Article

Microstructural Evolution, Thermodynamics and Kinetics of Mo-Tm₂O₃ Powder Mixtures during Ball Milling

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Abstract: The microstructural evolution, thermodynamics and kinetics of Mo-21%Tm₂O₃ (mass fraction, %) powder mixtures during ball milling were investigated using X-ray diffraction and transmission electron microscopy. Ball milling induced Tm₂O₃ to be decomposed and then dissolved into Mo crystal. The supersaturated nanocrystalline solid solution of Mo (Tm, O) was obtained after 96 h of ball milling. The elements of Mo, Tm and O were distributed uniformly in the ball-milled particles. Based on the semi-experimental theory of Miedema, a thermodynamic model was developed to calculate the driving force of phase evolution. There was no chemical driving force to form a crystal solid solution of Tm atoms in Mo crystal or an amorphous phase because the Gibbs free energy for both processes was higher than zero. For Mo-21%Tm₂O₃, it was mechanical work, not negative heat of mixing, that provided the driving force to form supersaturated nanocrystalline Mo (Tm, O) solid solution.

Key words: neutron absorber; Mo-based composites; ball milling; microstructure; thulium oxide

1. Introduction

Control rods are very important in both operating and accident conditions because nuclear reactivity needs to be controlled to safely operate a nuclear reactor [1]. The advantages of using gray control rods to control nuclear reactivity are that they greatly reduce the daily dealing with reactor coolant, distinctly simplify the chemical and tolerant control system and its operation, significantly decrease high expense for short time change, and also make the nuclear reactor safer [2-3]. The key to gray control rods controlling and adjusting nuclear reactivity is in their internal neutron absorbers, which can be prepared from bulk with elemental boron, Dy₂TiO₅ materials, Ag-In-Cd alloys, tungsten and silver and their alloys [4-7]. Initially, boron carbide and boron steels were used as neutron absorbers in Russian nuclear power water reactors such as VVER-1000, VVER-440 and RBMK-1000 [4]. However, large irradiation damage such as from swelling and cracking were induced by (n,a)-reactions, void formation and gas bubble growth. Ag-In-Cd alloys are usual strong neutron absorbers [6] and could be also used in gray control rods [7]. However, the nuclear reactivity worth of gray control rod assemblies drops to approximately 80% of its initial value after running five years, which no longer meets the control condition of mechanical shim. Dy₂TiO₅ pellets are used as neutron absorbers in Russian thermal reactors such as VVER-1000 RCCAs [5,8]. However, the neutron absorption ability decreased with increasing service time [9,10]. Other lanthanide elements such as terbium are needed as a reinforcing agent to balance the neutron absorption ability to keep the stability of the nuclear reaction.

According to nucleon calculation, nuclear property assessment and burnup analysis, thulium is an excellent candidate for use as a neutron absorber [7]. Natural thulium consists of one stable isotope, Tm¹⁶⁹. Both thulium decay products, Yb and Lu, and their offspring decay products can absorb neutrons. The absorption cross section of thulium is 105 barns, and the region of resonance absorption is 3.92-17.6 eV [9,10]. Theoretical calculation shows that the gray control rod assemblies of loaded neutron absorbers with thulium element still have a nuclear reactivity worth that is close to the initial value after a long time of burnup. Even if they are burned up for 20 years, the nuclear reactivity worth is changed only slightly. Therefore, it is important to develop and synthesize bulk material containing elemental thulium for use as neutron absorbers in gray control rods. According to in-service conditions in the nuclear reactor core, neutron absorbers should have high density, good thermophysical properties, sufficient physical efficiency and excellent irradiation resistance. Meanwhile, the size factor, fabricability of neutron absorber pellets in a practical application, and linear density of thulium in pellets must also be considered. Therefore, Mo-based Tm₂O₃ composites are designed and considered as excellent neutron absorbers based on the excellent characteristics of Mo component such as high melting temperature, high temperature strength, excellent thermal conductivity and outstanding irradiation resistance. Meanwhile, the thermal neutron absorption cross section of Mo (2.7 barns) is much smaller than that of Tm (105 barns), which has little effect on the behavior of nuclear characteristics of Tm element during long-time burnup. In addition, there is no the chemical reaction between Mo and Tm₂O₃ at the temperature below 1600 °C, which can ensure the stability of the phases in service. However, almost no studies relating to their microstructure, thermophysical properties, physical efficiency and irradiation behavior can be found.

