

1 Article

# 2 Modeling the Dynamic Concentration Profiles and 3 Efficacy of Type-I and Type-II Photopolymerization

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13 **ABSTRACT:** The kinetics and efficacy profiles of photoinitiated polymerization are theoretically  
14 presented. For the same dose, lower light intensity achieves a higher steady-state-efficacy (SSE) in  
15 type-I; in contrast, type-II has an equal SSE. Higher light intensity has a faster rising efficacy, due to  
16 faster depletion of photoinitiator (PS) concentration. However, type-II process is also affected by the  
17 available oxygen. Higher light intensity produces more efficient singlet oxygen, resulting a higher  
18 transient efficacy, in which all intensities reach the same SSE when oxygen is completely depleted.  
19 With external oxygen, type-II efficacy increases with time, otherwise, it is governed only by the light  
20 dose, i.e., same dose achieves same efficacy. Moreover, type-II has an efficacy follows Bunsen Roscoe  
21 law (BRL), whereas type-I follows non-BRL. The measured type-I efficacy and gelation profile are  
22 analyzed by our analytic formulas. Schematics of the photocrosslinking stage defined by the  
23 availability of oxygen is developed, where both type-I and -II coexist until the oxygen is depleted.  
24 The overall efficacy may be enhanced by resupply of PS or oxygen during the light exposure. The  
25 roles of light dose and PS concentration on the efficacy of photoinitiated polymerization should be  
26 are governed a new concept of a volume efficacy ( $V_e$ ), defined by the product of the crosslink (or  
27 gelation) depth (CD) and local [efficacy].

28 **Keywords:** polymerization; modeling; kinetic; singlet oxygen; polymerization efficacy; crosslinking

29

## 30 1. Introduction

31 External stimuli such as such as pH, temperature, hydrophobicity, the presence of ions or  
32 enzymes, and light-stimulation will cause the three-dimensional polymeric networks, called  
33 hydrogels, to be physically or chemically linked for “gelation”, or crosslinking [1,2]. These processes  
34 result in the increasing of polymer networks and viscosity of the medium, but decreasing solubility  
35 of polymers. Physically-assembled gels are built with polymer networks tied via hydrogen bonds,  
36 ionic interactions, hydrophobic associations, or agglomerations. Chemically-linked hydrogels are  
37 commonly prepared via a three-dimensional polymerization using a water-soluble monomeric  
38 polymer and a multi-functional cross-linker [2]. Photoinitiated polymerization and crosslinking  
39 provide advantageous means over the thermal-initiated polymerization, including fast and  
40 controllable reaction rates, spatial and temporal control over the formation of the material, and  
41 without a need for high temperatures or pH conditions [1,2].

42 Photodynamic therapy (PDT) offers biometrical applications in dermatology, orthopedics  
43 (tissue engineering), ophthalmology and cancer treatments in various parts of human body,  
44 including early stage (micro-invasive) lung cancer, lung tumors (endobronchial, mesothelioma), skin,

45 brain, colorectal and breast cancer, chronic skin diseases (psoriasis, vitiligo), and oral cavity (anti-  
46 bacterial, curing) [7-14]. Ophthalmology applications include age-related macular degeneration  
47 (AMD), coagulation of retina and corneal crosslinking for keratoconus and ulcers [15-17].

48 The kinetics of photoinitiated polymerization systems (PPS) have been studied by many  
49 researchers for uniform photoinitiator distribution or for the over simplified cases that the photolysis  
50 product becomes completely transparent after polymerization or constant light intensity [18-23].  
51 Kinetic modeling of PPS assumed either a constant light intensity (for thin polymers), or a  
52 conventional Beer-Lambda law for the light intensity [13, 18-20]. For more realistic systems, the  
53 distribution of the photoinitiator is non-uniform and the UV light may still be absorbed by the  
54 photolysis product besides the absorption of the monomer. To improve the efficiency and spatial  
55 uniformity (in the depth direction) particularly in a thick system (>1.0 cm), we have presented the  
56 numerical results using a focused light [24] and two-beam approach [25] for the case of uniform PS  
57 distribution; and analytic and computer modeling for the non-uniform case [26]. Optimal efficacy in  
58 light-activated biomedical systems and nonlinear laws versus linear Beer-Lambert law were also  
59 reported by Lin [27].

60 The kinetics and macroscopic modeling of PPS for anti-cancer have been reported by Zhu et al  
61 [13] and Kim et al [14], which, however, are limited to the type-II oxygen-mediated process. Lin  
62 reported the kinetic modeling for both type-I and type-II mechanism for the application in corneal  
63 collagen crosslinking (CXL) [16,17,28], where the temporal and spatial profiles of PS concentration  
64 and the CXL efficacy were also reported. Accelerated CXL has been clinically used for faster  
65 procedure (within 3 to 10 minutes) using higher light intensity of 9 to 45 mW/cm<sup>2</sup>, in replacing the  
66 conventional 3 mW/cm<sup>2</sup> which took 30 minutes [29]. However, much less efforts have been done for  
67 fast PPS in thick polymers using a high light intensity.

68 Photo-polymerization offers two major categories of biomedical applications: (a) photodynamic  
69 therapy (PDT) using light-initiated oxygen free radical; and (b) crosslinking (or gelation) of  
70 biomaterials using radical-substrate coupling for tissue engineering [1,2]. In general, both type-I and  
71 type-II reactions can occur simultaneously, and the ratio between these processes depends on the  
72 types and the concentrations of PS, substrate and oxygen, the kinetic rates involved in the process  
73 [16,17]. Kim et al [14] have focused on the type-II oxygen-mediated process, whereas we have  
74 previously focused on the non-oxygen-mediated type-I mechanism [27]. It is not yet fully understood,  
75 theoretically or experimentally, the interaction between these two mechanisms.

