

1 Article

# 2 $H^\infty$ Robust Control of a Large-Piston MEMS 3 Micromirror for Compact Fourier Transform 4 Spectrometer Systems

5 Huipeng Chen <sup>1</sup>, Mengyuan Li <sup>2,4</sup>, Yi Zhang <sup>3</sup>, Huikai Xie <sup>2</sup>, Chang Chen <sup>1</sup>, Zhangming Peng and  
6 Shaohui Su <sup>1,\*</sup>

7 <sup>1</sup> School of Mechanical Engineering, HangZhou DianZi University, Hangzhou 310018, China;  
8 lynx\_chen@163.com (H.P. Chen); chenchang@hdu.edu.cn (C. Chen)

9 <sup>2</sup> Department of Electrical and Computer Engineering, University of Florida, Gainesville, FL 32611, USA;  
10 limengyuan@ufl.edu (M.Y Li); hkxie@ece.ufl.edu (H.K. Xie)

11 <sup>3</sup> Shangdong University of Science and Technology, Qingdao 266590, China;  
12 zhangyigb@163.com (Y. Zhang)

13 <sup>4</sup> School of Information and Electronics, Beijing Institute of Technology, Beijing 100081, China

14 \* Correspondence: 13735515369@163.com; Tel.: +86-137-355-15369

15 **Abstract:** Incorporating linear-scanning MEMS micromirrors into Fourier transform spectral  
16 acquisition systems can greatly reduce the size of the spectrometer equipment, making portable  
17 Fourier transform spectrometers (FTS) possible. How to minimize the tilting of the MEMS mirror  
18 plate during its large linear scan is a major problem in this application. In this work, an FTS system  
19 has been constructed based on a biaxial MEMS micromirror with a large piston displacement of 180  
20  $\mu\text{m}$ , and a biaxial  $H^\infty$  robust controller is designed. Compared with open-loop control and PID  
21 closed-loop control,  $H^\infty$  robust control has good stability and robustness. The experimental results  
22 show that the stable scanning displacement reaches 110.9  $\mu\text{m}$  under the  $H^\infty$  robust control, and the  
23 tilting angle of the MEMS mirror plate in that full scanning range falls within  $\pm 0.0014^\circ$ . Without  
24 control, the FTS system cannot generate meaningful spectra. In contrast, the FTS yields a clean  
25 spectrum with an FWHM spectral linewidth of 96  $\text{cm}^{-1}$  under the  $H^\infty$  robust control. Moreover, the  
26 FTS system can maintain good stability and robustness under various driving conditions.

27 **Keywords:** electrothermal micromirror; robust control; bimorph actuator modeling; active tilting  
28 rejection; Fourier transform spectrometer

---

## 30 1. Introduction

31 Fourier transform infrared spectroscopy (FTIR)[1] is a technique that is used to obtain absorption  
32 or emission IR spectra of various matters and determine materials' compositions and concentrations  
33 in both laboratory and field environments. A Fourier transform spectrometer (FTS) is based on a  
34 Michelson interferometer consisting of a beam splitter, a photodetector (PD), and one movable mirror  
35 and one fixed mirror respectively in its two optical path arms. Conventional FTS systems are only for  
36 lab use as they are expensive and bulky largely due to the complex scanning mirror system [11].

37 Recently, FTS systems based on MEMS micromirrors begin to emerge, and such miniature FTS  
38 systems can enable real-time, in-field analysis in many environments such as national border  
39 checkpoints and in natural or manmade hazardous conditions[2-3]. For miniature FTS systems, the  
40 scanning characteristics of the moving MEMS micromirror are critical. Electrothermal micromirrors  
41 are more widely used in FTS than other types, such as piezoelectric, electromagnetic, and electrostatic  
42 micromirrors, because they can provide much larger linear scan range at low drive voltage [4-10].

43 In principle[11], for a FTS, the Fourier Transform of the interferogram yields a spectrum whose  
44 resolvable spectral linewidth is inversely proportional to the movable mirror's scan range, i.e., the  
45 larger the linear scan range of the movable mirror, the higher the achievable spectral resolution is.

46 However, limited by the fabrication process variations, the piston motion of MEMS micromirrors  
47 always comes with tilting. The tilting during the mirror's piston deteriorates the interferogram,  
48 resulting in low spectral resolution or even no recoverable spectrum. Thus, controlling the mirror's  
49 tilting has become the biggest challenge in practical MEMS based FTS systems.

50 Several tilt control methods have been developed to reduce the tilting of the movable  
51 micromirror of FTS system during scanning motion. Wu et al. [12] developed a FTS based on dual-  
52 reflective MEMS mirror. When the driving signal was not compensated or controlled, the tilting angle  
53 was  $0.7^\circ$  during the scan. By using a pre-shaped drive signal, the mirror tilting was reduced to  $0.06^\circ$ .  
54 S. R. Samuelson et al. [13] proposed a piston motion micromirror with a laddered inverted-series-  
55 connected (ISC) electrothermal actuator array and demonstrated a uncompensated tilt of  $0.25^\circ$  over  
56 its full displacement range. By designing a pair of ratio optimized drive signals, the tilting angle was  
57 reduced to  $0.004^\circ$  [14]. Wang et al. [11] reported a large-stroke electrothermal MEMS mirror with an  
58 original tilting angle  $0.3^\circ$ . In order to reduce the tilting angle, the micromirror was driven by an open-  
59 loop control using a pair of corrected ramp drive signals, and the final tilting angle was reduced to  
60  $\pm 0.002^\circ$ .

61 In view of the high sensitivity of open loop control to the environmental variations and  
62 disturbances, there have been numerous studies on designing of closed-loop controllers for  
63 improving the repeatability and stability of micromirrors [15-17], most of which are focused on  
64 electrostatic micromirrors. A study on closed-loop tilt control for a single-axis electrothermal  
65 micromirror has just recently been reported [28], where the tilt angle was controlled within  $\pm 0.0015^\circ$ .  
66 Although the closed-loop control algorithm developed in this study improved the robustness of the  
67 electrothermal micromirror scanning, there are still some practical problems that have not been  
68 considered and solved. First of all, the uniaxial electrothermal micromirror previously studied has  
69 no actuators in the orthogonal axis, and thus the jitter in the orthogonal axis is in an uncontrolled  
70 state. Secondly, the change of the characteristics of the electrothermal micromirror due to aging and  
71 any changes in the operating environment may cause the originally stable system to lose its stability.

72 To solve these problems, we firstly propose to use a dual-axis electrothermal micromirror  
73 instead of a uniaxial micromirror in the MEMS FTS system. Thus both directions of the micromirror  
74 can be controlled to avoid the instability caused by no control in one axis. Secondly, as the  
75 electrothermal micromirror scanning exhibits strong non-linearity and resonance characteristics, the  
76 device model was obtained through experiments rather than theoretical modeling. Thirdly, in order  
77 to ensure the robustness of the MEMS FTS system, a  $H^\infty$  robust control method is proposed and the  
78  $H^\infty$  controller is designed for the perturbation of the system parameters. As a comparison, a PID  
79 controller and a look-up table control are also implemented. The experimental results show that with  
80 the  $H^\infty$  control, not only the tilting angle is greatly reduced, but also it has the advantages of good  
81 anti-interference ability.

82 This paper is organized as follows. Section 2 introduces the electrothermal MEMS mirror and its  
83 experimental model. In Section 3, the design of a  $H^\infty$  robust controller is introduced in detail, and a  
84 PID controller and a look-up table driving curve are also presented. In Section 4, the experimental  
85 results with the  $H^\infty$  robust controller, PID controller and look-up table driving curve on the MEMS  
86 mirror and their application in the MEMS FTS system are analyzed and compared.

## 87 2. THE ELECTROTHERMAL MICROMIRROR

### 88 2.1. Device Description

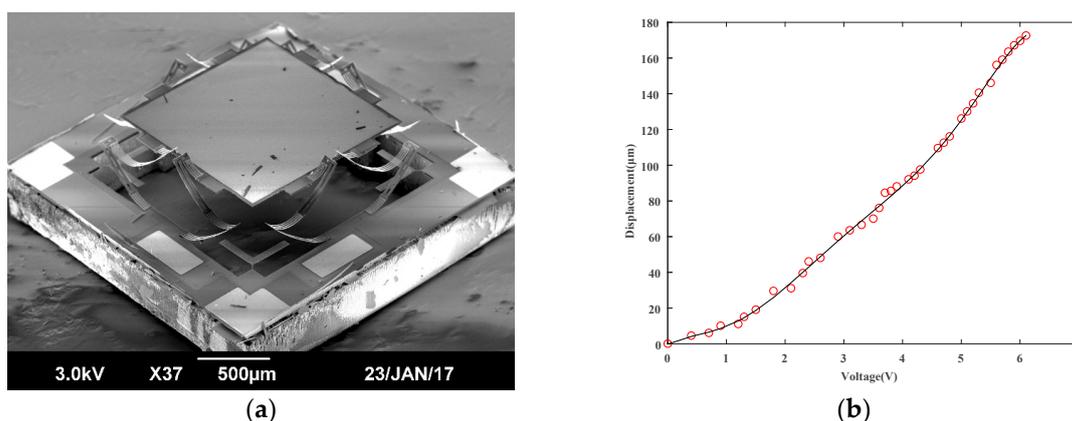
89 A bimorph is composed of two layers of different materials that have different coefficients of  
90 thermal expansion (CTEs), as shown in Figure 1(a). Al and SiO<sub>2</sub> are chosen as the two bimorph layers  
91 for their large CTE difference which can lead to large actuation. A Pt resistor is also integrated as a  
92 heater. When a current is injected into the Pt resistor, Joule heating will be produced, which increases  
93 the bimorph temperature and consequently causes the bimorph to bend due to the different CTEs of  
94 the two bimorph layers. The simple bimorph shown in Figure 1(a) also generates undesired large  
95 lateral shift upon actuation. Thus, a bimorph actuator design consisting of two segments of silicon-

96 backed rigid beams and three segments of Al/SiO<sub>2</sub> bimorphs has been proposed, as shown in Figure  
 97 1(b). This design can generate a lateral shift free (LSF) large vertical displacement[10,18].



