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2 **Orientating the future bio-macromolecular electron microscopy**

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Fei Sun ^{1,2,3}

4 ¹ Center for biological imaging, Institute of Biophysics, Chinese Academy of Sciences;

5 ² Center for Excellent, National Laboratory of Biomacromolecule, Institute of Biophysics, Chinese

6 Academy of Sciences; Datun Road 15, Chaoyang District, Beijing 100101, China.

7 ³ School of Life Science, University of Chinese Academy of Sciences.

8 Email: feisun@ibp.ac.cn

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1 ABSTRACT

2 With forty years of developments, bio-macromolecule cryo-electron microscopy has met its
3 revolution of resolution and is playing a very important role in structural biology study. According
4 to different specimen states, cryo-electron microscopy (cryo-EM) involves three specific techniques,
5 single particle analysis (SPA), electron tomography and sub-tomogram averaging, and electron
6 diffraction. All these three techniques have not realized their full potentials of solving structures of
7 bio-macromolecules and therefore need to be developed in the future. In this review, the current
8 existing bottlenecks of cryo-EM SPA are discussed with theoretical analysis, which includes air-
9 water interface during specimen cryo-vitrification, bio-macromolecular conformational
10 heterogeneity, focus gradient within thick specimen, and electron radiation damage. Besides,
11 potential solutions of these bottlenecks are proposed and discussed, which are worthy of further
12 investigations in the future.

13 KEYWORDS

14 Cryo-electron microscopy; Air-water interface; Conformational heterogeneity; Focus gradient;
15 Radiation damage.

16

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1 **1. Introduction**

2 Revealing the detailed 3D structure of bio-macromolecules is one of the important steps to
3 understand the life. In recent years, cryo-electron microscopy (cryo-EM) has completed its
4 revolution and been becoming one of the major biophysical techniques to study the 3D structures
5 of bio-macromolecules, especially of membrane protein complexes and supra macromolecular
6 assemblies, giving a big impact on our understanding of the life.

7 The development of cryo-EM technology started from the last century in 1970~1980s when the
8 significant electron radiation damage of biological specimen was discovered and the low-dose
9 illumination technique was proposed [1, 2], the cryo-vitrification method to fix the native structure
10 of biological specimen was established [3, 4], and the image analysis theory to process the low
11 signal-to-noise-ratio (SNR) low-dose cryo-EM micrographs of bio-macromolecules was developed
12 [5]. The potential of electron microscopy to determine the high resolution structure of bio-
13 macromolecules had been revealed by the 3D structure of purple membrane in 1975 [6] and
14 seriously discussed with a theoretical consideration in 1995 [7]. In recent years, with the
15 instrumental advances in electron optics and improved stabilities of electron microscopes, the
16 development of sophisticated image processing software [8-10] and the automation of data
17 collection [11-13], the improvement of phase plate technology [14, 15], especially the success of
18 the direct electron detector [16-18], cryo-EM technology has reached its revolution in resolution
19 [19] with the milestone work that the structure of transient receptor potential cation channel TRPV1
20 was determined to 3.4 angstrom in 2013 [17] and the structure of the glutamate dehydrogenase
21 (GDH) was solved to atomic level in 2016 [20] by cryo-EM. More and more supra bio-

1 macromolecular structures, e.g. spliceosome [21], photosynthetic complex [22], and mitochondrial
2 respiratory complex [23], have been solved to near atomic level, which could not be achieved by
3 using traditional structural biology approaches.

4 According to the states of the biological specimen and the experimental workflow, cryo-EM
5 technologies can be classified into three different techniques, cryo-EM single particle analysis (SPA),
6 cryo-electron tomography (cryo-ET) and sub-tomogram analysis (STA), and cryo-electron
7 diffraction (cryo-ED). Cryo-EM SPA is used to analyze the 3D structure of purified bio-
8 macromolecules in solution, which are cryo-vitrified into a thin ice layer. Ten of thousands of cryo-
9 EM images of the bio-macromolecule are needed to increase SNR and reconstruct a high resolution
10 structure [24]. Cryo-ET can reconstruct the native 3D structure of a local region within a cell or
11 tissue and the *in situ* structure of bio-macromolecules can be further analyzed by STA [25]. There
12 is no need to purify bio-macromolecules from cell or tissue while hundreds of tomograms of cryo-
13 lamella of cell or tissue are needed to obtain a high resolution *in situ* structure of bio-
14 macromolecules. Cryo-ED is a kind of technique to use cryo-electron microscope to analyze the
15 crystallized biological specimen. It collects the electron diffraction data from 2D crystals or 3D
16 nano-crystals of bio-macromolecules and then solves the high resolution structures using
17 crystallographic theories, where the technique to study 2D crystals was used to be called electron
18 crystallography [26] and the one to study 3D nano-crystals emerged recently and called as micro-
19 electron diffraction [27]

20 Cryo-EM SPA has been developed mature enough in past decades and been becoming the major
21 approach in the current structural biology study. The Nobel prize in chemistry in 2017 was also

1 given to Jacques Dubochet, Joachim Frank and Richard Henderson for their great dedicative work
2 in developing cryo-EM SPA technique. Cryo-ET STA has been developed quickly in recent years
3 and will become an important as well as a unique approach to study the *in situ* structure of bio-
4 macromolecules in the future. In this paper, I will discuss the current bottlenecks of cryo-EM SPA
5 and present some potential solutions in my personal view.

6 **2. Current technical bottlenecks in cryo-EM SPA**

7 There have been many good reviews to describe the technique of cryo-EM SPA including the
8 theory, workflow, image processing and applications [28-30]. In brief, cryo-EM SPA starts from the
9 cryo-vitrification of bio-macromolecular solution and needs to collect thousands of high quality
10 cryo-EM micrographs in a high-throughput way with a limited illumination dose ($20 \sim 60 \text{ e}/\text{\AA}^2$
11 normally), a defined magnification yielding a proper pixel size ($0.8 \sim 1.5 \text{ \AA}/\text{pixel}$ normally), and a
12 proper defocus range ($0.8 \sim 3.0 \text{ }\mu\text{m}$). The subsequent image processing includes micrograph
13 correction (motion and distortion correction, dose weighting) and evaluation, contrast transfer
14 function estimation, particle picking and sorting, 2D and 3D classification, orientation refinement
15 and reconstruction, and post-processing (map sharpening). Since the limited illumination dose
16 yields a very noisy raw image of bio-macromolecules that are embedded in vitreous ice, ten
17 thousands of particle images are needed to increase SNR. Thus, the basic assumption behind cryo-
18 EM SPA is that all the analyzed bio-macromolecules should have an identical structure and
19 conformation, which is actually not always true.

