

1 Article

2 One Universe, Many Spaces: A Non-Local, 3 Relativistic Quantum Spacetime

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6

7 **Abstract:** The nonlocality of entangled quantum mechanical systems is incompatible with the
8 standard interpretation of special relativity as a single 4D Minkowskian metric spacetime. The
9 difficulty is that the definition of a spacetime interval between any pair of events precludes any
10 form of nonlocal interaction, even the relatively benign non-signaling correlations. By an
11 application of the relativity principle, and the use of the space \longleftrightarrow time symmetry of the Lorentz
12 boost I propose here a reinterpretation of special relativistic spacetime. This new ontology consists
13 of a set of coexisting 3+1D spaces ('framespaces'), each containing unique content in the form of a
14 complex density. These spaces are related by the Lorentz boost, and coupled pairwise in a manner
15 dictated by the Lorentz transformation. The inter-space coupling acting on the spacetime content
16 gives rise to a nonlocal wave phenomenon, which is identified as quantum wave mechanics. The
17 interspace coupling strength is then inversely proportional to Planck's constant. The coexistence of
18 multiple spaces is interpreted as momentum superposition, implying that momentum is the
19 fundamental physical basis of quantum superposition. This new spacetime interpretation of
20 quantum mechanics has many consequences, including explanations of quantum non-locality, the
21 spacetime role of Planck's constant, quantum measurement as a symmetry-breaking process and
22 the redundancy of description of gauge theory.

23 **Keywords:** Special Relativity; Quantum Mechanics; Non-Locality; Planck's constant; EPR.

24

25 1. Introduction

26 There is famously no consensus on the spacetime meaning of quantum mechanics [1]. From
27 the Einstein-Podolsky-Rosen ('EPR') paper in 1935 [2], to Bell's theorem in the 1960s [3], to Aspect's
28 EPR experiment in the 1980s [4] and its recent refinements demonstrating kilometer range
29 correlations [5], it has become increasingly clear that entangled quantum mechanical systems are
30 inherently nonlocal [6,7]. Troublingly, this quantum non-locality is incompatible with the standard
31 interpretation of Minkowskian spacetime as a single 4D metric space, in which a spacetime interval
32 is definable between pairs of events [8]. Moreover, without a theory of quantum spacetime, the
33 fundamental quantum phenomenon of superposition has remained unreconciled with classical 4D
34 relativistic spacetime [9]. Here by appeal to spacetime symmetry we reinterpret relativistic
35 spacetime as a set of coexisting and coupled 3+1D spaces. Assigning a complex density to each space
36 gives rise to a nonlocal inter-space wave phenomenon, identifiable as quantum waves with Planck's
37 constant playing a critical spacetime role governing inter-space coupling. The relativistically
38 necessary coexistence of multiple spaces is identified as the origin of quantum superposition, and
39 offers a physical explanation for the redundancy of description found in gauge theories. One key
40 consequence, impacting the infamous quantum measurement problem, is that momentum is the
41 fundamental physical basis of quantum superposition. This 'many-spaces' proposal represents a
42 unified quantum-spacetime applicable to all atomic-scale quantum phenomena, without any need to
43 appeal to the inaccessible Planck scale. This proposal is both a reinterpretation and, due to the
44 spacetime superposition postulate, a modification of quantum theory.

45

46 It is sometimes suggested that relativistic quantum field theory (QFT) [10] adequately
47 integrates special relativity (SR) with quantum mechanics (QM). However, we argue that despite
48 the successes of QFT, a deeper understanding of quantum mechanics within flat spacetime is still
49 needed. Two issues highlight the friction between QFT and spacetime. Firstly, the nonlocality
50 inherent in the path integral formulation [11] illustrates that QFT has merely inherited, rather than
51 resolved, the quantum conflict with metric-based spacetime locality. Secondly, lacking a
52 spacetime role for Planck's constant, QFT has failed to lead to a theory of quantum gravity.
53 Although semiclassical gravity [12] does incorporate both QFT and gravity, and despite successes in
54 black hole thermodynamics [13], the semi-classical approach is of limited validity because
55 expectation values, rather than the full quantum state, act as the source of spacetime curvature.

