

1 *Review*

# 2 **A Trajectory-Based Method to Explore Reaction** 3 **Mechanisms**

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11 **Abstract:** The method tsscds, recently developed in our group, discovers chemical reaction  
12 mechanisms with minimal human intervention. It employs accelerated molecular dynamics,  
13 spectral graph theory, statistical rate theory and stochastic simulations to uncover chemical reaction  
14 paths and to solve the kinetics at the experimental conditions. In the present review, its application  
15 to solve mechanistic/kinetics problems in different research areas will be presented. Examples will  
16 be given of reactions involved in photodissociation dynamics, mass spectrometry, combustion  
17 chemistry and organometallic catalysis. Some planned improvements will also be described. The  
18 source code can be downloaded from: <http://forge.cesga.es/wiki/g/tsscds/HomePage>

19 **Keywords:** automated algorithm; molecular dynamics; graph theory; statistical rate theory; kinetics  
20 simulations.

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## 22 **1. Introduction**

23 Theoretical studies of reaction mechanisms can greatly benefit nowadays by leveraging the  
24 surge of automated methods developed in the last few years [1-58]. The idea of these new  
25 computational protocols is to substitute human intervention by less error-prone and less tedious  
26 automated algorithms. The automated methodologies range from chemical heuristics to the use of  
27 artificial forces to boost chemical reactions.

28 Our group has contributed with the development of a method called tsscds [43-47], which is  
29 based on accelerated molecular dynamics (MD), as are some others [29, 30]. In our trajectories, the  
30 bonds of the molecule(s) are broken/formed thanks to large amounts of energy placed in each normal  
31 mode/atom of the system [45]. The distinctive feature of tsscds compared to others is the primary  
32 target of the post-processing analysis: the search for transition states (TS) rather than minima.  
33 Additionally, having determined the TS of a given process, its rate can easily be determined using  
34 transition state theory (TST) [59-62]. Thus, finding the relevant TSs on a given potential energy  
35 surface (PES), as our method does, is a subject of fundamental importance in chemistry.

36 In tsscds, after completion of a trajectory, an algorithm named bond breaking/formation search  
37 (BBFS) [45] is employed to select good TS guess structures, which are then optimized using  
38 Eigenvector Following (EF) [63]. In particular, the adjacency matrix, which indicates whether pairs  
39 of atoms form a bond, is monitored along each trajectory to identify the atoms/bonds involved in all  
40 chemical reactions taking place. Then, for each of the selected candidates, a partial optimization is  
41 firstly carried out by freezing the atoms involved in the reaction. The partially-optimized structure is  
42 subsequently subjected to TS optimization using the EF algorithm. The resulting TSs are then  
43 connected with the minima using intrinsic reaction coordinate (IRC) calculations [64]. Finally, tsscds  
44 also features a Kinetic Monte Carlo [65] module that provides the desired kinetic information using  
45 the network of TSs and minima.

46 The method has been successfully employed to study reactions involved in combustion [66, 67],  
47 photolysis [68-70], mass spectrometry [71] and organometallic catalysis [43]. The aim of this review  
48 is to go over several examples where tsscds is employed to either discover new mechanisms and/or  
49 to explain the experiments. For detailed comparisons among different methods for exploring reaction  
50 space, the reader is referred to two recent reviews [58, 72]. Additionally, in the last section, some  
51 planned improvements to enhance the efficiency/efficacy or to expand the scope of tsscds will be  
52 described.

## 53 2. Method

54 The method tsscds has been recently put forward by one of the authors as an automated tool to  
55 discover reaction mechanisms [73, 74]. The basic idea behind tsscds is to run accelerated MD  
56 simulations with the aim to break/form bonds within a few hundred femtoseconds. The simulations  
57 are called “accelerated” because the molecules experience breakage or formation of new bonds very  
58 rapidly thanks to large amounts of vibrational energy placed in each normal mode of the system. In  
59 particular, a range of vibrational energies of ~20-50 kcal/mol per normal mode is initially employed.  
60 However, this range is automatically adjusted to attain at least 60% reactive trajectories in the MD  
61 simulations. Although the default option is to excite all vibrational modes of the system (using  
62 microcanonical normal mode sampling [75]), the user can decide to heat only one part of the system  
63 selecting a few normal modes to be initially excited. The latter option can be particularly useful for  
64 large systems.

65 The trajectory results are then analyzed with a post-processing algorithm (named BBFS), which  
66 identifies geometries with partly formed/broken bonds. Those structures serve as TS candidates in  
67 subsequent transition state optimizations. As detailed below, BBFS is based on the adjacency matrix,  
68 a Graph Theory object that has been employed in other successful automated methods like the one  
69 developed by Zimmerman [16]. Similar ideas have also been recently employed to analyze changes  
70 in conformations occurring in MD simulations [76].

71 Once the TSs are optimized, a reaction network is constructed by computing the intrinsic  
72 reaction coordinates (IRCs) [64] connecting TSs with intermediates [64]. The method employs two  
73 levels of theory: semi-empirical and *ab initio*/DFT. The semi-empirical calculations are performed to  
74 run the MD simulations and to obtain approximate TSs structures, while a higher level of theory is  
75 used to re-optimize the TSs and run IRC calculations. Two different electronic structure programs  
76 are employed: MOPAC2016[77, 78] and Gaussian09[79] for the semi-empirical and *ab initio*/DFT  
77 calculations, respectively.

78 Unlike other automated methods like GRRM [42], our methodology has been employed so far  
79 to study only the ground electronic state. This is in part due to the fact that, currently, the potential  
80 energy and gradients can only be calculated at the semiempirical level of theory.

81 The following is a description of the graph-theoretic tools and kinetic models employed in our  
82 method.

### 83 2.1. Graph Theory

84 A number of graph theoretic tools are employed at various stages of the procedure to find  
85 transition states (TS), screen their structures and construct a reaction network. Specifically, the time  
86 dependence of the adjacency matrix  $\mathbf{A}$  is employed to discriminate TS-like geometries along the  
87 trajectories. The elements of this matrix are defined as:

$$88 \quad a_{ij} = \begin{cases} 1 & \text{if } r_{ij} < r_{ij}^{\text{ref}} \\ 0 & \text{otherwise} \end{cases} \quad (1)$$

89 with  $r_{ij}$  being the distance between atoms  $i$  and  $j$ , and  $r_{ij}^{\text{ref}}$  a reference value that sets the upper  
90 limit for the bond length between the pair; in practice  $r_{ij}^{\text{ref}}$  is taken 20% greater than the sum of the  
91 covalent radii of  $i$  and  $j$  [45]. Thus, for an  $N$ -atom system,  $\mathbf{A}$  is a  $N \times N$  symmetric matrix with  
92 zeros on its diagonal.

