

1 *Invited Review*2

New Reentrant Insulating Phases in Strongly 3 Interacting 2D Systems with Low Disorder

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10 **Abstract:** The apparent metal-insulator transition (MIT) in two-dimension (2D) was discovered by
11 Kravchenko et al. [1] more than two decades ago in strongly interacting 2D electrons residing in a
12 Si-metal-oxide-semiconductor field-effect transistor (Si-MOSFET). Its origin remains unresolved.
13 Recently, low magnetic field reentrant insulating phases (RIPs), which dwell between the zero-field
14 ($B=0$) metallic state and the integer quantum Hall (QH) states where the Landau-level filling factor
15 $\nu > 1$, have been observed in strongly correlated 2D GaAs hole systems with large interaction
16 parameter r_s ($\sim 20-40$) and high purity. A new complex phase diagram was proposed, which includes
17 zero field MIT, low magnetic field RIPs, integer QH states, fractional QH states, high field RIPs and
18 insulating phases (HFIPs) with $\nu < 1$ in which the insulating phases are explained by the formation
19 of Wigner crystal. Furthermore, evidences of new intermediate phases were reported. All contribute
20 to the further understandings of the puzzle. This review article serves the purpose of summarizing
21 those recent experimental findings and theoretical endeavors, to foster future research efforts.

22 **Keywords:** 1; Metal-insulator transitions 2; Electronic transport in interface structures 3; Quantum
23 Hall effects.

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1. Introduction

26 Tremendous knowledge from more than fifty years of research has been accumulated regarding
27 the transport behaviors of electrons in varieties of 2D materials and systems. It has regained strong
28 interests in the field of physics and material science due to recent discoveries of topological insulator
29 [2] and new 2D materials [3].

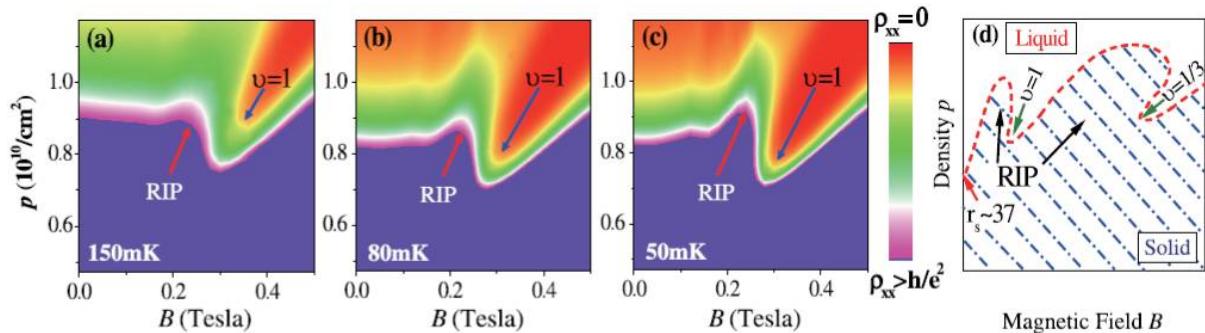
30 With different types and strengths of interactions, many fascinating phenomena can emerge at
31 various temperatures and energy scales, such as Fermi liquid, Wigner crystal [4], integer quantum
32 Hall (QH) effect [5], fractional QH effect [6], etc. The scaling theory of localization for noninteracting
33 2D systems [7] predicts that the electronic states are localized while temperature T approaches zero.
34 In low density and high mobility Si MOSFET, however, a metal-insulator transition (MIT) was
35 observed by Kravchenko et al. [1]. It implies that the strong correlation effects play a key role in the
36 MIT [8, 9], as the value of r_s (the ratio between Coulomb energy and kinetic energy, $r_s \equiv$
37 $1/[a_B^* \sqrt{\pi p}]$, $a_B^* = \hbar^2 \epsilon / m^* e^2$ is the effective Bohr radius) is much greater than one.

38 In this brief review, we will first introduce and focus on the recently discovered low magnetic
39 field RIPs [10-13], with the new phase diagram which connects the zero field MIT with RIPs, integer
40 QH states, fractional QH states, and high magnetic field insulating phases [14]. Resistivity,
41 capacitance and inductance measurements exploring the RIPs and phase diagram are presented and
42 discussed. Second, experimental observations of possible Wigner crystal melting and new
43 intermediate phases are shown and examined. Finally, relevant theoretical models that try to explain
44 the MIT are briefly discussed.