20 Starting from 2013 when the near atomic structure of TRPV1 was first solved [17], the number
21 of bio-macromolecular structures studied by cryo-EM SPA increased quickly. Till 2017, there are

1 already 5541 cryo-EM maps deposited in electron microscopy database (EMDB,
2 <http://www.ebi.ac.uk/pdbe/emdb>) [31] while there are only 1566 entries in 2012. Most of the
3 deposited maps were obtained by cryo-EM SPA. However, only a small portion of the map
4 (313/5541) can reach a resolution higher than 4 Å and only a few structures can be solved to the
5 resolution higher than 3 Å (**Figure 1**), which actually forms a barrier for cryo-EM SPA to be widely
6 applied into pharmaceutical industry.

7 Several bottlenecks are still there in cryo-EM SPA from sample preparation, data collection to
8 image processing, which forms the barrier to reach a better resolution. In cryo-vitrification during
9 sample preparation, the existence of air-water interface increases the possibility of disassociation
10 and denature of bio-macromolecules, which prohibited many fragile macromolecular complexes to
11 be studied by cryo-EM SPA. The existence of intrinsic conformational flexibility of bio-
12 macromolecules rules out the basic assumption upon the identical structure and conformation in
13 cryo-EM SPA, which has been restricting to approach high resolution. During cryo-EM imaging,
14 the SNR from the current instrument and hardware is still not enough to study the high resolution
15 3D structures of bio-macromolecules when their molecular weight is smaller than 60 kDa. The phase
16 plate and direct detection detector technologies can be improved better to increase the current cryo-
17 EM SNR. The cryo-electron microscope can be made more stable, easy to use and dedicative for
18 the cryo-EM SPA workflow with much improved throughput. When the size of bio-macromolecule
19 becomes large or the thickness of biological material increases, the existence of the Ewald sphere
20 effect and focus gradient will limit the approachable resolution of cryo-EM imaging in the current
21 theoretical framework, which needs to be solved in both experimental and image processing

1 procedures. The ultimate bottleneck of cryo-EM SPA is the physical nature of radiation damage of
2 bio-macromolecules during cryo-EM imaging, which can not be avoided and is the key limitation
3 of achieving atomic resolution in cryo-EM SPA.

4 Glaeser has made a very good discussion of the above bottlenecks and given thoughtful
5 perspectives [32]. Here, I would like to focus on cryo-vitrification, conformational heterogeneity,
6 thick sample, and electron radiation damage, further discuss these limitations in theory, and propose
7 new solutions to solve these bottlenecks.

8 **2.1 Cryo-vitrification and air-water interface**

9 The current cryo-vitrification method was originally invented by Dubochet's group in 1984 [3],
10 which is called plunge freezing. The EM grid coated with a carbon supporting film is pre-treated
11 with plasma cleaner and then nipped by a fine tweezer and mounted onto the plunge freezer. 3~5 μ l
12 bio-macromolecule solution is dropped onto the supporting film. After a few seconds of incubation,
13 most of the liquid is blotted with filter paper, leaving a thin layer of the solution (30 ~ 50 nm
14 thickness) on the grid. Subsequently, the grid is then quickly plunged into liquid ethane that is pre-
15 cooled by liquid nitrogen, resulting the vitrification of the thin solution layer. Finally, the bio-
16 macromolecules are embedded in the vitreous ice with their native structure preserved. The
17 procedure of this cryo-vitrification method did not change much since its invention and is still now
18 widely used in the world for cryo-EM SPA. Several vendors made commercialized instruments
19 (Thermo Fisher Vitrobot, Leica EM GP and Gatan CP3) for cryo-vitrification, which increase the
20 throughput and reproducibility by controlling experimental parameters (e.g. humidity, temperature,
21 blotting time) accurately.

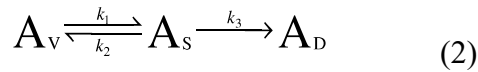
1 However, more and more labs found that the homogeneity of the bio-macromolecule became
2 significantly decreased from cryo-vitrified sample when compared with the negatively stained
3 sample. The worst case is that although the specimen shows even distribution with homogenous size
4 and shape from the negative stain electron micrographs, few particles could be identified/recognized
5 from the cryo-EM micrographs (**Figure 2A**). The reason of this phenomena has been fully discussed
6 by Glaeser and colleagues [33, 34] and now widely recognized to be the effect from the air-water
7 interface (**Figure 2B**).

8 During plunge freezing, the thin solution layer after filter paper blotting results a very large
9 surface-to-volume ratio ($\sim 20 \text{ um}^{-1}$) in comparison with its original value of ($\sim 2.0 \text{ e-3 um}^{-1}$). Thus
10 the bio-macromolecule in the thin layer of the solution has large opportunity to reach the surface,
11 which has been observed from the recent Cryo-ET study that $\sim 90\%$ of particles were absorbed to
12 the air-water interface [35]. The movement of the bio-macromolecule $\sqrt{x^2}$ follows the Brownian
13 motion law according to the following equation,

$$14 \quad \frac{x^2}{t} = \frac{k_B T}{3\pi\eta r} \quad (1)$$

15 where, t is the time of the motion, r is the radius of the bio-macromolecule, η is the coefficient
16 of viscosity of the solution (for water at 10°C , it is $1.308 \text{ mPa}\cdot\text{s}$), k_B is the Boltzmann constant,
17 and T is the temperature in Kalvin. For the thickness of 50 nm of the layer and a 20 nm diameter
18 of the bio-macromolecule, the averaged time (at $T = 283 \text{ K}$) for the bio-macromolecule reaching
19 the surface can be estimated as $\sim 20 \text{ ms}$. The time can be even shorter ($\sim 6 \text{ ms}$) for a thinner ice (40
20 nm thickness) and smaller bio-macromolecule (10 nm diameter). As a result, upon the formation of

1 the thin layer of the solution, the bio-macromolecules can quickly approach the surface called the
2 air-water interface, which has the large possibility to induce denaturation of the bio-macromolecule
3 [34]. The thermodynamics of this procedure can be further described as below (**Figure 2B**),



4
5 where A_V represents the concentration of bio-macromolecule in the solution, A_S represents the
6 concentration of bio-macromolecule bound to the air water interface, and A_D represents the
7 concentration of denatured bio-macromolecule. Thus, the speed of bio-macromolecule denaturation
8 is determined by the reaction constants k_1 , k_2 , and k_3 . The molecular dynamics simulation to
9 study the absorption behavior of lysozyme onto the hydrophobic surface of graphite suggested that
10 the secondary structures of lysozyme disappear at 10 ns after bound to the surface [36]. Thus, the
11 order of k_3 can be estimated of $1 \times 10^7 \sim 1 \times 10^8 \text{ s}^{-1}$, the denaturation of the bio-
12 macromolecule can occur less than 1 μs .

