

Article

Features of Structure and Absorption in the Jet-launching Region of M87

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Abstract: M87 is the best available source for studying the AGN jet-launching region. To enrich our knowledge of this region, with quasi-simultaneous observations using VLBA at 22, 43 and 86 GHz, we capture the images of the radio jet in M87 on a scale within several thousand R_s . Based on the images, we analyze the transverse jet structure and obtain the spectral-index distribution in the jet. We find that the absorption is enhanced at the collimation regions and the interruption of the jet. The external medium, which may not be uniform, may have contributed a lot to the absorption in the jet-launching region. Additionally, the temporal morphology of the jet in its launching region may be largely affected by the local, short-lived kink instability growing in itself.

Keywords: galaxies: individual (M87); galaxies: jets; galaxies: active; radio continuum: galaxies

1. Introduction

M87 is a nearby ($D \sim 16.7$ Mpc, [1]) FRI radio galaxy which harbors a central black hole with a mass up to $7.2 \times 10^9 M_{\text{sun}}$ [2]. Profiting from its proximity, the angular resolution of ground-based Very Long Baseline Interferometry (VLBI) now is down to a few times of the angular size of schwarzschild radii (R_s) of the super massive black hole (SMBH) in M87. In recent years, more and more facts of the vicinity of the SMBH in M87 are unveiled by high angular-resolution VLBI observations especially at millimeter wavelength. e.g. the accurate position of SMBH in M87 is constrained by Hada et al.[3]; the central compact radio source in M87 is imaged directly by EHT Collaboration 2019 [4].

M87 is a prototype to study the AGN jet-launching region as well as the vicinity of the SMBH. The structure of the jet in its launching region has been determined with VLBI monitoring, that smears out the variable details by stacking multi-epoch observations at 22, 43 and 86 GHz [5–7]: the limb-brightened jet structure starts at a projected distance down to $7R_s$ to the SMBH, with an apparent opening angle wider than 100° [7,8]; as the limb-brightened jet propagates outward, it has to experience multiple expansions followed with subsequent recollimations [6,7] i.e. "collimation regimes" and finally reaches an equilibrium parabolic expansion in several thousand R_s , and keeps the shape until $10^5 R_s$ [9,10]. Besides the motion along the jet, the jet flow rotates around the axis clockwise, and the toroidal component of magnetic field is found close to the core by VLBI polarimetry [6]. Additionally, the counterjet is proved to be existent and with a limb-brightened structure as well [6].

In this work, we present images of the radio jet in M87 on a scale within 10 mas, i.e. a few thousand R_s , captured by Very Long VLBI Array (VLBA) at 22, 43 and 86 GHz almost simultaneously. Based on these images, we analyze the transverse jet structure and obtain the spectral-index distribution in

Table 1. Observations

Date	Frequency (GHz)	Telescopes	On-source Time (Min)	Beam Size ^a (mas × mas, deg)	I_{peak} ^b (mJy beam ⁻¹)	σ ^c (mJy beam ⁻¹)
2015-03-28	22.72	VLBA, -MK	124	0.93×0.45, 12.6	1054	0.360
	43.12	VLBA, -MK	124	0.56×0.23, 16.7	726	0.224
2015-04-03	86.28	VLBA, -HN, -NL, -SC	289	0.24×0.12, -17.7	501.4	0.217

Notes:

^a The Synthesized beam FWHM size of the images.

^b The peak flux density on the images.

^c The root-mean-square noise of the images.

31 the jet. Further, we make discussions on the structural features and the absorption features and the
 32 association between them. We also discuss the cause of the temporal jet morphology at the observed
 33 epoch. Throughout this work, we use cosmological parameters $H_0 = 67.8 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_M = 0.308$,
 34 $\Omega_\Lambda = 0.692$, then 1 pc corresponds to 11.3 mas, and $R_s \approx 7 \times 10^{-4} \text{ pc} = 7.9 \mu\text{as}$.

35 2. OBSERVATIONS AND DATA REDUCTION

36 In 2015 March 28, we observed M87 at 22 GHz ($\lambda = 1.3\text{cm}$) and 43 GHz ($\lambda = 7\text{mm}$) with 9
 37 antennas of VLBA (Mauna Kea was absent due to the "bad diskpack" problem). To optimize the $u - v$
 38 coverage, the two observing frequencies were cycled in each of the observing runs. Six days later,
 39 we observed M87 at 86 GHz ($\lambda = 3\text{mm}$) with 7 antennas of VLBA (Hancock and Saint Croix have
 40 no 86 GHz receiver while NL was "disc failure"). An automatic reference pointing control was used
 41 for the antenna-pointing. 5-minute scans on the nearby bright source 3C 273 and 3C279 as fringe
 42 finders, delay and bandpass calibrators, were inserted during the observations, every 35 minutes for
 43 22/43 GHz and 75 minutes for 86 GHz. The received signals were sampled with 2-bit quantization
 44 and recorded with an aggregate data rate of 2048 Mbit s⁻¹. See details of observations in Table 1.