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2. New Reentrant Insulating Phases at Low Magnetic Fields

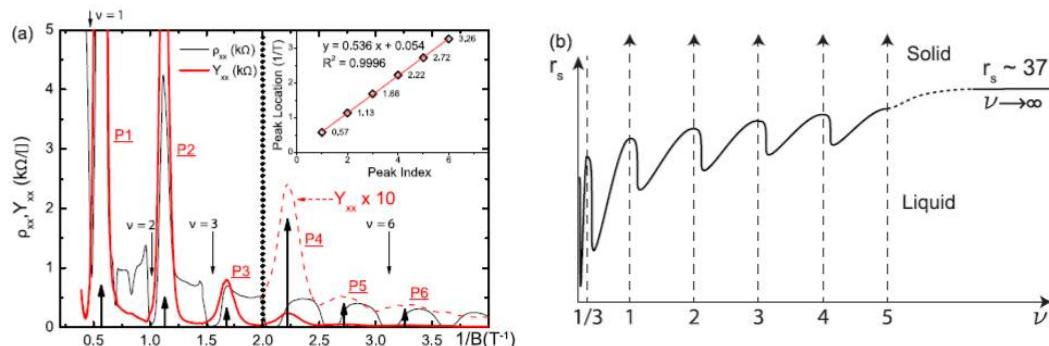
46 Recently, a new reentrant insulating phase at Landau-level filling factor $\nu > 1$ was first
 47 discovered by Qiu et. al. [10, 11] in ultra clean GaAs quantum well samples. Figure 1 is taken from
 48 reference 10, which shows the observed phase diagram at 150, 80 and 50 mK.



49 **Figure 1.** Longitudinal resistivity map plotted in the density (p) – magnetic field (B) plane at $T = 150$
 50 (a), 80 (b) and 50 (c) mK for dilute 2D holes in a 10nm wide high mobility GaAs quantum well. The
 51 RIP phase ($\rho_{xx} > h/e^2$) becomes more prominent at lower temperature. (d) The proposed phase
 52 diagram includes MIT, RIP, integer and fractional QH states. Figure is adapted from reference 10.

53 The formations of WC were widely believed to be the origin of the HFIP and RIPs where the
 54 Landau-level filling factor $\nu < 1$ [14]. Experiment evidences were given from transport,
 55 thermodynamic compressibility, and microwave transmission measurements [14]. Most recent study
 56 [15] observed tunneling resonance feature that was attributed to the vibrations of WC. However,
 57 theoretically, liquid-WC transition was predicted to appear both at $B = 0$ and $\nu > 1$ around $r_s \sim 30
 58 [16, 17], which had not been shown in experimental measurements until the study in Reference 10.
 59 Therefore, the observation of the new RIP in reference 10 provided a consistent phase diagram with
 60 theory [17], which further suggests that the zero-field MIT is a liquid-WC transition.$

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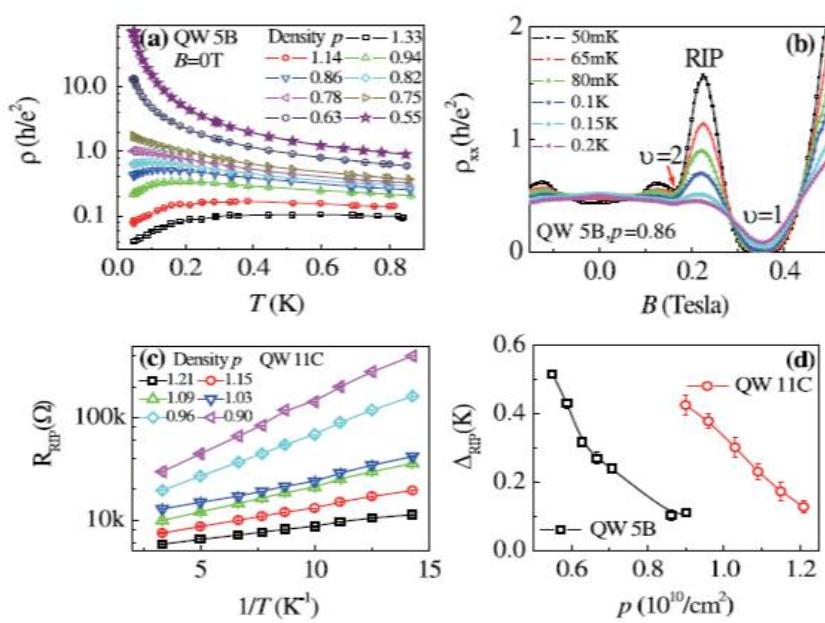
63 **Figure 2.** (a) Observation of multiple RIPs between Landau-level filling factors 1, 2, 3 and 4 in a dilute
 64 2D hole system in GaAs with $r_s \sim 20$. (b) The modified phase diagram with multiple RIPs. Figure is
 65 taken from reference 12.