13 During plunge freezing, the time normally takes seconds from the completion of blotting to
14 plunging into liquid ethane, which are much longer than the one for bio-macromolecules
15 approaching to the air-water interface and denatured. Therefore, it can be explained why it was
16 difficult to obtain good cryo-EM micrographs for those fragile bio-macromolecules that are easily
17 denatured (or dissociated) although they can be well captured in negative stain electron micrographs
18 (**Figure 2A**).

19 The existence of air-water interface has become one of the most important bottlenecks for cryo-
20 EM SPA to obtain high resolution structures of many bio-macromolecules. To overcome this

1 bottleneck, people have developed multiple ways to reduce the air-water interface by utilizing an
2 additional ultrathin ($\sim 2\text{nm}$) carbon film [34] or developing affinity grids [37, 38]. Adding surfactant
3 into the solution could be also useful to form a ‘cover-slip’ at the air-water interface and therefore
4 protect the bio-macromolecules from denaturation [39]. In addition, the new instruments using
5 automatic robotics have been recently developed, e.g. Spotiton [40] and ‘grid writer’ [41], which
6 avoid paper blotting and can minimize the time between the formation of thin layer and the plunge
7 freezing to ~ 500 ms, which therefore greatly reduce the possibility of the molecule approaching to
8 the surface. In the recent study, a new Spotiton robot was reported and the procedure was optimized
9 to reduce the exposure time in the air-water interface to $100 \sim 200$ ms, which effeciently reduced
10 the number of particles reaching the air-water interface and improved the final reconstructed map
11 [42].

12 According to the equations (1) and (2), there are also other ways to reduce the effect of air-water
13 interface. Using a high concentration of bio-macromolecules during vitrification may increase the
14 coefficient of viscosity and thus increase the time of approaching to the surface. Meanwhile, the
15 high concentration can also have a chance to saturate the air-water interface with the denatured
16 molecules and thus allow enough number of native molecules left in the solution. To be noted that,
17 the high concentration could increase the difficulty in the following particle picking and image
18 processing. Adding a proper chemical reagent to reduce the interaction between the molecule and
19 the air-water interface would be an alternative way. In such case, the reaction constant k_1 in
20 equation (2) is much smaller than k_2 , and thus large portion of molecules are kept in the solution.
21 Using chemical cross-linker to increase the structural stability of bio-macromolecules would be an

1 additional way that the denaturing reaction constant k_3 is reduced.

2 **2.2 Conformational heterogeneity**

3 Cryo-EM SPA assumes that all the bio-macromolecules analyzed in the electron micrographs
4 have an identical structure and conformation. However, this assumption is not rigorously true in
5 most cases because of the unavoidable thermodynamics of bio-macromolecules and would become
6 much worse when there are heterogeneities existing in the sample.

7 There are two kinds of heterogeneities of bio-macromolecules, the composition heterogeneity
8 and the conformational heterogeneity. The composition heterogeneity refers to the specimen
9 composed of a mixture of molecules with different ligand-bound states [43], a mixture of
10 macromolecular complexes with different subunits stoichiometry, a mixture of macromolecular
11 assembly with different symmetries, or even in a worst case a mixture of the target molecule and
12 contaminations. The conformational heterogeneity refers to the specimen containing the target bio-
13 macromolecules but in different functional states, and can be further divided into two cases, the
14 heterogeneity with discrete conformations [44] and the one with continuous conformations [45].

15 The existence of heterogeneity will add difficulty in image processing of SPA and prevent
16 achieving high resolution structure. The composition heterogeneity can be possibly and efficiently
17 solved by improving biochemical preparation procedure, e.g. more specific affinity chromatography.
18 In addition, recent image processing algorithms have been well developed by applying sophisticated
19 statistics tools, e.g. principal component analysis [46, 47], maximum log likelihood [48] and
20 Bayesian method [49]. Utilizing these advanced image processing tools, we are able to perform

1 efficient image classification to separate bio-macromolecular particles with different compositions
2 and solve the composition heterogeneity [50].

3 The conformational heterogeneity reflects the functional and thermodynamics nature of bio-
4 macromolecules, which could not be easily improved by conventional biochemical approaches. For
5 the case of discrete conformations, the current image classification algorithms can work very well
6 if SNR of the particles are large enough to discriminate the different conformations. However, for
7 the case of continuous conformations, it would be difficult to improve the reconstruction resolution
8 by carrying out 3D classification approaches unless incredible large number of particles are
9 collected. There have been a few image processing approaches developed to try to solve the
10 heterogeneity of continuous conformations, including local optimization refinement [51], masked
11 refinement [52], multi-body refinement [53], particle segmentation on micrograph [54], normal
12 mode method [55] and manifold-embedding method [56]. When using local optimization or masked
13 refinement approach, we assumed that the bio-macromolecular particle can be divided into a number
14 of rigid parts and the relative orientations and positions of different rigid parts contribute to the
15 flexibility of the molecular complex. This assumption worked in many cases to improve the
16 resolution of bio-macromolecular flexible modules [21]. However, it is not always true and the
17 internal conformational changes of different modules should be also considered in many cases.

18 The recent proposed normal mode and manifold-embedding methods would be good solutions
19 to study the intrinsic conformational dynamics of bio-macromolecules directly from the raw cryo-
20 EM images of bio-macromolecular particles. The normal mode method first performs atomization
21 of cryo-EM map and then calculates various normal modes of clustered pseudo atoms. The specific

1 modes are selected to simulate cryo-EM maps with continuous conformations, which are then
2 compared with the raw cryo-EM images [55, 57, 58]. The normal mode method has been
3 successfully applied to study the structure of transcription pre-initiation complex recently [59].

4 The manifold-embedding method maps each projection of bio-macromolecule into a point of
5 hyperspace ($N \times N$ dimension, N is the size of projection image). All the points of bio-
6 macromolecules with different orientations and conformations will build a manifold in this
7 hyperspace. The dimension of this manifold is determined by the degree of freedom of bio-
8 macromolecular motion including rotation (3 freedoms), shift (2 freedoms) and conformational
9 changes (various freedoms). Manifold-embedding is a kind of mathematical approach to estimate
10 the degree of freedom of the manifold and decompose these freedoms into different principal
11 coordinates. After decomposition, a specific coordinate can be selected and sorted to reconstruct the
12 conformational changes of bio-macromolecules [56, 60-62].

13 The recent published review [63] has discussed various image processing algorithms to solve
14 the conformational heterogeneity, especially the continuous conformation problem. Besides the
15 image processing approaches, here, I would like to propose another biochemical approach to reduce
16 the conformational heterogeneity of bio-macromolecules.

