

1 Entropy, Information and Symmetry: Ordered is 2 Symmetrical

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10 **Abstract:** Entropy is usually understood as the quantitative measure of “chaos” or
11 “disorder”. However, the notions of “chaos” and “disorder” are definitely obscure. This
12 leads to numerous misinterpretations of entropy. We propose to see the disorder as an
13 absence of symmetry and to identify “ordering” with symmetrizing of a physical system;
14 in other words, introducing the elements of symmetry into initially disordered physical
15 system. We demonstrate with the binary system of elementary magnets that introducing
16 elements of symmetry necessarily diminishes its entropy. This is true for 1D and 2D
17 systems of elementary magnets. Imposing symmetry does not influence the Landauer
18 Principle valid for the addressed systems. Imposing the symmetry restrictions onto the
19 system built of particles contained within the chamber divided by the permeable
20 partition also diminishes its entropy.

21 **Keywords:** Entropy; symmetry; ordering; Landauer principle.

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23 1. Introduction

24 The notion of entropy, introduced by Rudolf Clausius [1-2], is considered to be one
25 of the important, but also most abstract and least visualisable quantities of physics [3].
26 Rudolf Clausius (1867) coined the term entropy from the Greek word *Entrophein* for
27 transformation and change [4]. The classical thermodynamics macroscopic definition of
28 the entropy change within a reversible process is $\Delta S = \frac{\Delta Q_{rev}}{T}$, where ΔQ_{rev} is the heat
29 flow which took place under the reversible process, and T is the temperature [5]. The
30 statistical definition of the entropy is:

$$31 \quad S = k_B \ln \Omega, \quad (1)$$

32 Where Ω is the multiplicity of a certain state, i.e. the number of different configurations
33 that a system defined by macroscopic variables could assume (or in other words, the
34 number of states accessible to the system) [6-8]. The widespread illustrative interpretation
35 of entropy is “the measure of disorder” in macroscopic systems built from, a large

36 particles [8]. However equating entropy with disorder was criticized recently [4, 9, 10].
 37 Indeed, the multiplicity has a well-defined statistical meaning; whereas, “disorder” is an
 38 obscure notion, deeply rooted in the human psychology, namely in the human perception
 39 of ordered and disordered patterns [11-13]. We propose the hypotheses that “ordering”
 40 may be strictly related to the symmetry, and may be quantified by symmetry. In turn,
 41 “chaos” or disorder is an absence of symmetry. It seems, that the idea that just a symmetry
 42 is an adequate characteristics of ordering for various patterns was suggested in ref. 14
 43 and elaborated further in refs. 15-16. We demonstrate that introducing symmetry into
 44 physical systems necessarily orders them and decreases their entropy. Symmetry
 45 considerations play a key role in modern science [17-19], giving rise to the conservation
 46 laws in physics, and being dominating in quantum theory [20], crystallography [21] and
 47 condensed-matter physics. We show that they are crucial for constituting
 48 thermodynamics and informational properties of physical systems.

49 2. Symmetry and entropy of binary systems

50 Consider first a binary 1D system illustrated in **Figure 1A**. We assume that there are
 51 N separate and distinct states fixed in a space and aligned as shown in **Figure 1A** [8].
 52 Attached to each site is an elementary magnet, which can point only up or down. Thus, a
 53 binary system is created. The total number of arrangements of the N magnets is $\Omega = 2^N$
 54 if no external restrictions are imposed on the system and all of the states are equally
 55 accessible to the system. Thus, the entropy of the system is given by:

$$56 \quad S = k_B N \ln 2 \quad (2)$$

57 If the addressed binary system is in the thermal equilibrium with a thermal
 58 reservoir T the modulus of the term of the Helmholtz Free Energy $|F_S|$ arising from
 59 entropy contributions is supplied by:

$$60 \quad |F_S| = k_B T N \ln 2, \quad (3)$$

61 thus, illustrating the Landauer principle [23-25]. Indeed, if erasing of one bit of
 62 information is performed under re-orientation of the single elementary magnet, the
 63 energy cost of such an erasure equals $k_B T \ln 2$ in a strict accordance with the Landauer
 64 Principle [23-25].