93 Additionally, a weighted adjacency matrix  $\mathbf{A}^w$  is also employed in tsscds, whose off-diagonal  
94 elements are defined as:

$$95 \quad a_{ij}^w = \frac{1 - (r_{ij}/r_{ij}^{\text{ref}})^n}{1 - (r_{ij}/r_{ij}^{\text{ref}})^m} \quad (2)$$

96 Values of 6 and 12 have been employed in previous work for  $n$  and  $m$ , respectively [44]. Matrix  $\mathbf{A}^w$   
97 contains information on the 3D geometry of the molecule,[80] and its eigenvalues and eigenvectors  
98 can be employed to construct the so-called SPRINT coordinates [80]. An important property of these  
99 coordinates is their invariance with respect to translation, rotation and permutation of atoms, which  
100 makes them good molecular descriptors in trajectory-based methods. SPRINT coordinates are  
101 employed in tsscds to remove redundant structures.

102 Another matrix employed to determine the number of fragments in the system is the Laplacian,  
103 which is defined as:

$$104 \quad \mathbf{L}^{(w)} = \mathbf{D} - \mathbf{A}^{(w)} \quad (3)$$

105 where  $\mathbf{D}$  is the so-called degree matrix [44], whose elements are defined as:

$$106 \quad d_{ij} = \begin{cases} \text{deg}(v_i) & \text{if } i = j \\ 0 & \text{otherwise} \end{cases} \quad (4)$$

107 where the degree  $\text{deg}(v_i)$  of an atom counts the number of contacts. The superscript  $(w)$  on  $\mathbf{L}$  and  
108  $\mathbf{A}$  indicates that the corresponding matrix can either be weighted or not. For a non-weighted graph,  
109 the lowest eigenvalue of the Laplacian  $\lambda_1$  is always zero, and the total number of zero eigenvalues  
110 determines the number of fragments of the system. For a weighted graph, an upper threshold for  $\lambda_1^w$   
111 is employed to identify fragmented structures [44]. The smallest non-zero eigenvalue is called the  
112 spectral gap, which is a measure of the degree of fragmentation of the structure. Thus, a small value  
113 of the spectral gap is associated with structures presenting non-covalent bonds (like van der Waals  
114 complexes), which are usually of no interest in chemical dynamics and kinetics.

115 The invariance of the SPRINT coordinates upon atom permutation is very important for the  
116 analyses of trajectories, where scrambling of atoms is frequent, as stated above. However, since the  
117 identity of each atom is absent in the adjacency matrix, SPRINT coordinates are identical for two  
118 structures where two non-equivalent atoms swap positions. For that reason, another type of  
119 molecular descriptor, based on a modified (weighted or not) adjacency matrix, is employed in tsscds.  
120 This new matrix, denoted as  $\mathbf{A}_Z^{(w)}$ , contains the atomic numbers  $Z_i$  of the atoms on the diagonal:

$$121 \quad a_{z,ij}^{(w)} = \begin{cases} a_{ij}^{(w)} & \text{if } i \neq j \\ 1 + \frac{Z_i}{10} & \text{if } i = j \end{cases} \quad (5)$$

122 The expression for the diagonal elements is chosen to provide values comparable to the off-  
123 diagonal ones. Most importantly, the eigenvalues of this new matrix are only invariant with respect  
124 to the permutation of like atoms, and it is widely employed in tsscds.

## 125 2.2. Kinetics simulations

126 The kinetics module of tsscds calculates rate constants for all the elementary steps and solves  
127 the set of first-order differential equations that describe the time evolution of all species (usually  
128 known as chemical master equation).

129 The rate constants can either be obtained as a function of temperature or energy. In the former  
130 case, transition state theory is employed [59-62]:

$$131 \quad k(T) = \sigma \frac{k_B T}{h} \left( \frac{RT}{p_0} \right)^{\Delta n} e^{-\frac{\Delta G^\ddagger}{RT}} \quad (6)$$

132 where  $\sigma$  is the reaction path degeneracy,  $T$  is the temperature,  $h$  is Planck's constant,  $\Delta G^\ddagger$  is the  
133 free energy of activation,  $p_0$  is 1 bar and  $\Delta n = 1$  (0) for bimolecular (unimolecular) reactions. The

134 reaction path degeneracy is calculated as  $\sigma = \frac{m^{TS}}{m}$ , where  $m$  and  $m^{TS}$  are the number of optical  
135 isomers of the reactant and transition states, respectively [81].

136 By contrast, the microcanonical rate constants are computed according to RRKM theory [81]:

$$137 \quad k(E) = \sigma \frac{W^{TS}(E)}{h\rho(E)} \quad (7)$$

138 where  $W^{TS}(E)$  is the sum of states at the TS,  $\rho(E)$  is the density of states at the reactant, and  $E$  is  
139 the excitation energy of the system. The sums and densities of states are evaluated by direct count of  
140 the harmonic vibrational states using the Beyer-Swinehart algorithm.

141 Once all state-to-state rates are determined, the chemical master equation is solved using Kinetic  
142 Monte Carlo simulations [65].

### 143 3. Overview of the applications of tsscds

144 The tsscds methodology has been employed in our lab to elucidate reaction mechanisms  
145 involved in photodissociation dynamics, mass spectrometry, combustion and organometallic  
146 catalysis, and in this section, several examples of each type are reviewed.

#### 147 3.1. Photodissociation dynamics

148 The dissociation of molecules can be promoted by using a laser source, which is known as  
149 photodissociation. Although many photodissociations take place in excited states, important  
150 mechanisms may occur in the ground electronic state following internal conversion. One of the  
151 quantities of interest is the product yield, which is usually determined in the experiments. The  
152 understanding of the dissociation channels in organic compounds has greatly benefited from the  
153 interplay between photolysis experiments and computational studies [70, 82-95].

154 In this section, we summarize the results obtained with our automated method for systems that  
155 have also been studied in photodissociation experiments, highlighting the most important  
156 conclusions. In particular, the dissociation channels of formaldehyde, formic acid, vinyl cyanide,  
157 acrolein, acryloyl chloride and methyl cyanofornate were studied with our tsscds methodology.

158 Formaldehyde was employed as a benchmark system to test tsscds. The system had been  
159 previously studied with other automated methods like the scaled hypersphere search [33] and the  
160 global reaction route mapping (GRRM) [35]. The results obtained with all algorithms are comparable,  
161 and the kinetically-relevant stationary points are found using any procedure.