17 According to the Boltzmann's distribution law, the number of bio-macromolecules at a specific
18 state, N_c , is proportional to $e^{-\frac{E_c}{k_B T}}$, where E_c is the Gibbs energy of the bio-macromolecule in this
19 state c , and k_B is the Boltzmann constant (8.62×10^{-5} eV/K), T is the temperature of the bio-
20 macromolecule solution. Suppose the lowest Gibbs energy of the state is E_L (normally this state is

1 called steady state and it is not degenerative, i.e., only one conformation corresponds to this state),
2 and the highest Gibbs energy of the state is E_H (normally this state is degenerative, i.e., only
3 multiple conformations correspond to this state), the ratio of numbers of molecules between these
4 two states can be determined as,

$$5 \quad r(T) = \frac{N_L}{N_H} = e^{\frac{E_H - E_L}{k_B T}} = e^{\frac{\Delta E}{k_B T}} \quad (3)$$

6 The Gibbs energy difference ΔE between two states of bio-macromolecules is 40~90 meV [45, 62].
7 Then the ratio of the numbers can be estimated as, $r(298K) = 4.7 \sim 33.2$ (room temperature),
8 $r(277.6K) = 5.3 \sim 43.0$ (4 degree), $r(253.6K) = 6.2 \sim 61.4$ (-20 degree), $r(193.6K) =$
9 $11.0 \sim 219.9$ (-80 degree). Therefore, if we could use chemical cross-linker to fix the steady state
10 at the low temperature (e.g. -20 or -80 degree) before vitrification, there will be more populations
11 of homogeneous bio-macromolecular particles of the steady state in the cryo-EM images, which
12 provides an alternative approach to solve the conformational heterogeneity. To utilize this approach,
13 we need to add glycerol or other cryo-protectants into the bio-macromolecular solution in order to
14 keep the solution in liquid state at low temperature, so that the thermodynamics equilibrium can be
15 reached. Then the chemical cross-linker is added to the cooled solution to allow the cross-linking
16 reaction occur. A good cross-linker needs to be screened and optimized to allow an efficient and fast
17 reaction at low temperature.

18 **2.3 Thick specimen and focus gradient**

19 The current image processing procedure of cryo-EM SPA assumes that the specimen is thin
20 enough so that the dynamic scattering, the Ewald sphere and focus gradient effects can be neglected.

1 For 300 kV accelerate voltage and 100 nm thickness of vitrified bio specimen, the dynamic
2 scattering effect can be still neglected because it is smaller than the mean free path (~ 350 nm) of
3 300 keV electron for the vitrified bio specimen [64]. However, the Ewald sphere effect has limited
4 the resolution to 3.8 angstrom according to the formula of $\sqrt{t * \lambda / (2 * 0.7)}$, where t is the
5 specimen thickness and λ ($= 0.02 \text{ \AA}$) is the wavelength of 300 keV electron [65]. The 100 nm
6 focus gradient will induce a phase error of $\frac{\pi}{2}$ at the resolution of 6.3 angstrom according to the
7 formula of $\Delta\chi = \pi\lambda\Delta Zs^2$, where $\Delta\chi$ is the phase error and ΔZ is the focus gradient [66]. Thus,
8 when the size of bio-macromolecule or the vitrified ice is thick, the Ewald sphere, especially the
9 focus gradient will take effects and should be corrected to improve the resolution. In cryo-EM, the
10 Ewald sphere and focus gradient effects are combined and corrections of these two effects are
11 actually equivalent [65].

12 Theory of Ewald sphere correction has been well developed [65] and various algorithms have
13 been implemented into difference programs and tested with simulated data [67, 68] and recent
14 experimental data [69]. A recent block-focused algorithm was proposed [70] and proved efficient to
15 solve the structure of a gigantic herpesvirus capsid (~ 125 nm in diameter) to 3.1 \AA [71]. The effect
16 of focus gradient was also carefully discussed recently using simulated data [72]. Here, I will discuss
17 the focus gradient effect in a different way and provide a new approximation to solve and correct
18 this effect.

19 As shown in **Figure 3**, suppose the thickness of the specimen is D , the underfocus of the
20 proximal side is f_0 , and that of the distal side is $f_0 + D$, the averaged underfocus of the specimen
21 is $f_a = f_0 + D/2$. For the thin specimen, the final image can be simply formulated as,

$$1 \quad I(x, y) = p(x, y) \otimes PSF(f_a, x, y) \quad (4)$$

2 where, $p(x, y)$ is the projection of the structural density $f(x, y, z)$ of the specimen,

$$3 \quad p(x, y) = \int_{f_0}^{f_0+D} f(x, y, z) dz \quad (5)$$

4 and $PSF(f_a, x, y)$ is the point spread function of the objective lens, and is the Fourier
5 transformation of contrast transfer function,

$$6 \quad CTF(f_a, s_x, s_y) = \text{Sin}\chi(f_a, s) = \text{Sin}\left(\frac{\pi}{2} \lambda C_s^3 s^4 - \pi \lambda f_a s^2\right) \quad (6)$$

7 For simplicity, here we do not consider astigmatism and amplitude contrast and therefore $s^2 =$
8 $s_x^2 + s_y^2$.

9 However, when the thickness D is large enough, equation (4) has to be corrected by dividing
10 the specimen into a series of thin specimen ($n=0, 1, 2, \dots, N-1$), thus the final image of a thick
11 specimen can be formulated as,

$$12 \quad I(x, y) = \sum_{n=0}^{N-1} \Delta z * f(x, y, f_0 + n\Delta z) \otimes PSF(f_0 + n\Delta z, x, y) \quad (7)$$

13 When $N \rightarrow +\infty$, we have,

$$14 \quad I(x, y) = \int_{f_0}^{f_0+D} f(x, y, f_0 + z) \otimes PSF(f_0 + z, x, y) dz$$

$$15 \quad = \int_{-\frac{D}{2}}^{\frac{D}{2}} f(x, y, f_a + z) \otimes PSF(f_a + z, x, y) dz \quad (8)$$

16 Considering $D (\sim 100 \text{ nm}) \ll f_a (1000 \sim 2500 \text{ nm})$, the point spread function can be expanded
17 at $z = 0$ and approximated as,

$$1 \quad PSF(f_a + z, x, y) = PSF(f_a, x, y) + z * \frac{\partial PSF}{\partial z} \Big|_{z=0} \quad (9)$$

2 Combining (8) and (9), we have,

$$3 \quad I(x, y) = I^0(x, y) + I^1(x, y) \quad (10)$$

$$4 \quad I^0(x, y) = \int_{-\frac{D}{2}}^{\frac{D}{2}} f(x, y, f_a + z) dz \otimes PSF(f_a, x, y) \quad (11)$$

$$5 \quad I^1(x, y) = \int_{-\frac{D}{2}}^{\frac{D}{2}} f(x, y, f_a + z) z dz \otimes \frac{\partial PSF}{\partial z} \Big|_{z=0} \quad (12)$$