45 The initial data calibration was made in AIPS (Astronomical Image Processing System) software
 46 package. For 22 and 43 GHz, we followed the standard procedure of VLBA data reduction. For 86 GHz,
 47 the procedure is slightly different: the global fringe-fitting was first performed on scans of 3C273
 48 and 3C279 with point source models to derive time evolutions of the residual delay, rate, and phase
 49 for each IF separately; the derived residual delay varies slowly with time, so the solutions of 3C 273
 50 and 3C279's delay were used as a first-order approximation for M87's residual delay; then the global
 51 fringe-fitting was performed twice on M87, the first was performed with a point source model, the
 52 solution interval was set to 30 s and the threshold of signal to noise ratio (SNR) was set to 3.0, to avoid
 53 false signals, a tight delay and rate search window (10 nsec and 20 mHz for delay and rate respectively)
 54 were used, the IFs are averaged to increase the SNR by a factor of 2.8; then we produce an initial image,
 55 and perform the second one with the CLEAN models obtained from the initial map to increase the
 56 fringe detection rate (the detection rate was increased slightly for 0.2%).

57 After calibration, multi-source data are split into single-source files and imported into DIFMAP
 58 for self-calibration and imaging. Self-calibration/Imaging loops are performed for multiple iterations
 59 to obtain images with high dynamical ranges. The final image was produced in DIFMAP with natural
 60 weighting.

61 3. Results

62 We show the resulting VLBI images in Figure 1 and Figure 2 as contour plot overlapping with
 63 pseudo-color image. We list the parameters of the images (e.g. synthesized beam size, peak flux
 64 density and root-mean-square noise σ) in Table 1. Some structural features in the jet are apparently
 65 seen, e.g. on the 22 GHz image, the jet shrinks locally between 1-2 mas from the center of the image,
 66 and the northern limb is interrupted between 5 and 7 mas from the center of the image; on the 43 GHz

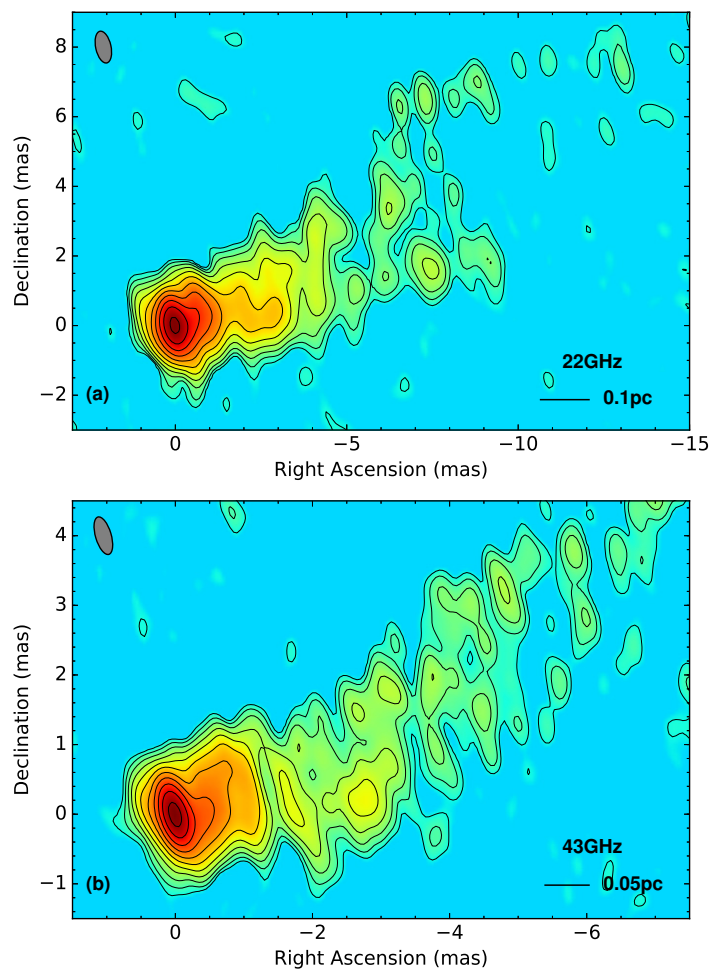


Figure 1. 22 and 43 GHz contour plots overlapping with pseudo-color image of M87, with the synthesized beam shown at the top-left corner of each panel. Panel (a) shows image at 22 GHz with a synthesized beam of 0.93×0.45 mas at a position angle of 12.6° ; contours in this image are from $(-1, 1, 2, 4, 8, \dots) \times 0.86$ mJy beam $^{-1}$. Panel (b) shows image at 43 GHz with a synthesized beam of 0.56×0.23 mas at a position angle of 16.7° ; contours in this image are from $(-1, 1, 2, 4, 8, \dots) \times 0.54$ mJy beam $^{-1}$.

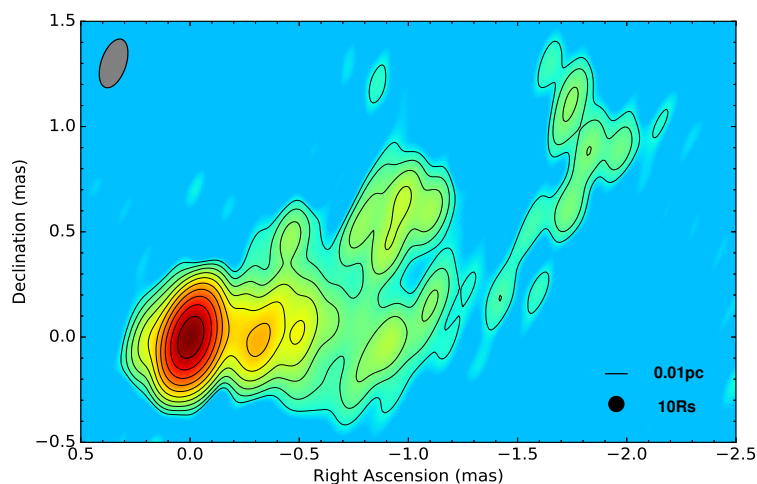


Figure 2. 86 GHz contour plot overlapping with pseudo-color image of M87; contours are from $(-1, 1, 2, 4, 8\dots) \times 0.60 \text{ mJy beam}^{-1}$. The synthesized beam of $0.24 \times 0.12 \text{ mas}$ at a position angle of -17.7° is shown at the top-left corner.