76 In this study, we will further study the oxygen-mediated type-II mechanism, and compare the  
77 significant different kinetics and efficacy of type-I and type-II by numerical simulations and analytic  
78 formulas. Dynamic profiles of singlet-oxygen concentration, polymerization efficacy, and the cell  
79 viability are produced by numerical solutions of macroscopic equations. Depending on the rate  
80 constant, type-I and type-II mechanism may coexist to achieve a higher efficacy than type-II dominant  
81 case, which is limited by the available oxygen. We will also explore the new strategy to enhance the  
82 overall efficacy by resupply of PI or oxygen during the light exposure.

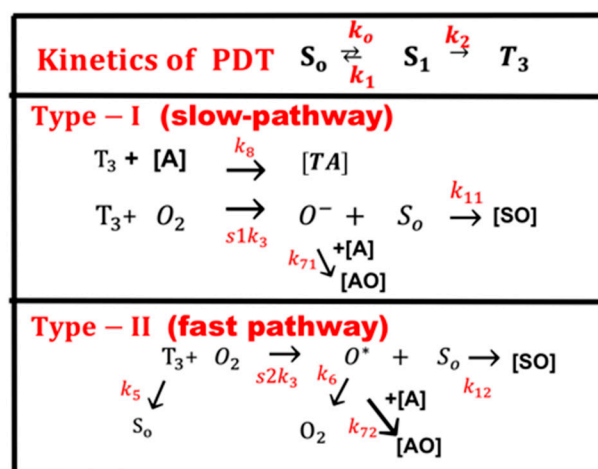
## 83 2.Methods and Modeling systems

### 84 2.1. Photochemical kinetics

85 The kinetics of corneal collagen crosslinking (CXL) shown by Fig. S-1 (shown in Supporting  
86 Information) was previously reviewed by Lin [17] using UVA (365 nm) initiated riboflavin solution  
87 as the sensitizer. Fig. 1 shows the 4 coupled dynamic equations defined as  $dI/dz$ ,  $d[C]/dt$ ,  $d[T]/dt$  and  
88  $d[O_2]/dt$ , for, respectively, light intensity, PS concentration, triplet state, and oxygen molecular. The  
89 three pathways are revised for a more general polymer system as shown in Fig. 1 and briefly  
90 summarized as follows. In type-I pathway, the excited PS triple-state ( $T_3$ ) can interact directly with  
91 the substrate (A); or with the ground state oxygen ( $O_2$ ) to generate a superoxide anion ( $O^-$ ), which  
92 further reacts with oxygen to produce reactive oxygen species (ROS). In comparison, in type-II  
93 pathway,  $T_3$  interacts with ( $O_2$ ) to form a reactive singlet oxygen ( $O^*$ ). In general, both type-I and  
94 type-II reactions can occur simultaneously, and the ratio between these processes depends on the

95 types and the concentrations of PS, substrate and oxygen, the kinetic rates involved in the process  
96 [16].

97



98

99 **Fig. 1** The kinetics of PDT, where  $[S_0]$ ,  $[S_1]$  and  $[T_3]$  are the ground state, singlet excited state, and  
100 triplet excited state of PS molecules. Three pathways are shown for both type-I and type-II process.  
101 Ground state oxygen may couple to  $T_3$  to form either singlet oxygen  $[O^*]$ , or other reactive radicals  
102  $[O^-]$ . In type-I pathway,  $T_3$  can interact directly with the collagen substrate (A); or with the oxygen  
103 ( $O_2$ ) to generate a superoxide anion ( $O^-$ ); in type-II pathway,  $T_3$  interacts with the ground oxygen ( $O_2$ )  
104 to form a singlet oxygen ( $O^*$ ) [16].

105 These factors also influence the overall photopolymerization efficacy, particularly the PS triplet  
106 state quantum yield ( $q$ ) and its concentration. Furthermore, the specific protocols and the methods of  
107 PS instillations prior to and during the photopolymerization also affect the short and long term  
108 outcomes. The overall photopolymerization efficacy is proportional to the time integration of the light  
109 intensity,  $I(z, t)$  and the PS and oxygen concentration,  $C(z, t)$ , and  $[O_2]$ . The efficacy reaches a  
110 saturated (steady) state when  $C(z, t)$  or  $[O_2]$  is depleted by the light, where higher intensity depletes  
111  $C(z, t)$  and  $[O_2]$  faster and therefore reaches a lower steady-state efficacy [26,29].

112 Referring to the kinetic pathways shown by Fig. 1, a set of quasi steady-state macroscopic kinetic  
113 equation for the PI ground-state,  $C(z, t)$ , and the ground state oxygen molecule  $[O_2]$  is constructed  
114 [14,16]

$$\frac{\partial C(z, t)}{\partial t} = -b[g + g']C \quad (1.a)$$

$$\frac{\partial [O_2]}{\partial t} = -sbCG + P \quad (1.b)$$

$$\frac{\partial I(z, t)}{\partial z} = -A'(z, t)I(z, t) \quad (1.c)$$

$$A'(z, t) = 2.3[(a' - b')C(z, t) + b'C_0(z) + Q] \quad (1.d)$$

115 where  $b=aqI(z, t)$ ;  $a=83.6wa'$ ;  $w$  is the light wavelength;  $a'$  and  $b'$  are the molar extinction coefficient  
116 of the initiator and the photolysis product, respectively;  $Q$  is the absorption coefficient of the  
117 monomer and the polymer repeat unit. Typical values are [4,9]:  $a'=0.2$  to  $0.3$  (1/mM/cm),  $b'=0.1$  to  $0.15$   
118 (1/mM/cm), and  $C_0=1.0$  to  $3.0$  mM; and for a UV light at 365 nm,  $a=0.00305a'$  (1/cm).