High-energy ball milling is a non-equilibrium solid-state alloying technology for powder mixture preparation that can be used to synthesize novel materials impossible to obtain by conventional technology. It is well known that ball milling of powder mixtures can generate equilibrium and non-equilibrium structures including supersaturated solid solution, nanocrystalline powders, metastable compounds and amorphous solids [11-14]. The grain size and the agglomeration degree of mixture particles have major influence on both the packing achievement in the green bodies and the final sintered density and grain size [15]. The nanocrystalline powders have excellent sintering ability, low-sintering temperature and a high tendency for density increase [16], which can improve the properties and performance of sintered bulks.

Therefore, based on the defects of neutron absorbers such as bulk with elemental boron, Dy₂TiO₅ materials and Ag-In-Cd alloys, and the advantages of Mo and Tm₂O₃ components, the Mo-Tm₂O₃ composite will be an excellent candidate as neutron absorber. In the present work, ball-milling technology was first used to prepare Mo-Tm₂O₃ powder mixtures. The microstructural evolution behavior, thermodynamics and kinetics of Mo-Tm₂O₃ powder mixtures during ball milling were investigated. The corresponding mechanism was also analyzed and discussed. The current experimental results of ball-milled powder mixtures can provide sufficient support for subsequent bulk preparation.

2. Experiments

The powders of Mo (99.9% purity) and Tm_2O_3 (99.6% purity) with an average particle diameter of 60 μ m and 5 μ m, respectively, were used as raw materials. The raw powders of Tm_2O_3 and Mo were purchased from Beijing HWRK Chem Co., LTD of China and Jiangxi KETAI

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Advanced Material Co., LTD of China, respectively. Ball milling of Mo-21% Tm₂O₃ (mass fraction, %) powder mixtures was carried out at room temperature on a SFM-1 high-energy planetary ball mill manufactured by Shenyang Kejing Auto-instrument Co., Ltd of China. The mass fraction of Mo-21% Tm₂O₃ is equal to the molar fraction of Mo-6.2 at% Tm₂O₃. Stainless steel balls of 5 mm in diameter were used as milling media. The ball-to-powder mass ratio was 10:1 and the rotational speed was 500 rpm. No more than one weight percent stearic acid was added in powder mixtures as a process control agent to prevent excessive cold welding and aggregation amongst powder particles. During ball milling, a 5 min interval stopping was used after running 55 min to prevent excess generated heat that had an obvious effect on the ball-milling procedure. In order to avoid contamination, the powder mixtures were firstly milled at low speed for a few minutes, so a powder coating covered the stainless steel balls and vial wall. Meanwhile, the powder mixtures used for microstructure analysis were extracted from loose powders in the steel can, not from powders stuck on the stainless ball surface and the steel can wall surface.

After ball milling over different time periods, small amounts of powder mixtures were taken from the container and characterized by X-ray diffraction (XRD) on a Rigaku D/max-3C X-ray diffractometer with Cu K α radiation ($\lambda = 0.1540598$ nm). The grain size was calculated using Suryanarayana and Grant Norton's formula [17],

$$B_r \cos \theta = \frac{R\Lambda}{L} + \eta \sin \theta \tag{1}$$

where, K is a constant (with a value of 0.9); λ is the wavelength of the X-ray radiation; L and η are the grain size and internal strain, respectively; and θ is the Bragg angle. B_r is the full width at half-maximum (FWHM) of the diffraction peak after instrumental correction and can be calculated from the following equation:

$$B = B_r + B_s \tag{2}$$

where, B and Bs are the FWHM of the broadened Bragg peaks and the standard sample's Bragg peaks, respectively.

The microstructure of the ball-milled powders was analyzed using transmission electron microscopy (TEM) on a JEM-2100 instrument. The ball-milled mixtures were first put in the ethyl alcohol, and then adequately dispersed by ultrasonic vibration. A carbon-coated copper grid was used to collect the dispersed powders in the ethyl alcohol and then dried by ultraviolet lamp. After that, the prepared samples were observed by TEM instrument.