98 **Figure 1.** (a) Bimorph structure. (b) Large displacement electrothermal bimorph actuator design.

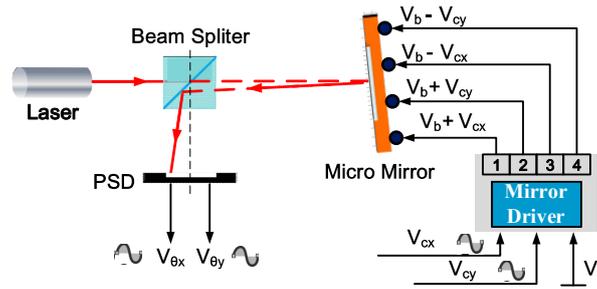
99 Figure 2(a) shows an SEM of a fabricated LSF MEMS mirror, where the initial elevation of the  
 100 pop-up mirror plate is about 180 μm. This elevation is caused by the thin film intrinsic stress and  
 101 thermal residual stress generated during fabrication. There are two LSF bimorph actuators attached  
 102 on each side of the mirror plate. When a drive voltage is applied on all four bimorph actuators, the  
 103 measured static piston response is shown in Figure 2(b), where the vertical displacement reaches up to  
 104 180 μm. The allowable drive voltage range to generate stable vertical motion is from 0 to 6.6 V.  
 105 Increasing the voltage further would eventually burn out the bimorphs due to overheating. Although  
 106 the optimized drive signal ratio [14] or open-loop compensation method [11] can reduce the tilting  
 107 during the piston scanning, the residual tilt angle is relatively large due to time-varying  
 108 characteristics of thermal bimorph actuation.



109 **Figure 2.** (a) SEM of an electrothermal MEMS mirror. (b) Measured vertical displacement versus  
 110 applied voltage on bimorph actuators.

## 111 2.2. Dynamic Model of the Micromirror

112 The micromirror system studied in this paper has a complex structure and multi-input and  
 113 multi-output characteristics. The analytical model is difficult to obtain. Therefore, in this work, a  
 114 frequency domain experimental method is used to obtain the system model for the control system  
 115 design. The frequency response of the mirror tilting motion was measured in an experimental setup  
 116 as illustrated in Figure 3.



117

118

Figure 3. Setup to measure open-loop frequency response of micromirror tilting.

119

The mirror is first biased in the linear region of the scan displacement-voltage characteristics. In both directions  $x$  and  $y$ , the AC voltages  $v_{cx}$  and  $v_{cy}$  are superimposed on the DC bias  $V_b$  via a unit-gain driver. The driver generates two pairs of the voltage outputs,  $V_b + V_{cx}$  and  $V_b - V_{cx}$ ,  $V_b + V_{cy}$  and  $V_b - V_{cy}$ , which are applied to the four bimorph actuators in a differential fashion.  $V_b$  is used to bias the mirror plate at a certain displacement and  $V_{cx}/V_{cy}$  excites the mirror to tilt at the vertical position set by  $V_b$ . To test the frequency response, the frequency of  $V_{cx}$  and  $V_{cy}$  are swept and the actual tilt angles in both directions of  $x$  and  $y$  are measured by tracking the light beam reflected from the mirror plate using a position sensing detector (PSD). The PSD output voltages,  $V_{\theta_x}$  and  $V_{\theta_y}$ , are proportional to the mirror tilt angles in both directions of  $x$  and  $y$ .

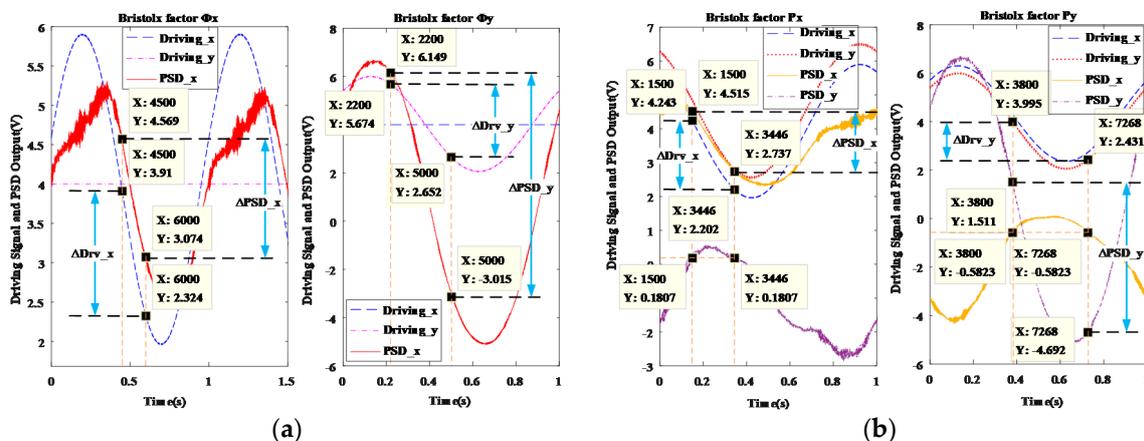
128

The system has two outputs, so it needs to analyze and judge the coupling relationship. If the coupling is strong, decoupling is needed. There are a variety of methods to evaluate the degree of coupling of multivariable systems [19-26], among which the most widely used one is the static and dynamic Relative Gain Array (RGA) theory [24] proposed by Bristol.

132

For a multivariable control system, Bristol defines a first amplification factor  $\Phi_k$  and a second amplification factor  $P_k$ . The first amplification coefficient  $\Phi_k$  means that, in the system of the mutual coupling, the channel gain between a driving signal of  $X/Y$  direction  $Drv_j(j=x,y)$  and a PSD signal of  $X/Y$  direction  $PSD_i(i=x,y)$  under the conditions of  $Drv_j$  observed with a change of  $\Delta Drv_j$  and other manipulated variables  $Drv_r(r \neq j, r=x,y)$  unchanged, the  $\Phi_k$  of  $x$  and  $y$  direction is:  $\Phi_x = \Delta PSD_x / \Delta Drv_x |_{Drv_y=const}$ ,  $\Phi_y = \Delta PSD_y / \Delta Drv_y |_{Drv_x=const}$ . The second amplification factor  $P_k$  refers to the change of the  $PSD_i(i=x,y)$  obtained by fixing the other  $PSD_r(r \neq i, r=x,y)$  and changing only  $Drv_j(j=x,y)$ , and the two change coefficient is the second amplification factor  $P_k$  in the static state. That is:  $P_x = \Delta PSD_x / \Delta Drv_x |_{PSD_y=const}$ ,  $P_y = \Delta PSD_y / \Delta Drv_y |_{PSD_x=const}$ . When  $0.8 < \lambda_k = \Phi_k/P_k < 1.2$ , the influence of other channels on the channel ( $Drv_j \rightarrow PSD_i$ ) is small and can be used as the main channel [24].

142



143

144

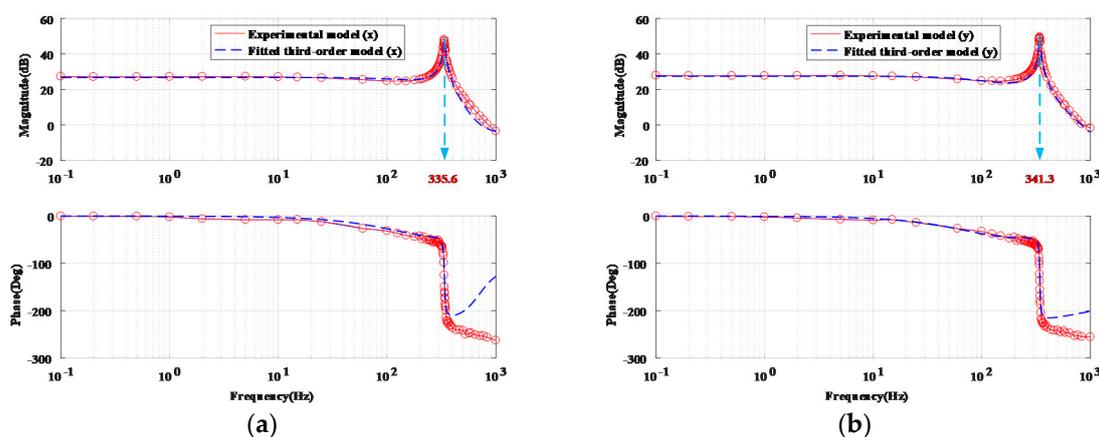
Figure 4. (a) Signals for calculating Bristol factors  $\Phi$  of  $X/Y$  direction (b) Signals for calculating Bristol factors  $P$  of  $X/Y$  direction.