66 Later, the existence of low field RIPs ($\nu > 1$) was confirmed by Knighton *et al.* [12, 13], see Figure
 67 2 adopted from reference 12. Moreover, three RIPs were observed, which are seen between $\nu = 1, 2, 3$
 68 and 4. This observation implies a phase diagram that the 2D WC or low field RIP can alternate with
 69 integer QH states to take lower energy down to low fields where $\nu > 1$.

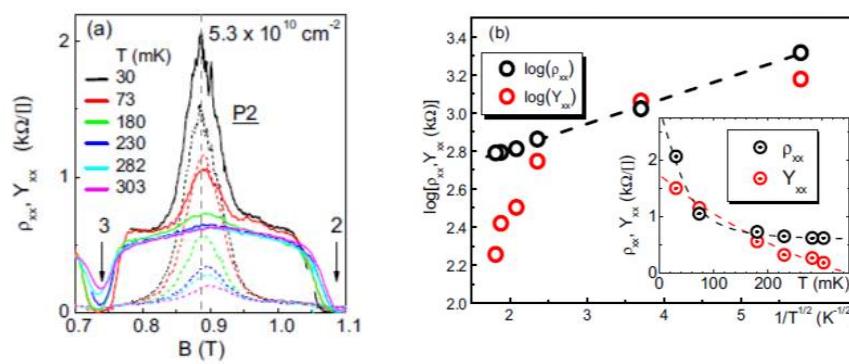
70 The reported resistance, capacitance and inductance characteristics of the RIPs are discussed
 71 below.

72 *2.1. Resistivity*

73 Both Qiu *et al.* [10-11] and Knighton *et al.* [12-13] reported clearly insulating behavior
 74 ($d \rho_{xx}/dT < 0$) of the RIPs that are far stronger than the SdH oscillation amplitude. The resistivity
 75 values of the RIP between $v=1$ and 2 are all above the quantum resistivity h/e^2 . However, there are
 76 some discrepancies in the amplitude and temperature dependency, which could be caused by the
 77 differences in sample structures and qualities. Figure 3 presents the result from reference 10, which
 78 used thermal activation model ($R_{RIP} \propto \exp(\Delta_{RIP}/2T)$) to fit the temperature dependence of the RIP
 79 peak resistance. On the other hand, Knighton *et al.* [12-13] fitted the data to Efros-Shklovskii variable-
 80 range-hopping model $\rho_{xx} \propto \exp(\sqrt{T^*/T})$. Figure 4 shows their data and fitting results. It is worth to
 81 point out that the resistivity value of the RIPs at $v > 2$ in Knighton *et al.*'s work was much lower than
 82 h/e^2 (Figure 4a) and the data do not have high dynamic range to warrant a reliable fitting to the model.
 83 Therefore, the mechanism (thermal activation vs variable-range-hopping etc.) of temperature
 84 dependent resistivity in the RIPs remains to be seen.
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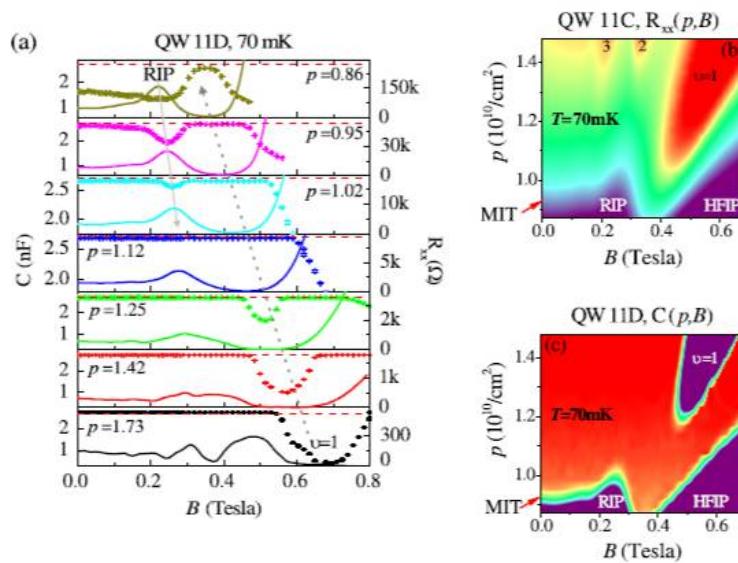
86
 87 **Figure 3.** (a) 2D MIT at zero magnetic field in a dilute 2D hole system in 10 nm wide GaAs quantum
 88 well, $p_c \sim 0.8 \times 10^{10}/cm^2$ (b) $\rho_{xx}(B)$ with $p = 0.86 \times 10^{10}/cm^2$ (c) Arrhenius plot of the RIP peak
 89 resistance at various hole densities. (d) Fitted thermal activation gap. From Reference 10.