6 Next, we consider Fourier transformation of the structural density $f(x, y, z)$ of the specimen,

$$7 \quad F(s_x, s_y, s_z) = \iint \int_{-\frac{D}{2}}^{\frac{D}{2}} f(x, y, f_a + z) * e^{-i\pi(xs_x + ys_y + zs_z)} dz dx dy \quad (13)$$

8 Thus, we have the projection theorem,

$$9 \quad F(s_x, s_y, 0) = \iint \int_{-\frac{D}{2}}^{\frac{D}{2}} f(x, y, f_a + z) dz * e^{-i\pi(xs_x + ys_y)} dx dy$$

$$10 \quad = \iint p(x, y) * e^{-i\pi(xs_x + ys_y)} dx dy \quad (14)$$

11 and the following relationship,

$$12 \quad \frac{\partial F}{\partial s_z} \Big|_{s_z=0} = \int f(x, y, f_a + z) * e^{-i\pi(xs_x + ys_y + zs_z)} * (-i\pi z) dx dy dz \Big|_{s_z=0}$$

$$13 \quad = \iint \frac{\pi}{i} \int_{-\frac{D}{2}}^{\frac{D}{2}} f(x, y, f_a + z) z dz * e^{-i\pi(xs_x + ys_y)} dx dy \quad (15)$$

14 As a result, combining (14) and (15) and utilizing the convolution theorem, the Fourier
15 transformation of equation (10) becomes,

$$1 \quad \hat{I}(s_x, s_y) = F(s_x, s_y, 0) * CTF(f_a, s) + \frac{i}{\pi} \frac{\partial F}{\partial s_z} \Big|_{s_z=0} * \frac{\partial CTF(f_a+z, s)}{\partial z} \Big|_{z=0} \quad (16)$$

2 Combining (6) and (16), we have,

$$3 \quad \hat{I}(s_x, s_y) = F(s_x, s_y, 0) * \text{Sin } \chi(f_a, s) - i\lambda \left(\frac{\partial F}{\partial s_z} \right)_{s_z=0} * \text{Cos } \chi(f_a, s) * s^2 \quad (17)$$

4 From (17), it is clear to see, for the thick specimen, the Fourier transformation of cryo-EM image
 5 contains an additional term $i\lambda \left(\frac{\partial F}{\partial s_z} \right)_{s_z=0} * \text{Cos } \chi(f_a, s) * s^2$. This term has a minor contribution at
 6 the low resolution but will interfere significantly with the first term and induce Thon ring fading out
 7 at the high resolution.

8 Equation (17) provides at least two ways to solve and correct the focus gradient effect. For the
 9 first way, we can take two cryo-EM images of the same specimen with different underfocus, yielding
 10 two equations. The underfocus parameters can be accurately estimated by fitting Thon ring and
 11 using the low resolution data. Then the structural factor $F(s_x, s_y, 0)$ can be solved directly. The
 12 concern of reduced SNR from additional radiation damage of the second exposure has been
 13 discussed carefully and can be eliminated [65].

14 The second way does not need two experimental images but apply an iterative algorithm to solve
 15 the structural factor. Initially, the second term of (17) is neglected and normal cryo-EM SPA
 16 procedure is performed to get the first round of structure. Then $\frac{\partial F}{\partial s_z}$ can be computed and will be
 17 used to update the structural factor $F(s_x, s_y, 0)$ according to (17). Therefore, the structure of cryo-
 18 EM map can be reconstructed again with an improved resolution. This procedure can be iterated
 19 several rounds until the convergence reaches.

1 **2.4 Beam induced motion and radiation damage**

2 Cryo-EM of bio-macromolecules embedded in vitreous ice has suffered from beam induced
3 motion (BIM) for many years. When accelerated high-energy electrons interact with the specimen,
4 the electrons from the specimen will be scattered and become secondary electrons coming out of
5 the specimen, leaving positively charged annulus at the illumination area [73]. This is called
6 charging effect ('Berriman effect') from electron beam illumination. Since the thin layer of vitreous
7 ice is an insulator, the positive charges can not be quickly compensated from the environment and
8 thus induce the subsequent physical effects. First, the irregular and metastable structure of the
9 vitrified ice can easily respond the internal electrostatic repulsion stress from the positive charges
10 and thus result a global mechanical deformation. Such global mechanical deformation is much
11 significant at the beginning of electron illumination [16] and was observed to have a dome-like
12 shape [74]. The mechanical deformation of the ice layer during electron illumination will result a
13 blurred cryo-EM image and weaken the high resolution information significantly. Second, the
14 positive charged annulus will induce additional phase shift of electron beam just like a microscopic
15 electrostatic lens, which will induce an additional contrast loss and blurring of the final image [73],
16 which, however and fortunately, has the minimal effect on the high resolution cryo-EM SPA
17 technique according to the recent study [75].

18 Besides charging effect that causes BIM, electron beam illumination will induce another more
19 severe effect, which is called radiation damage (or radiolysis). When the secondary electron is
20 kicked out, the chemical covalent bond of molecule will be broken, generating many free-radicals.
21 The free-radicals can migrate quickly and react with the adjacent molecules. As a result, the structure

1 of bio-macromolecule can be damaged effectively from the electron beam illumination and the high
2 resolution structural information will be damped significantly when the illumination dose increased
3 [2, 76]. Thus, electron radiation damage must be carefully considered and the electron dose should
4 be carefully controlled for high resolution cryo-EM SPA work.

5 The electron radiation damage of vitreous ice embedded bio-macromolecules will further cause
6 the ‘bubbling’ effect that is routinely observed in cryo-EM experiments. Since both water molecules
7 and bio-macromolecules contain abundant hydrogen atoms, electron radiolysis will generate large
8 amount of hydrogen free-radicals and these hydrogen free-radicals will react subsequently to form
9 hydrogen molecule [77]. At the interfaces between bio-macromolecules and ice or the support
10 carbon film and ice, hydrogen molecules are frequently accumulated into a high concentration and
11 form the hydrogen gas pocket, which was observed as the ‘bubbling’ effect [78]. The generation of
12 hydrogen gas pocket will produce additional mechanical stress within the cryo-vitrified specimen
13 and thus become another factor of BIM. By using a very low electron radiation dose-rate, the
14 accumulation of hydrogen gas can be effectively decreased and thus the ‘bubbling’ effect can be
15 alleviated [78]. To be noted, the ‘bubbling’ effect used to be developed as a kind of technique to
16 study the internal nucleic structure of virus [79].

17 Recently, another kind of BIM effect, called beam-induced Brownian motion, was proposed
18 and studied [80]. This effect describes a pseudo-Brownian motion of vitreous ice embedded bio-
19 macromolecules, which is generated from the beam-induced movement of water molecules.
20 Fortunately, the experimental data and simulation study by Henderson and coworkers suggested that
21 this beam-induced Brownian motion has the minimal effect on the current cryo-EM SPA work unless

1 the size of bio-macromolecule is small and the target resolution goes beyond 2 angstrom [80].