3. Results and discussion

3.1. Phase evolution and microstructure analysis

The XRD patterns of Mo-21% Tm₂O₃ powder mixtures milled for different times are shown in Fig. 1. The diffraction peaks of the Mo and Tm₂O₃ phase are broadened significantly and reduced in intensity with increasing milling time. Meanwhile, a small position shift of Mo peaks to a low diffraction angle is detected, which indicates the large size Tm atoms are dissolved into Mo crystal structure.

The diffraction peaks indicate Tm₂O₃ crystal structure disappears after 24 h of ball milling, which is due to the destruction of the Tm₂O₃ crystal structure by ball milling. Only a diffraction hill indicating Tm₂O₃ amorphization can be observed in the XRD spectrum as shown in Fig. 1 (b). Meanwhile, the intensity of the diffraction hill decreases with increasing milling time. The O atoms being dissolved into Mo crystal structure. This result is similar with Zhang's investigation that showed the intensity of Y₂O₃ diffraction peaks almost disappeared after 52 h of ball milling in a Co-based ODS alloy, and this was attributed to Y₂O₃ decomposition and dissolution into the Co crystal structure during ball milling [18]. The effect of ball milling is more evident in the Tm₂O₃ phase than the Mo phase. In addition, no diffraction peaks of new a phase among Mo, Tm and O elements are observed, which indicates that no intermetallic compounds are formed during ball milling. Ball milling does not induce chemical reaction between Mo and Tm₂O₃, and the amount of intermetallic compound formed between Mo and Tm₂O₃ is also negligible.

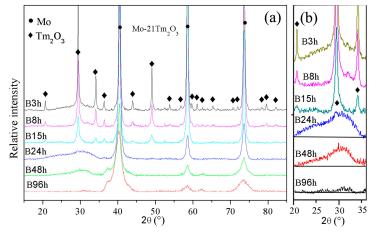


Fig. 1 (a) and (b) X-ray patterns of Mo-21% Tm₂O₃ powder mixtures milled for different times at 20 of 15° to 85° and 20° to 36°, respectively.

The broadening of X-ray diffraction peaks is associated with the refinement of grain size and lattice distortion during ball milling. Fig. 2 (a) shows the variation of the Mo lattice parameter with increasing ball-milling time. It is demonstrated that the lattice parameter increases from 0.31474 to 0.31564 nm as the ball-milling time increases from 0 to 96 h. Actually, in the initial stage, the Mo lattice parameter almost remains constant. After further ball milling, the Mo lattice parameter increases with increasing ball-milling time. The reason for the Mo lattice parameter incrementation can be attributed to the dissolution of Tm and O atoms. The supersaturated nanocrystalline solid solution of Mo (Tm, O) is formed after 96 h of ball milling in this system. This result differs from Raghavendra's research, which showed that the diffraction peaks of the Fe and ZrO₂ phase could be still observed after 100 h of ball milling in a Fe-15wt%ZrO₂ system [19]. However, Toualbi reported the diffraction peak disappearance of the Y₂O₃ phase in a Fe-9Cr-10%wt Y₂O₃ system during ball milling was due to the dissolution of a very small amount of yttria particles into matrix to form solid solution and to the amorphization of a large amount of Y₂O₃ particles [20].

Fig. 2(b) shows the grain size of the ball-milled powders with increasing ball-milling time. It indicates that the refinement rate of crystallite size is roughly logarithmic with ball-milling time. In fact, the grain size was calculated for Mo phase, not for Tm₂O₃ phase, because Tm₂O₃ peaks completely disappeared. It can be seen that ball milling results in a fast decrease of Mo grain size in the initial stage and a constant value in the later stage. The average grain size of the Mo phase is approximately 74 nm and 35 nm after ball milling of 0 and 3 h, respectively. The decrease rate is approximately 13 nm/h. After further ball milling from 3 to 24 h, the average grain size decreases from 35 to 17 nm. The decrease rate is approximately 0.85 nm/h. However, the average grain size

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only decreases from 17 to 8 nm after ball milling from 24 to 96 h. The decrease rate is only 0.125 nm/h, which indicates that further refinement of grain size occurs slowly after extended milling time. In particular, Mo grain size almost remains constant at 8 nm from 48 to 96 h of ball milling. The lattice distortion calculated from the X-ray broadening exhibits an increasing manner. The total lattice distortion is approximately 1.5% after 96 h of ball milling.