65 Now let us restrict the possible configurations of elementary magnets by introducing
 66 the symmetry axis, shown with the dashed line in **Figure 1B**. After introducing the
 67 symmetry axis only the symmetric configurations of the elementary magnets are available,
 68 as depicted in **Figure 1B**. The total number of arrangements of the N magnets is $\Omega = 2^{\frac{N}{2}}$.
 69 Hence, the entropy of the symmetrized, ordered binary system is given by:

$$70 \quad S = k_B \frac{N}{2} \ln 2 \quad (4)$$

71 It is seen that introducing symmetry orders the binary system, and consequently
72 decreases the entropy. It seems plausible to equate “ordering” with symmetrizing of the
73 system. If this identifying is accepted, the entropy is indeed the measure of disorder, in
74 other words, an absence of symmetry. Obviously the axis of symmetry may be replaced
75 by the center of symmetry [17].

76 The modulus of the entropy related term of the Helmholtz Free Energy $|F_S|$ for the
77 symmetrized system shown in **Figure 1B** equals $|F_S| = k_B T \frac{N}{2} \ln 2$. Thus, the Landauer
78 Principle remains untouched. Indeed, in this case erasing of one bit of information is
79 performed under simultaneous re-orientation of two symmetric elementary magnet, and
80 the energy cost of such an erasure again equals $k_B T \ln 2$ in an accordance with the
81 Landauer Principle [23-25].

82 Consider now the 2D system of elementary magnets depicted in **Figure 1C**.
83 Generalization of the aforementioned approach for 2D systems is straightforward. If the
84 system is not symmetrized (ordered) and all of the arrangement of elementary magnets
85 are equally available, the entropy of the system is still given by
86 $S = k_B N \ln 2$. If the system is symmetrized and it possesses two axes of symmetry, as
87 shown in **Figure 1C**, its entropy is $S = k_B \frac{N}{4} \ln 2$. Additional elements of symmetry in this
88 2D case (i.e. the axes or centers of symmetry) result is the larger decrease in the entropy
89 of system. This supports the idea that ordering (understood as symmetrizing) necessarily
90 decreases the entropy.

91 Consider one more example shown in **Figure 2**. Four molecules of two different kinds
92 (i.e. blue and red ones) are located within a chamber divided equally by the permeable
93 partition. The particles of the same kind are considered as indistinguishable. Consider the
94 states, within which two molecules may be located simultaneously on the one side of the
95 partition. These states may be realized by 3 various arrangements (recall, that the particles
96 of the same sort are identical and indistinguishable). Hence, the entropy of such a system
97 is $S = k_B \ln 3$. Now let us introduce “ordering” in our system. Impose the demand that
98 the particles should be necessarily located symmetrically relative to the axis OO' passing
99 over the partition as depicted in **Figure 2**. Introducing symmetry diminishes the
100 multiplicity to $\Omega = 1$. Consequently, the entropy of the symmetrized system of particles
101 equals to zero. The generalization of the suggested approach to the system of $2N$ particles
102 is straightforward.

103 It is noteworthy, that the notion of the Voronoi entropy fails to quantify the ordering
104 of points constituting the 2D pattern [26-28]. It was demonstrated that the symmetric
105 patterns may be characterized by the values of the Voronoi entropy being very close to
106 those inherent to random ones [29].

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108 Conclusions

109 Misinterpretations of entropy are ubiquitous among scientists and non-scientists [4,
110 10]. The widespread interpretation of entropy is “the measure of disorder” [4, 10].
111 However this understanding of entropy was criticized recently and equating of entropy
112 to “disorder” or “chaos” was considered as a misconception [4, 10]. Indeed, “disorder”
113 and “chaos” are obscure notions, actually rooted in the human psychology [11-13], and
114 hardly could be accurately defined. We demonstrate how this misconception may be
115 avoided. The misunderstanding is eliminated, if the “ordering” is identified with
116 symmetrizing of the system. Thus, “chaos” or “disorder” is understood as an absence of
117 symmetry. We demonstrate that introducing of elements of symmetry orders the system
118 and consequently diminishes its entropy. The idea is illustrated with the binary system
119 built from elementary magnets. 1D and 2D exemplifications of the binary systems are
120 treated. We also demonstrate that introducing of symmetry does not influence the
121 Landauer Principle for the binary systems. We also consider the system built of particles
122 contained within the chamber divided by the permeable partition. Again imposing the
123 symmetry restrictions diminishes the entropy of the system.