162 The study of the dissociation channels of formic acid ( $\text{CO}_2\text{H}_2$ ) with tsscds revealed the existence  
163 of a new TS for the water-gas shift reaction (WGSR:  $\text{CO} + \text{H}_2\text{O} \rightarrow \text{CO}_2 + \text{H}_2$ ) [45]. By contrast, GRRM  
164 predicted three consecutive steps for the shortest path of the WGSR [35]. The discovery of the new  
165 TS is a consequence of the highly non-IRC [96] nature of the trajectories employed in tsscds [45]; in  
166 other words, IRC jumps are not uncommon events [97]. This exemplifies one of the advantages of  
167 using trajectory-based methods to discover new reactions: we are not restricted to unimolecular  
168 reactions and the only constrain to discover new processes is the molecular formula of the system.  
169 Additionally, the large amounts of vibrational energy put in the normal modes enhances  
170 configurational space sampling in tsscds, which permits the exploration of all types of reactions.

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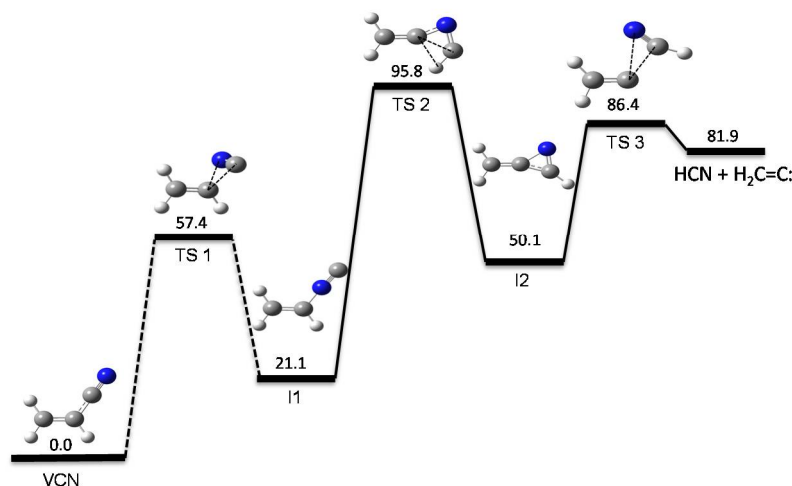
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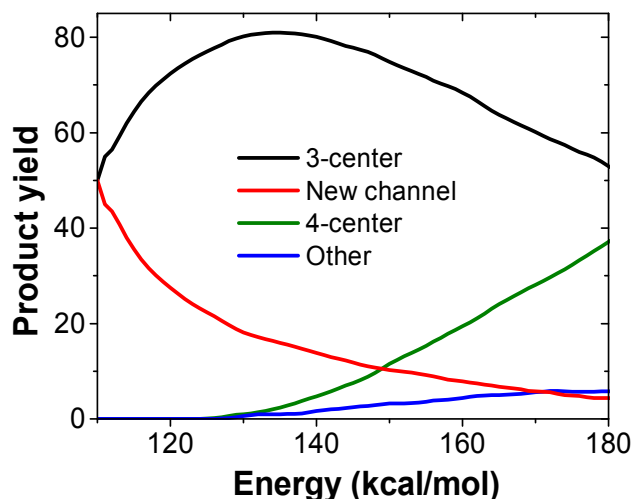
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181 **Figure 1.** New HCN elimination mechanism from VCN obtained with tsscads. The numbers are  
 182 relative energies (including the zero-point vibrational energy) with respect to VCN, calculated at the  
 183 CCSD(T)/6-311++G(3df,3pd)//CCSD/6-311+G(2d,2p) level of theory with the vibrational frequencies  
 184 obtained using CCSD/6-311+G(2d,2p) numerical Hessians.

185 Our automated computational study on the dissociation of vinyl cyanide (VCN) [70] provides a  
 186 HCN/HNC branching ratio in nearly perfect agreement with experiments for an excitation energy of  
 187 148 kcal/mol [98]. Besides the traditional 3-center and 4-center elimination mechanisms found in  
 188 many HX eliminations from CH<sub>2</sub>=CHX systems, a new HCN elimination pathway involving three  
 189 TSs was discovered in the tsscads study. The new mechanism involves three TSs and two  
 190 intermediates and is shown in Figure 1.



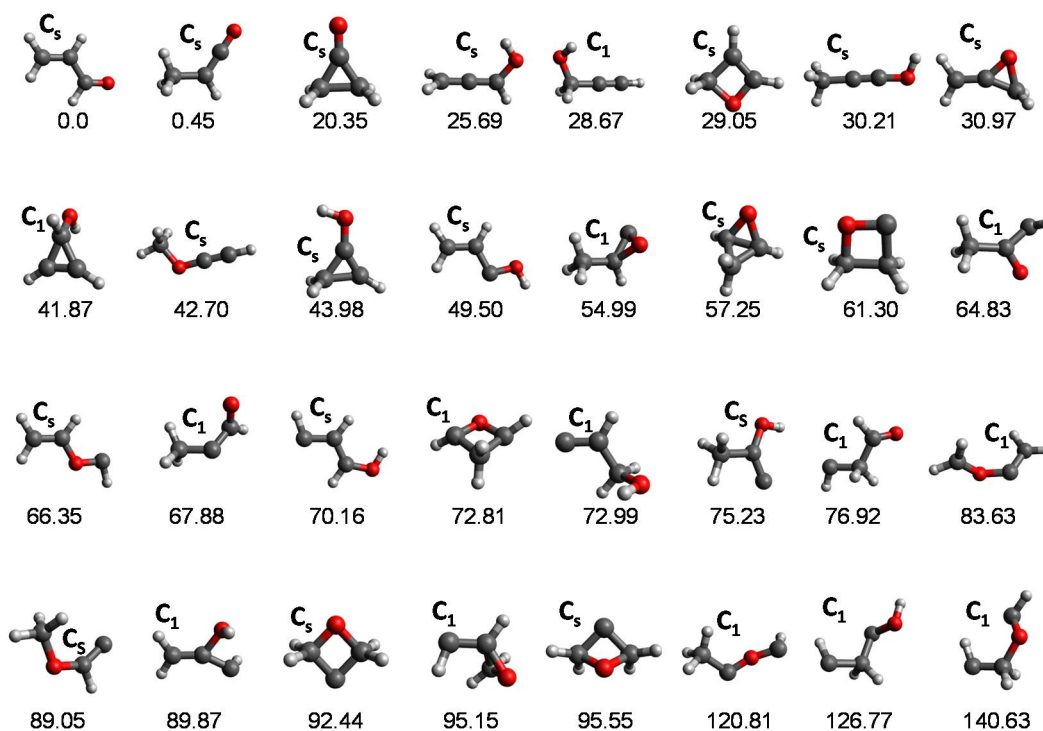
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192 **Figure 2.** Kinetic simulation results of the different HCN elimination channels from VCN.