Table 2. Core Component

Frequency (GHz)	Flux density ^a (mJy)	Radius ^b (as)	Theta ^c (deg)	Major Axis ^d (as)	Axial ratio ^e	Phi ^f (deg)	Brightness Temperature ^g (10^{10}K)
22	1605 ± 195	19.9 ± 2.1	-80.8 ± 6.1	479.2 ± 1.1	0.84	43.4	3.1
43	1016 ± 122	3.6 ± 0.4	-142.8 ± 5.8	261.2 ± 0.2	0.62	20.2	2.3
86	580 ± 58	3.6 ± 0.3	-44.7 ± 4.4	73.0 ± 0.2	0.52	31.8	5.0

Notes:

^a Integral flux density of Gaussians.

^b Distance of Gaussians to the center of the images.

^c Position angle of Gaussians referring to the north.

^d Size of Major axis of Gaussians.

^e Axial ratio of Gaussians.

^f Position angle of major axis of Gaussians referring to the north.

^g Brightness temperature of Gaussians.

67 image, the jet bends southward between 1-2 mas from the center of the image; on the 86 GHz image,
 68 the northern limb looks rather ragged, while the shape of its southern counterpart is far better defined;
 69 emission at the northeast of the most luminous region which is probably the hint of the counter jet, is
 70 seen on images at all three frequencies .

71 For the further analysis, we fit the most luminous region on the images as an elliptical Gaussian
 72 with procedure MODELFIT in DIFMAP, and refer it as "core" hereafter. We list the fitting results of
 73 the parameters in Table 2 with their uncertainty estimated as in Lee et al. [11]. We also estimate the
 74 brightness temperature of the core as in Güijosa & Daly [12]. The brightness temperature of the core
 75 is estimated to be at the magnitude of 10^{10}K for all three frequencies. This is the same as given in
 76 Hada et al. [8] and Kim et al. [7], which is nearly one order of magnitude lower than the equipartition
 77 brightness temperature.

78 3.1. Transverse Jet Structure

79 In this part, we use intensity slices to extract details of the jet and to analyze the transverse jet
 80 structure evolving with core distance d . To improve the angular resolution in the direction transverse
 81 to the jet with the least distortion of the images, in advance of slicing, we restored the images with
 82 circular Gaussian beams whose diameter equal to the geometric mean of the length of major and minor