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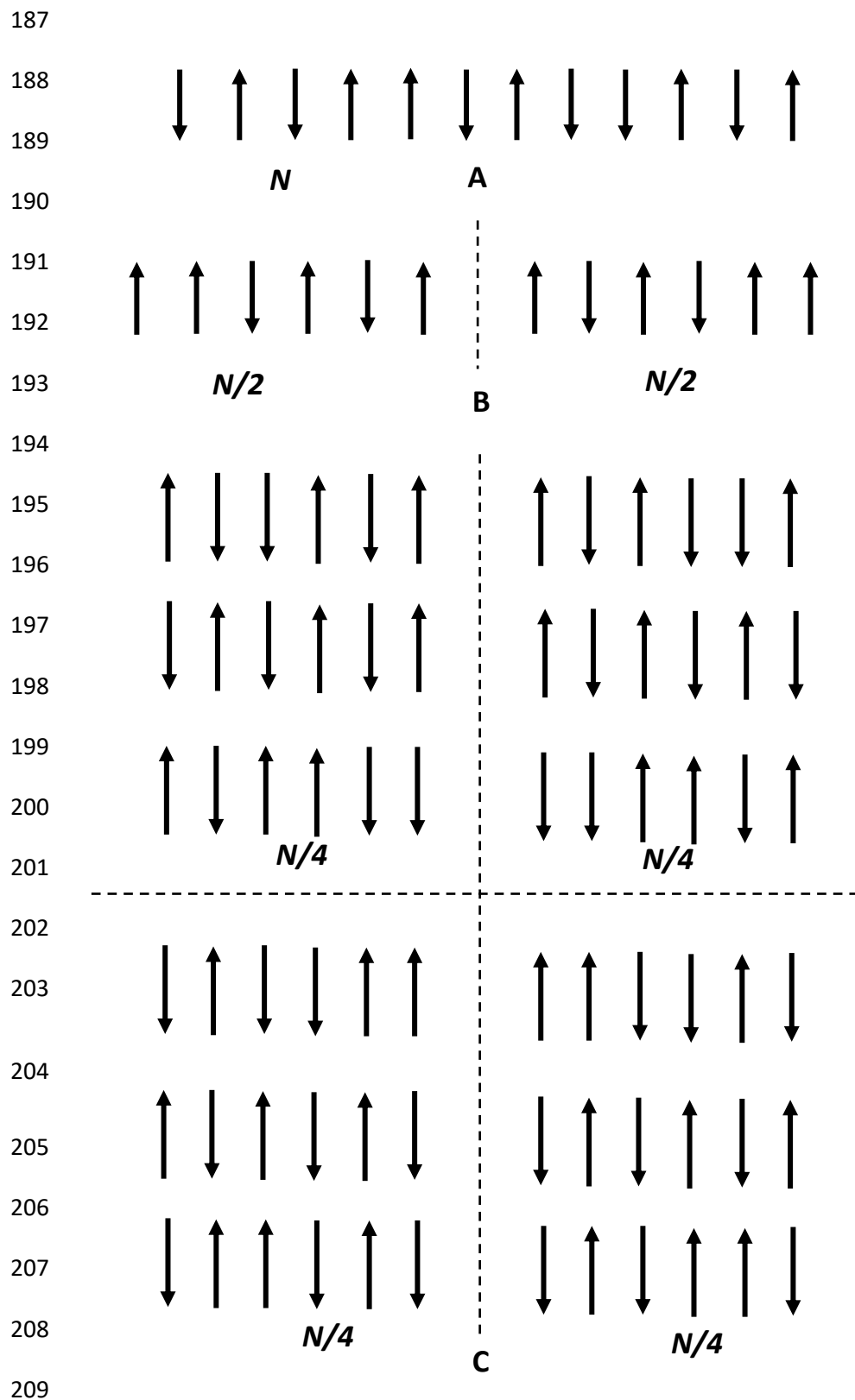
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210 **Figure 1.** A The binary 1D system of N elementary magnets is shown. All of up/down
 211 arrangements of the magnets are available. B The axis of symmetry shown with the
 212 dashed line restricts the number of available configurations of magnets. C 2D binary
 213 system of elementary magnets is shown. Axes of symmetry shown with dashed lines
 214 restrict the available arrangements of the magnets.