119 For type-I,  $g=k_8[A]G_0/k_3$ ,  $G_0=1/([O_2] + k + K')$ ; and for type-II,  $g'=K'(C+d)G(z)$ , with  $G(z)=[O_2]G_0$ ,  
120  $k=(k_5+k_8[A])/k_3$ ;  $K'=k_{12}/(k_6+k_{12}(C+d)+k_{72}[A])$ ;  $d$  is a low concentration correction related to the  
121 diffusion of singlet oxygen [14].  $[A]$  is the substrate concentration.  $q$  is the triplet state  $[T]$  quantum  
122 yield given by  $q=k_2/(k_1+k_2)$ ;  $s=s_1+s_2$ , with  $s_1$  and  $s_2$  are the fraction of  $[O_2]$  converted to the singlet

123 oxygen and other ROS, respectively, in type-I and type-II [16]. All other rate constants,  $k_j$ ,  $k_{ij}$  are  
 124 defined by the reaction paths shown in Fig. 1.

125 In Eq. (1.b) we have included the light intensity in the polymer given by a time-dependent Beer-  
 126 Lambert law [27]. We have also included in Eq. (1.b) the oxygen source term  $P(z, t) = p(1 - [O_2]/[O_0])$ ,  
 127 with a rate constant  $p$  to count for the situation when there is an external continuing supply, or nature  
 128 replenishment (at a rate of  $p$ ), besides the initial oxygen,  $[O_0]$ , in the polymer.

129 We note that Eq. (1) was also presented by Kim et al [14] for the anti-cancer kinetics. However,  
 130 they have assumed a constant UV intensity, i.e.,  $A'(z, t)$  is a constant in Eq. (1.d). They also ignored  
 131 the contribution from the type-I term,  $k_s[A]$ , since type-II is dominant in their anti-cancer process.  
 132 Most of the previous model have also ignored the dynamic of UV intensity given by Eq. (1.c) and the  
 133 depth-dependent profile of PI and light intensity [18-20]. Exact solutions of Eq. (1) require numerical  
 134 simulations. For analytic formulas, we will use an effective  $A(z, t)$  or its mean value, such that  $A'(z,$   
 135  $t)$  becomes time-independent in solving Eq. (1). The effective absorption is given by [28], The effective  
 136 absorption is given by [28],  $A' = 2.3 \times 0.5(a' + b') (1 - 0.25z/D) C_0 + Q$ .

### 137 2.2. Concentration Profile, Crosslink (or gelation) Time and Depth

138 In solving Eq. (1), we will choose initial profile (at  $t=0$ )  $I_0(z) = I_0(1 - 0.25z/D)$  for the light intensity;  
 139 and  $C_0(z) = C_0F(z)$ , with  $F(z) = 1 - 0.5z/D$ , for the PS concentration distribution; where  $D$  is the PS  
 140 concentration distribution depth; and when  $D \gg 1.0$  cm,  $F=1$  representing a flat (or uniform) PS  
 141 distribution. Analytic solution of Eq. (1) is available for the type-I process and under certain  
 142 approximation. For  $g \gg g'$ , or for the case that type-I is dominant over type-II, we obtain an  
 143 approximate solution,  $C(z, t) = C_0F(z)\exp(-bgt)$ , assuming  $b$  and  $g$  are time-independent, or taking  
 144 their averaged value.  $A'(z, t)$  in Eq. (1.d) has an initial value ( $A_1$ , with  $b'=0$ ) and steady state value  
 145 ( $A_2$ , with  $C=0$ ), given by:  $A_1 = 2.3[a'C_0F' + Q]$ , and  $A_2 = 2.3[b'C_0F' + Q]$ , with  $F'(z) = 1 - 0.25z/D$  being the  
 146 integration of  $F(z)$  over  $z$ ; the mean value is given by  $A = 0.5(A_1 + A_2)$ . We have also developed  
 147 numerically fit value [9],  $A = 2.3[m'b'C_0F + Q]$ , with  $m=0.8$  to  $1.0$  depending on the value of  $a'$  and  $b'$ .

148 A crosslink (or gelation) time  $T^*$  may be defined by when the PS concentration  $C(z, t=T^*)$   
 149  $= C_0(z)\exp(-M)$ , with  $M = 4$ , or  $C(z, t)$  is depleted to 0.018 of its initial value. We obtain an analytic  
 150 formula  $T^*(z) = T_0 \exp(Az)$ , where  $T_0$  is the surface depletion time given by  $T_0 = M/(bg)$ , which is inverse  
 151 proportional to the light initial intensity, since  $b = aqI(z)$ .  $T^*$  may be also defined by the level of  
 152 photopolymerization efficacy, or the crosslink time ( $T_c$ ), to be discussed later. The strong depletion  
 153 of  $C(z, t)$  will also affect the time-dependent profiles of the intensity,  $I(z, t)$ , which in general, will  
 154 not follow the conventional Beer-Lambert law (BLL), and should be governed by a generalized, time-  
 155 dependent BLL first presented by Lin [27,28].

156 Another important parameter is called crosslink (or gelation) depth ( $z^*$ ) may be defined by the  
 157 depth having PS concentration  $C(z^*, t=T^*)$  reduced to a low value of  $C'$  which typically is ( $1/e^4$ , or  
 158 0.18%) of its initial value (at  $t=T^*$ ). Therefore, it is given by (for the case of flat distribution or  $F=1$ )  
 159  $z^* = (1/A)\ln(bE_0/M)$ , with  $M = \ln(C_0/C')$ , which is proportional to the light fluence (or dose),  $E_0$ . In  
 160 general, for  $F < 1$  (with  $D < 1.0$  cm) and  $A$  is  $z$ -dependent,  $z^*$  needs numerical calculation.