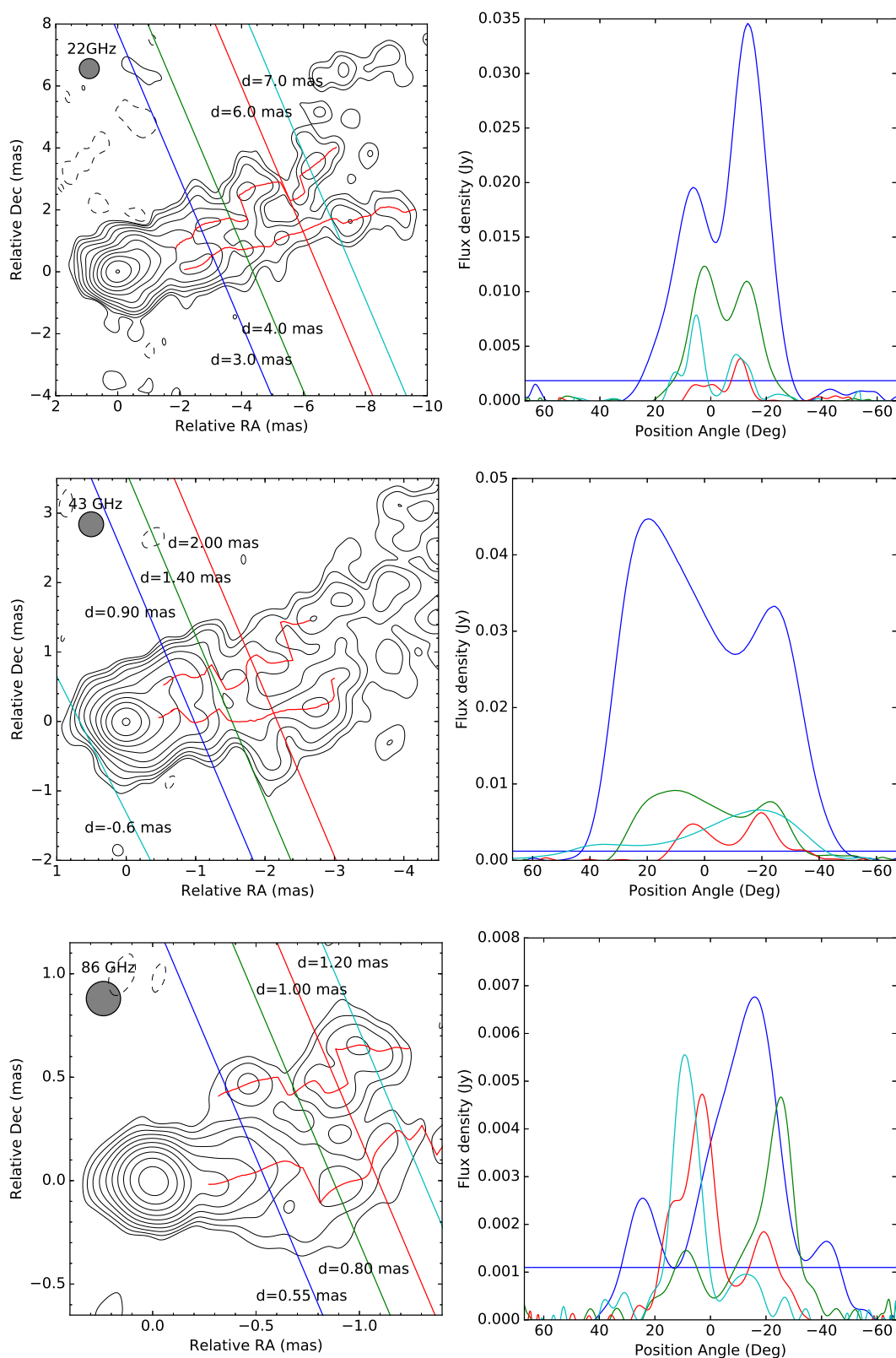


Figure 3. The left panels shows the images restored with circular beams as contours. The restored beams are shown at the top-left corner of each panel. Locations of sample slices are presented as lines superimposed on the contours. The ridgeline is presented as red curves superimposed on the contour plots. The right panels present sample slices, showing the intensity as functions of position angles PAs referring to the jet axis. The northern limb is at positive values, while its southern counterpart is at negative values. The blue horizontal lines represent 5σ of the restored images.

axis of their synthesized beams, which is 0.65, 0.36 and 0.17 mas for the image at 22, 43 and 86 GHz respectively. The restored images are shown as contours in the left panels of Figure 3. Then we slice the restored images in a direction perpendicular to the approximate overall jet axis, which is -67° respect to the north (referred as "jet axis" hereafter). For the 22 GHz image, we made slices between $d = 2.00$ and 10.00 mas, at an interval of 0.05 mas; for the 43 GHz image, we made slices between $d = 0.10$ and 3.00 mas, at an interval of 0.05 mas; for the 86 GHz image, we made slices between $d = 0.10$ and 1.40 mas, at an interval of 0.01 mas. The right panels of Figure 3 present the sample slices, showing the intensity as a function of position angles referring to the jet axis (PAs, the northern limb is at positive values, while its southern counterpart is at negative values). Locations of these sample slices are presented as lines superimposed on the contours in the left panels of Figure 3.

In our intensity slices, at 43 and 86 GHz, the northern limb represents as a well defined hump since $d = 0.65$ and 0.50 mas respectively, while its southern counterpart since $d = 0.45$ and 0.24 mas respectively. The counter jet emission also seems bifurcated on the restored 43 GHz image, and a slice at $d = 0.60$ mas shows this bifurcated structure clearly (see the cyan line in the second row of Figure 3). We also recover the ridgeline of the limb-brightened jet in M87 with the intensity slices. The ridgeline is presented in the left panels of Figure 3 as red curves superimposed on the contour plots. As our result shows, both limbs are tortuous within $d = 1.30$ mas; as going outward, the northern limb is still tortuous, while its southern counterpart becomes smooth.

We present PAs of the intensity peak of both limbs as a function of d in Figure 4. The general trend is both limbs approach to the jet axis gradually as going outward, but there are some local structures seen on both limbs. We also notice that, if the structure measured at higher frequency is moved outward for 0.1 or 0.2 mas, then it looks quite consistent with that measured at the adjacent lower frequency. It does not seem like a consequent of core-shift between frequencies, since the core-shift between adjacent frequencies in our observations is only up to 0.03 mas according to the relation obtained by Hada et al. [3]. We consider it as an effect brought by different angular resolutions of different frequencies. Within $d = 1.30$ mas, at 86 GHz, the northern limb is deflected northward at $d = 1.05$ mas, while at 43 GHz, the corresponding deflection is seen at $d = 1.30$ mas; at 86 GHz, the southern limb veers southward at $d = 0.62$ and 1.18 mas respectively, while at 43 GHz the corresponding changes are seen at $d = 0.75$ and 1.30 mas respectively. Beyond $d = 1.30$ mas, the northern limb becomes fragmented, and its orientation changes abruptly at least at three locations ($d = 1.45, 2.55$ and 4.55 mas), a wiggle is seen between $d = 1.70$ and 2.30 mas, and an interruption is seen around $d = 6.00$ mas; on the contrary, the southern limb evolves smoothly as going outward, only a mild wiggle is seen between $d = 3.00$ and 4.00 mas, and its PA tends to be stable at $\sim -10^\circ$ since $d = 4.00$ mas.