98 **Figure 4.** (a) Real (solid) and imaginary (dashed) parts of the magnetoresistance for RIP between $\nu =$
 99 2 and 3 for 2D holes in a 20nm wide GaAs quantum well. (b) Real and imaginary components of the
 100 magnetoresistance plotted in a semi-log scale. From the reference 12.

101 **2.2. Capacitance Measurement**

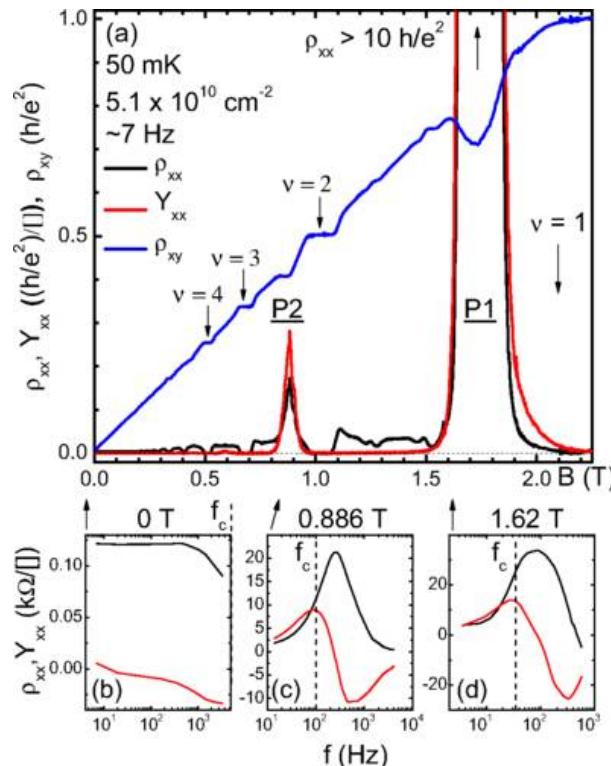
102 Qiu *et al.* [10-11] also reported the thermodynamic compressibility study through capacitance
 103 measurement in the RIP, shown in Figure 5. It is found that the RIP tends to be incompressible, like
 104 the zero-field insulating phase and the HFIP. The phase diagram from capacitance measurement
 105 matches well with that from the transport measurement. The observation suggests the possibility of
 106 the same origin for them, i.e. the liquid-WC transition. We also point out that in the capacitance
 107 measurement of a 2D heterostructure with a single-gate configuration, the geometric gate capacitance
 108 usually dominates over the 2D system's quantum capacitance which is related to the compressibility.
 109 It is only when the 2D system becomes very incompressible and the 2D system's quantum capacitance
 110 is greatly reduced, the measured gate capacitance starts to show observable deviation from the
 111 geometric capacitance. More intricate methods such as the penetration field measurements in 2D
 112 structures with both top and bottom gate [18] will be desirable to further study the compressibility
 113 of RIPs and their connection or competition with the zero field MIT or integer QHs to a better
 114 precision.