56 We argue here that the de facto adoption of a classical-style spacetime ontology, consisting of a
57 single objective, unobservable invariant reality [14], has critically hindered the conceptual
58 unification of spacetime and quantum mechanics. By contrast, a spacetime ontology consisting of
59 multiple subjective observable realities would be much closer aligned to that of the quantum world.
60 Such an approach is not without foundation, bearing in mind that in addition to invariants, relativity
61 theory also encompasses observer-dependent covariant quantities. Accordingly we propose a
62 relativistic spacetime ontology in which the direct observational experience of simple inertial
63 observers is treated as primary. This emphasis on observables over invariants might be captured by
64 the phrase: 'reality is relative'. There are some communalities in approach between this and
65 relational quantum mechanics [15] and also with a 'relationalist' interpretation of spacetime [16].

66 2. A Many-Spaces, Non-Local Spacetime

67 We introduce here a 'many-spaces' interpretation of spacetime, illustrated in Fig.1. Although
68 there is a great deal that might be said in critique of conventional single-space 4D spacetime
69 concepts and in support of this approach, in the interests of brevity the bulk of that discussion will
70 be deferred, but we will outline the arguments in this section.

71 Consider the direct observations made by a set of inertial observers, all in relative motion.
72 Following the non-detection of the ether by Michelson and Morley, each observer - unable to detect
73 any self-motion - directly observes itself at rest within a 3+1D space. Each observer *Alice* will assert
74 that its own directly experienced 3+1D space is physically real. No other observer is in a position to
75 deny that, for Alice at least, Alice's space is real. In this paper we will focus on the interpretation of
76 these multiple realities or views.

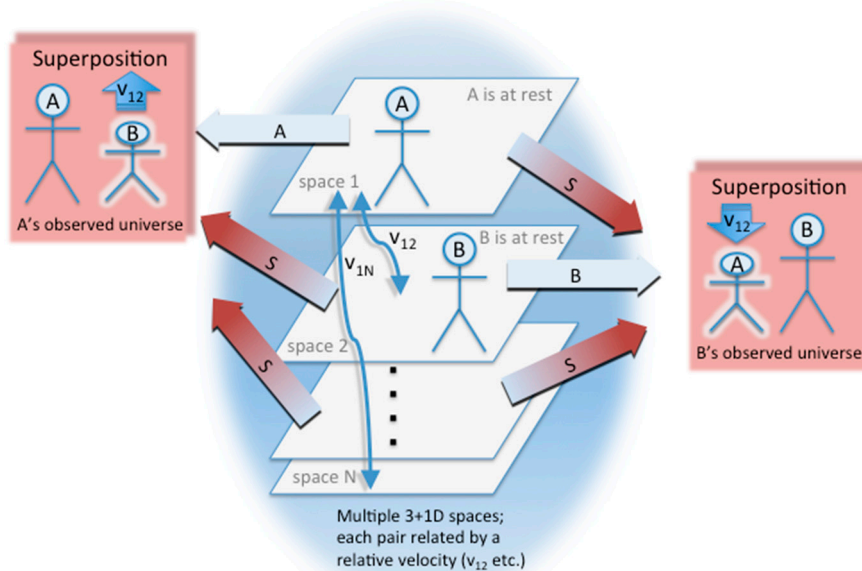
77 It is true that Minkowski successfully showed that within the confines of classical physics a
78 4-dimensional manifold can be used to encompass and relate all the differing views. We will take the
79 position here that despite all the benefits and insights of the 4-dimensional scheme, it does not
80 invalidate Alice's direct experience. The direct experiences of all observers A,B,C and D do not
81 dissolve into a 4-dimensional mist; they must persist. In other words, spacetime ontology cannot be
82 simply a 4-dimensional manifold, but must respect the existence of multiple observed realities.