### 161 2.3. Efficacy Profiles

162 The normalized photo-polymerization efficacy defined by  $C_{eff} = 1 - [A]/[A]_0 = 1 - \exp(-S)$ , with  $S$ -  
 163 function for type-I ( $S_1$ ) and type-II ( $S_2$ ), and the overall efficacy is given by  $C_{eff} = 1 - \exp[-(S_1 + S_2)]$ . The  
 164 type-I efficacy may be further expressed by rate equation of conversion of collagen monomers  $[M]$  to  
 165 polymers, where the NOM term of Eq. (1.a),  $g = k_s[A]G_0/k_3$ , is replaced by an overall rate constant ( $K$ )  
 166 including all polymerization chain reactions. The  $S$  functions are given by [13,16]

$$S_1 = \int_0^t [\sqrt{a q g K I C} + (f s_1 a q K') I(z, t) G] dt \quad (3.a)$$

for type-I, and

$$S_2 = \int_0^t f s_2 a q K' I(z, t) G dt \quad (2.b)$$

167 for type-II.

168 The first term in Eq. (3.a) relates to the direct coupling of triplet state [T] with the substrate [A]  
 169 under hypoxic conditions or any other non-oxygen-mediated (NOM) reactions; and the second term  
 170 relates with the reactive oxygen species (ROS)-mediated reactions (in type-I).  $f$  is the fraction of all  
 171 ROS (including singlet oxygen) interacting with acceptors [A], or the oxygen-mediated (OM)  
 172 reactions in type-I and type-II.  $s_2$  and  $s_1$  are the fraction of  $[O_2]$  interacting with [T] to produce singlet  
 173 oxygen (in type-II) and other ROS (in type-I), respectively.

#### 174 2.4. Analytic formulas

175 For analytic formulas, we will use the mean value of  $A(z)$  such that  $I(z, t) = I_0 \exp(-Az)$ , and  $C(z,$   
 176  $t) = C_0 \exp(-Bt)$ , with  $B = bg = aqgI(z)$ , Eq. (3.a), for the case that  $g' \ll g$ , the type-I, S- function is given by  
 177 [16,28]

$$S_1 = \sqrt{4KC_0F \exp(Az) / (aqgI_0)} E1 \quad (4.a)$$

$$E1 = [1 - \exp(-0.5Bt)] \quad (3.b)$$

178 which is a nonlinear function of the light dose (E) given by its Taylor expansion  $S_1 =$   
 179  $\sqrt{(aqI_0KC_0F) \exp(-Az)} t [1 - 0.5aqgE + \dots]$ , which follows Bunsen-Roscoe reciprocal law (BRL)  
 180 only for small time with the first term kept. In contrast, type-II efficacy, given by the time integral of  
 181 [IC] follows the BRL [26]. A crosslink time (T) may be defined by Eq. (4.b) when  $E1 = 0.87$ , or  $0.5BTc = 2$ ,  
 182 which gives us  $Tc = 4/B = 4/(bg) = [4/(aqgI_0)] \exp(Az) = T_0 \exp(Az)$ , with the surface crosslink time given  
 183 by  $T_0 = 4/(aqgI_0) = 1000/I_0$ , for  $aqg = 0.004$ . We note that the crosslink time equals to the depletion time  
 184 ( $T^*$ ), when  $M = 4$ , and it also defines the gelation time, or crossover time.

185 Type-II process is much more complex than type-I and requires numerical solutions to be shown  
 186 later. For analytic results for type-II dominant case (with  $g' \gg g$ ), we assume an approximated oxygen  
 187 concentration given by  $[O_2] = [O_0] - m'btC_0$ , with  $m'$  being a fit parameter; and the similar functional  
 188 forms for  $C(z,t)$  and  $I(z,t)$  as type-I, the time integral of Eq. (3.b) gives us, for  $k \ll [O_0]$  and  $P = 0$ ,

$$189 \quad S_2 = (f sag) I'(z) C_0 [(1 - k/[O_0])(E1 + dt) - HO] \quad (5)$$

190  
 191  
 192 with  $I'(z) = 0.5(I_1 + I_2) = 0.5I_0 \exp(-A'z)$ , is a mean light intensity, and HO is a high-order term. Eq. (5)  
 193 shows that the type-II efficacy is an increasing function of  $I_0 C_0 [O_0]$ ; and  $S_2$  has a transient state (with  
 194  $E1 = t$ ) proportional to the light dose,  $I_0 t$ ; and steady-state is only dose-dependent (for the case of  $P = 0$ )  
 195 to be justified numerically later.  
 196

#### 197 2.5. Nonlinear scaling law

198 As predicted by our  $S1$  formula, the efficacy at transient state (for small dose) is proportional to  
 199  $tI_0^{0.5}$ , however, at steady-state, it is a nonlinear increasing function of  $[C_0/I_0]^{0.5}$  or  $[t/E_0]^{0.5}$ . This nonlinear  
 200 scaling law predicts the clinical data more accurate than the linear theory of Bunsen Roscoe law (BRL)  
 201 [26]. Accelerated PPS based on BRL, therefore, has undervalued the exposure time (t) for higher  
 202 intensity using the linear scaling of  $t = [E_0 / I_0]$ , rather than  $t = [E_0 / I_0]^{0.5}$ , based on our nonlinear law.  
 203 To achieve the same efficacy, higher PS concentration requires higher light intensity; and for the same  
 204 dose, higher light intensity requires a longer exposure time.

205 The BRL is based on the conventional Beer-Lambert law for light intensity without PS depletion,  
 206 such that  $I(z)$  is time-independent, and  $C(z,t) = \text{constant} = C_0F$ , therefore,  $S_1 =$   
 207  $\sqrt{4KC_0FE_0 \exp(Az)}$  which is a linear function of the dose  $E_0 = (tI)$ .