124 **Figure 5.** (a) Capacitance (symbol) and resistance (line) Vs. perpendicular magnetic field at several
 125 hole densities in a 10nm wide GaAs quantum well system. (b) The phase diagram viewed in the
 126 longitudinal resistance map at 70mK. (c) The phase diagram viewed in the capacitance map at 70mK.
 127 From Reference 10.

128 **2.3. Inductance**

129 Inductance of 2D systems is a rarely studied topic. Knighton *et al.* [12-13] reported the inherent
 130 inductive behavior of the RIPs, illustrated in Figure 6. The inductance behaviors are different between
 131 the RIPs at low magnetic field with $\nu > 1$ and the RIPs at high magnetic field where $\nu < 1$. There has
 132 not been much (or any) theoretical studies on the inductive behavior of correlated 2D electron
 133 systems. Thus the understanding of these anomalous inductance observations is limited and awaits
 134 further theoretical investigation.



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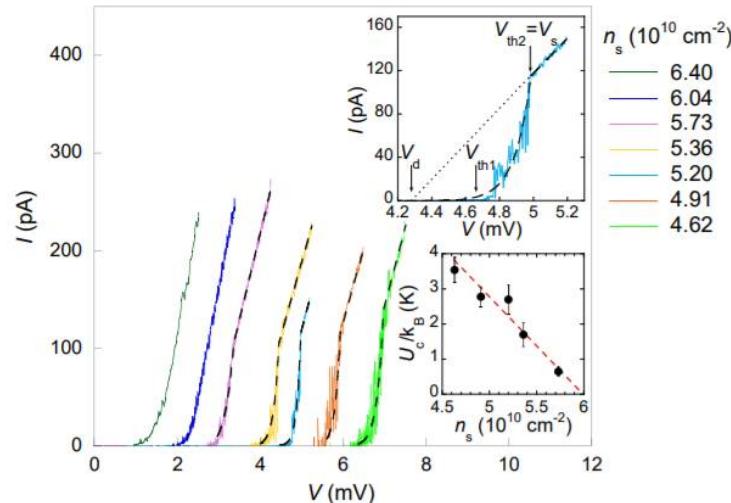
136 **Figure 6.** (a) Magnetoresistance of 2D holes in GaAs showing two RIs at $v>1$ (labeled as P1 and P2).
 137 (b-d) The longitudinal resistivity ρ_{xx} and inductive signal Y_{xx} vs. frequency at various magnetic fields,
 138 showing the clear inductive effect in the RIs. Figure taken from Reference 12.

139 In summary, all the measured properties of the RIs indicate that the RIs share the same origin
 140 with the zero-field insulating phase and the RIs at higher magnetic field $v < 1$. The phase diagram
 141 of very high mobility 2D p-GaAs systems with high r_s in the density-perpendicular magnetic field
 142 plane is consistent with the liquid-WC transition phase diagram in *clean* 2D systems.

143 3. Possible transport evidences for intermediate phases and Wigner crystal melting

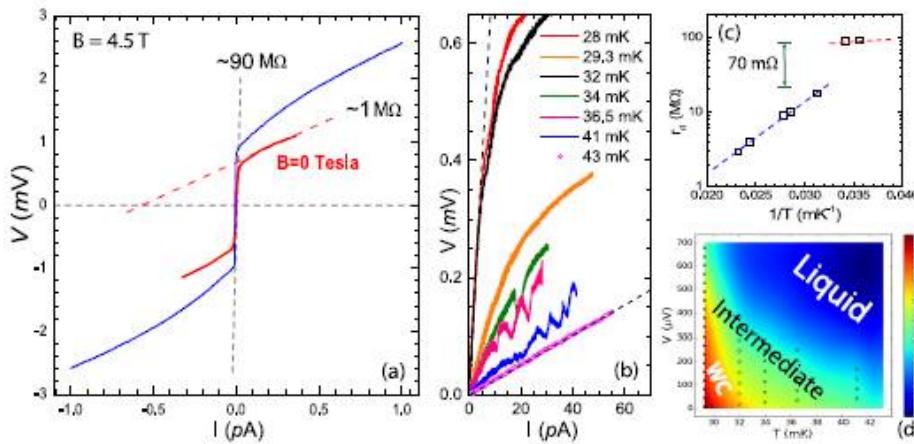
144 Since the discovery of the RIs at lower magnetic field $v > 1$ implies a liquid to Wigner crystal
 145 transition in clean 2D systems with high r_s , then an important question arises: what type of phase
 146 transition is the experimentally observed 2D metallic liquid to insulator (whereas the insulator is
 147 either the zero field insulator or the RIs at low magnetic fields). Given the well-known theoretical
 148 results that there is no long range order in 2D solid and the possible existence of various phases
 149 intermediate between 2D WC and liquid [19-32], are there any experimental evidences for
 150 intermediate phases when the WC melts into liquid?