2 The existence of BIM and electron radiation damage has been aware of the key bottleneck of
3 high resolution cryo-EM SPA for many years until the emerging of the direct electron detector (DED)
4 [81]. The high detective quantum efficiency (DQE) of the DED camera [82] allows to retain the
5 high resolution weak signal under low dose electron radiation that is important to reduce radiation
6 damage of bio-macromolecules. Equally importantly, the CMOS architecture of DEM camera
7 enables a high frame rate to record a single exposure into a movie, which can be utilized to correct
8 BIM efficiently by applying appropriate image processing algorithms [16, 74, 83]. In addition, by
9 using dose fractionation and damage compensation algorithms, the movie mode of DED camera can
10 further allow to use a high illumination dose to take cryo-EM images for a better contrast [76].

11 Besides the direct electron detector that can correct BIM during image processing procedure,
12 additional efforts have been also performed to alleviate BIM effect and these include spot scan
13 imaging [84, 85], paraxial charge compensator [86], and development of various supporting films
14 such Cryo-Mesh grid, graphene film and pure gold grid [87-89]. The pure gold grid was proved to
15 have sufficient mechanical stiffness and good conductivity, which can therefore reduce BIM [88]
16 and has been successfully applied in many high resolution cryo-EM SPA applications [90].

17 Overall, the efforts of developing direct electron detector, motion correction algorithms and
18 new types of supporting films in recent years have significantly reduce the effects from BIM and
19 electron radiation damage. More and more bio-macromolecular structures are solved to near atomic
20 resolution (3 ~ 4 angstrom) by cryo-EM SPA approach and there were a few cases reaching sub-2
21 angstrom resolution [20, 69]. However, the electron radiation damage effect is still existing and will

1 be the most important bottleneck of cryo-EM SPA in the future to achieve atomic resolution.
2 Previous studies have showed that the high resolution (~ 3 angstrom) information of vitreous
3 embedded biological specimen starts falling off after a low dose ($3 e / \text{\AA}^2$) of electron radiation [2].
4 The severe mechanical deformation of ice layer at the first dose fractionated frames could not be
5 corrected by image processing algorithms [16]. As a result, while the first couple of frames with less
6 radiation damage contains atomic resolution information, this information can not be restored due
7 to a large BIM and thus these frames have to be discarded in the subsequent image processing.

8 In the future, there will be two potential ways to further alleviate the electron radiation damage
9 effect. The first possible way is to utilize quantum entanglement effect of electrons to reduce the
10 shot noise of electron beam from the normal scale $\sim 1/N^{1/2}$ to the Heisenberg limit $\sim 1/N$ [32,
11 91]. As a result, we could utilize an even low electron dose ($1 e / \text{\AA}^2$) to capture a good image with
12 enough SNR and less radiation damage [91]. The other possible way is to consider the time scale of
13 electron radiation damage. If we could take cryo-EM image before the specimen damage occurs,
14 we thus could have an opportunity of obtaining the damage-free and high resolution structure of
15 bio-macromolecule. This idea has been proved in the field of serial femtosecond X-ray
16 crystallography with the term of ‘diffraction before damage’ [92]. In cryo-EM of bio-
17 macromolecules, it is important to estimate the time scale of specimen damage from electron
18 radiation and then verify the possibility of the term of ‘imaging before damage’.

19 For 300kV transmission electron microscope, the accelerated electron gains a high speed $v =$

20 $c \sqrt{1 - \frac{1}{(1+E/E_0)^2}} = 0.78c = 2.3 \times 10^8 \text{ m/s}$, where the relativistic effect has to be considered, the

1 static energy of electron is $E_0 = m_0c^2 = 511keV$, the kinetic energy is $E = 300keV$ and the light
2 speed in vacuum is $c = 3 \times 10^8 m/s$. Then, the time for an electron traveling across a specimen
3 with the thickness of $d = 100 nm$ is $\Delta t = \frac{d}{v} = 0.33 fs$. Thus, we could estimate the time scale of
4 the interaction between the specimen and the high energy electron is $\sim 1 fs$. Then there will be
5 many damage events occurred in the biological specimen, which can be divided into two processes,
6 the primary damage process and the secondary damage process [93, 94]. The primary damage
7 process includes chemical bond breaking, ionization, and production of secondary electrons and free
8 radicals. Previous studies suggested that the time scale of the primary damage process is $1 \sim 10 ps$
9 and such damage does not influence the electron microscopic image appreciably [95], because the
10 positions of atoms do not change noticeably at this time scale. The only detectable damage in the
11 electron microscopic image occurs in the secondary process, which initiates from the transition of
12 free radicals and includes subsequent cascade reactions induced by free radicals and productions of
13 new chemical bonds. During the second process, the positions of atoms in the specimen have
14 changed significantly, resulting appreciable damage effect in the final electron microscopic image.
15 The time scale of the second process depends on the rate of free radical transition, which can be
16 estimated as below.

17 As discussed above, the abundant free radicals generated from vitreous ice embedded bio-
18 macromolecules are hydrogen free radicals H^* . The transition of H^* follows the Fick's laws of
19 diffusion,

$$20 \quad J = vC = -D \cdot \frac{dC}{dx} \quad (18)$$

21 where, v is the transition rate of H^* , C is the local concentration of H^* , the diffusion coefficient

1 can be calculated according to Stokes-Einstein relationship,

$$2 \quad D = \frac{k_B T}{6\pi\eta r} \quad (19)$$

3 where, k_B is the Boltzmann constant ($1.38\text{E-}23$ J/K), T is the temperature of vitrified specimen,
 4 η is the viscosity of the vitreous ice, and r ($\sim 10^{-10}\text{m}$) is the radius of H^* . Combining (18) and
 5 (19), we can calculate the transition rate of H^* as,

$$6 \quad v = \frac{1}{c} \cdot \frac{k_B T}{6\pi\eta r} \cdot \frac{dC}{dx} \quad (20)$$

7 Suppose, initially, the free radicals H^* are concentrated in a small cubic region with the size of 1
 8 nm^3 , thus we could have the following estimation,

$$9 \quad \frac{1}{c} \cdot \frac{dC}{dx} = \frac{1}{c} \cdot \frac{\Delta C}{\Delta x} \sim \frac{1}{\Delta x} \sim 10^9 \text{m}^{-1} \quad (21)$$

10 Thus, at the temperature ($T = 90$ K) of cryo-EM experiments, the transition rate of H^* can be
 11 estimated as,

$$12 \quad v \sim \frac{1.38 \times 10^{-23} \times 90}{6 \times 3.14 \times 10^{-3} \times 10^{-10}} \times 10^9 = 0.7 \text{ m/s} \quad (22)$$

13 Then the time for H^* traveling 0.3 nm , the averaged distance to reach adjacent groups and then
 14 perform radical reaction, can be estimated as,

$$15 \quad \Delta t \sim \frac{0.3 \text{ nm}}{0.7 \text{ m/s}} = 0.4 \text{ ns} \quad (23)$$

16 To be noted that, the estimation in (22) utilizes the water viscosity at room temperature, $\eta_{H_2O} =$
 17 $10^{-3} \text{ Pa} \cdot \text{s}$, where the viscosity of the vitreous ice η_{ice} at 90 K should be much larger (e.g. more

1 than ten times) than the water viscosity at room temperature. Therefore, the time scale for the second
2 damage process is in the nanoseconds ~ 10 ns.