In the following analysis, we move the jet structure measured at 86 GHz outward for 0.15 mas, and the structure measured at 22 GHz inward for 0.25 mas to eliminate the effect brought by different angular resolutions. We define the apparent jet opening angle ψ_{app} as the difference of PAs of two limbs. The panel (a) of Figure 5 shows ψ_{app} as a function of d . The parabolic collimation profile, $\psi_{app} \propto r^{0.58}$, determined by Asada & Nakamura [9] and Hada et al [10], is also presented for reference. The overall jet generally evolves around the parabolic profile, but some local structural features are clearly seen. So we divide the jet into six segments, and each segment includes an expansion and a following recollimation: segment 1 includes a quick expansion between $d = 0.75$ and 0.90 mas, followed by a gradual recollimation between $d = 0.90$ and 1.30 mas; segment 2 includes an expansion between $d = 1.30$ and 1.45 mas, followed by a quick recollimation between $d = 1.45$ and 1.50 mas; segment 3 includes a gradual expansion between $d = 1.50$ and 2.00 mas, followed by a gradual recollimation between $d = 2.00$ and 2.55 mas; segment 4 includes a quick expansion between $d = 2.55$ and 2.60 mas, followed by a slow recollimation between $d = 2.60$ and 3.35 mas; segment 5 includes an expansion between $d = 3.35$ and 3.65 mas, followed by a gradual recollimation between $d = 3.65$ and 4.25 mas; segment 6 includes an expansion between $d = 4.25$ and 4.60 mas, followed by a recollimation beyond $d = 4.60$ mas. We define the apparent offset of jet center δ_{app} as the arithmetic average of the PAs of two jet limbs as well. The panel (b) of Figure 5 shows δ_{app} as a function of d , and the jet axis is presented as

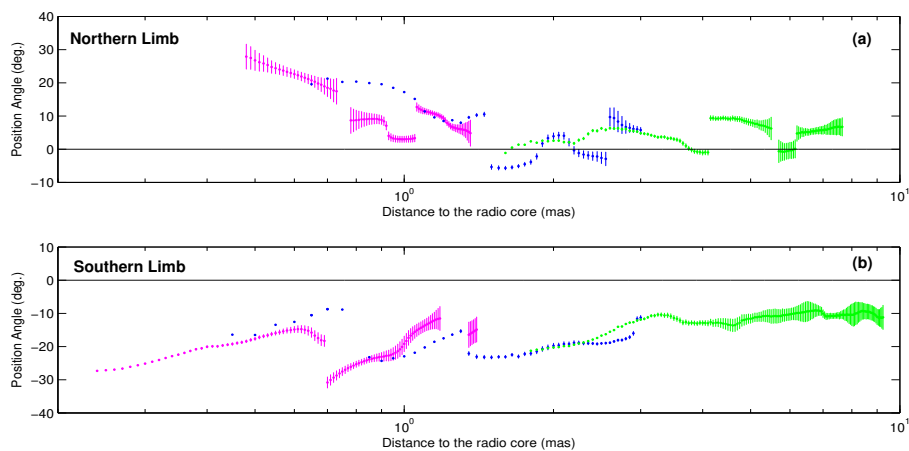


Figure 4. (a) Position angles of the northern limb in M87 as functions of core distance d . (b) Position angles of the southern limb in M87 as functions of core distance d . The colors represent the measured frequencies, magenta for 86 GHz, blue for 43 GHz and green for 22 GHz. The uncertainty of PA is estimated as $\frac{20^\circ}{\text{SNR}}$, in which 20° is a typical width of the jet limb, and SNR is the ratio between the intensity peak and the σ of images. Only peaks with intensity above 5σ are presented to make sure the result is reliable

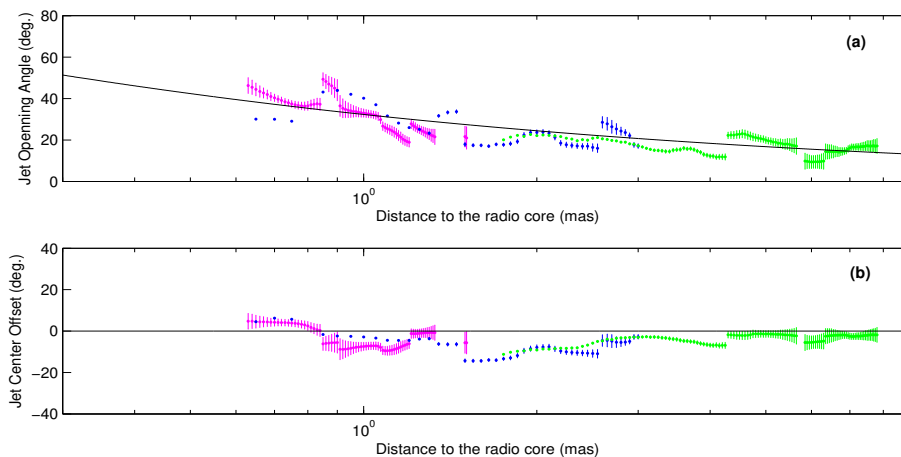


Figure 5. (a): The apparent opening angle ψ_{app} as functions of core distance d . Parabolic collimation profile determined by Asada & Nakamura is shown for reference. (b): The jet center offset δ_{app} as functions of core distance d . The jet axis is shown for reference. The colors represent the measured frequencies: magenta for 86 GHz, blue for 43 GHz and green for 22 GHz. The uncertainties of ψ_{app} and δ_{app} are estimated as $\sqrt{\Delta\text{PA}_{\text{North}}^2 + \Delta\text{PA}_{\text{South}}^2}$. Only peaks with intensity above 5σ are presented to make sure the result is reliable.