83
84 The other important caveat of Minkowski's classical scheme is that the move to a 4-dimensional
85 metric space, with a meaningful concept of locality, has unfortunately failed for quantum systems.
86 Quantum nonlocality violates the locality so prized by Minkowski. In retrospect we can see
87 Minkowski's contribution was a regressive move, tethering Einstein's theory to 19th Century
88 absolutists notions, and blinding it to the multiplicity that would become so central to the quantum
89 world.

90 We will take the position here that for a set of coexisting observers, there must be a
91 corresponding set of coexisting physical 3+1D spaces, with each pair related by a relative velocity.

92
93 In this picture, the Lorentz boost represents a transformation between physically distinct 3+1D
94 spaces, and not, as commonly supposed, a mere relabeling exercise of a 4D spacetime [14]. The
95 entirety of this spacetime consists of many coexisting spaces, each containing unique content. These
96 spaces ('framespaces') can be thought of as reference frames made real: pairs are related by the

97 Lorentz boost (like reference frames), however each contains unique content (like spaces). This
 98 expanded spacetime ontology does not conflict with established experimental or theoretical results
 99 of classical relativity, but we will show that it does provide a much more coherent framework for
 100 quantum systems.
 101



102 **Figure 1.** Diagrammatic representation of the many-spaces spacetime for classical systems. The many
 103 3+1D spaces, each pair related by a relative velocity (e.g. v_{12}), are represented diagrammatically as a
 104 stack of 2D planes. There is no motion within any one 3+1D space, and classical relatively moving
 105 inertial observers A and B each reside at rest within one of these spaces. As shown to left and right,
 106 and as indicated by the arrows, each observer's reality is constructed from a superposition of many
 107 3+1D spaces. However, A and B each experience a different version of the superposition of the
 108 underlying multiple spaces. For the classical relativistic case shown, each observes the other as in
 109 motion and Lorentz contracted. Although for classical systems, the familiar special relativistic
 110 4-vector calculus may still be applied, the corresponding ontology is no longer that of a single 4D
 111 spacetime.

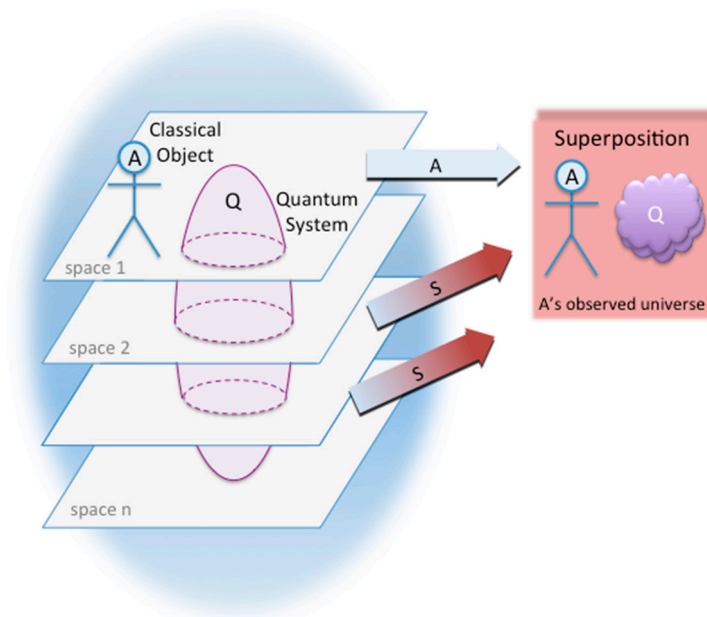
112 *Non-Locality*

113 This many-spaces ontology is able to very naturally accommodate both local classical and
 114 nonlocal quantum systems without any conflict or paradox, as follows. A classical system consists of
 115 a set of objects (inertial observers), each pair possessing a definite relative velocity (Fig.1). This is
 116 interpreted as each object residing (at rest) wholly within its own space. There is no motion within
 117 any one space, however relatively moving objects enter into the universe of an observer by the
 118 superposition of all spaces. Therefore even in classical systems superposition plays a role.