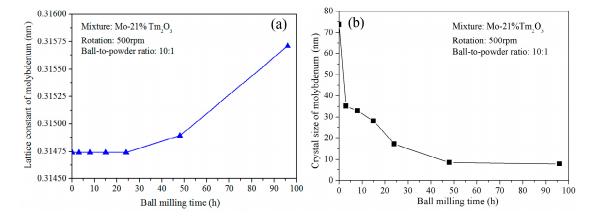


Fig. 2 (a) Curve of lattice parameter of Mo vs. ball-milling time; (b) Curve of crystalline size of Mo vs. ball-milling time

Fig. 3 is SEM image, bright field TEM image, the corresponding selected area electron diffraction (SAED) pattern and energy dispersive spectrometer (EDS) results of Mo-21% Tm₂O₃ powder mixtures milled for 48 h. SEM image in Fig. 3(a) shows the morphology of the ball-milled powder mixtures, which indicates that many particles with small size aggregates to form a large size particle. The size of particle is over one micrometer. A spherical particle with a size of approximately 500 nm can be observed in the TEM image. Some defects such as dislocations, voids and crystal boundaries can also be observed in this particle, which were induced by heavy deformation of powder mixtures during ball milling. These defects enhance the diffusivity of solute atoms into the Mo matrix. The SAED pattern taken from the region marked letter 'A' in Fig. 3 (b) is shown in Fig. 3 (c). After analyzing and indexing the ring-shaped SAED pattern, it is indicated that this SAED pattern belongs to the Mo phase, not the Tm₂O₃ phase. Meanwhile, it is also demonstrated that the Mo component has already been nano-crystallized, which is also confirmed by the calculation results of Mo crystallites derived from Bragg diffraction peaks in Fig. 1. Although the size of the particles is still in the micrometer as shown in Fig. 3(a), the grain size is up to nanometer level. The EDS spectrum taken from the ball-milled particle is shown in Fig. 3(d), which includes Mo, O and Tm element peaks. Similar EDS results are detected in other particles. Although the SAED results show that the particle belongs to an Mo crystal, Tm and O elements are also detected in this particle. Therefore, it is indicated that Tm₂O₃ either uniformly distributes on the Mo particle surface in amorphous form or dissolves into the Mo crystal structure in the form of Tm and O atoms. However, few diffraction halos related to the amorphous phase can be found in this SAED pattern as show in Fig. 3(c), which indicates that Tm₂O₃ does not mainly exist in the amorphized form in this particle.

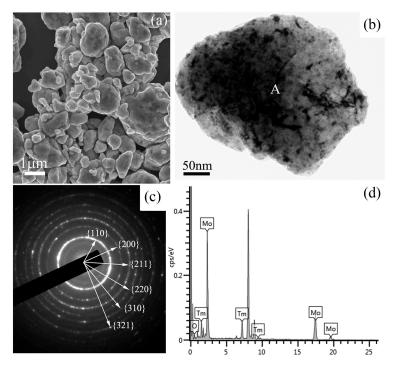


Fig. 3 Microstructure analysis results of Mo-21% Tm₂O₃ powder mixtures milled for 48 h, (a) SEM image showing the morphology of the ball-milled powders; (b) Bright field TEM image; (c) Selected area electron diffraction pattern from region 'A' in (a) showing Mo phase; (d) EDS analysis results from region 'A' in (a) showing the diffraction peaks of Tm, Mo, O and Cu elements. Cu peaks comes from copper grid.

Fig. 4 (a) is high angle annual dark field STEM image of Mo-21% Tm₂O₃ powder mixtures milled for 96 h, which shows several small particles agglomerated to form a large-sized particle. Fig. 4 (b), Fig. 4 (c) and Fig. 4 (d) are Mo L, Tm L and O K element mappings of the particles in Fig. 4 (a), respectively. It can be seen that Mo, Tm and O atoms are uniformly distributed in the particles. Every location contains Mo, Tm and O elements, which further demonstrates that Tm₂O₃ dissolves into Mo crystal structure after 96 h of ball milling because the diffraction hill indicating Tm₂O₃ amorphization disappears completely after 96 h of ball milling as shown in Fig. 1 (b).