151 Evidences of intermediate phases have been presented by several groups [33-35]. Qiu *et al.*
 152 studied how the low field RIP is suppressed by either increasing the temperature or carrier density
 153 and found evidences for the 2D holes in p-GaAs transforming into a mixture state of incipient RIP
 154 and metallic liquid [36]. Other researchers took a different approach and used an increasing voltage
 155 (or electric field) applied to the sample to probe the breakdown of a presumed 2D WC deep in the
 156 insulating state where the resistivity is very high. In such state where the WC is presumably well
 157 formed, Brussarski *et al.* [34] and Knighton *et al.* [35] found non-linear I-V curves at low temperatures,
 158 shown in Figure 7 and Figure 8. Different threshold voltages were seen, suggesting two-stage phase
 159 transition instead of direct phase transition, although different explanations were given by the two
 160 groups. More experimental data and theoretical models are needed to further elucidate the situation.



161

162 **Figure 7.** Voltage-current characteristics of 2D electrons in Si-MOSFET in a possible WC state. V-I
 163 curves of different electron densities are shown to show the depinning of WC. Figure taken from
 164 Reference 34.



165

166 **Figure 8.** Voltage-current characteristics of 2D holes in GaAs in a possible WC state to show the two-
 167 stage melting of WC. (a) dc IV at 28 mK. (b) IVs at different temperatures. (c) Temperature dependence of
 168 $r_d(T)|_{V \rightarrow 0}$ (d) Suggested phase diagram. From Reference 35.

169 4. Discussion and Outlook

170 Although various recent transport experiments and findings point to the existence of 2D Wigner
 171 crystal and relevance of liquid-WC transition in the observed RIPs and peculiar voltage-current
 172 characteristics in strongly correlated 2D carrier systems, due to the limitations of transport data and
 173 various challenging aspects of the experiments (ultra-low temperatures, low noise-low level
 174 measurements, and stringent sample quality requirements), information about how the 2D Wigner
 175 crystal transforms into metallic liquid transition and the nature of intermediate phase are still limited
 176 and further experiments are required to obtain more in-depth understandings.

177 On the theoretical side, from the beginning of the discovery of 2D MIT, a number of weak-
 178 interaction based models were proposed to explain the metal-insulator transition or cross-over
 179 behavior in the electrical resistivity such as screening, potential fluctuation, percolation [36-38]. In
 180 these theories, conventional Anderson localization (weak-localization in the metal-side and strong
 181 localization in the insulator side of the MIT) is still relevant. From extensive prior studies on the
 182 Anderson insulator to 2D QH transition, it is expected the zero field insulator would directly
 183 transition to an integer QH state upon the application of perpendicular magnetic fields [39]. This is