119 Transit of content between spaces corresponds to acceleration, so an inter-space coupling
 120 mechanism is required. It is important to realize that in special relativity, the concept of locality (i.e.
 121 the ability to define a spacetime interval) is completely dependent upon the operation of the Lorentz
 122 boost velocity transformation. Locality is therefore contingent upon the existence of definite relative
 123 velocities. Because classically relative velocities are indeed well defined, classical systems behave
 124 locally.

125 However, the situation is quite different for quantum systems. A quantum system in
 126 momentum superposition lacks definite relative velocities, i.e., coexists within multiple spaces, see
 127 Fig.2. Lacking well-defined relative velocities, such a system behaves non-locally. To allow unitary
 128 quantum evolution, again, an inter-space coupling will be required. Next we show that this picture

129 is more than mere metaphor because the mechanics of a system coexisting within multiple spaces is
 130 indeed quantum.
 131



132 **Figure 2.** Diagrammatic representation of the many-spaces spacetime for a quantum system. Here a
 133 classical object A is shown together with a quantum system Q which is depicted as penetrating
 134 multiple planes, representing its coexistence within multiple 3+1D spaces. On the right hand side, the
 135 observed universe of classical object A is created from the superposition of the content of the
 136 multiple spaces, as indicated by the arrows labeled 'S'. In this superposition, the quantum system Q
 137 is represented as a cloud, indicating quantum properties such as non-locality.

138

139 3. Superposition, Coupling and Wave Behaviour

140 A radical yet compelling interpretation is to equate *coexistence* with *superposition*, i.e. the many
 141 spaces coexist in superposition. This implies that coexistence of spaces is the origin of quantum
 142 superposition, and reinterprets spacetime as an inherently quantum mechanical superposed set of
 143 3+1D spaces, rather than a single 4D classical space. This is a major step in the argument: the claim is
 144 that superposition is a fundamental ingredient in the structure of spacetime, not merely an
 145 expression of 'inexplicable quantum weirdness'. This is an attempt to take superposition seriously in
 146 a spacetime context. Does this claim stand up? We show below that it does.

147 Equating framespace-coexistence with quantum-superposition immediately leads to the
 148 identification of momentum as the fundamental basis of quantum superposition. These spaces are
 149 real, contain unique content, and interact. Specifically, we postulate that each space F_i contains a
 150 complex-valued density function: $\rho_i(x_i, ct_i)$. Spaces are related by a relative velocity therefore the
 151 inter-space interaction must be pairwise and must conform to the two fundamental relativistic
 152 observational conditions [17], namely the simultaneous observation of length ('simultaneous-length'),
 153 and the colocal observation of duration ('colocal-duration'). The duality between these conditions is
 154 a consequence of the space \leftrightarrow time symmetry of the Lorentz boost
 155 [17].

156 *Observation and Coupling between Spaces*

157 Due to the Relativity of Simultaneity and its dual the Relativity of Colocality [17], observation
 158 possesses an inherent directionality in which the observational condition (simultaneity or colocality

159 respectively) – applying to the observer only – physically distinguishes between observer and observed.
 160 This is a relativistic answer to the question: ‘What is an observer?’

161 To be physically meaningful, it is necessary that observation result in physical change for the
 162 observer. Accordingly, we propose an active form of inter-space coupling, arising from the Lorentz
 163 boost’s mixing of spatial and temporal axes, and therefore only in force for non-zero relative
 164 velocity. We show below that this results in wave behavior with coupling strength governed by
 165 Planck’s constant. For brevity, in the following we present one case that illustrates this ‘coupling
 166 postulate’ (see appendix for other cases).