145 As shown in Figure 4(a), X/Y direction uses 2V~6V sine waves to drive, and Y/X direction uses  
 146 4V constant value to drive, get the  $\Phi$ :  $\Phi_x=1.495/1.586$ ,  $\Phi_y=9.164/3.022$ . As shown in Figure 4(b), X/Y  
 147 direction uses 2V ~ 6V sine to drive, Y/X direction uses 2.5V ~ 6.5V sine to drive, get the P:  
 148  $P_x=1.778/2.041$ ,  $P_y=1.347/0.462$ . And finally get the following results: (1)  $\lambda_x = \Phi_x/P_x = 1.082 \in$   
 149  $(0.8,1.2)$ ; (2)  $\lambda_y = \Phi_y/P_y = 1.04 \in (0.8,1.2)$ .  $\lambda_x$  and  $\lambda_y$  indicating that the control of the x, y  
 150 direction can be used as the main control channel, which can independently be controlled.

151 As illustrated in Figure 3, the frequency responses of the x-scan and y-scan can be expressed as:

$$152 \quad G_{mx}(s) = \frac{V_{\theta x}(s)}{V_{cx}(s)}, \quad G_{my}(s) = \frac{V_{\theta y}(s)}{V_{cy}(s)} \quad (1)$$

153 A sinusoidal voltage driving signal with an amplitude of 0.1 V and a dc bias of 4.0 V was used  
 154 to obtain frequency response. The measured frequency responses are shown in Figure 5 (a) and (b),  
 155 where the resonant frequencies in x and y direction are 335.6 Hz and 341.3 Hz, respectively.



156 **Figure 5.** Experimental frequency response and fitted micromirror model (a) The micromirror model  
 157 of X direction. (b) The micromirror model of Y direction.

158 Based on the experimental data shown in Figure 5, we used the system identification toolbox in  
 159 Matlab to identify the system, and selected the best-fit results as the system model for the feedback  
 160 controller design, i.e.,

$$161 \quad G_{mx}(s) = \frac{4915s^2 + 1.68 \times 10^7 s + 9.46 \times 10^{10}}{s^3 + 1042s^2 + 4.54 \times 10^6 s + 4.36 \times 10^9} \quad (2)$$

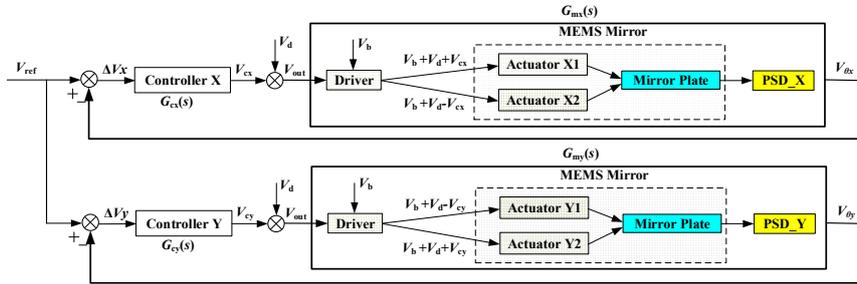
$$162 \quad G_{my}(s) = \frac{1.92 \times 10^7 s + 6.16 \times 10^{10}}{s^3 + 614.4s^2 + 4.66 \times 10^6 s + 2.61 \times 10^9} \quad (3)$$

163 This transfer function model was tested under room temperature and one atmospheric pressure.  
 164 For different micromirrors or operating conditions, the model parameters can be measured and fitted  
 165 by using the above identification method.  
 166

### 167 3. Design of Robust Controller

168 When the MEMS micromirror based FTS system is placed in different environments and  
 169 operating conditions, the system parameters will change and thus the system model will need to be  
 170 changed accordingly. So it is a great challenge to design a controller that can ensure the stability and  
 171 response characteristics of the system even when there exist external disturbances or internal  
 172 structure deterioration. PID control and look-up table control methods are simple and easy to  
 173 implement, but these control methods cannot adapt well to changes in system parameters, causing  
 174 large errors or even system instability. Therefore, in this work, a controller design with an  $H^\infty$  hybrid  
 175 sensitivity control algorithm is proposed, so that the FTS system can not only be applied to working  
 176 environments other than the laboratory, but also be adapt to the measurement errors or the  
 177 parameters' drift caused by the aging of the system. In order to evaluate the performance of the  $H^\infty$   
 178 control method, a PID controller and a look-up controller are also implemented for the FTS system  
 179 as a comparison.

180 A block diagram of the micromirror tilting control loop is shown in Figure 6, where  $G_{cx}(s)$  and  
 181  $G_{cy}(s)$  denote the feedback controllers to be designed, the error  $\Delta V_x$  and  $\Delta V_y$  are the measures of the  
 182 residual tilting of the MEMS mirror plate, and  $V_d$  is the equivalent disturbance input voltage  
 183 intended to evaluate the robustness of the closed-loop system. The system models  $G_{mx}(s)$  and  
 184  $G_{my}(s)$  have been experimentally obtained as given in Eqs. (2) and (3).

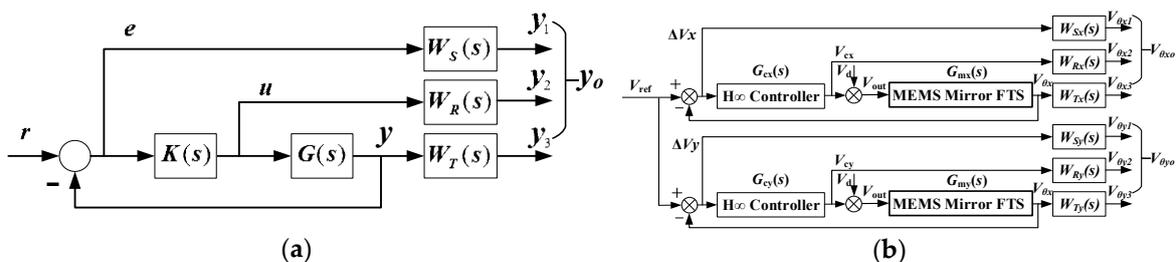


185  
186 **Figure 6.** Block diagram of the closed-loop micromirror tilt control system.

187 **3.1. Design of  $H^\infty$  robust Controller**

188  $H^\infty$  robust control theory [25] is a control theory in  $H^\infty$  space (Hardy space) that can yield robust  
 189 controllers by optimizing the infinite norm of certain performance index, solving the problems of a  
 190 robust control model that has a certain range of uncertainties and external interference signals  
 191 existing in a system.

192 To control a system with both interference and uncertainty, the  $H^\infty$  mixed sensitivity design  
 193 method can be employed [26]. The control system structure diagram of the weighted  $H^\infty$  mixed  
 194 sensitivity method is shown in Figure 7(a), and the same form of the FTS system control structure is  
 195 shown in Figure 7(b), where  $r$  is the PSD reference input  $V_{ref}$ ,  $e$  is the PSD error  $\Delta V_x$  or  $\Delta V_y$ ,  $u$  is  
 196 the  $H^\infty$  controller output,  $V_{cx}$  or  $V_{cy}$ ,  $y_o = [y_1, y_2, y_3]$  is the output of the MEMS mirror in the FTS  
 197 system,  $V_{\theta x0} = [V_{\theta x1}, V_{\theta x2}, V_{\theta x3}]$  or  $V_{\theta y0} = [V_{\theta y1}, V_{\theta y2}, V_{\theta y3}]$ , and  $W_{S/Sx/Sy}(s)$ ,  $W_{R/Rx/Ry}(s)$   
 198 and  $W_{T/Tx/Ty}(s)$  are three weight functions added in order to respectively suppress the interference  
 199 and limit controlled quantity and thus improve the system performance.



200 **Figure 7.** (a) Structure diagram of the  $H^\infty$  mixed sensitivity control system. (b) The block diagram of the  
 201 FTS system  $H^\infty$  controller.