184 in stark contrast to the observation of low field RIPs between the zero field insulator and integer QH
185 state [10-12]. It appears that it is difficult to use weak-interaction based models to reconcile the
186 observations of various RIP, QH and HFIP states and their connection to the liquid-WC transition
187 when the behavior of strongly interacting high r_s 2D systems is examined beyond the zero magnetic
188 field to finite perpendicular magnetic fields. In contrast to weakly interaction theories, many other
189 theories emphasize the importance of strong correlations and relevance of WC physics in the systems
190 showing 2D MIT [9, 25, 26, 32] and therefore may be further compared with the experimental results.
191 These theories are based on a number of different approaches: analytical mean-field models [25, 26],
192 quantum Monte Carlo simulations [29], or dynamic mean field theory (DMFT) [31, 32]. In the mean-
193 field theories by Kivelson and Spivak [25, 26], the various spectacular transport behavior in correlated
194 2D systems showing MIT are attributed to the Fermi liquid to WC transition where intermediate
195 states are unavoidable [30]. In the intermediate 'micro-emulsion' states (e.g. WC bubbles in a Fermi
196 liquid background), it is the interplay or transformation between Fermi liquid and WC components
197 tuned by temperature or magnetic field that dictates the transport behavior and gives rise to the
198 resistivity change of the system. It seems that the most relevant micro-emulsion phase to the
199 experimentally observed RIPs and intermediate phases is the scenario where WC bubbles co-exists
200 in a Fermi liquid. Whether other micro-emulsion states (Fermi liquid bubbles in WC, 1D ordered
201 stripes) exist in experiments requires further research and more theoretical developments are desired
202 to establish more quantitative predictions on the experimental systems. New theoretical approaches
203 based on hydrodynamics seem quite promising and are currently being developed [40, 41]. In
204 addition to the mean-field models by Kivelson and Spivak, strong interaction and Wigner-Mott
205 transition based theoretical studies led by Dobrosavljevic and collaborators [31, 32] may also be
206 relevant to the experimental findings. Modern DMFT was applied to study the MIT in 2D carriers
207 with high r_s and the early approaches of Wigner and Mott were reconciled. Based on this 'Wigner-
208 Mott' transition scenario, DMFT calculations are able to explain many detailed behavior in the
209 electrical transport and charge ordered intermediate phases similar to charge density wave (CDW)
210 are predicted to form before the system enters WC. It is worth to note that in the DMFT theory, both
211 metallic CDW and insulating CDW are found [32]. It will be very interesting to see whether such
212 CDW states exist in experiments.

213 On the experimental side, besides the previously mentioned compressibility and inductance
214 measurements that require further advancements and understanding, other striking effects found in
215 the zero field 2D MIT are very worth to explore in relation to the RIPs. One particular case is the effect
216 of an in-plane magnetic field. In the zero field MIT, it was established that an in-plane magnetic field
217 causes large positive magneto-resistance and drives the system toward the insulating phase [8, 9, 42,
218 43]. Recent in-plane field magneto-transport experiments revealed that the resistivity of dilute 2D
219 electrons in Si-MOSFET in the insulating state are the same for zero field and in the presence of an
220 in-plane field that polarizes the spins [43], a behavior different from the metallic phase of the 2D MIT
221 [8, 9, 42]. It would be thus very insightful to study the effect in-plane magnetic field induced spin-
222 polarization effect on the RIPs and examine how the in-plane field affects the transition between the
223 RIPs and metallic liquid. In addition to spin-polarization effect, understanding the physics of
224 correlated 2D systems at 'high temperatures' comparable with the Fermi temperature is also an
225 interesting topic. In this 'semi-quantum' regime, the correlated electron fluid is expected to show
226 unique behavior in viscosity [26]. While there were many experiments done at low temperatures,
227 transport studies in this 'semi-quantum' regime where $T-T_F$ is limited and worth exploring further
228 [45]. In addition to transport, new techniques such as thermopower measurement are strongly
229 desired to shed new light on the RIPs.

230 In summary, the recent observations of RIPs in low magnetic fields where $v > 1$ and the
231 connections between low field RIPs and high field RIP and HFIP with $v < 1$ point to the formation of
232 WC as the origin of the low field RIPs and the 2D MIT being driven by the liquid-WC transition.
233 There are also transport evidences for possible new intermediate phases when the insulating phase
234 (either the insulator in zero field or the RIPs at low magnetic fields) is destructed by an increasing
235 voltage, temperature or carrier density. Further experimental and theoretical progresses in these

236 fronts are needed and expected to yield many more exciting new insights on the long standing
237 problems of 2D WC-liquid transition and MIT.

238

239 **Funding:** This research was funded by NSF grant number DMR-0906415 and DMR-1607631.

240 **Acknowledgments:** We thank Vladimir Dobrosavljevic and Jian Huang for helpful discussions.

241 **Conflicts of Interest:** The authors declare no conflict of interest.

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