133 a black solid line for reference. As the Figure shows, the jet bends southward since $d = 0.75$ mas, the
 134 offset reaches its maximum at $d = 1.45$ mas (where it is most clearly seen on the 43 GHz image), where
 135 the northern limb veers south, and gradually returns to zero around $d = 3.00$ mas.

136 By comparing the behaviors of each limbs with ψ_{app} evolving with d , we find that, structural
 137 features including the expansion and recollimation in segment 2, 3, 4 and 6, result from behaviors
 138 of the northern limb only, while the features in segment 1 and 5 result from behaviors of both limbs.
 139 By further comparing with earlier works, we notice that the gradual recollimations in segment 1
 140 between $d = 0.90$ and 1.30 mas and in segment 5 between $d = 3.65$ and 4.25 mas, spatially associate with
 141 recollimations in the first and second collimation regime (at $d = 1.1$ and 3.7 mas respectively) suggested
 142 by Walker et al. [6]. So they could possibly be the representations of these persistent recollimation
 143 features at our observed epoch.

144 3.2. Spectral-Index Distribution

145 In this part, we present the spectral-index distribution in the jet calculating between 22 and
 146 43 GHz. To align the VLBI images at different frequencies, we employ an algorithm which is first
 147 introduced by Croke & Gabuzda [13]. It is based on the premise that the positions of optically thin
 148 regions to synchrotron radiation are not affected by absorption effects, then images can be aligned by
 149 finding out the maximum of 2D cross-correlation of the optically thin regions. This method is now
 150 commonly used [14–16] and was once employed by Park et al. [17] to obtained the RM map of jet in
 151 M87.

152 The practical steps are listed as follow: (1) First, we produce VLBI images using the fully
 153 self-calibrated dataset. A same $u - v$ range is used to minimize the difference of baseline coverages
 154 between 22 and 43 GHz. Both images are restored with a circular beam with a diameter of 0.65 mas
 155 (the same as used in 3.1). (2) Second, we produce an initial spectral-index map by simply aligning
 156 the geometric center of two images. With the initial map, we roughly determined the optical thick
 157 region which will be masked in the next step. (3) Third, we import the two images into an interactive
 158 Python program, VIMAP [18], based on the algorithm introduced by Croke & Gabuzda [13]. A circular
 159 mask with a diameter of 1 mas is used to cover the core region, and a rectangle box with relative
 160 RA range from 2 to -4 mas and DEC range from -2 to 3 mas is used to choose a optically thin region.
 161 The VIMAP gives the values of 2D-correlation after shifting and find out the shifting corresponding
 162 to the maximum (see Figure 6). In our case, shifting corresponding to the maximum is zero. (4) The
 163 final spectral-index map is produced in AIPS with task COMB. Only regions with intensity $> 5\sigma$
 164 are used to make sure the result is reliable and a noise map showing the distribution of errors is also
 165 produced. The 22-43 GHz spectral-index map superimposed on the intensity contour plots of 22 GHz
 166 are presented in panel (a) of Figure 7 with its noise map in the panel (b).

167 We also produce a 43-86 GHz spectral-index map with the same procedure. Considering there is
 168 a gap of six days between the observations of two frequencies, this spectral-index map is produced
 169 only for reference. In calculating the 2D cross-correlation, a circular mask with a diameter of 0.5 mas
 170 and a rectangle box with a RA range from 1 to -2 mas and a DEC range from -1 to 0.5 mas are used.
 171 The 43-86 GHz spectral-index map with its noise map are presented in the lower panel of Figure 7.

172 On the 22-43 GHz spectral-index map, inverted-spectrum with index $\alpha > 0$ ($S_{obs} \propto \nu^\alpha$) are
 173 detected between $d = 1.30$ and 1.60 mas on the edge of the northern limb, and between $d = 1.30$ and
 174 1.80 mas on the edge of the southern limb. Some less confident inverted-spectrum are seen on the edge
 175 of northern limb around $d = 3.80$ mas. Strongly inverted spectrum with $\alpha > 2$ is detected in the northern
 176 limb between $d = 5.00$ and 7.00 mas where the limb is interrupted. These inverted-spectrum features
 177 indicate absorptions in the jet in M87. Interestingly, on the map, we see a pattern-like distribution along
 178 the jet, the regions with steep spectrum interlace with regions with flatter spectrum (the pattern-like
 179 distribution is also seen on the 43-86 GHz spectral-index map).