202 The closed-loop transfer function  $S(s)$  from the input  $r$  to the error  $e$  is

$$203 \quad S(s) = [I + G(s)K(s)]^{-1} \quad (4)$$

204 Where  $S(s)$  is called the sensitivity function, which is the most important indicator to determine  
 205 the size of the PSD signal tracking error. The lower the sensitivity  $S(s)$  is, the smaller tilting of the  
 206 FTS system is. The closed-loop transfer function  $T(s)$  from input  $r$  to output  $y_o$  is

$$207 \quad T(s) = I - S(s) \quad (5)$$

208 We can use the P-K structure of the  $H^\infty$  standard problem. The generalized object  $P(s)$  is given by

$$209 \quad P(s) = \begin{bmatrix} P_{11} & P_{12} \\ P_{21} & P_{22} \end{bmatrix} = \begin{bmatrix} W_S & -W_S G \\ 0 & W_R \\ 0 & W_T G \\ I & -G \end{bmatrix} \quad (6)$$

210 The closed-loop transfer function matrix from input  $r$  to output  $y_o$  is

$$T_{ry_o}(s) = LFT(P(s), K(s)) = P_{11} + P_{12}K(I - P_{22}K)^{-1}P_{21} = [W_s S, W_R R, W_T T]^T \quad (7)$$

To obtain the  $H^\infty$  mixed sensitivity controller is to optimize the  $\|T_{ry_o}(s)\|_\infty$  by finding a controller  $K(s)$  so that the closed-loop transfer function of the FTS is internally stable and satisfies:

$$\inf_K \|T_{ry_o}(s)\|_\infty = \gamma_{min} \Rightarrow \|T_{ry_o}(s)\|_\infty < \gamma (\gamma > \gamma_{min}) \quad (8)$$

From the frequency domain point of view, the classical  $H^\infty$  control algorithm is essentially a system loop forming method. The  $H^\infty$  mixed sensitivity control strategy directly performs closed-loop gain shaping on the closed-loop function such as the sensitivity function  $S(s)$  or the complementary sensitivity function  $T(s)$ , thus eliminating the large peaks that may occur in open-loop gain shaping. The ideal S/T curve is given in Figure 8.

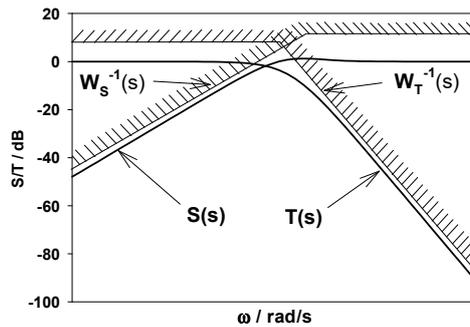


Figure 8. Ideal S/T curve.

As show in Figure 9, in order to make the FTS system meet the shape of the ideal S/T curve, according to the FTS system characteristics, through multiple iterative tests, we set the weight functions  $W_{S_x}(s) = (0.001s + 200)/(s + 1)$ ,  $W_{S_y}(s) = (0.001s + 265)/(s + 4.1)$ ,  $W_{T_x}(s) = (0.66s + 1)/(0.0025s + 100)$  and  $W_{T_y}(s) = (0.86s + 3.6)/(0.0025s + 100)$  to guarantee the desired low-pass characteristics, and set the weight function  $W_{R_x}(s) = W_{R_y}(s) = 1 \times 10^{-4}$  to adjust the output of the controller to ensure that the PSD output changes at the millivolt level. The robust controllers  $G_{cx}(s)$  and  $G_{cy}(s)$  in the x and y directions are obtained as follows,

$$G_{cx}(s) = \frac{1.15 \times 10^8 s^4 + 4.71 \times 10^{12} s^3 + 5.30 \times 10^{15} s^2 + 2.13 \times 10^{19} s + 1.99 \times 10^{22}}{s^5 + 1.29 \times 10^{10} s^4 + 1.75 \times 10^{14} s^3 + 6.96 \times 10^{17} s^2 + 2.52 \times 10^{21} s + 2.52 \times 10^{21}}$$

$$G_{cy}(s) = \frac{2.42 \times 10^8 s^4 + 9.81 \times 10^{12} s^3 + 7.06 \times 10^{15} s^2 + 4.57 \times 10^{19} s + 2.52 \times 10^{22}}{s^5 + 1.15 \times 10^7 s^4 + 6.64 \times 10^{13} s^3 + 1.29 \times 10^{18} s^2 + 3.44 \times 10^{21} s + 1.41 \times 10^{22}}$$

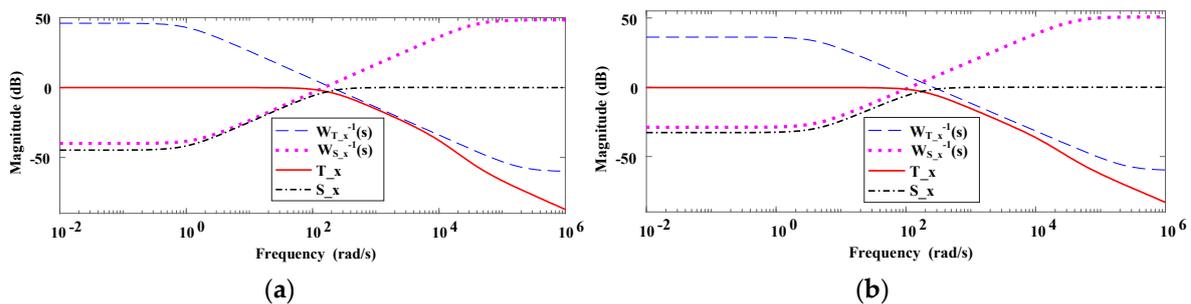


Figure 9. S/T curve. (a) S/T curve of X direction. (b) S/T curve of Y direction.

### 3.2. Design of the PID Controller

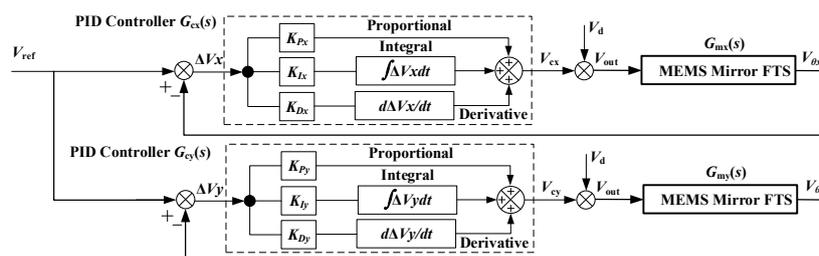
Figure 10 shows the PID controller designed respectively in two directions: x and y. The two control loops can be independently controlled and adjusted for optimal conditions of the final controller to be determined is that the x and y direction achieve the best.

A PID controller is a kind of linear controller that tries to minimize the error between the reference value and the actual output value. The errors in x and y directions are given by

$$\Delta V_x = V_{ref} - V_{\theta_x}, \quad \Delta V_y = V_{ref} - V_{\theta_y} \quad (9)$$

Its control law is

$$G_{cx}(s) = \frac{V_{cx}}{\Delta V_x} = K_{Px} \left[ 1 + \frac{1}{K_{Ix}s} + K_{Dx}s \right], \quad G_{cy}(s) = \frac{V_{cy}}{\Delta V_y} = K_{Py} \left[ 1 + \frac{1}{K_{Iy}s} + K_{Dy}s \right] \quad (10)$$



242

243

**Figure 10.** The block diagram of our PID controller.

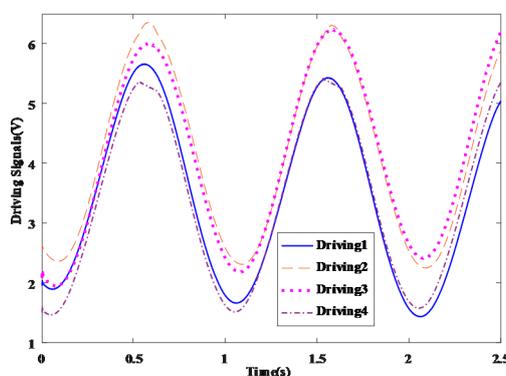
244 In the above formula,  $K_P$  is the proportional coefficient,  $K_I$  is the integral time constant, and  $K_D$   
 245 is the differential time constant. In the PID controller of the correction link,  $K_P$  adjusts the system  
 246 error. Once the error occurs, the controller produces control effect to reduce the error.  $K_I$  is used to  
 247 eliminate the static difference and improve the system's no difference.  $K_D$  adjusts the rate of change  
 248 of the error, accelerates the speed of movement of the system and reduces the adjustment time.

249 Even without the system model,  $K_P$ ,  $K_I$  and  $K_D$  can be experimentally determined by multiple  
 250 trials. In this work, the PID controller is obtained without using the above FTS system model.