208 Our nonlinear law, as shown by Eq. (4) predicts that, for the same dose, higher intensity depletes  
 209 the PS concentration faster and reach a lower steady-state efficacy. Further discussion will be shown  
 210 later. As shown by our S-formula, diffusion depth (D) also plays important role. Larger D will achieve

211 higher efficacy as shown by the PS distribution function,  $F(z)=1-0.5z/D$ , which is an increasing  
 212 function of  $D$ , and  $F=1.0$  for the flat (uniform) distribution case. The above features have been  
 213 clinically shown in corneal crosslinking [26], but not yet for other PPS.

### 214 2.6. Volume efficacy

215 The new concept of a volume efficacy ( $V_e$ ), first introduced by Lin [30], is defined by the product  
 216 of the crosslink (or gelation) depth (CD) and local [efficacy], or  $V_e=z[1-\exp(-S)]$ , where  $z$  is the  
 217 polymerization (or crosslink) depth (PD) given by  $z=(1/A)\ln(b'/E_0)$ , with  $b'=b/M$ , and  $E_0$  is the light  
 218 dose. For a type-II process, the local (at a specific depth  $z$ ) efficacy is defined by  $\text{Eff}=1-\exp(-S)$ , with  
 219  $S$  function is given by Eq. (3). For a polymer thickness of  $z_0$ , the normalized  $V_e$  is given by  $V'=V_e/z_0$ .  
 220 For type-II application of anti-cancer for a cell sample thickness of  $z_0$ , the cell viability (normalized  
 221 by  $z_0$ ) is given by  $CV=1-V_e/z_0$ . We should note that both CD and  $S$  are increasing function of the light  
 222 dose, however, they are competing functions with respect to the PS concentration. Higher  $C_0$  offers  
 223 higher  $S$  (or local efficacy) but it has a smaller depth (due to the larger absorption, or larger  $A$ -value).  
 224 The general feature of  $V_e$  may be stated as follows: increasing light dose (for a fixed  $C_0$ ) offers both  
 225 higher [local efficacy] and [depth], and  $V_e$ ; however, increasing  $C_0$  (for a fixed light dose) suffers a  
 226 shallow depth, although the [local efficacy] increases. Therefore, there is an optimal  $C_0$  for maximum  
 227  $V_e$ . Numerical results and application for cell viability in anti-cancer and corneal crosslinking will be  
 228 presented elsewhere.

## 229 3. Results and Discussions

### 230 3.1. Concentration profiles

231 Numerical results of Eq. (1) are shown in Fig. 3 for the PS concentration dynamic profiles, for a  
 232 type-I dominant case, with  $g' \ll g$  in Eq. (1.a). One may see that depletion of PS starts from the polymer  
 233 surface, and gradually into the volume ( $z > 0$ ). We note that the PS concentration profile is an  
 234 increasing function of  $z$  for the uniform case. In contrast, the non-uniform case shows a decreasing  
 235 function of  $z$ .

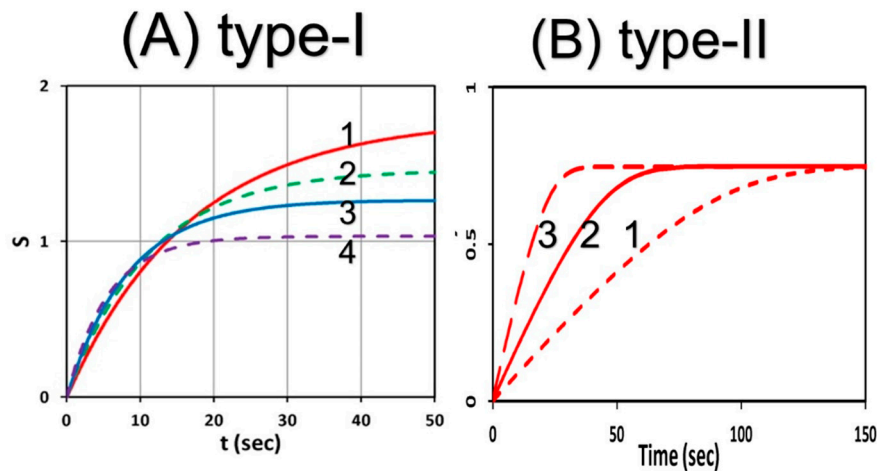
### 236 3.2. Efficacy profiles

237 Using typical values of:  $a'=0.2(1/\text{mM}/\text{cm})$ ,  $b'=0.1(1/\text{mM}/\text{cm})$ ,  $Q=0.1(1/\text{cm})$ ,  $q=0.5$ ,  $aqg=0.012$   
 238  $\text{cm}^2/\text{mW}$ ; the mean  $A(z)=0.35C_0F(z)+0.23$ ,  $B=(0.006I_0)\exp(-Az)$ , with  $C_0$  in mM,  $I_0$  in  $\text{mW}/\text{cm}^2$ . Eq. (4)  
 239 gives a normalized  $S$ -function defined by  $S=S_1/S_0$ , where  $S_0=[4K/(aqg)]^{0.5}$  is a proportional constant,  
 240 such that  $S = E1\sqrt{(C_0F/I_0)\exp(Az)}$ .  
 241 In the follow figures, we will show the normalized  $S$ -function for  $S_0=4$ , or  $K/(aqg)=4$ . In addition, the  
 242 transient factor  $E1$  is based on  $aqg=0.012$ , or  $B=(aqg)I(z)=0.012I(z)$ , and  $K=4(aqg)=0.048$ .