193 Although alternative routes for HX elimination were also found for other ethylene analogues,  
 194 those pathways involved high-energy TSs and were not competitive with the conventional 3-center  
 195 and 4-center channels. This was the first time a new HX elimination channel competes with the well-  
 196 known 3-center and 4-center processes in the dissociation of CH<sub>2</sub>=CHX species.

197 Figure 2 shows the product yields as a function of excitation energy obtained in our kinetic  
 198 simulations from VCN. As seen in the figure, at low excitation energies (<150 kcal/mol) the new  
 199 channel (red) is more important than the 4-center channel (green) and accounts for half of the HCN  
 200 eliminations when the excitation energy is 110 kcal/mol.





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**Figure 3.** Minima obtained by tsscads for the  $C_3H_4O$  system. The structures are arranged in ascending order of their relative energies (shown at the bottom of each structure), which are obtained at the CCSD(T)/6-311+G(3df,2p)//B3LYP/6-311G(d,p) level of theory. Conformers are not included in the figure and only the lowest lying of each family is displayed.

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The tsscads methodology was also employed to study the dissociation of acrolein (ACRL,  $C_3H_4O$ ), which comprises many different fragmentation channels involving more than 250 transition states and 66 minima [44]. This system was studied with an enhanced procedure (now fully integrated in the method) consisting in the initialization of the MD simulations from multiple minima. In this new procedure the method works in an iterative manner. In the first iteration all MD simulations start from a starting structure, but once some TSs and intermediates are located, subsequent iterations utilize not only the starting equilibrium structure but also the newly generated intermediates to initialize the MD simulations. Compare to a single-minimum initialization, the use of multiple minima to start the dynamics ensures a better sampling of the PES of the system.

The potential energy surface of the  $C_3H_4O$  system is very complex and the 32 equilibrium structures (not including conformers) shown in Figure 3 were found with tsscads, with ACRL being the global minimum. To exemplify the importance of automated reaction discovery methods, we compare our results with those obtained by Chin et al. [99], who manually located equilibrium structures and TSs. Using the same levels of theory as in our study, Chin et al. only found 6 of the 66 minima obtained with tsscads. Most importantly, the relative product abundances obtained with tsscads at 148 kcal/mol (the energy corresponding to the experimental wavelength of 193 nm) are much closer to the experimental results than the computational results of Chin et al., as seen in Table 1.

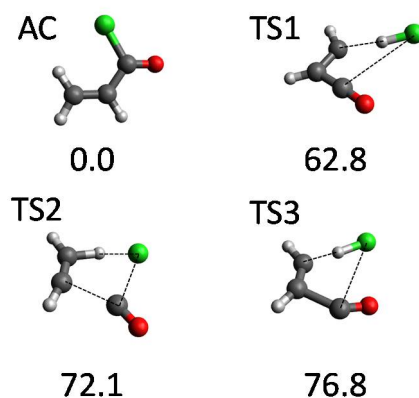
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**Table 1.** Relative product abundances obtained by different computational studies and experiment in the photodissociation of ACRL at 193 nm.

Channel	Chin et al. [99]	tsscads	Exp [100]
H <sub>2</sub> O	0.01	0.03	0.07
CH <sub>2</sub> O	0.65	0.20	0.07
H <sub>2</sub>	0.09	0.19	0.00
CO	1.00	1.00	1.00
H <sub>2</sub> +CO+HCCH	6.82	1.49	1.10

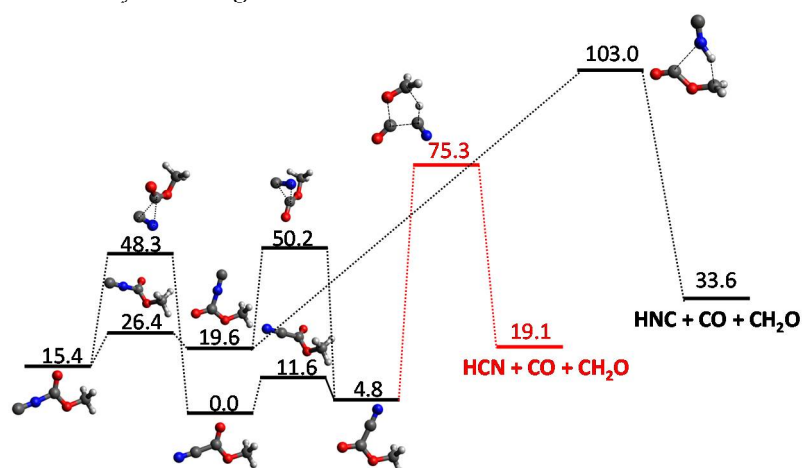
225 Another system that attracted our attention was acryloyl chloride (AC). Overall, around 700  
 226 stationary points were found using our tsscds strategy. Of all possible dissociation channels from AC,  
 227 experiments focus on the HCl dissociations. The use of our automated procedure led to the discovery  
 228 of the three new HCl dissociation TSs [69] displayed in Figure 4; the figure also shows the AC  
 229 equilibrium structure. The highest-energy TSs (TS2 and TS3) correspond to three-body dissociations  
 230 leading to acetylene, carbon monoxide and hydrogen chloride, and they only become important at  
 231 high excitation energies. By contrast, HCl elimination over TS1 is predominant at the experimental  
 232 conditions (148 kcal/mol) [101], showing again that tsscds is capable of finding competitive pathways.



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234 **Figure 4.** Structure of AC minimum and the three new TSs found with tsscds for the HCl elimination  
 235 from AC. Numbers are relative energies in kcal/mol (including the zero-point vibrational energy)  
 236 with respect to AC, calculated at the CCSD(T)/6-311+G(3df,2p)//B3LYP/6-311+G(2d,2p) level of  
 237 theory.