### 251 3.3. Design of Look up Table Controller

252 The look-up table is actually generated from the FTS system with the PID controller described  
 253 above. The output voltage signals for the four actuators X1, X2, Y1 and Y2 are acquired and the data  
 254 are plotted in Figure 11, where Driving1, Driving2, Driving3 and Driving4 are the driving voltages  
 255 of the four actuators respectively. Taking the Driving1 as the reference, the other three signals are  
 256 fitted to obtain the following results:

$$257 \text{drv1} = \text{drv1}; \text{drv2} = 1.00045 * \text{drv1} + 0.6799; \text{drv3} = 0.9887 * \text{drv1} + 0.6837; \text{drv4} = 0.9842 * \text{drv1} + 0.0038.$$



258

259

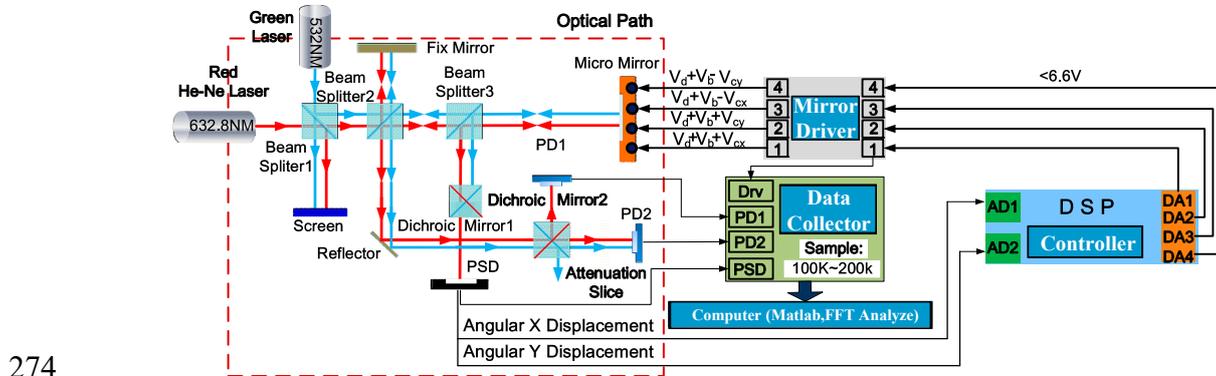
**Figure 11.** The block diagram of our PID controller.

## 260 4. Experimental Results and Discussion

### 261 4.1. Experimental Setup

262 An experimental setup, as illustrated schematically in Figure 12, is used to evaluate the tilting  
 263 performance of the closed-loop controlled electrothermal micromirror with the proposed control  
 264 schemes. This setup is actually a Michelson interferometer based Fourier transform spectrometer  
 265 (FTS), which is composed of an MEMS mirror to be controlled, a red He-Ne laser (632.8 nm) source  
 266 (LS-R), a green laser (532 nm) source (LS-G), three beam splitters (BS), two dichroic mirrors, a position  
 267 sensitive detector (PSD), two photodiodes (PD1 and PD2), a high speed data collector, an MEMS  
 268 driver to drive the four bimorph actuators in both x and y directions, and a digital controller realized  
 269 with a 32-bit digital signal processor (DSP) and high speed A/D, D/A converters. Here, the controller

270 sampling frequency is set at 10 kHz benefiting from the high-speed DSP (TMS320F28335) with a  
 271 powerful floating-point unit. As the frequency of the PD signals may reach 9~10 kHz, the data  
 272 collector sampling rate is set to 200 kHz. A computer is used to reconstruct the spectrum via Fast  
 273 Fourier Transform (FFT) from the raw interferogram signals from PD1 and PD2.



274

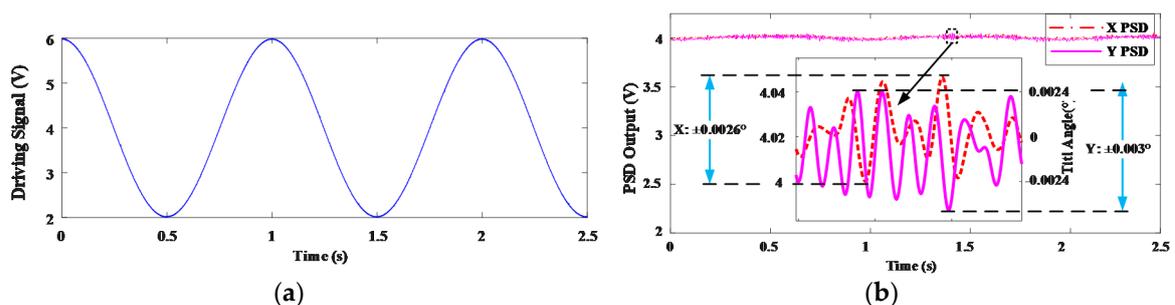
275 **Figure 12.** The Schematic Diagram of FTS Experiment System.

276 The schematic diagram of the optical path is shown in the upper part of Figure 12. The laser  
 277 beams from LS-R and LS-G are combined by the first beam splitter (BS1) and directed into the second  
 278 beam splitter (BS2), where the combined light beam is then split into two beams which are reflected  
 279 back respectively from a fixed mirror (FM) and the MEMS mirror (MM) through the third beam  
 280 splitter (BS3) and then re-combined as a single beam by the BS2. Under the combined action of the  
 281 BS3 and the first dichroic mirror (DM1), the red laser light is received by the PSD. After that, the  
 282 second dichroic mirror (DM2) only allows the red laser to pass through to the first photodiode (PD1)  
 283 and the reflector directs the combined light to the second photodiode (PD2) and through an  
 284 attenuation slice to reduce the excessive light intensity. The red laser here is introduced as the  
 285 reference light for spectrum calibration to overcome the variable velocity of the mirror [14]. The green  
 286 laser (532 nm) combined with the red laser is used as the testing light to be measured. As the  
 287 employed MEMS micromirror has a low thermal cut-off frequency of less than 5 Hz [11,27], we use  
 288 0.2, 0.5, 1.0 and 2.0 Hz drive signals to carry out the experiments, and mainly use the results of the  
 289 1.0 Hz drive signal to perform the comparison analysis. When the drive frequency is increased, the  
 290 micromirror has lower response to the high frequency drive, so that the amplitude decreases and the  
 291 optical path difference (OPD) of the FTS system becomes smaller, but this is not discussed in this  
 292 paper.

#### 293 4.2. Tilting Control

294 The PID control, look-up table control and  $H^\infty$  control described above were respectively  
 295 applied to the MEMS FTS system to control the tilting of the MEMS mirror under the drive signal  
 296 with an amplitude of 4V and a frequency of 1Hz. The results are described one by one.

##### 297 1)PID control



298 **Figure 13.** Measured tilt angle responses for PID FTS system. (a)Driving Signal. (b) PSD signals of X and  
 299 Y two directions.

300 Using the PID controller design shown in Figure 10 and tuning  $K_P$ ,  $K_I$  and  $K_D$ , we used the  
 301 driving signal shown in Figure 13(a) and obtained the PSD output signals of the x and y two  
 302 directions as shown in Figure 13(b), where the tilt angle in x direction is reduced to  $\pm 0.0026^\circ$  from  
 303  $0.453^\circ$  when no control was used, and the tilt angle in y direction is reduced to  $\pm 0.003^\circ$  from  $1.786^\circ$   
 304 when no control was used, both close to the optimal tilt angle range.

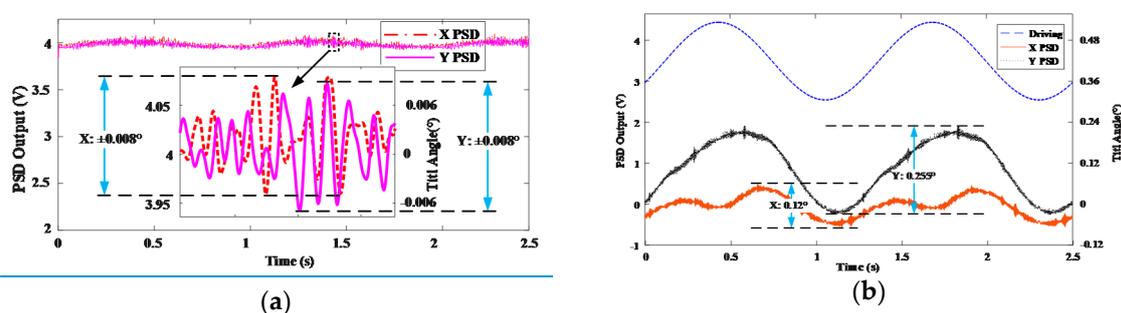
305 When changing the drive signal frequency, the tilt angles in both x and y directions vary slightly,  
 306 as shown in Table 1. As the PID parameters are optimized for the 1Hz, 2V ~ 6V sine wave drive signal,  
 307 the tilting increases due to the limited robustness of the PID controller when using other drive  
 308 frequencies.

309 **Table 1.** Tilt angle variations when the drive signal frequency changes..

| Tilting Angle | Drive signal frequency (Hz) |                    |                    |                    |
|---------------|-----------------------------|--------------------|--------------------|--------------------|
|               | 0.2                         | 0.5                | 1                  | 2                  |
| X direction   | $\pm 0.0028^\circ$          | $\pm 0.0029^\circ$ | $\pm 0.0026^\circ$ | $\pm 0.0029^\circ$ |
| Y direction   | $\pm 0.0031^\circ$          | $\pm 0.0033^\circ$ | $\pm 0.003^\circ$  | $\pm 0.0035^\circ$ |

310  
 311 2) Look-up table control

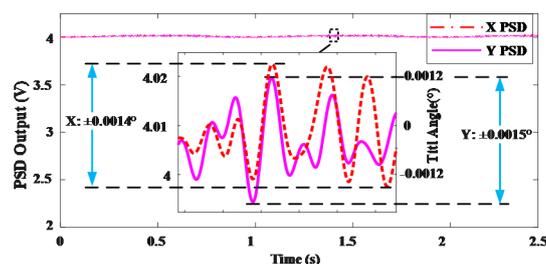
312 The first drive signal is set as a 1Hz, 2V~6V sinusoidal signal, as shown in Figure 13(a), and the  
 313 other three drive signals are generated according to the method described in section 3.3. The  
 314 corresponding output signals of the PSD in x and y directions are shown in Figure 14(a), where the  
 315 tilt angles in x and y directions both are  $\pm 0.008^\circ$ , which are greatly reduced compared to the no-  
 316 control case. However, the robustness of the look-up table control method is poor. As shown in Figure  
 317 14(b), when the drive voltage is changed to a 1Hz, 2.5V ~ 4.5V sine wave, the tilt angles in x and y  
 318 directions are increased to  $\pm 0.12^\circ$  and  $\pm 0.255^\circ$ , respectively.