167 *Non-local Waves Behavior*

168 We have postulated that each framespace F_i contains a complex-valued density function:
 169 $\rho_i(x_i, ct_i)$. A simultaneous-length observation made from space F_1 results in a change of observer
 170 density $\delta\rho_1$ equal to the product of a coupling strength (written as ib), the density ρ_2 in the
 171 observed space F_2 , and a temporal interval $c\delta t_2$ in the observed space F_2 :

$$172 \quad \delta\rho_1 = ib\rho_2(c\delta t_2) \quad (1)$$

173 where the temporal interval arises from the relativity of simultaneity $\delta x_1 = -(1/\beta\gamma)c\delta t_2$.
 174 Equation (1) in differential form becomes:

$$175 \quad \frac{\partial\rho_1}{\partial x_1} = -i(b\beta\gamma)\rho_2 \quad (2)$$

176 Analogously to Eq.(2), a colocal-duration observation from space F_2 yields:

$$177 \quad \frac{\partial\rho_2}{c\partial t_2} = +i(a\beta\gamma)\rho_1 \quad (3)$$

178 The parameters a and b are respectively space-to-time and time-to-space inter-space coupling
 179 parameters. By writing a complex density as $\rho = \rho^R + i\rho^I$, the following pair of equations for ρ_2^I
 180 and ρ_1^R can be derived from Eq.(3) and Eq.(2) respectively:

$$181 \quad \frac{\partial\rho_2^I}{c\partial t_2} = (a\beta\gamma)\rho_1^R \quad \text{and} \quad \frac{\partial\rho_1^R}{\partial x_1} = (b\beta\gamma)\rho_2^I \quad (4)$$

182 This equation pair is of the form $\partial U/\partial t_2 = \omega V$ and $\partial V/\partial x_1 = kU$, using the substitutions $U =$
 183 ρ_2^I , $V = \rho_1^R$, $\omega = a\beta\gamma c$ and $k = b\beta\gamma$. Following Bohm [18], U and V are the real and imaginary parts
 184 of a wave function $\psi = U + iV = e^{i(kx_1 - \omega t_2)}$, so in this case we have:

$$185 \quad \psi = \rho_2^I + i\rho_1^R = e^{i\beta\gamma(bx_1 - act_2)} \quad (5)$$

186 This wave function represents a superposition of content within different spaces and so
 187 inherently has nonlocal properties. This form of wave function can satisfy either a Schrodinger type
 188 wave equation:

$$189 \quad i\partial\psi/\partial t_2 = -(ac/b^2\beta\gamma)\partial^2\psi/\partial x_1^2 \quad (6)$$

190 or a 1D wave equation:

$$191 \quad (b/ac)^2\partial^2\psi/\partial t_2^2 = \partial^2\psi/\partial x_1^2. \quad (7)$$

192 These wave equations contain dimensions from different spaces, so are also nonlocal. This
 193 represents a wave-like interaction between framespaces. In summary, since multiple spaces are not
 194 related by a metric, the wave behavior arising from interaction between these spaces is also
 195 non-local. This emergence of nonlocal wave behavior is of course is highly suggestive of quantum
 196 behavior.

197 *Matter Waves*

198 This nonlocal complex wave mechanics has arisen purely from relativistic arguments and an
 199 inter-space coupling postulate, (i.e. without any specifically wave or quantum postulates). The
 200 striking parallels demand the identification of these relativistic waves as indeed quantum waves.
 201 Making this identification by invoking the quantum mechanical formulae for energy: $KE = \hbar\omega =$
 202 $mc^2(\gamma - 1)$, and momentum: $p = \hbar k = \gamma m\beta c$, of a quantum particle of rest mass m , allows the
 203 evaluation of the coupling parameters a and b as:

$$204 \quad a = \frac{mc(\gamma-1)}{\hbar\beta\gamma} \quad \text{and} \quad b = \frac{mc}{\hbar}. \quad (8)$$

205 A pair of equations describing quantum spacetime wave behavior for matter may then be
 206 written:

$$207 \quad \frac{\partial \rho_2}{c \partial t_2} = \frac{imc(\gamma-1)}{\hbar} \rho_1 \quad \text{and} \quad \frac{\partial \rho_1}{\partial x_1} = -\frac{imc\beta\gamma}{\hbar} \rho_2 \quad (9)$$

208 To respect the symmetry between spaces there is also a companion pair with the space labels
 209 reversed, (see Appendix). The full evolution of a quantum system is described by a set of these
 210 equations between all pairs of participating spaces. It is notable that quantum (\hbar, ρ_i) and
 211 many-space spacetime (x_i, t_i, β_{ij}) quantities are directly related. This results in quantum evolution
 212 being expressed explicitly in non-local terms, in contrast to the classical local single-space spacetime
 213 Schrodinger equation formulation.