In fact, the Tm₂O₃ component is brittle powder, which is fragmented during ball milling. The particle size reduces continuously as a consequence of the energy provided by ball milling. However, the Mo component belongs to ductile powder, which is repeatedly flattened, cold welded, fractured and rewelded by the force of the impact. In the initial stage of ball milling, the microstructural evolution of this ductile-brittle Mo-Tm₂O₃ system is that the ductile Mo metal powder particles get flattened by the ball-powder-ball collisions, while the brittle Tm₂O₃ oxide get fragmented. These fragmented Tm₂O₃ particles tend to become occluded by the ductile Mo constituents and trapped in the ductile Mo particles. With further ball milling, Mo particles get work hardened, fractured, and refined. The newly created surface of the Mo component has a large number of dangling bonds that have high chemical activity, which makes the fined Mo particles aggregate together to form a large-sized particle and also combined with dangling bonds of the fined Tm₂O₃ particles at an atomic level. Meanwhile, fragmented Tm₂O₃ particles will be continuously fined to smaller-sized particles with increasing ball-milling time and will be decreased to several nanometers. In particular, as the size of Tm₂O₃ is decreased to 1~2 nanometers, the ratio of atoms on the particle surface is very large. With increasing ball-milling time, Tm_2O_3 is transformed gradually from crystal structure to amorphous phase as shown in Fig. 1(b). The amorphous atom arrangement is disordered over long distance, but ordered over a very short distance, which could induce Tm and O atoms to attach on a Mo particle surface in a state of dissociation. Tm and O atoms can easily diffuse into the Mo crystal structure. Meanwhile, heavy deformation is introduced continually into Mo particles. This is manifested by the presence of a variety of crystal defects such as dislocations, vacancies, stacking faults, and an increased number of crystal defect boundaries [21], which can enhance the diffusivity of Tm and O atoms into the Mo matrix. Moreover, the refined microstructural features decrease Tm and O atom diffusion distances. Additionally, the rise in temperature during ball milling further promotes diffusion. Consequently, Tm and O atoms dissolve into the Mo crystal structure. The longer the ball-milling time is, the greater the amount of dissolution of Tm atoms. After 96 h of ball milling, the supersaturated nanocrystalline solid solution of Mo (Tm, O) is formed.

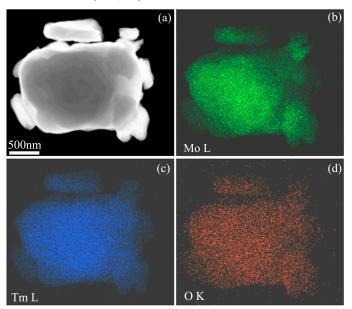


Fig. 4 STEM image of Mo-21% Tm₂O₃ powder mixtures milled 96 h, (a) High angle annual dark field image; (b), (c) and (d) Mo L, Tm L and O K element mappings, respectively

3.2. Thermodynamics analysis

As mentioned above, long time ball milling induces the Tm₂O₃ component to decompose and then dissolve into Mo crystal structure. Therefore, thermodynamics will be used to analyze the dissolution behavior of Tm into Mo crystal structure.

Based on Miedema semi-experimental theory [22], the Gibbs free energy of a Mo-Tm solid solution is estimated as:

$$\Delta G^s = \Delta H_m^s - T \Delta S^s \tag{3}$$

where ΔH_m^s and ΔS^s are mixing enthalpy and mixing entropy, respectively. With the assumption of an ideal solution, ΔS^s can be expressed as:

$$\Delta S^{s} = -R(x_{i} \ln x_{i} + x_{i} \ln x_{i}) \tag{4}$$

where x_i and x_j are the molar fraction of element A and B, $x_i + x_j = 1$. R is the gas constant and T is the reaction temperature.