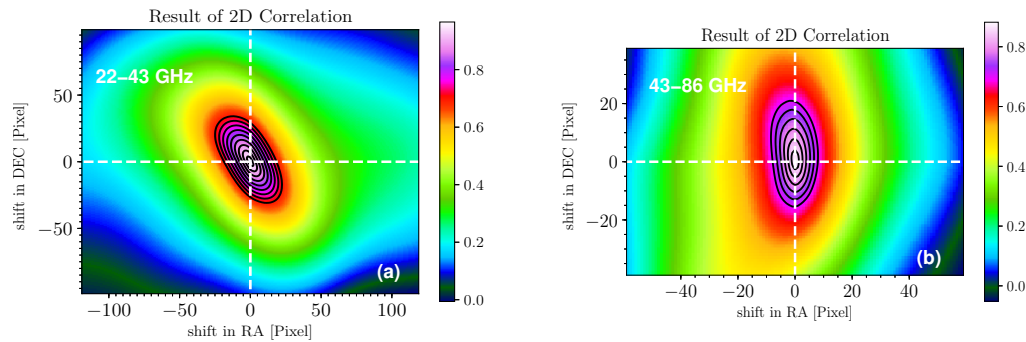


Figure 6. Panel (a) shows 2D cross-correlation of the optically thin regions of 22 and 43 GHz images. Panel (b) 2D cross-correlation of the optically thin regions of 43 and 86 GHz images.

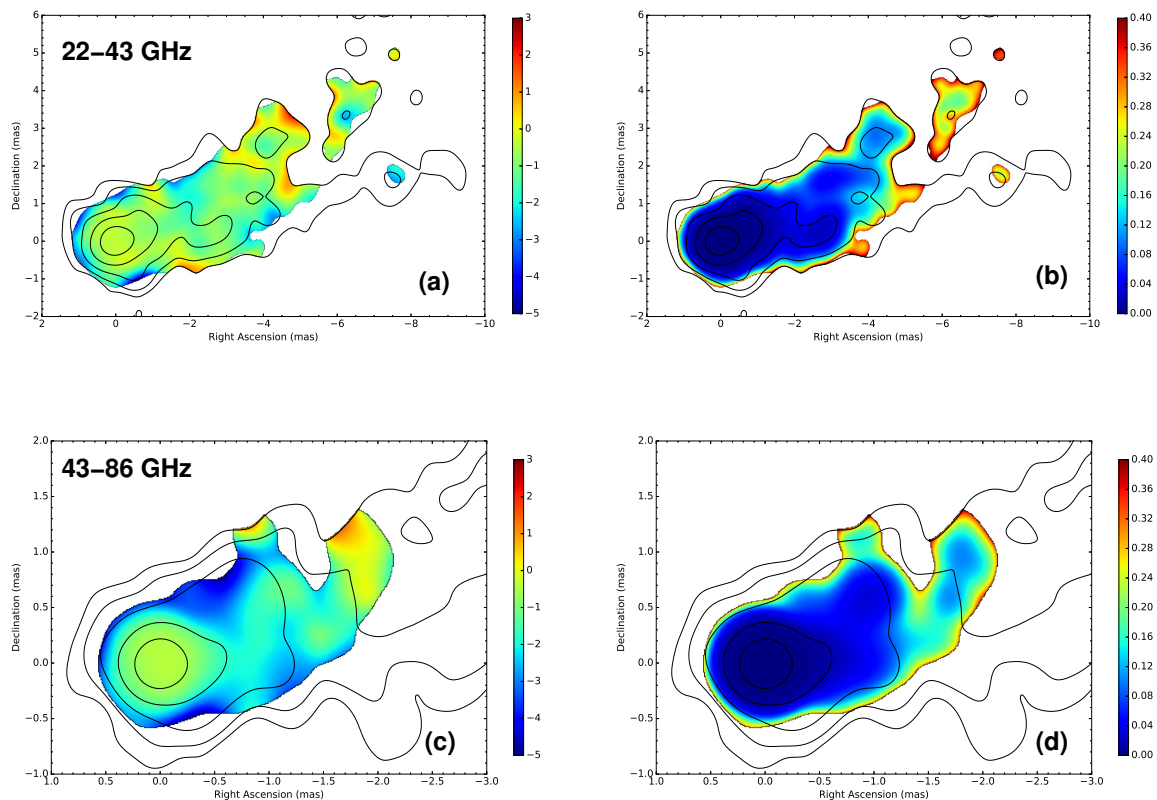


Figure 7. The pseudo-color maps of spectral-index distribution with color bars superimposed on the total intensity contour plots. Panel (a) shows 22-43GHz spectral-index map superimposed on the contour plot of 22 GHz. Panel (c) shows the 43-86 GHz spectral-index map superimposed on the contour plot of 43 GHz. Their noise maps showing distribution of uncertainty are also presented in panel (b) and (d) respectively.

180 4. DISCUSSION

181 4.1. Jet Recollimation with Absorption

182 As we have mentioned in 3.1, recollimations between $d = 0.90$ and 1.30 mas and between $d = 3.65$
 183 and 4.25 mas could possibly be the representations of the persistent recollimations suggested by Walker
 184 et al.[6] at our observed epoch. By comparing our obtained transverse jet structure and the 22-43 GHz
 185 spectral-index map, we find these two recollimations are spatially close to the inverted-spectrum
 186 features. Collimations associating with inverted-spectrum features used to appear on the 22-43 GHz
 187 spectral-index map in Ly et al. [19], at locations between $d = 1$ and 2 mas.

188 The collimations are nonequilibrium behaviors in the AGN jet-launching region before
 189 the hydrodynamic or magnetohydrodynamic equilibrium is achieved, and they reminds us the
 190 reconfinement nodes in the simulations of a jet propagating in the external medium with pressure
 191 decreasing slower than the thermal pressure of jet [20–27]. As to the external medium, since these
 192 features are between 300 - $2000R_s$ from the central engine (consider the viewing angle of the jet in M87
 193 as 17° , as used in Walker et al.[6]), in the hot accretion flow models like ADIOS (see Yuan & Narayan
 194 [28] for an overview), the external medium interacting with an AGN jet on this scale is believed to be
 195 "winds", i.e. the moderately magnetized, non-relativistic un-collimated, extremely hot and generally
 196 fully ionized gas outflows launched from the accretion disk.