319 **Figure 14.** (a) Measured tilt angle responses for Look-up Table FTS system using the same driving as PID  
 320 control. (b) Measured tilt angle responses for Look-up Table FTS system using different driving.

321 3)  $H^\infty$  robust Control

322 The  $H^\infty$  controller design given in section 3.1 is applied to control the MEMS mirror in the FTS  
 323 system with a 1Hz, 2V ~ 6V sine drive signal, as shown in Figure 13(a). The tilt angles of the x and y  
 324 direction are reduced down to  $\pm 0.0014^\circ$  and  $\pm 0.0015^\circ$ , respectively, as shown in Figure 15, which are  
 325 within the optimal tilt angle range. When the frequency of the drive signal is changed to 0.2Hz, 0.5Hz  
 326 or 2Hz, the maximum tilt angles are all smaller than  $0.0018^\circ$ . This result indicates the robustness of  
 327 the  $H^\infty$  controller.



328

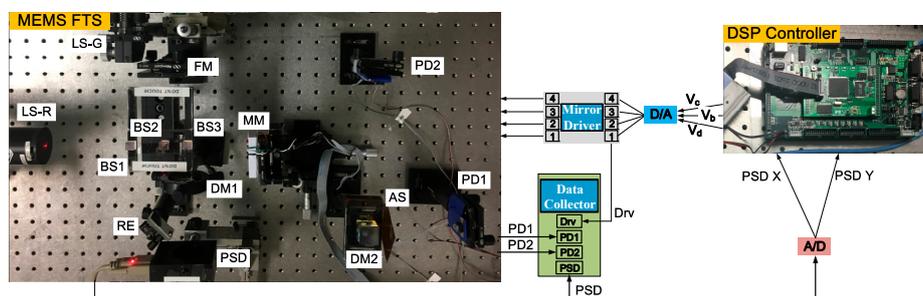
329

**Figure 15.** Measured tilt angle responses for FTS system with  $H^\infty$  control.

### 330 4.3. FTS Application

331 The FTS system as illustrated in Figure 12 has been built with the MEMS mirror. Figure 16 shows  
 332 a picture of the implemented system. The interferograms of the reference light and the unknown light  
 333 are picked up concomitantly by two photodetectors, PD1 and PD2, and then digitized by a data  
 334 acquisition module. The usable OPD scan range is only a fraction of the total MEMS mirror scan  
 335 range. Here, under the sinusoidal drive of 1Hz and 2V~6V, the OPD of the micromirror is  $221.8\mu\text{m}$ .  
 336 Theoretically, the spectral resolution is inversely proportional to the usable OPD scan range in which  
 337 the amplitude of the interferogram fringes does not have significant loss [1]. Then the Piecewise Cubic  
 338 Hermite Interpolating Polynomial (PCHIP) interpolation is employed to convert the testing light  
 339 interferogram data into an evenly sampled interferogram data of the testing light in the spatial  
 340 domain. After that, the spatial domain interferogram is transformed into a spectrogram via FFT and  
 341 Mertz phase correction [28]. The corresponding spectrum of the testing light source can be recovered  
 342 finally.

343



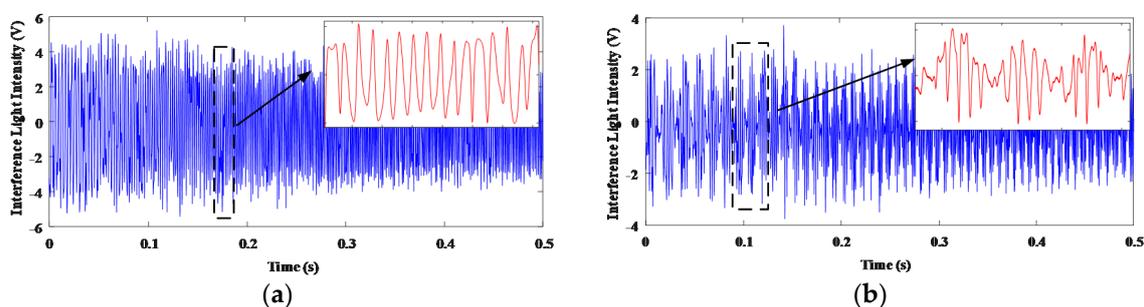
344

**Figure 16.** Experimental setup for MEMS FTS control system.

345 In the following, the experimental results of the FTS system with the MEMS under respective  
 346 look-up table control, PID control, and  $H^\infty$  control are compared and analyzed. In all four cases, the  
 347 MEMS was driven by a 1Hz and 2V~6V sinusoidal voltage signal.

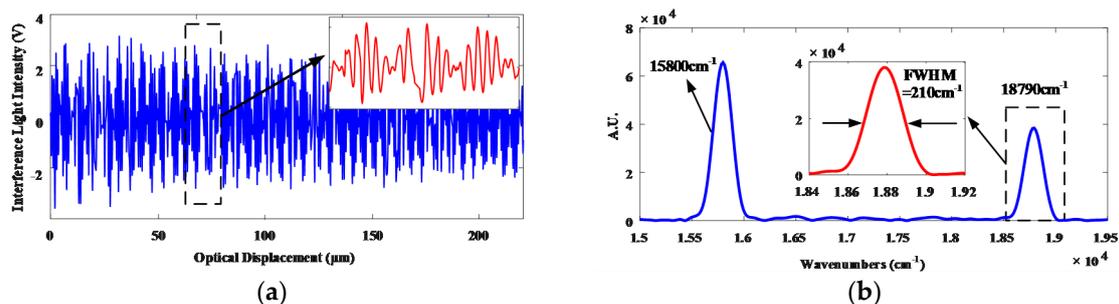
#### 348 1) Look-up table control

349 Figure 17 shows the interferograms of the reference light and the testing light when a look-up  
 350 table control was applied to the MEMS mirror. As the look-up table control is not robust, the quality  
 351 of the interferogram signals under the look-up table control is slightly worse than the PID control,  
 352 there are some burr noise.



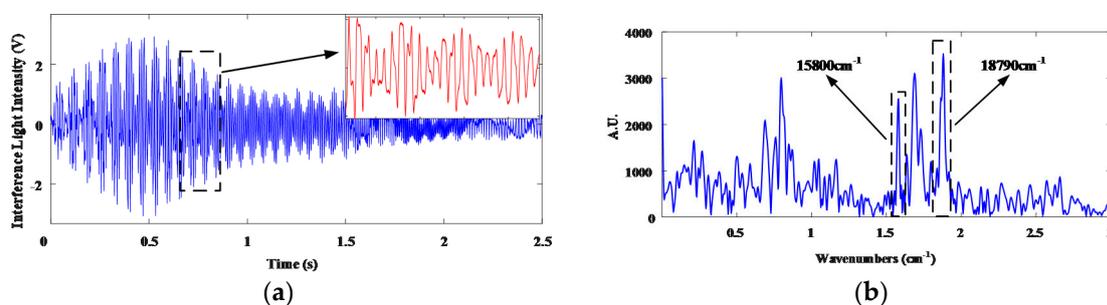
353 **Figure 17.** The interferogram signals acquired in time domain with Look-up table control using the same  
 354 driving as PID control (a) Reference light. (b) Testing light.

355 Figure 18(a) shows the reconstructed interferogram of the testing light in spatial domain with  
 356 look-up table control and the corresponding spectrum. Figure 18(b) shows the spectrum recovered  
 357 of the testing light in the system with Look-up table control. The measured FWHM resolution is  $210\text{ cm}^{-1}$ ,  
 358 corresponding to  $5.94\text{ nm}$ , Error is slightly larger.



359 **Figure 18.** Under Look-up table control using the same driving as PID control (a) The reconstructed  
 360 interferogram of the testing light in spatial domain testing light. (b) Spectrum recovered of the testing  
 361 light.