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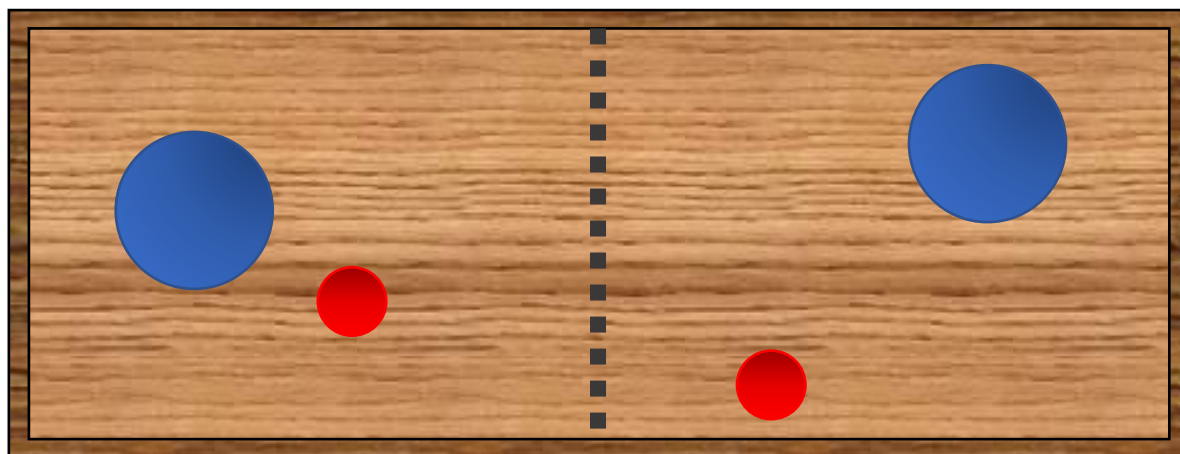
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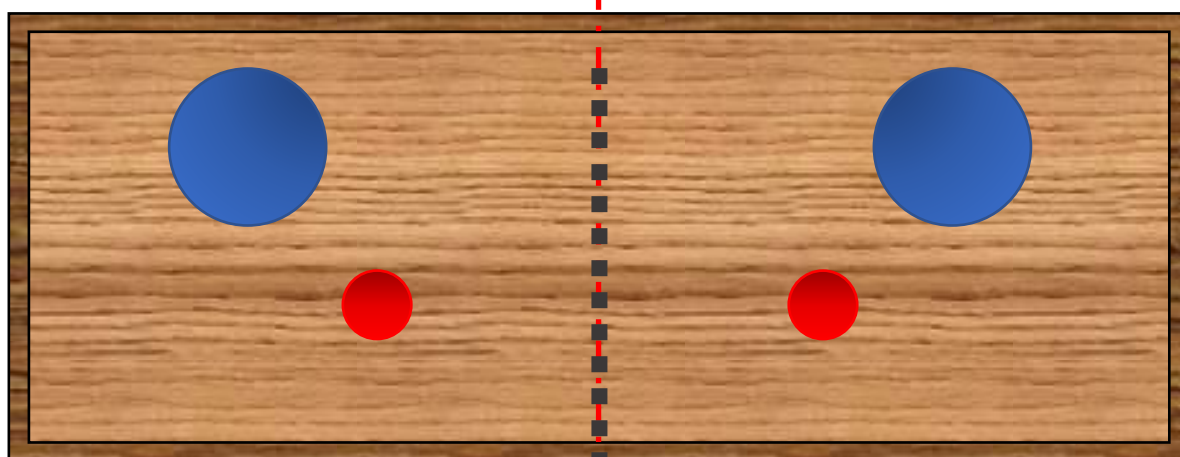
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B O'

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235 **Figure 2. A** Particles of two kinds (blue and red) are located within the chamber divided
 236 equally by the permeable partition. The arrangements at which two particles are
 237 simultaneously located at one side of the partition are permitted. **B** Particles of two
 238 kinds (blue and red) are located within the chamber divided equally by the permeable
 239 partition. Only the arrangements symmetric relative to axis OO' shown with the red
 240 dashed line are permitted.

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