238 Finally, with the aim of exploring possible sources of HCN and HNC in astrophysical  
 239 environments, the dissociation channels of methyl cyanoformate (MCF) were probed with tsscds,  
 240 excited state calculations and photolysis experiments [68]. In particular, time-resolved infrared  
 241 spectroscopy measurements indicate that both HCN and HNC are formed after the 193-nm  
 242 photolysis of MCF [68]. The excited state calculations suggest that most of the dissociations take place  
 243 in the  $S_2$  excited state leading to  $\text{CH}_3\text{O} + \text{NCCO}$  via a Norrish type I reaction, in agreement with  
 244 experiment. However, our calculations are also consistent with cascading internal conversion from  
 245  $S_2$  to produce vibrationally excited ground state MCF.



246

247 **Figure 5.** Relevant HCN and HNC pathways in the ground-state PES of methyl cyanoformate for an  
 248 excitation energy of 148 kcal/mol. Relative energies (in kcal mol<sup>-1</sup>) include ZPE contributions and  
 249 were obtained by CCSD(T)/6-311++G(3df,3pd)//MP2/6-311+G(2d,2p) calculations.

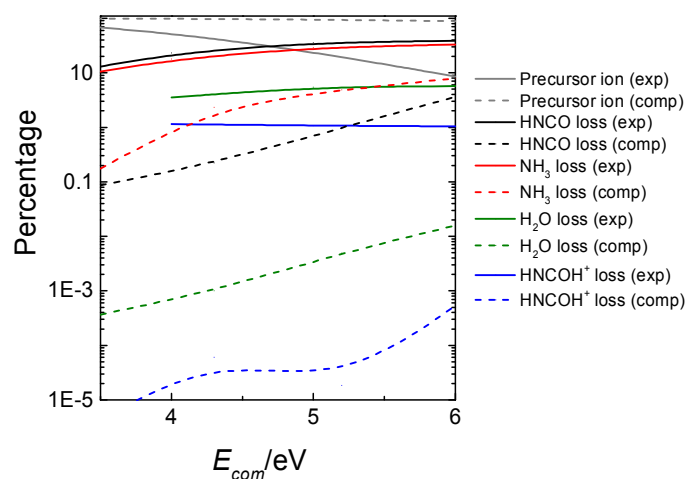
250 To study the dissociation of vibrationally excited MCF molecules in the  $S_0$  electronic state, tsscds  
 251 was employed. Our approach assumes that, after the internal conversion process, intramolecular  
 252 vibrational redistribution is fast enough to ensure RRKM behavior. With the tsscds procedure several  
 253 HNC and HCN mechanisms are found, and Figure 5 shows the kinetically-relevant ones at 148  
 254 kcal/mol. The kinetic simulations predict a HNC/HCN branching ratio of 0.01, which is in  
 255 semiquantitative agreement with that determined in the experiments ( $\approx 0.07$ ). The work provides  
 256 further insights into the intriguing observation of overabundance of HNC in astrophysical  
 257 environments.

### 258 3.2. Mass spectrometry

259 The prediction of mass spectra remains much of a challenge for the community of computational  
 260 chemists. The common computational approaches employed for such endeavor include statistical  
 261 rate theory calculations, MD simulations and electronic structure calculations [102-116]. Our  
 262 automated method is very useful in this regard and can easily be coupled with MD simulations of  
 263 collisions to generate theoretically-based mass spectra as described below.

264 In particular, tsscds was employed to simulate mass spectrometry (MS) experiments of  
 265 protonated uracil,  $[\text{uracil}]H^+$ . Our computational results indicate that the decomposition of  $[\text{uracil}]H^+$   
 266 involves more than one thousand stationary points and 751 elementary reactions [71]. Branching  
 267 ratios for the different fragmentation channels can be automatically obtained from tsscds. However,  
 268 these fractions are a function of the ion's internal energy and cannot be directly compared with MS  
 269 experiments, where the collision energy in the center-of-mass framework ( $E_{com}$ ) is employed instead.  
 270 For that reason the tsscds results were combined with collisional dynamics simulations [71], which  
 271 provide the fraction of  $E_{com}$  transferred to the ion's internal energy.

272 The resulting computationally-predicted product abundances (dashed lines) are compared in  
 273 Figure 6 with the experimental ones (solid lines). As seen in the figure, for the predominant  
 274 dissociation channels, the computationally-predicted product abundances are in qualitative  
 275 agreement with experiment, and formation of HNCO (black),  $NH_3$  (red),  $H_2O$  (green) and  $HNCOH^+$   
 276 (blue) are the major channels. Discrepancies with experiment can be attributed to the possible  
 277 existence of well-known non-statistical behavior in many collision-induced dissociations [103, 117],  
 278 which cannot be captured with our statistical model.



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 280 **Figure 6.** Experimental (exp) and calculated (comp) intensities of precursor and fragment ions  
 281 produced in the fragmentation of protonated uracil.

### 282 3.3. Combustion chemistry

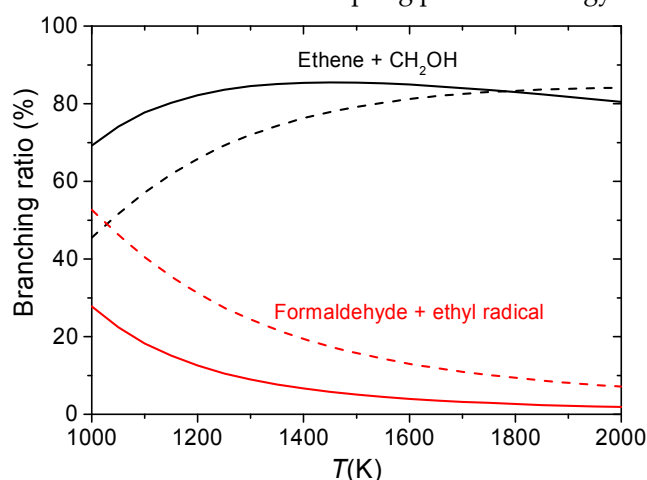
283 Modeling the combustion reactions of oxygenated fuels is of great interest due to their potential  
 284 use as alternatives to conventional petroleum-based fuels. To investigate combustion mechanisms, it  
 285 is important to use kinetic models and perform computer simulations as a complement to



286 experimental determinations, due to the tremendous complexity of these chemical processes. In  
287 general, different approximations are employed in combustion simulations to handle the complicated  
288 mechanisms. One of these simplifications consist of considering only the lowest energy rotamers of  
289 the involved species, which can lead to large errors in the calculation of rate coefficients.

