43].

also agrees with a recent LVP study by Ishii *et al.* [36]. Such an excellent agreement despite the fact that our LVP measured temperature using a W/5%Re-W/26%Re thermocouple, Ishii *et al.* [36]’s LVP measured using a Pt/Pt-13%Re thermocouple, and our LHDAC measured using spectroradiometry. This result now enables direct comparison of the experimental results at the  $P-T$  conditions of the deep mantle transition zone for temperature. However, our study reveals that pressure measurement is the main source for the discrepancy between LVP studies and also, between LVP and LHDAC results. This result now calls for efforts in improving pressure measurement techniques in LVP and LHDAC.

**Author Contributions:** conceptualization, S.-H.S. and K.L.; methodology, V.P.; software, S.-H.S.; validation, B.K. and J.D.; formal analysis, B.K. and J.D.; investigation, B.K. and J.D.; resources, S.-H.S.; data curation, B.K., J.D., and S.-H.S.; writing—original draft preparation, B.K., J.D., and S.-H.S.; writing—review and editing, K.L. and S.-H.S.; visualization, B.K. and J.D.; supervision, S.-H.S.; funding acquisition, S.-H.S.

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**Conflicts of Interest:** The authors declare no conflict of interest.

## Abbreviations

The following abbreviations are used in this manuscript:

<i>P</i>	Pressure
<i>T</i>	Temperature
LVP	Large-volume press
LHDAC	Laser-heated diamond anvil cell
Ak	Akimotoite
Mj	Majorite
Bm	Bridgmanite

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**Sample Availability:** Samples of the compounds ..... are available from the authors.