The Gibbs free energy, ΔG^s , can be rewritten as

$$\Delta G^{s} = \Delta H_{m}^{s} + RT(x_{i} \ln x_{i} + x_{j} \ln x_{j})$$

$$\tag{5}$$

The enthalpy, ΔH_m^s , of a solid solution formation can be written as [23]:

$$\Delta H_m^s = \Delta H_{chem} + \Delta H_{Elast} + \Delta H_{Struct}$$
 (6)

where ΔH_{Chem} , ΔH_{Elast} and ΔH_{Struct} are the chemical, elastic and structural contributions due to mixing of the two different atoms, the atom size mismatch, and taking into account the difference in valence and crystal structure of solute and solvent, respectively. Compared with the first two items, the structural contributions have only a minor effect [24], which will be ignored here. Thus ΔH_m^s will be rewritten as:

$$\Delta H_m^s = \Delta H_{chem} + \Delta H_{Elast} \tag{7}$$

 ΔH_{Chem} can be expressed by:

$$\Delta H_{Chem} = x_i x_j \left(f_i^i \Delta H_{Sol}^{iinj} + f_i^j \Delta H_{Sol}^{jini} \right)$$
 (8)

where ΔH_{Sol}^{iinj} is the solution enthalpy of i in j given by:

$$\Delta H_{Sol}^{iinj} = \frac{2PV_i^{2/3}}{\left(n_{ws}^{-1/3}\right)_i + \left(n_{ws}^{-1/3}\right)_j} \left[-\left(\Delta\phi^*\right)^2 + \frac{Q}{P} \left(\Delta n_{ws}^{1/3}\right)^2 \right]$$
(9)

and f_i^J is the parameter of adjacent atom in solution solid and can be expressed as:

$$f_i^j = C_i^S \left[1 + \delta \left(C_j^S . C_i^S \right)^2 \right] \tag{10}$$

$$C_i^S = \frac{x_i V_i^{2/3}}{\left(x_i V_i^{2/3} + x_j V_j^{2/3}\right)}$$
(11)

where $\Delta \phi^*$, V and n_{ws} are the work function, molar volume and electron density of the constituents, respectively. P and Q are empirical constants having the same value for widely different metal combinations. δ is taken to be five for short-range order and eight for long-rang order, respectively.

 ΔH_{Elast} can be expressed as:

$$\Delta H_{\text{Elast}} = x_i x_i \left(x_i \Delta E_e^{\text{iinj}} + x_i \Delta E_e^{\text{jini}} \right)$$
 (12)

where ΔE_e^{iinj} is the size mismatch contribution to the enthalpy of a solution of i and j per mol i and can be estimated as [25]:

$$\Delta E_e^{iinj} = \frac{2K_i u_j (V_j^* - V_i^*)^2}{3K_i V_j^* + 4u_j V_i^*}$$
 (13)

where K, u, V^* are the bulk modulus, shear modulus, and atomic volume, respectively. V^* can be expressed as:

$$(V_i^*)^{2/3} = V_i^{2/3} \left(1 + c f_j^i \Delta \phi_{AB} \right)$$
 (14)

where *c* is constant.

The Gibbs free energy of formation of an amorphous phase can be estimated by [26]:

$$\Delta G^{Amorp.} = \Delta H^{Amorp.} - T\Delta S^{Amorp.} + x_i \Delta G_i^{Amorp.-Cryst.} + x_j \Delta G_i^{Amorp.-Cryst.}$$
(15)

where $\Delta H^{Amorp.}$ and $\Delta S^{Amorp.}$ are the enthalpy and entropy of mixing of the amorphous phase.

 $\Delta G_i^{Amorp.-Cryst.}$ is the difference in Gibbs free energy between the amorphous and crystalline phase of the pure element at room temperature and can be calculated according to the formula proposed by Thompson:

$$\Delta G_{i(j)}^{Amorp,-Cryst.} = \frac{2T\Delta H_f(T_m - T)}{T_m(T + T_m)}$$
(16)

where ΔH_f and T_m are the enthalpy of fusion and melting temperature, respectively.

 $\Delta H^{Amorp.}$ contains only the chemical contribution due to the amorphous structure, which can be calculated by:

$$\Delta H^{Amorp.} = \Delta H_{Chom} \tag{17}$$

$$\Delta S^{Amorp.} = -R(x_i \ln x_i + x_i \ln x_i)$$
(18)

The calculated Gibbs free energy changes on forming an amorphous phase and crystalline solid solution in the Mo-Tm system is shown in Fig. 5. The temperature for calculating the Gibbs free energy is taken to be 373 K [11]. The parameters used for calculation are listed in Table 1 [27,28]. The calculated Gibbs free energy of both crystalline solid solution and amorphous phase are higher than zero for all compositions. Therefore, there is no chemical driving force to form a crystalline solid solution or an amorphous phase. From Fig. 5, it can be seen that, with molybdenum content in the range of 0~9.2 at.% or 81.26~100 at.%, the Gibbs free energy for the formation of an amorphous phase is higher than that of a crystalline solid solution. The Gibbs free energy curve of the crystalline solid solution intersects with that of the amorphous phase at the 9.2 at.% and 81.26 at.% Mo content.