243 Fig. 2 compares the efficacy  $S$ -function profiles for: (A) type-I, and (B) type-II, for various light  
 244 intensity,  $I_0 = (9,18,30,45) \text{ mW}/\text{cm}^2$ , for curves (1,2,3,4). The important features demonstrated by Fig.  
 245 4 and 5 are summarized as follows:

- 246 (i) For the same dose, lower light intensity achieves a higher steady-state-efficacy (SSE) in type-  
 247 I, as also shown by Eq. (4); in contrast to type-II, which has an equal SSE, as also shown by  
 248 Eq. (5).
- 249 (ii) In both type-I and type-II, higher light intensity has a faster rising efficacy, due to faster  
 250 depletion of PS concentration. However, type-II process is also affected by the available  
 251 oxygen. Higher light intensity produces more efficient singlet oxygen, resulting a higher  
 252 transient efficacy, as shown by Fig. 4-(B), in which all intensities reach the same SSE when  
 253 oxygen is completely depleted, as shown by Fig. 4-(B).
- 254 (iii) As shown by Fig. S-2 (shown in Supporting Information), for the same dose, lower light  
 255 intensity always achieves higher efficacy in type-I; in contrast, type-II efficacy is governed

256 only by the light dose (for the case of no external oxygen); i.e., same dose achieves same  
 257 efficacy. Moreover, type-II has an efficacy follows BRL, whereas type-I follows non-BRL  
 258 [27,17]. These numerical results are also predicted by our analytic formulas, Eq. (4) and (5).



259

260 **Fig. 2** Comparing the efficacy  $S$ -function temporal profiles for: (A) type-I, and (B) type-II, for various  
 261 light intensity,  $I_0 = (9,18,30,45)$  mW/cm<sup>2</sup>, for curves (1,2,3,4).

262 As shown by  $S$ -formulas, Eq. (5), for the anti-cancer type-II PDT efficacy  $S \sim [O_2]C$ , which requires  
 263 both PS concentration  $C$  and  $[O_2]$ . Therefore, resupply of PS or oxygen would enhance the generation  
 264 of singlet oxygen radicals, and improve the anti-cancer efficacy via type-II PDT. Resupply of PS or  
 265 oxygen during the light exposure would enhance the overall efficacy; this new strategy has been  
 266 proposed in type-I corneal crosslinking [13], but not yet in gelation of thick polymers. These  
 267 theoretically predicted features have been only partially proven clinically for corneal crosslinking  
 268 [27]. Therefore, further experimental studies are highly desired in polymer systems.

### 269 3.3. Analysis of experiments

270 Our formula, Eq. (4), predicts that the type-I steady-state- $S$  is proportional to the square-root of  
 271 the PS concentration ( $C_0$ ). Therefore, the crosslink efficacy, defined by  $Eff=1-\exp(-S)$ , is also an  
 272 increasing function of  $C_0$ . This feature has been clinically reported by O'Brart et al in corneal  
 273 crosslinking [30], where the PS is riboflavin solution initiated by a UVA light at 365 nm. The role of  
 274 PS concentration was also shown by Table 1 of Holmes et al [31], where (for LAP) increasing the PS  
 275 concentration from 0.1% to 0.5% (w/v) in the thiol-ene mixture resulted in a 15-fold increase in the  
 276 storage modulus. This increasing feature may be analyzed by our  $S$ -function, Eq. (4), given by a  
 277 steady-state formula  $S'=S_0 [C_0 \exp(Az)]^{0.5}$ . For  $C_0$  increase five times, and for  $Az=0.9$ , we calculate  
 278  $S$ -function increases a factor of  $2.23 \times \exp(2Az) = 14.5$ , which is comparable to the increase of storage  
 279 modulus; noting that when  $C_0$  is 5 times,  $A$  is also 5 times, given by  $A=2.3mb'C_0F(z) + Q$ , if  $Q \ll 0.1$   
 280 (1/cm).

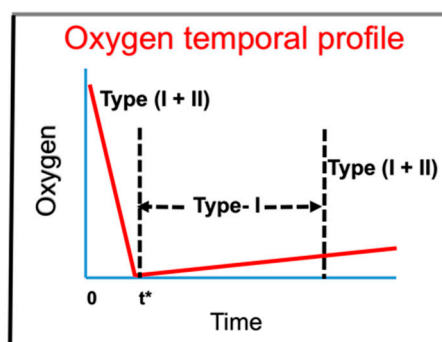
281 Fig. 3 of Holes et la [31] shows the gelation kinetic profiles which has a strong similarity as our  
 282 Fig. 2 (red curve). The storage modulus was found to increase with time and UV exposure until a  
 283 plateau was reached within 300 s, indicating no further elastic properties (complete chemical gel).  
 284 The plateau-time corresponds to our crosslink time ( $T^*$ ) defined before. Similarly, the measured data  
 285 of Shih et al [11] (in their Fig. 1) showed crosslinking of thiol-norbornene PEG- peptide hydrogels  
 286 (initiated by a visible light). Also has a strong similarity as our Fig. 2 (red curve), except the time scale  
 287 which depends on the types of PS and light used in the process. Unfortunately, Shih et al [11] and  
 288 Holes et al [31] did not measure the profiles for different light intensity, as shown in our Fig. 2, to  
 289 justify our predicted feature that higher light intensity is less efficient in gelation. However, our  
 290 predicted feature has been clinically demonstrated in ophthalmic system for CXL [27].

291 Anti-cancer via oxygen-mediated type-II mechanism has been clinically studied [32,33], where  
 292 the cytotoxic effect of photodynamic therapy (PDT) to tumor tissue is resulted by the generation of  
 293 singlet oxygen. Efficacy of PDT is mainly influenced by: the concentration of PS drug accumulated  
 294 into the cells, molecular oxygen in tissue, the light dose, intensity and dose (fluence) [34]. High  
 295 concentrations of singlet oxygen can lead to necrotic cell death. In contrast, low concentrations lead  
 296 to cell survival and increase the metabolism; whereas medium singlet oxygen concentrations lead to  
 297 initiation of apoptosis or autophagy [32]. Therefore, the threshold light dose and singlet oxygen dose  
 298 play the important role in PDT for anti-cancer. The singlet oxygen threshold dose, and the dose-  
 299 dependence cell viability curves of human cancer cells of K562 and Hela after red-light irradiation of  
 300 Radachlorin were reported in vivo by Klimenko et al [32]. They showed that the cell viability, defined  
 301 by  $CV = \exp(-S)$ , is lower for higher  $C_0$  and/or light intensity ( $I_0$ ). Moreover, with  $p > 0$ , external oxygen  
 302 offers lower CV, or better cell killing. The threshold dose  $E^* = I_0 t^*$  (or time  $t^*$  for a fixed light intensity)  
 303 to reach a cell viability  $CV^* < 25\%$ , is higher for smaller  $C_0$  and/or  $I_0$ . These features are in consistent  
 304 with our numerical results for type-II PDT to be shown later.