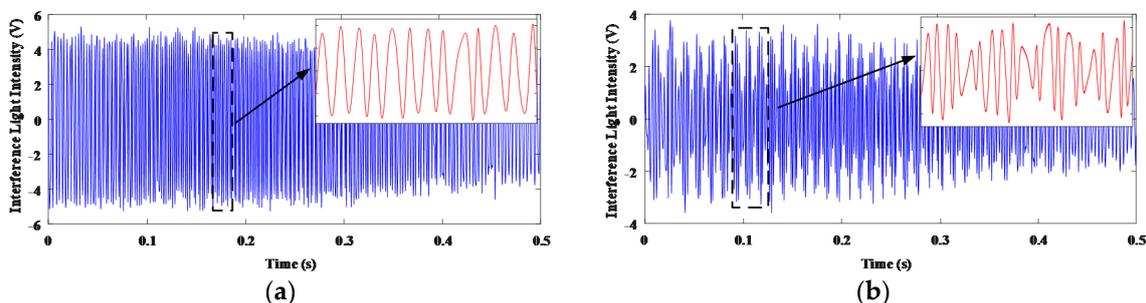
362 When the drive signal is 1Hz, 2.5V ~ 4.5V sinusoidal drive, as shown in Figure 19, tilting is  
 363 increased a lot, and the noise is also greatly increased. The resulting spectral quality is very poor, and  
 364 the spectrum of the test laser is mixed in the noise and is difficult to distinguish.



365 **Figure 19.** Under Look-up table control using different driving (a) The interferogram signals of testing  
 366 light acquired in time domain. (b) Spectrum recovered of the testing light.

367 2) PID control

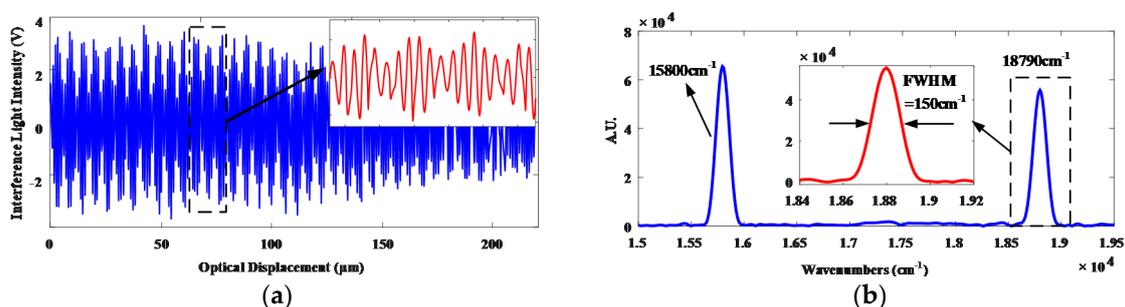
368 Figure 20 shows the interferograms of the reference light and the testing light, acquired when  
 369 the MEMS mirror was under PID control. The quality of the interferogram signals under PID control  
 370 is improved significantly. The distortion of the interferograms and the loss of the fringe contrasts  
 371 are much reduced. The envelope of the reference interferogram signal has only small variations,  
 372 indicating that there is only a small residual tilting of the MEMS mirror left in its full scan.



373 **Figure 20.** The interferogram signals acquired in time domain with PID control (a) Reference light. (b)  
 374 Testing light.

375 Figure 21 shows the reconstructed interferogram in spatial domain of the testing light and its  
 376 corresponding spectrum after FFT. It is obvious that under PID control the quality of the spectrum  
 377 is greatly improved. The spectral peaks of the testing light, which is a combination of the He-Ne laser  
 378 and the green laser, are detected accurately at 15800  $\text{cm}^{-1}$  and 18790  $\text{cm}^{-1}$ , or 632.9 nm and 532.2 nm  
 379 in wavelength, respectively. As shown in Figure 25(b), the measured FWHM resolution is 150  $\text{cm}^{-1}$ ,  
 380 corresponding to 4.24 nm at 532 nm wavelength. Since the Gaussian window is employed for

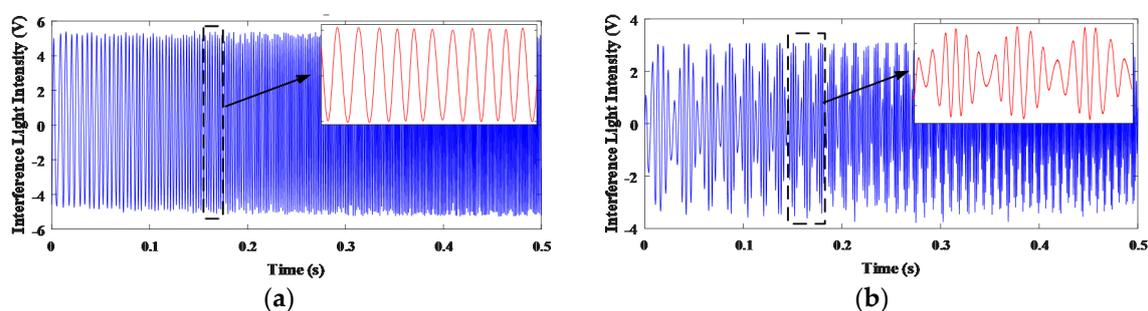
381 apodization, the theoretical value of the full-width at half-maximum (FWHM) resolution is given by  
 382  $2.0/\text{OPD}$  [1], or  $90.17 \text{ cm}^{-1}$  for this FTS, corresponding to  $2.55 \text{ nm}$  at  $532 \text{ nm}$  as the OPD is  $221.8 \mu\text{m}$ .  
 383 This deviation is believed to be caused mainly by the residual tilting of the MEMS mirror.



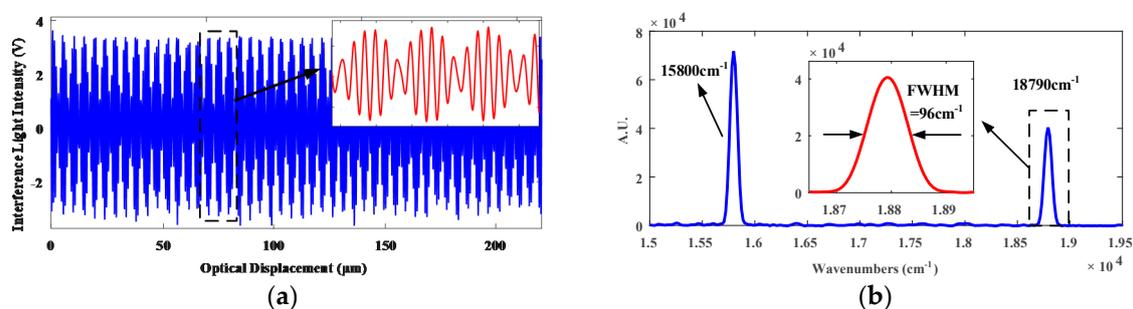
384 **Figure 21.** Under PID control (a) The reconstructed interferogram of the testing light in spatial domain  
 385 testing light. (b) Spectrum recovered of the testing light.

386 3)  $H^\infty$  robust control

387 Figure 22 shows the interferograms of the reference light and the testing light when the  $H^\infty$   
 388 control was applied to the MEMS mirror. Figure 23 shows the reconstructed interferogram of the  
 389 testing light in spatial domain with  $H^\infty$  control and the corresponding spectrum. Under the  $H^\infty$   
 390 control, the measured FWHM resolution is  $96 \text{ cm}^{-1}$ , corresponding to a spectral resolution  $2.71 \text{ nm}$   
 391 at  $532.2 \text{ nm}$ , which is in good agreement with the theoretical calculation. By changing the frequency  
 392 of the drive signal to  $0.2, 0.5$  and  $2.0 \text{ Hz}$ , the reconstructed spectra and measured FWHM's are  
 393 basically the same as those driven at  $1.0 \text{ Hz}$ , which further proves that the designed  $H^\infty$  robust  
 394 controller is robust.



395 **Figure 22.** The interferogram signals acquired in time domain with  $H^\infty$  control (a) Reference light. (b)  
 396 Testing light.



397 **Figure 23.** Under  $H^\infty$  control (a) The reconstructed interferogram of the testing light in spatial domain  
 398 testing light. (b) Spectrum recovered of the testing light.

## 399 5. Conclusions

400 In this paper, the methods of controlling the tilting of the MEMS mirror in an MEMS based  
 401 Fourier transform spectrometer are studied and experimentally verified. Based on the study of the

MEMS mirror's biaxial coupling relationship, a  $H^\infty$  mixed sensitivity controller is designed to suppress the tilting of the mirror plate in both x and y directions for the purpose of maintaining pure piston motion. Compared with the PID control and look-up table control, the  $H^\infty$  control has better robustness. Experimental results demonstrate that the residual tilting is as small as  $0.0014^\circ$  under the  $H^\infty$  control. In the built MEMS FTS system, the OPD generated by the MEMS mirror reaches  $221.8 \mu\text{m}$  and a spectral resolution of  $96 \text{ cm}^{-1}$ , or  $2.71 \text{ nm}$  at  $532 \text{ nm}$ , has been achieved. Compared with the previous work, the mirror tilting is reduced dramatically by the  $H^\infty$  control, the robustness and anti-interference capability of the FTS system are also improved.

Furthermore, considering that the LSF bimorph-based MEMS mirror employed in this study has large thermal response time ( $\sim 100 \text{ ms}$ ) which will limit the scan speed, our future work will focus on design, fabrication and control of a large-stroke electrothermal micromirror with faster thermal response [29].