290 In a recent paper, our group analyzed the influence of multiple conformers and paths in the  
291 evaluation of rate constants and relative abundances of products formed in the thermal  
292 decomposition of 1-propanol radicals using different methodologies including tsscds [66].  
293 Specifically, the most relevant pathways reported in the literature [118-124] are obtained with tsscds,  
294 except for the barrierless dissociation leading to propene + OH, since the present version of tsscds  
295 cannot handle this type of reactions. Of significance, an important number of reactant and TS  
296 conformers, not described in the previous studies, are obtained with tsscds.

297 A conformational reaction channel (CRC) was defined in our study [66] as the group of all the  
298 paths that connect the conformers of a given reactant with the corresponding TS conformers. The  
299 influence of these conformers on the rate constants and branchings ratios was investigated in detail  
300 [66]. To study such influence, the output of tsscds (families of CRCs) was fed into a computer program  
301 to treat torsional anharmonicity named Q2DTOR (also developed in our group) [125]. The results  
302 obtained with tsscds and Q2DTOR were finally employed to calculate variational transition state  
303 theory (VTST) [126-128] rate constants for all the CRCs. The multipath (MP) approach within VTST  
304 was employed [128-132], where the rate constant of a given CRC is calculated using contributions  
305 from all the conformers and paths. For comparison purposes the simplest one-well (1W) approach is  
306 also considered; in the 1W method only the most stable conformers of reactant and TS are considered.  
307 As seen in Figure 7, the product abundances obtained in the temperature range 1000-2000 K  
308 are greatly influenced by the selected approach (MP vs 1W), particularly for the major products: ethene  
309 + CH<sub>2</sub>OH and formaldehyde + ethyl radical [66]. Our results show the importance of using automated  
310 codes for discovering reaction mechanisms and sampling potential energy surfaces.



311

312 **Figure 7.** Branching ratios obtained in the kinetics simulations starting from one of the isomers of 1-  
313 propanol (only the two major mechanisms are shown). The solid and dashed lines correspond to the  
314 MP and 1W results, respectively.

315 Very recently, Fenard et al. developed a detailed kinetic model of the low-temperature oxidation  
316 of tetrahydrofuran (THF) based on theoretically-calculated rate constants [67]. The reaction pathways  
317 involved in these processes were probed with our automated software tsscds [67] using CBS-QB3 as  
318 the choice for the high-level of electronic structure. The rate constants were determined using TST  
319 with a tunneling correction using an Eckart potential.

320 The predictions from the model developed by Fenard et al. are overall in good agreement with  
321 the different experimental measurements. Namely, it reproduces ignition delay times obtained in a  
322 rapid-compression machine and in a shock tube, as well as numerous product mole fractions  
323 measured in a jet-stirred reactor.

## 324 3.4. Organometallic catalysis

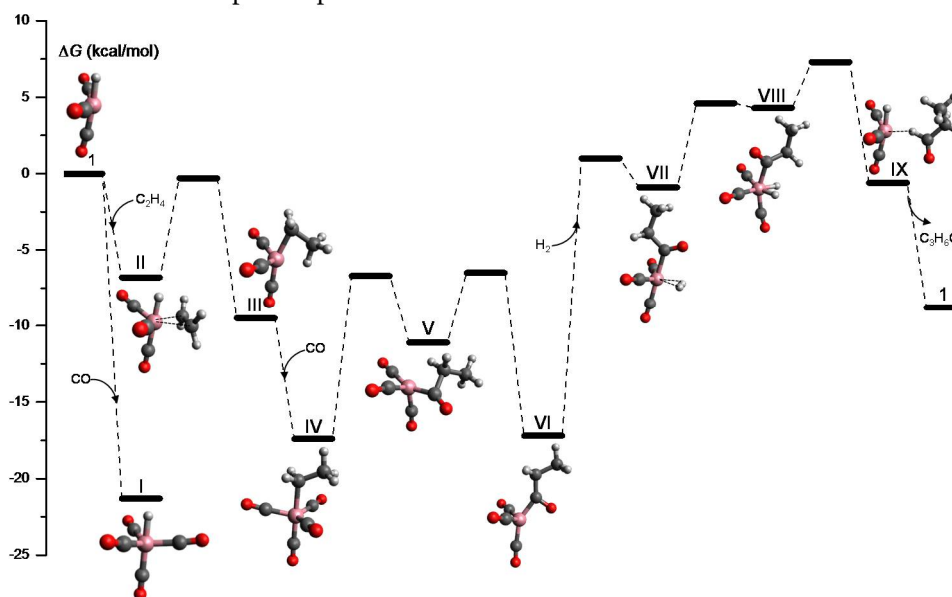
325 Computational studies of organometallic catalysis are becoming increasingly more important  
326 because they can help elucidate reaction mechanisms, characterize catalytic intermediates,  
327 supplement experimental studies, and also because of their predictive power [127, 133-136].

328 However, the traditional workflow of most computational studies consists of using chemical  
329 intuition in the design of reaction routes and construction of guess TS structures. In recent years the  
330 appearance of powerful automated computational methods to study homogenous catalysis [27, 43,  
331 137-139] very much eased the tedious work of manual searches.

332 To exemplify the use of tsscds in organometallic catalysis, the cobalt-catalyzed  
333 hydroformylation of ethylene was chosen [43]. Very briefly, the first step in our computational study  
334 was to generate all combinations of the catalyst  $\text{Co}(\text{CO})_3$  with any of the starting materials (CO,  $\text{H}_2$   
335 and ethylene), which in this case amounts to eight. Each of these combinations has fewer atoms than  
336 the overall system and they were named sub-systems in our original paper [43]. Standard tsscds is  
337 then run in each sub-system to build the reaction networks. Finally, the full reaction network is  
338 obtained after merging the individual results for each sub-system.

339 Figure 8 shows the tsscds-calculated free energy profile for the formation of propanal ( $\text{C}_3\text{H}_6\text{O}$ ),  
340 which is the predominant channel; the level of theory employed was B3LYP/6-31G(d,p). As pointed  
341 out in the original paper, this is not the best electronic structure method for this system and it was  
342 only selected for comparison purposes. Additionally, we simulated the reactivity in the gas phase  
343 because, for this system, solvent effects are unimportant [43, 136].

344 The mechanism shown in Figure 8 was obtained in an automated manner, and agrees with the  
345 one predict by Heck and Breslow in the 1960s [140] and with more recent mechanistic studies [136].  
346 This is a very interesting result as we needed to make no assumptions in our automated calculations.  
347 Additionally, our method predicts that hydrogenation of ethylene is a side reaction that can be  
348 predominant under low CO partial pressures.



349

350 **Figure 8.** Free energy profile for the Co-catalyzed hydroformylation of ethylene obtained in our tsscds  
351 study using DFT calculations [136].