305 Above examples demonstrate that our formulas predict very well the measured results, at last  
 306 the overall trends. However, the accuracy of our formulas will require accurate measurement of the  
 307 parameters involved, such as the rate constant (K), the quantum yield (q), the molar extinction  
 308 coefficient of the initiator ( $a'$ ), the photolysis product ( $b'$ ), and the monomer and the polymer repeat  
 309 unit (Q) et al. In addition, further experimental measurements should also include the roles of PS  
 310 concentration and light intensity. Our group has also worked on the in vitro measurement of cell  
 311 viability, which may be empirically analyzed by our formulas, and results will be published  
 312 elsewhere.

### 313 3.4. The role of oxygen

314 By solving the kinetic coupled Eq. (1) for oxygen, we have previously developed a schematic for  
 315 type-I and type-II mechanisms in CXL, which has a very short oxygen depletion time ( $t^*$ )  
 316 approximately 10 to 30 seconds [16,17]. For thick polymer system,  $t^*$  has a wide range of 50 to 500  
 317 seconds, depending on the PI, light intensity and kinetic rate constants. For example, in anti-cancer  
 318 process,  $t^*$  is few minutes and type-II process is predominant [14], whereas in CXL (with  $t^* < 20$   
 319 seconds), type-I is predominant. Therefore, we re-plot a schematic shown in Fig. 3 for a more general  
 320 PPS without specifying  $t^*$ . In the transient stage (with  $t < t^*$ ), both type-I and -II coexist until the oxygen  
 321 is completely depleted; then type-I dominates before the oxygen is resupplied or replenished. In  
 322 general, both type-I and type-II reactions can occur simultaneously, and the ratio between these  
 323 processes depends on the types and the concentrations of PS, substrate and oxygen, and the kinetic  
 324 rates involved in the process [17].



325

326 **Fig. 3** Schematics of the oxygen profiles during the photocrosslinking process; in the transient stage,  
 327 both type-I and -II coexist until the oxygen is depleted. (Figure revised from previous schematics  
 328 [17]).

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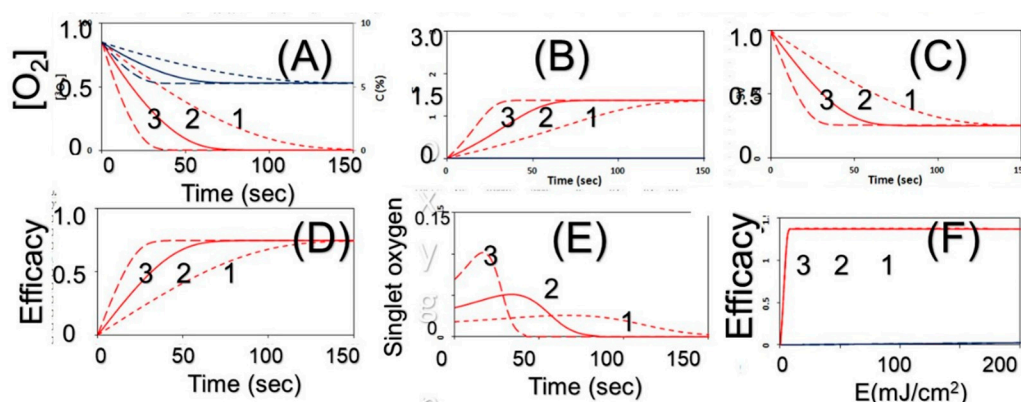
## 330 3.5. Numerical results

331 We will now explore the roles of each of the parameters involved in type-I and type-II  
 332 mechanisms, where the macroscopic coupled Eq. (1) will be numerically solved for various  
 333 parameters of: initial values,  $C_0$ ,  $I_0$ ,  $[O_0]$ ,  $[A]$ ; rate constants  $k'=k_s/k_3$ ,  $k=(k_5+k_s[A])/k_3$ ,  $K'$ , small signal  
 334  $d$ , and the oxygen source term ( $p$ ). The  $S$  functions defined by Eq. (2) will then calculated to obtain  
 335 the associate singlet-oxygen concentration, efficacy,  $Eff=1-\exp(-S)$ , and the cell viability,  $CV=\exp[-$   
 336  $(S_1+S_2)]$ . Typical values are used in our calculations [13,14]: fixed  $[O_0]=7.3(\mu M)$ ,  $k=11.9(\mu M)$ ,  $d=33$   
 337  $(\mu M)$ ; and others will have ranges to show their roles:  $I_0=(50,100,200) \text{ mW/cm}^2$ ,  $C_0=(6,8.5,10) \mu M$ ,  $[A]=$   
 338  $(50,100) \mu M$ ,  $k'=(0.0001, 0.05) 1/s$ . Above various parameters allow us to investigate the roles of  $[A]$ ,  
 339  $I_0$ ,  $C_0$ , the interaction of type-I and type-II (via  $k'$ ) and the oxygen source term ( $p$ ).