**Acknowledgments:** We gratefully acknowledge the financial support of this research by the following projects: The Natural Science Foundation of Zhejiang Province, China (Grant No.LY15E050023) and National Natural Science Foundation of China (Grant No.51475129, No.51405117, No.U1509203 and No.U170920062).

**Author Contributions:** H.P. C., S.H. S. and Z.M. P. designed the study and wrote the paper, Y. Z. was responsible for the experimental design while M.Y. L. and C. C. processed the data. H.K. X. guided the project and also revised the final version of the paper.

**Conflicts of Interest:** The authors declare no conflict of interest

## References

1. P.R. Griffiths; J. A. De Haseth. Theoretical background. in *Fourier Transform Infrared Spectrometry*, 2nd ed. Hoboken, NJ; Wiley: USA, 2007; pp. 30–50 and 110–112, 9780471194040.
2. H. Xie; Z.-M. Qi. MEMS-based optical chemical sensors. in *Semiconductor Device-Based Sensors for Gas, Chemical, and Biomedical Applications*, S. J. Pearton, Ed. Boca Raton, FL; CRC Press: USA, 2011, pp. 268–275, 978-1-4398-1387-4.
3. Handheld and Portable FTIR Spectrometers for the Analysis of Materials: Taking the Lab to Sample. Available online: <http://www.americanlaboratory.com>(accessed Sep. 30, 2015).
4. U. Wallrabe; C. Solf; J. Mohra; J. G. Korvink. Miniaturized Fourier transform spectrometer for the near infrared wavelength regime incorporating an electromagnetic linear actuator. *Sensors & Actuators A Physical* **2005**, 123-124, 459-467, 10.1016/j.sna.2005.05.014.
5. T. Sandner; T. Grasshoff; E. Gaumont; H. Schenk; A. Kenda. Translatory MOEMS actuator and system integration for miniaturized Fourier transform spectrometers. *Journal of Micro/nanolithography Memes & Moems* **2012**, 13(1), 011115-1-13, 10.1117/1.JMM.13.1.011115.
6. N. Das; J. Sin; D. O. Popa; H. E. Stephanou. Design and manufacturing of a Fourier transform microspectrometer. in Proc. 8<sup>th</sup> IEEE Nanotechnol, Aug. 2008, pp. 837–840.
7. W. Jung et al. Three-dimensional optical coherence tomography employing a 2-axis microelectromechanical scanning mirror. *IEEE Journal of Selected Topics in Quantum Electronics* **2005**, 11(4), 806-810, 10.1109/JSTQE.2005.857683.
8. C. Duan et al. A  $45^\circ$ -tilted 2-axis scanning micromirror integrated on asilicon optical bench for 3D endoscopic optical imaging. in Proc. 28<sup>th</sup> IEEE MEMS, Jan. 2015, pp. 948-951.
9. O. Solgaard; A. A. Godil; R. T. Howe; L. P. Lee, Y.-A. Peter; H. Zappe. Optical MEMS: From micromirrors to complex systems. *J. Microelectromech. Syst* **2014**, 23(3), pp. 517–538, 10.1109/JMEMS.2014.2319266.
10. L. Wu; H. Xie. A millimeter-tunable-range microlens for endoscopic biomedical imaging applications. *IEEE Journal of Quantum Electronics* **2010**, 46(9), 1237-1244, 10.1109/JQE.2010.2048304.

- 450 11. W. Wang; J. Chen; A. S. Zivkovic; Q. A. A. Tanguy; H. Xie. A compact Fourier transform spectrometer on  
451 a silicon optical bench with an electrothermal MEMS mirror. *Journal of Microelectromechanical Systems* **2016**,  
452 *25*(2), 347-355, 10.1109/JMEMS.2016.2522767.
- 453 12. L. Wu; A. Pais; S. R. Samuelson; S. Guo, and H. Xie. A mirror tilt-insensitive Fourier transform spectrometer  
454 based on a large vertical displacement micromirror with dual reflective surface. in Proc. 15<sup>th</sup>  
455 TRANSDUCERS, Jun. 2009, pp. 2090–2093.
- 456 13. S. R. Samuelson; H. Xie. A large piston displacement MEMS mirror with electrothermal ladder actuator  
457 arrays for ultra-low tilt applications. *Journal of Microelectromechanical Systems* **2014**, *23*(1), 39-49,  
458 10.1109/JMEMS.2013.2290994.
- 459 14. W. Wang; S. R. Samuelson; J. Chen; H. Xie. Miniaturizing Fourier transform spectrometer with an  
460 electrothermal micromirror. *IEEE Photonics Technology Letters* **2015**, *27*(13), 1418-1421,  
461 10.1109/LPT.2015.2423637.
- 462 15. Y. Ma; S. Islam; Y. J. Pan. Electrostatic torsional micromirror with enhanced tilting angle using active  
463 control methods. *IEEE/ASME Transactions on Mechatronics* **2011**, *16*(6), 994-1001,  
464 10.1109/TMECH.2010.2066283.
- 465 16. H. Chen; W. J. Sun; Z. D. Sun; J. T. W. Yeow. Second order sliding mode control of a 2D torsional MEMS  
466 micromirror with sidewall electrodes. *Journal of Micromechanics & Microengineering* **2013**, *23*(1), 015006-1-9,  
467 10.1088/0960-1317/23/1/015006.
- 468 17. H. Chen; M. Pallapa; W. J. Sun; Z. D. Sun; J. T. W. Yeow. Nonlinear control of an electromagnetic polymer  
469 MEMS hard-magnetic micromirror and its imaging application. *Journal of Micromechanics &*  
470 *Microengineering* **2012**, *24*(4), 57-63, 10.1088/0960-1317/24/4/045004.
- 471 18. L. Wu; H. Xie. A Large vertical displacement electrothermal bimorph microactuator with very small lateral  
472 shift. *Sensors & Actuators A Physical* **2008**, *145–146*(7), 371-379, 10.1016/j.sna.2007.10.068.
- 473 19. Ferrara; Antonella; Lombardi; Cesare. Interaction control of robotic manipulators via second-order sliding  
474 modes. *International Journal of Adaptive Control & Signal Processing* **2007**, *21*(8-9), 708-730, 10.1002/acs.969.
- 475 20. Yu, Haoyong; Huang, Sunan; Chen, Gong; Pan, Yongping; Guo, Zhao. Human-Robot Interaction Control  
476 of Rehabilitation Robots with Series Elastic Actuators. *IEEE Transactions on Robotics* **2015**, *31*(5), 1089-  
477 1100, 10.1109/TRO.2015.2457314.
- 478 21. Li, Xiang; Pan, Yongping; Chen, Gong; Yu, Haoyong. Adaptive Human-Robot Interaction Control for  
479 Robots Driven by Series Elastic Actuators. *IEEE Transactions on Robotics* **2017**, *33*(1), 169-182,  
480 10.1109/TRO.2016.2626479.
- 481 22. Sivaranjani; Rakkiyappan; Cao; Alsaedi. Synchronization of nonlinear singularly perturbed complex  
482 networks with uncertain inner coupling via event triggered control. *Applied Mathematics & Computation*  
483 **2017**, *311*, 283-299, 10.1016/j.amc.2017.05.007.
- 484 23. Vromant, Pieter; Weyns, Danny; Sam, Malek; Andersson, Jesper. On interacting control loops in self-  
485 adaptive systems. Proceedings - International Conference on Software Engineering, p 202-207, 2011,  
486 SEAMS'11 - Proceedings of the 6th International Symposium on Software Engineering for Adaptive and  
487 Self-Managing Systems, Co-located with ICSE 2011.
- 488 24. E.H. Bristol. On a New Measure of Interaction for Multivariable Process Control. *Automatic Control IEEE*  
489 *Transactions on* **1966**, *11*(1), 133-134, 10.1109/TAC.1966.1098266.
- 490 25. Doyle J C; Glover K; Khargonekar P P et al. State-space solutions to standard H<sub>2</sub> and H<sub>∞</sub> control  
491 problems. *IEEE Transactions on Automatic Control* **1989**, *34*, 831–847, 10.1109/9.29425.
- 492 26. K. Glover; J.C. Doyle. State-space formulae for all stabilizing controllers that satisfy an H<sub>∞</sub> norm bound  
493 and relations to risk sensitivity. *Systems and Control Letters* **1988**, *11*(3), 167–172 , 10.1016/0167-  
494 6911(88)90055-2.
- 495 27. F. T. Han; W. Wang; X. Zhang; H. Xie. Modeling and control of a large-stroke electrothermal MEMS mirror  
496 for Fourier transform microspectrometers. *Journal of Microelectromechanical Systems* **2016**, *25*(4), 750-760,  
497 10.1109/JMEMS.2016.2565388.
- 498 28. C. D. Porter; D. B. Tanner. Correction of phase errors in Fourier spectroscopy. *International Journal of Infrared*  
499 *& Millimeter Waves* **1983**, *4*(2), 273-298, 10.1007/BF01008607.
- 500 29. X. Zhang; L. Zhou; H. Xie. A fast, large-stroke electrother mal MEMS mirror based on Cu/W. bimorph.  
501 *Micromachines* **2015**, *6*(12), 1876-1889, 10.3390/mi6121460.