352 With the full reaction network constructed, the kinetics simulation module of tsscds can provide  
353 a rate law for the hydroformylation reaction when a range of different initial conditions for each  
354 species is employed. The kinetics calculations consist of transition state theory calculations [59-62] for  
355 the thermal rate constants at 423 K, and subsequent Monte Carlo simulations using different initial  
356 conditions of the reactants. Table 2 shows the orders of the catalyst and starting materials for the  
357 hydroformylation reaction obtained experimentally [141], with tsscds [43], using a kinetic model

358 based on highly-accurate electronic structure calculations by Harvey and co-workers [136], and  
 359 obtained from another automated method by Habershon [27].

360 As seen in Table 2, tsscads agrees rather well with experiment and with the results obtained by  
 361 Harvey and co-workers [136]. Moreover, tsscads agrees much better with experiment than the other  
 362 automated method does [27] (last column of Table 2), despite the fact that both employ the same  
 363 alkene, initial conditions for the kinetics, and level of theory for the electronic structure calculations.

364 **Table 2.** Orders of the hydroformylation reaction with respect to the catalyst and starting materials.

Species	Exp [141]	tsscads [43]	Harvey [136]	Habershon [27]
H <sub>2</sub>	0.6	0.4	0.5	1
CO	<0	<0	<0	<0
catalyst	0.8	0.5	0.5	1
alkene	1	1	1	0.55

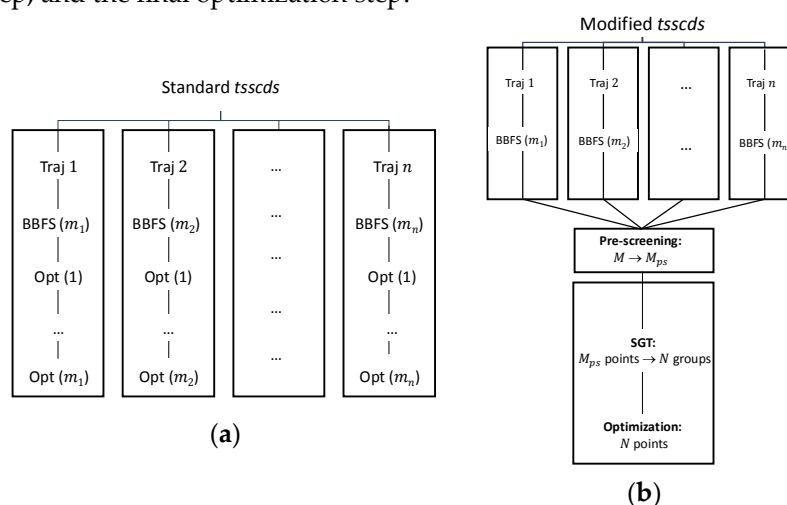
#### 365 4. Improvements

366 In this section we describe some improvements we plan to implement in the near future. They  
 367 include: the use of Spectral Graph Theory, implementation of knowledge-based methods,  
 368 implementation of rare event acceleration MD simulations, interface with other electronic structure  
 369 codes, reparametrization of semiempirical methods, and the study of condensed phase reactions.

##### 370 4.1. Use of Spectral Graph Theory to minimize the number of Hessian calculations

371 In standard tsscads, every single structure obtained after the BBFS analysis is subjected to TS  
 372 optimization [45]. As seen in Figure 9(a), for a trajectory  $i$ , BBFS selects  $m_i$  TS candidates, which  
 373 results in  $M = \sum_{i=1}^n m_i$  optimizations, where  $n$  is the total number of trajectories. On the one hand,  
 374 these  $M$  optimizations are the most CPU-time consuming step of the procedure as they involve  
 375 Hessian calculations, while the integration of the trajectories only requires gradients. On the other  
 376 hand, a number of those optimizations are repeated. This is so because trajectories visit more often  
 377 those areas of the configurational space around the kinetically most relevant TSs, leading to multiple  
 378 optimizations of those structures.

379 The workflow of the enhanced procedure is shown in Figure 9(b). Briefly, instead of carrying  
 380 out the optimizations for every single structure selected by the BBFS algorithm (as in the original  
 381 implementation), the new procedure will run the MD simulations and store at once the  $M$  structures  
 382 for the analysis of all trajectory data. This analysis will consist of a pre-screening, a Spectral Graph  
 383 Theory (SGT) step, and the final optimization step.



384 **Figure 9.** (a) Original tsscads showcasing an example with  $n$  different trajectories resulting in a total  
 385 number of  $M = \sum_{i=1}^n m_i$  optimizations. (b) Modified tsscads showcasing the same example as in panel  
 386 (a) with  $n$  different trajectories resulting in a total number of  $N$  optimizations.

387 Upon completion of the MD simulations, a pre-screening of the  $M$  structures will be performed  
388 based on the eigenvalues of the Laplacian matrix [44]. As pointed out above, the lowest eigenvalues  
389 of this matrix indicate the degree of fragmentation of the molecular system. We aim here to discard  
390 highly fragmented structures, i.e., TSs connecting van der Waals complexes, usually of negligible  
391 relevance in a kinetics study. In the SGT step the remaining points will be partitioned into  $N$  groups  
392 according to the eigenvalues of a TS adjacency matrix, calculated as the average of the reactant and  
393 product adjacency matrices. Finally, we will select the closest point (geometry) to the centroid of each  
394 cluster for optimization. With this new scheme the gain in efficiency can easily be quantified as the  
395 reduction in the number of optimizations from  $M$  to  $N$ .

#### 396 4.2. Implementation of knowledge-based mechanism generators

397 A number of reaction discovery methods are based on the so-called chemical heuristics [23, 48-  
398 50]. In these methods, molecules are typically represented as graphs, in pretty much the same way as  
399 in tsscds. Then, by applying transformations, based on encoded rules or principles inspired by  
400 organic chemistry, to the reactant molecule graph, reactions, products and intermediates can readily  
401 be obtained. Compared to MD-based methods, heuristic-based methods are less CPU-time  
402 demanding.

403 Our idea will be to combine a heuristic-based bias in the MD simulations alongside with our  
404 BBFS algorithm to obtain TSs. In particular, having defined a set of encoded rules based on chemical  
405 knowledge, every single MD simulation will suffer a different bias, aimed to trigger a particular  
406 reaction mechanism. In this way, the problem of multiple optimizations of a given TS mentioned  
407 above would be minimized, if not completely avoided. The bias (analytical) potentials will be added  
408 on top of the semiempirical potential to steer the dynamics towards a particular intermediate or  
409 product.