340 Fig. 4 shows the calculated temporal profiles of: (A) oxygen (red curves) and PS concentration  
 341 (blue curves), (B)  $S_2$ -function, (C) cell viability, (D) efficacy vs. time, (E) singlet-oxygen, and (F)  
 342 efficacy vs. light dose ( $E_0$ ) for a small  $[A]=50 \mu M$  and  $k'$  value (0.0001), or type-II is dominant, whereas  
 343 Fig. S-3 (in Supporting Information) shows the profiles for a higher  $[A]=100 \mu M$  which leads to a lower  
 344 efficacy. Fig. S-4 (in Supporting Information) shows profiles with external oxygen supply, or  $p>0$ ,  
 345 where the type-II efficacy increases due to the resupply of oxygen, comparing Fig. (D) of Fig.4 and  
 346 Fig.S-4. The role of substrate  $[A]$  is shown by Fig. S-5 (in Supporting Information) that higher  $[A]=100$   
 347  $\mu M$  has a lower type-II efficacy, but higher type-I efficacy. Fig. S-6 (in Supporting Information) shows  
 348 the comparison of type-II  $S$ -profiles for the case without ( $p=0$ ), and with external oxygen ( $p>0$ ). We  
 349 note that the efficacy is governed by the light dose only when  $p=0$  and independent to light intensity;  
 350 in contrast to  $p>0$ , which also shows the intensity-dependence for the transient state.

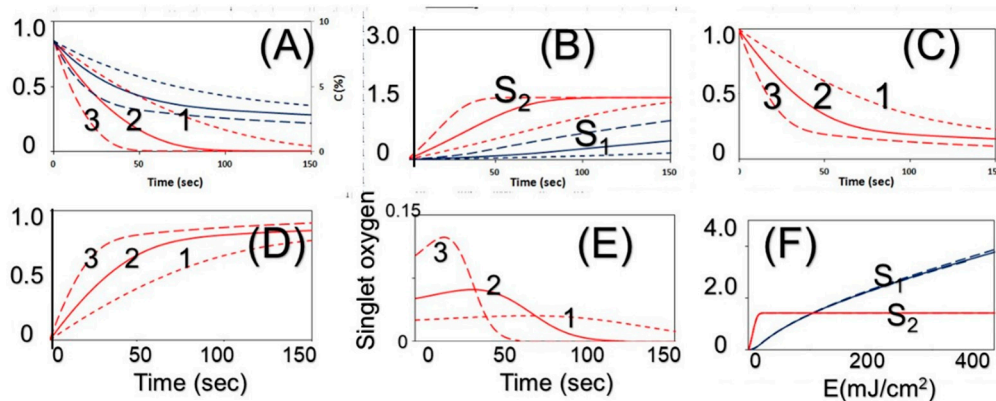
351 In contrast to Fig. 4 with a small  $k'=0.0001$  (for type-II dominant), Fig. 5 shows profiles for a  
 352 higher  $k'=0.05$ , in which type-I and type-II coexist. Comparing to Fig. 4(A), Fig. 5(A) shows more  
 353 depletion of PS concentration,  $C(z, t)$ , due to the combined type-I and type-II, which also shows  
 354 higher efficacy and lower cell viability, as shown by Fig. (C) and (D). We note that, as shown by Fig.  
 355 5 (B) and (F),  $S_2$  reaches its steady state, when oxygen is completely depleted; in contrast,  $S_1$ , is an  
 356 increasing function of PS concentration and does not require oxygen.

357 The singlet oxygen threshold dose, and the dose-dependence cell viability curves after red-light  
 358 irradiation of Radachlorin were reported in vivo by Klimenko et al [32], in which their Fig. 5 may be  
 359 compared with our Fig. (C) of Fig. S-4 (in Supporting Information). Our group has also worked on  
 360 the in vitro measurement of cell viability, which may be empirically analyzed by our numerical data  
 361 (results will be published elsewhere). Resupply of PS or oxygen during the light exposure would  
 362 enhance the overall efficacy. This new strategy has been proposed in type-I corneal crosslinking by  
 363 Lin [29] and type-II anti-cancer by Lin et al [35].



364

365 **Fig. 4** The calculated temporal profiles of: (A) oxygen (red curves) and PS concentration (blue curves),  
 366 (B)  $S_2$ -function, (C) cell viability, (D) efficacy vs. time, (E) singlet-oxygen, and (F) efficacy vs. light dose  
 367 ( $E_0$ ), for various light intensity of 50, 100, 200 mW/cm², (for curves 1,2,3), without external oxygen  
 368 source ( $p=0$ );  $d=33$ ,  $k'=0.0001$  (for type-II dominant) and for substrate  $[A]=50 \mu M$ .



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Fig. 5 Same as Fig. 4, but for a higher  $k'=0.05$  and with external oxygen source ( $p=0.15$ ), in which type-I and type-II coexist. Fig (B) and (F) show both  $S_1$  (blue curves) and  $S_2$  (red curves).

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#### 4. Conclusion

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For the same dose, lower light intensity achieves a higher steady-state-efficacy (SSE) in type-I; in contrast to type-II, which has an equal SSE. Type-II process is also affected by the available oxygen. Higher light intensity produces more efficient singlet oxygen, resulting a higher transient efficacy, in which all intensities reach the same SSE when oxygen is completely depleted. With external oxygen, type-II efficacy increases with time, otherwise, it is governed only by the light dose, i.e., same dose achieves same efficacy. Moreover, type-II has an efficacy follows Bunsen Roscoe law (BRL), whereas type-I follows non-BRL. The photopolymerization dynamics may be defined by the availability of oxygen, where both type-I and -II coexist until the oxygen is depleted. The roles of light dose and PS concentration on the efficacy should be governed a new concept of a volume efficacy ( $V_e$ ), first introduced by Lin [30], defined by the product of the crosslink (or gelation) depth (CD) and local [efficacy].

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**Conflicts of Interest:** Jui-Teng Lin is the CEO of New Vision Inc.

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