214 The two-way nature of the Born probability rule arises from this pairwise interaction between
 215 spaces, itself a consequence of the two-way nature of relative velocity (i.e. there is no concept of a
 216 three-way relative velocity).

217 *Electromagnetic Waves*

218 As a second case, by postulating equal coupling: $a = b = p/\hbar\beta\gamma$, the non-dispersive waves at
 219 velocity $\omega/k = c$, occur and may be identified as electromagnetic waves.

$$220 \quad \frac{\partial \rho_2}{c \partial t_2} = \frac{ip}{\hbar} \rho_1 \quad \text{and} \quad \frac{\partial \rho_1}{\partial x_1} = -\frac{ip}{\hbar} \rho_2 \quad (10)$$

221 In this case, the relative velocity between the spaces disappears from the equations. In both
 222 cases the coupling between spaces is inversely proportional to Planck's constant. Quantum behavior
 223 is therefore completely integral to this spacetime picture.

224 **4. Discussion**

225 To conclude, a new relativistic-quantum unification has been achieved without violation of
 226 established principles. Planck's constant plays a fundamental role in an inherently quantum
 227 mechanical spacetime, in which relativistic and quantum concepts are interdependent, rather than in
 228 conflict. The 'many spaces' spacetime describes the coexistence of a set of (3+1)-dimensional spaces,
 229 tightly but not rigidly coupled by a 'quantum glue'. Coexistence of spaces corresponds to
 230 superposition of momentum, which therefore emerges as the fundamental physical basis of
 231 quantum superposition. This implies that momentum measurement is fundamentally different from
 232 the quantum measurement of other observables. The spacetime approach provides general insights
 233 into the nature of quantum measurement. Although classical relativistic observers disagree on
 234 simultaneity and colocality, such separable observers are free to 'agree to disagree'. However in a
 235 non-separable quantum system disagreement becomes unsustainable. Experimental interactions can
 236 then demand the answer to an 'unanswerable question' for which the unitarily evolving quantum
 237 system itself possesses no unique answer. A response can only be provided by a symmetry-breaking
 238 process for which no predictable deterministic outcome is possible. This momentum-based
 239 symmetry-breaking description of the reduction of the quantum state is a new proposal, with
 240 different experimental predictions from other interpretations.

241 To return to issues of locality: EPR non-locality arises because two quantum-entangled particles
 242 in relative motion are a single entity residing within multiple different spaces. Coexistence across
 243 multiple spaces implies non-locality because the concept of definite relative velocity is lost, and with
 244 it the ability of the Lorentz transformation to calculate definite spacetime intervals. A further
 245 consequence of the coexistence of multiple observers is a necessary redundancy of physical
 246 description, which may be directly related to the origin and ubiquity of gauge invariance. In
 247 summary, the recasting of superposition from an abstract principle into a specific physical
 248 phenomenon has far reaching consequences for the interpretation of quantum mechanics and for
 249 quantum measurement. The integration of quantum and spacetime concepts results in a quantum
 250 spacetime providing solutions to many hitherto paradoxical phenomena.

251

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254 design of the study; in the collection, analyses, or interpretation of data; in the writing of the manuscript, and in
255 the decision to publish the results.