#### 410 4.3. Implementation of rare-event acceleration MD methods

411 One of the shortcomings of tsscds is the fact that chemical reactions are triggered by using very  
412 high energies in the MD simulations. While this approach was successfully employed to tackle  
413 different problems, it is biased towards the entropically favored reaction pathways. To alleviate this  
414 drawback of the method we propose to replace the current MD strategy by the rare-event acceleration  
415 method named Boxed Molecular Dynamics (BXD) [142]. BXD has its roots in work done by one of us  
416 and D. Shalashilin more than a decade ago [143]. It introduces several reflective barriers in the phase  
417 space of a MD trajectory along a particular collective variable. Those boundaries are employed to  
418 push the dynamics along the collective variable into regions of phase space which would be rarely  
419 sampled in an unbiased trajectory. However, the use of BXD constrains in configuration space suffers  
420 from the same "entropic" bias mentioned above.

421 A generalization of BXD has been very recently put forward by D. R. Glowacki and co-workers  
422 [144]. They show that the BXD bias can also be introduced along the potential energy ( $E$ ) of the  
423 system, which is referred to as BXDE. By scanning through potential energy "boxes", the energetic  
424 "windows" at which different chemical reaction channels switch on or off can be identified. The  
425 software design of tsscds is highly modular, which means that interfacing it with BXDE only requires  
426 little effort, like the need of compatible input/output geometry formats in both codes and the use of  
427 extra keywords in tsscds.

#### 428 4.4. Interface with other electronic structure codes

429 At present tsscds has been only interfaced with the MOPAC2016 [145] and the G09 [79] electronic  
430 structure packages. The MD simulation employs gradients calculated at the semiempirical level of  
431 theory, and the optimization step is carried out at both the semiempirical level with MOPAC2016  
432 and using higher levels (ab initio/DFT) with G09. Although we plan to reparametrize a semiempirical  
433 Hamiltonian for use in organometallic catalysis (see below), we do not want to be limited to this low-

434 level electronic structure calculations. Therefore, we will use the ASE package[146] to interface tsscads  
435 with other electronic structure codes like NWChem [147] or ORCA [148].

#### 436 4.5. Reparametrization of semiempirical methods

437 The application of the tsscads method relies on the use of semiempirical Hamiltonians for  
438 exploring potential energy surfaces. For this reason, it is important that the semiempirical method  
439 provides a reasonably accurate representation of the system under investigation. Although  
440 significant improvements in these methods have been made over the last years [149], there are still  
441 known limitations, which claim for further developments and more accurate parametrizations. Two  
442 important limitations concern the non-covalent interactions for large systems and ligand dissociation  
443 energies for transition metal complexes. In both cases, the performance of the semiempirical methods  
444 is, in general, quite poor. Our goal is therefore to improve the description of both non-covalent  
445 interactions and transition metal complexes in PM7.

446 Regarding non-covalent interactions, we aim to develop an analytical correction for PM7. To this  
447 end, we will consider a set of small molecules, which are representative of the most important  
448 functional groups. All pairs of molecules will be considered to calculate interaction energies at three  
449 levels of theory: coupled-cluster (CC), DFT and PM7. For every pair, various orientations will be  
450 considered, each one emphasizing a different two-body interaction.

451 Then, sums of two-body Buckingham potentials (supplemented with damping functions for the  
452 dispersion) will be fit to the CC, DFT and PM7 interaction energies using our genetic algorithm  
453 program GAFit [150]. Finally, the resulting potentials  $V_{\text{fit,CC}}$ ,  $V_{\text{fit,DFT}}$  and  $V_{\text{fit,PM7}}$  will be employed to  
454 build corrections  $V_X^{\text{corr}}$  to the PM7 interaction energies:

$$455 \quad V_X^{\text{corr}} = V_{\text{fit,X}} - V_{\text{fit,PM7}} \quad (8)$$

456 where X is either CC or DFT. Whereas the  $V_{\text{DFT}}^{\text{corr}}$  correction term will be employed to validate this  
457 methodology as explained below, the highly-accurate  $V_{\text{CC}}^{\text{corr}}$  correction will be used once the  
458 validation succeeds.

459 The correction will be added to the PM7 energy  $V_{\text{PM7}}$  so that the PM7 Hamiltonian corrected for  
460 non-covalent (*nc*) interactions would read:

$$461 \quad V_{\text{PM7,X}}^{\text{nc}} = V_{\text{PM7}} + V_X^{\text{corr}} \quad (9)$$

462 The strategy of using small representative molecules and sums of two-body functions was  
463 successfully employed in the development of intermolecular potentials for interactions of protonated  
464 peptides and silyl ions with perfluoroalkane self-assembled monolayers [151, 152]. Nevertheless, this  
465 strategy will be validated for the new functional groups by running DFT calculations for large  
466 systems. This will allow us to compare the DFT-calculated energies with those obtained with  
467  $V_{\text{PM7,DFT}}^{\text{nc}}$ .

468 The semiempirical methods, and particularly PM6 and PM7, do not perform well for transition-  
469 metal complexes [153]. Our strategy here will be to reoptimize the PM7 Hamiltonian as in previous  
470 studies of our group (e.g., see ref. [68]). We will select popular transition metals and ligand molecules  
471 used in organometallic catalysis, and will carry out high-level ab initio calculations for our own  
472 benchmark database. To gain flexibility in the parametrizations, we will consider the possibility of  
473 defining "atom types" for the ligand atoms, depending on the functional groups, in much the same  
474 way as that done for the parametrization of the hpCADD NDDO Hamiltonian [154].

#### 475 4.6. Study of condensed phase reactions

476 Our method is not limited to gas phase reactions. Although currently it only handles reactions  
477 in the gas phase, its modular design allows for a smooth adaptation of tsscads to deal with condensed  
478 phase reactions. For instance, to study solvent effects, the easiest way would be to use an implicit  
479 model, which in practice would only entail adding the appropriate keywords to the templates  
480 employed for the different electronic structure programs.



481 On the contrary, if one wants to use explicit solvent molecules, the MD module must be changed  
482 or substituted. At present, the MD module is a modified version of DRC routine in MOPAC2016,  
483 which includes different strategies for enhanced sampling, as detailed in the tutorial of tsscds [47].  
484 To include solvent molecules in the MD simulations, one possibility would be to use CHARMM [155]  
485 or to adapt DRC.

486 Finally, if the interest is a gas surface reactions, VENUS [156] would be the choice to run the MD  
487 simulations because the authors have vast experience using this program.

488  
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498

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