3 The above estimation suggests that if we could take a cryo-EM exposure within 10 ns, the
4 appreciable electron radiation damage during the second process can be almost avoided in the final
5 recorded micrograph. The recent developed ultrafast transmission electron microscopy (UEM) [96-
6 99] has actually provided an opportunity to test this kind of idea. There are two operation modes of
7 UEM, the stroboscopic mode with picosecond temporal resolution and the single-pulse mode with
8 nanosecond temporal resolution [97]. The stroboscopic mode is useful for ultrafast electron
9 diffraction experiments and suitable to study the reversible process of the material. However, the
10 electron radiation damage of bio-macromolecules is irreversible. Thus, to achieve the concept of
11 ‘imaging before damage’, it is necessary to develop the cryo-ultrafast transmission electron
12 microscopy (cryo-UEM) that is operated in the single-pulse mode. Although there have been some
13 reports of using UEM to observe the biological specimens [100-102], all these studies were
14 performed in the stroboscopic mode and utilized the dehydrated specimen, which should not be
15 relevant to biological functions. There is still a big space to develop and improve the single-pulse
16 UEM technology. We are looking forward to the future maturation of cryo-UEM that will bring the
17 bio-macromolecular electron microscopy into a new era.

18 **3. Conclusions**

19 Cryo-EM SPA has become the most important technique of bio-macromolecular electron
20 microscopy. The era of studying the structures of bio-macromolecules by using Cryo-EM SPA is
21 just beginning. In the near future, we will see more and more sophisticated bio-macromolecular

1 complexes whose structures are solved into near atomic resolution, making significant implications
2 to their biological functions. However, just as Henderson [103] and Glaeser [32] pointed, cryo-
3 electron microscopy has not yet realized its full potential. In the future, with a better cryo-
4 vitrification technique to avoid air-water interface problem, with a better camera and new type of
5 microscope to further alleviate electron radiation damage effect and with some novel image
6 processing algorithms and experimental techniques to solve the focus gradient problem as well as
7 the conformational heterogeneity issue, Cryo-EM SPA will expand its full ability to solve the atomic
8 resolution of bio-macromolecules.

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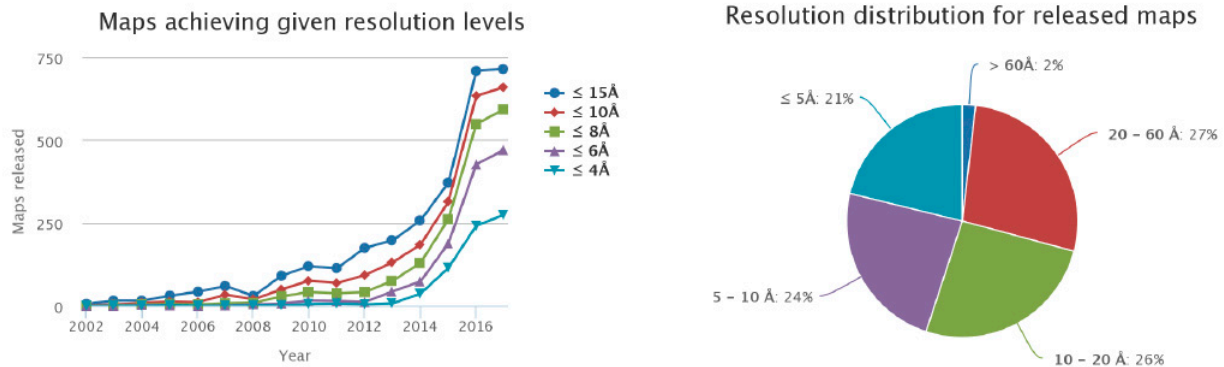
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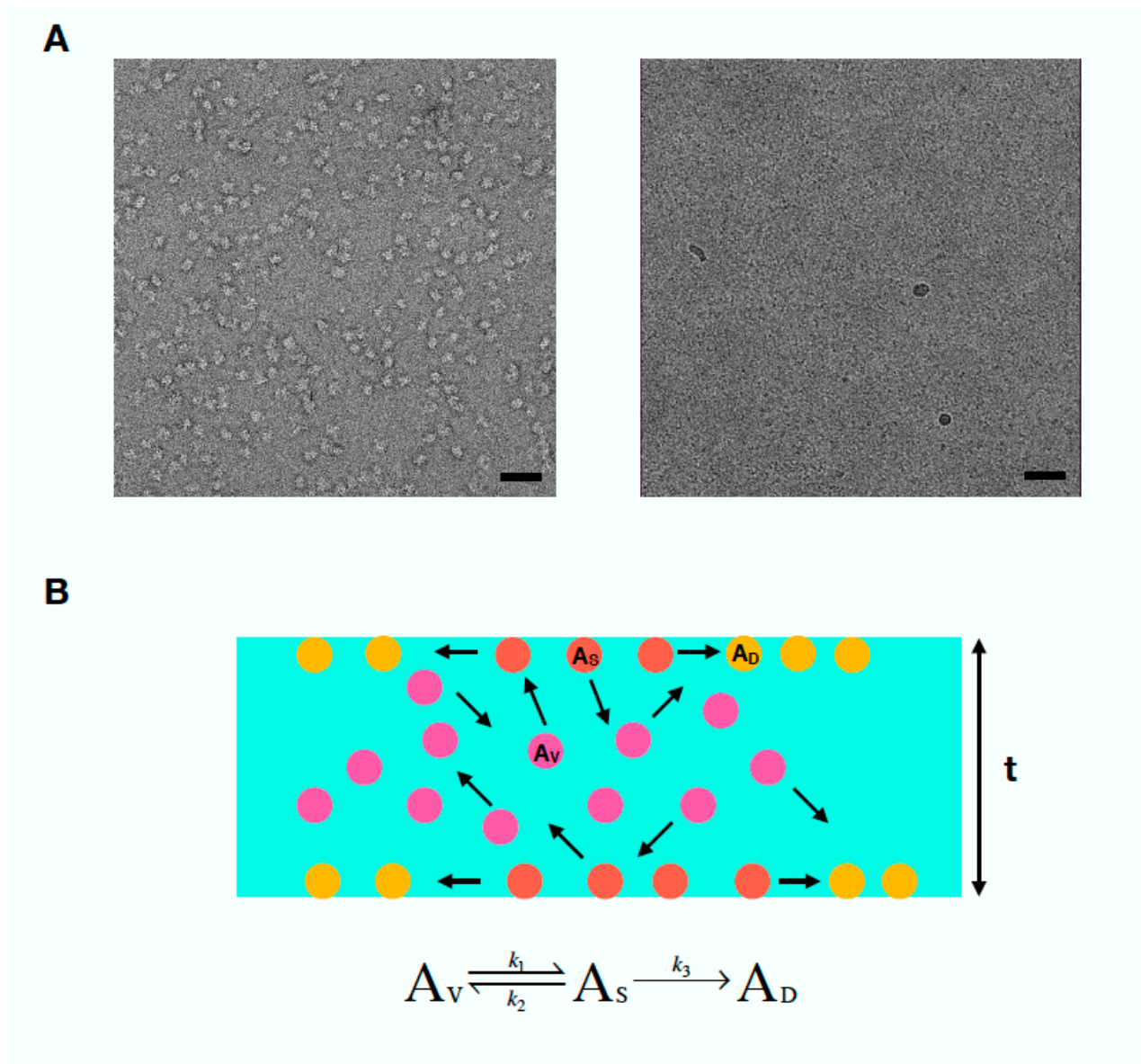
Figures