197 Stationary recollimation shock at the reconfinement node is a prospection of the simulations
 198 mentioned above. Since the shocks can enhance the magnetic fields locally [25–27], thus increase
 199 the opacity of Synchrotron Self-Absorption, SSA could be a possible mechanism responsible for the
 200 spectral turnover at the recollimations. But it could not be the sole mechanism, since in a SSA-only
 201 case, according to Eq. 2 of Marscher[29], for a region with intensity of $0.02 \text{ Jy Beam}^{-1}$ (a typical value
 202 of intensity of jet in M87 between 300 - $2000R_s$), turnover frequency higher than 43 GHz requires a
 203 local magnetic field of 10^5 G , which is definitely unreasonable. So other mechanisms like free-free
 204 absorption from the external medium must have taken part in as well. The pattern-like spectral-index
 205 distribution along the jet might indicate that the external medium in which the jet propagates is not
 206 uniform.

207 4.2. Jet Interruption with Absorption

208 As we have mentioned in 3, an interruption of the northern limb between 5 and 7 mas from the
 209 center of the image is seen at 22 GHz . On the restored 22 GHz image, the peak of the northern limb
 210 around $d = 6.0$ mas is below 5σ . According to the results of 22 GHz KaVA monitoring and 15 GHz
 211 MOJAVE survey on M87, at many epochs in a period longer than 20 years, an interruption is seen in
 212 the northern limb somewhere between $d = 5$ and 10 mas, thus the jet in M87 shows a single-ridgeline
 213 morphology in this region, e.g. at 22 GHz , 2014 March 2nd and 2014 May 3rd [5]; at 15 GHz , 1998
 214 November 1st, 2004 August 9th, 2010 February 11th, 2019 August 27th, et al.

215 Inverted spectrum with α up to 2.2 is detected around the interruption on our 22 - 43 GHz
 216 spectral-index map, suggesting a strong absorption feature in this area. Interestingly, we find that the
 217 inverted spectrum is also seen at a similar location around $d = 7$ mas on the 22 - 43 GHz spectral-index
 218 map in Ly et al.[19]. Since α is very close to the upper limit of SSA (if we set the cutoff to 3σ , the
 219 α is almost 3), free-free absorption may probably be the main absorption mechanism here. Since
 220 this feature is between 2000 - $4000R_s$ from the central engine, "winds" are still believed to be the main
 221 absorption medium. With our intensity slices, we estimate the opacity of this absorption feature: on the
 222 restored 22 and 43 GHz images, the intensity peak of the northern limb at $d = 6.0$ mas is $1.5 \text{ mJy beam}^{-1}$
 223 and $2.3 \text{ mJy beam}^{-1}$ respectively. If emissions were optically thin here, assuming the optically thin
 224 spectral index is -2 , then the intensity peak in this region should be $8.8 \text{ mJy beam}^{-1}$ at 22 GHz . If
 225 $S_{\text{obs}} = S_{\text{pre}}e^{-\tau}$, then this implies an opacity of $\tau \sim 1.8$ at 22 GHz .

226 4.3. Temporal Features and Short-lived Instability

227 As we have mentioned in 3.1, on our 86 GHz image, the northern limb looks rather ragged.
 228 The analysis of transverse structure shows the northern limb starts at a larger core distance than its
 229 southern counterpart, and it has a tortuous and fragmented structure, which means its orientation
 230 changes several times abruptly, resulting in structural features in the jet segment 2, 3, 4 and 6. These
 231 phenomenon have never been reported by works based on long-term mm-VLBI monitoring, so we see
 232 them as temporal features and representations of the variability of the AGN jet in its launching region.

233 It has been proved by mm-VLBI observations, that the flow in the jet in M87 is rotating clockwise
 234 and with helical magnetic fields [6]. A local, internal kink instability can grow over short timescales in
 235 a jet like this and lead to oscillations [24,30–33]. At the observed epoch, a southward jet oscillation
 236 is found between $d = 0.75$ mas and 3.00 mas, and most clearly seen at $d = 1.45$ mas. So, the
 237 temporal features could possibly be attributed to the observed southward jet oscillation caused
 238 by the kink instability. In the regions nearby the northern edge of the jet in M87 was compressed
 239 by the oscillation, thus local structure of the northern limb was damaged, and finally brought this
 240 fragmented morphology of the northern limb and these temporal structural features at the observed
 241 epoch.

242 5. CONCLUSIONS

243 With the multi-frequency quasi-simultaneous VLBI observations on M87, we have enriched our
 244 knowledge of AGN jet-launching region: the absorption in this region is substantial and could be
 245 enhanced at the collimation regions which represent the nonequilibrium jet behaviors. The external
 246 medium which believed to be "winds", may not be uniform, and greatly contribute to the absorption in
 247 this region. For a rotating jet with helical magnetic field, like the one in M87, its temporal morphology
 248 in the launching region may be largely affected by the local, short-lived kink instability growing in
 249 itself.

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327 **Sample Availability:** Samples of the compounds are available from the authors.