256

257

258 **Appendix**

259 Consider the possible interactions between a pair of spaces. Between the two complex
260 functions $\rho_1(x_1, ct_1)$ and $\rho_2(x_2, ct_2)$, representing content in two coupled spaces (F_1, F_2), there are
261 four possible measurement interactions:
262

| Observer | Observation | | |
|----------------|---------------------|------------------------------------------------------------------|------|
| F ₁ | Simultaneous-Length | $\frac{\partial \rho_1}{\partial x_1} = -i(b\beta\gamma)\rho_2$ | (A1) |
| F ₂ | Colocal-Duration | $\frac{\partial \rho_2}{c\partial t_2} = +i(a\beta\gamma)\rho_1$ | (A2) |
| F ₁ | Simultaneous-Length | $\frac{\partial \rho_2}{\partial x_2} = +i(b\beta\gamma)\rho_1$ | (A3) |
| F ₂ | Colocal-Duration | $\frac{\partial \rho_1}{c\partial t_1} = -i(a\beta\gamma)\rho_2$ | (A4) |

263 Table A1: Four possible measurement interactions between two spaces. Note that these are
264 symmetric between spaces: swapping space labels $1 \leftrightarrow 2$ implies $\beta \rightarrow -\beta$.

265 *Observational Compatibility*

266 As shown by the four rows in Table A1, there are four measurement modes: F₁-SL, F₁-CD, F₂-SL,
267 F₂-CD. However not all of these measurement modes are compatible. In a given space, the
268 simultaneous-length (SL) and colocal-duration (CD) observational conditions are incompatible, so
269 only one may apply. The condition for length measurement (simultaneity) precludes the condition
270 for measurement of duration (colocality). Therefore, the F₁-SL and F₁-CD modes are incompatible, as
271 are the F₂-SL and F₂-CD modes. The relativity of simultaneity forbids simultaneity in both spaces.
272 Similarly, the relativity of colocality forbids colocality in both spaces. Therefore F₁-SL and F₂-SL are
273 incompatible, as are F₁-CD and F₂-CD. The only two compatible combinations are therefore the pair
274 F₁-SL & F₂-CD, and the pair F₁-CD & F₂-SD.

275 *Real Densities*

276 It is useful to express the complex densities as a pair of real densities. Each of Eqns.(A1)-(A4),
277 containing complex density functions, may be split into two equations with real density functions,
278 using $\rho = \rho^R + i\rho^I$. Illustrating this for Eq.(A1) (the F₁ observation of simultaneous-length) we
279 have:

$$280 \quad \frac{\partial \rho_1}{\partial x_1} = -i(b\beta\gamma)\rho_2 \quad (A1)$$

$$281 \quad \frac{\partial(\rho_1^R + i\rho_1^I)}{\partial x_1} = -i(b\beta\gamma)(\rho_2^R + i\rho_2^I) = +(b\beta\gamma)(\rho_2^I - i\rho_2^R)$$

282 which splits into the independent pair:

$$283 \quad \partial \rho_1^R / \partial x_1 = +(b\beta\gamma)\rho_2^I \quad (A1a)$$

$$284 \quad \partial \rho_1^I / \partial x_1 = -(b\beta\gamma)\rho_2^R \quad (A1b)$$

285

286 This procedure generates eight equations with real quantities, as shown:

287

| | | | |
|-----------------------|----------------------------------------------------------------------|--------------------------------------------------------------------|-----------------------------------------------------------------------|
| F ₁ Obs SL | $\frac{\partial \rho_1}{\partial x_1} = -i(b\beta\gamma)\rho_2$ (A1) | $\partial \rho_1^R / \partial x_1 = +(b\beta\gamma)\rho_2^I$ (A1a) | $\partial \rho_1^I / \partial x_1 = -(b\beta\gamma)\rho_2^R$ (A1b) |
| | | F ₂ Obs CD | $\frac{\partial \rho_2}{c\partial t_2} = +i(a\beta\gamma)\rho_1$ (A2) |
| F ₁ Obs SL | $\frac{\partial \rho_2}{\partial x_2} = +i(b\beta\gamma)\rho_1$ (A3) | | |
| | | F ₂ Obs CD | $\frac{\partial \rho_1}{c\partial t_1} = -i(a\beta\gamma)\rho_2$ (A4) |

288 Table A2: As shown in the text, pairs of these equations can be combined to exhibit non-local wave
 289 behaviour. One example is the pair of equations (A1a) and (A2b), shown highlighted.

290

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