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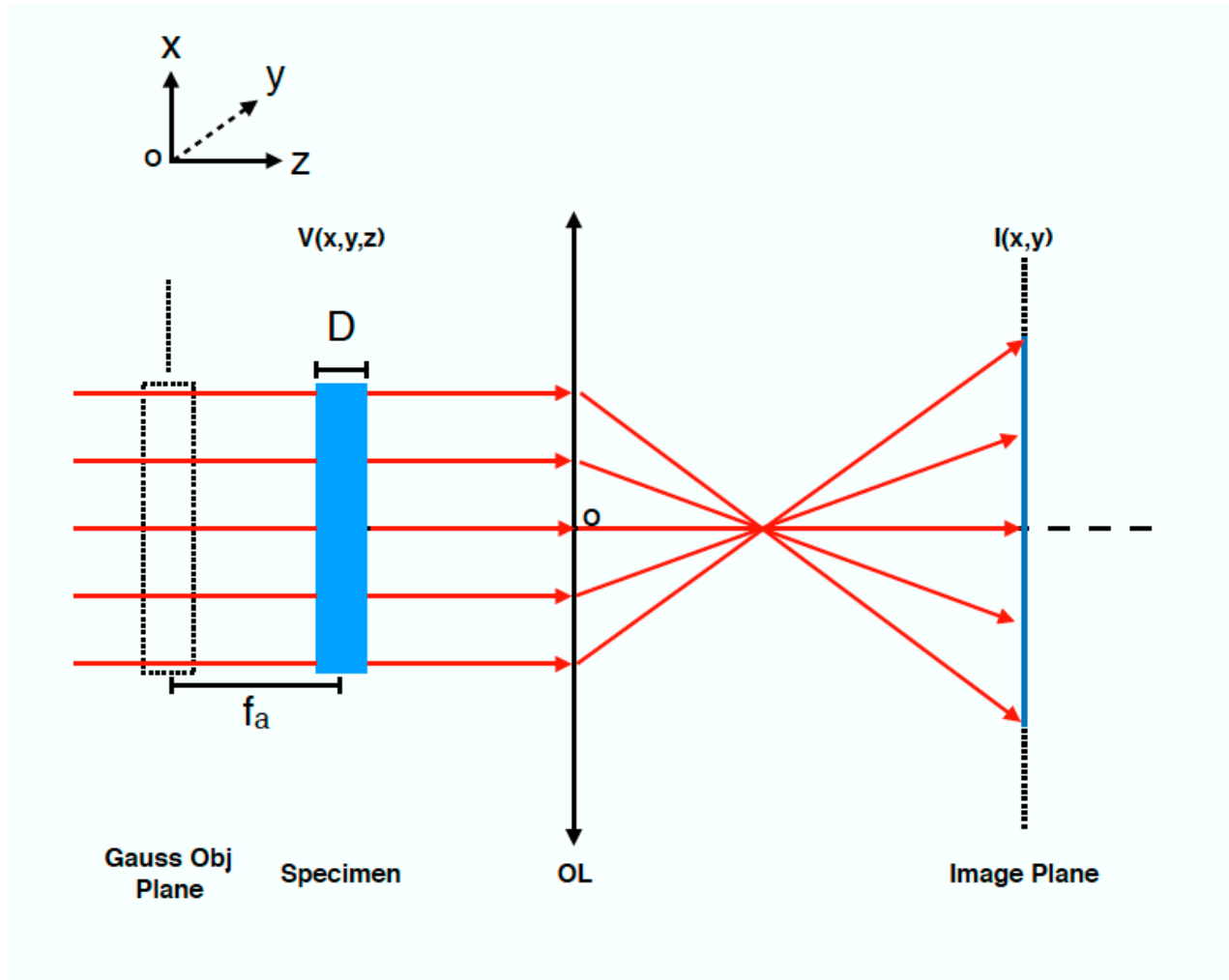
4 **Figure 1. Statistics of the electron density maps deposited in EMDB (Electron Microscopy**
5 **Data Bank).** The annual numbers of released maps with different reported resolutions are plotted
6 from 2002 to 2017 (left). And the distribution of all the released maps till 2017 is statistically plotted
7 vs different resolutions (right). Both panels were generated using the tool of EMDB
8 (<http://www.ebi.ac.uk/pdbe/emdb>).

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2 **Figure 2. Air-water interface effects during specimen cryo-vitrification.** A. The bio-
 3 macromolecule specimen exhibits homogenous distribution and good shape in negative electron
 4 microscopy (left), but is prone to degrade and hard to observe in cryo-electron microscopy (right).
 5 Scale bar, 50 nm. **B.** A diagram showing the physical process during specimen cryo-vitrification.
 6 Bio-macromolecules in native state (A_V) are colored in pink, those (A_S) absorbed to air-water
 7 interface are colored in orange, and that (A_D) denatured in yellow. The thickness (t) of the solution
 8 layer after blotting is between 30 ~ 100 nm.



1

2 **Figure 3.** A diagram showing the image formation of a thin specimen with the averaged
 3 **underfocus f_a .** Gauss Obj Plane, the plane where the idea objective locates according to Gauss
 4 image formation formula. D, the thickness of the specimen. OL, objective lens. $V(x,y,z)$, the density
 5 function of the specimen. $I(x,y)$, the image function at the image plane.