Table 1 Thermodynamics parameters for calculation [27,28]

Element	P(KJ V ⁻² cm ⁻¹)	Q(KJ cm)	n ^{1/3} (cm ⁻¹)	$\Phi(V)$	K(10 ¹⁰ N m ⁻²)	$\mu(10^{10}N\;m^{\text{-}2})$	V(cm ³ mol ⁻¹)	Tm(K)	Hf (KJ/mol)
Mo	12.35	132.54	1.77	4.65	26.12	12.56	9.4	2890.2	37.48
Tm	12.35	132.54	1.23	3.12	44.5	30.5	18.124	1818	16.8

In the present work, the molar percentage of Tm atoms is 11.68 at.%. The Gibbs free energy for the formation of an amorphous phase and a crystalline solid solution is 16.36 KJ/mol and 13.5 KJ/mol, respectively. The Gibbs free energy of amorphization is higher than that of crystalline solid solution formation. Therefore, the formation of crystalline solid solution is energetically preferred over the amorphous phase. However, the Gibbs free energy of the formation of crystalline solid solution and amorphization are both higher than zero for this chemical composition. There is no chemical driving force to form a crystalline solid solution and an amorphous phase. Meanwhile, there is also no chemical driving force for the transformation from crystalline to amorphous at this chemical composition. To more deeply understand the mechanism for microstructural evolution of Mo-Tm₂O₃ powder mixtures, especially in regard to the dissolution of Tm atoms into Mo crystal structure, the reaction kinetics during ball milling must be considered.

In addition, although the calculated Gibbs free energy of both crystalline solid solution and amorphous phase are higher than zero for all compositions as shown in Fig.5, these calculated values are limitations and cannot completely represent actual situation. Because the studied system is Mo-Tm₂O₃ powder mixtures, but the calculated Gibbs free energy comes for the Mo-Tm system. The Mo-Tm system originated from Mo-Tm₂O₃ system during the ball milling process is a little bit different with the Mo-Tm system of the equilibrium state. However, ball milling induced the decomposition of the Tm₂O₃ component followed by solvation into the Mo crystal structure, the investigation of Gibbs free energy using Mo-Tm system to take the place of Mo-Tm₂O₃ system is also reasonable.

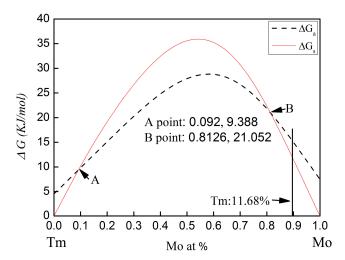


Fig. 5 Gibbs free energy of the Mo-Tm system vs. molybdenum content. ΔG_s and ΔG_a are the Gibbs free energy of the crystalline solid solution and amorphization, respectively

3.3. Dynamics analysis

The atomic radius of the Mo atom and Tm atom is 0.19 nm and 0.222 nm, respectively [29]. The difference in atomic radius is approximately 15%. At high pressure, the radii of the Mo atom

and Tm atom will reach the same value. Under this condition, the Tm atom will easily dissolve into Mo crystal structure. The required pressure is given by the calculation process as follows:

$$R_{\text{Mo}} (1-\varepsilon_{\text{Mo}}) = R_{\text{Tm}} (1-\varepsilon_{\text{Tm}}) \tag{19}$$

where $\varepsilon_{\text{Mo}} = q^*/E_{\text{Mo}}$, $\varepsilon_{\text{Tm}} = q^*/E_{\text{Tm}}$. Here, R is the initial atom radius; q^* is the pressure; E is the bulk elastic modulus and ε is the amount of strain. Therefore, equation (19) can be rewritten as follows:

$$R_{\text{Mo}} \left(1 - q^* / E_{\text{Mo}} \right) = R_{\text{Tm}} \left(1 - q^* / E_{\text{Tm}} \right) \tag{20}$$

where $R_{\text{Mo}} = 0.19$ nm, $R_{\text{Tm}} = 0.222$ nm, $E_{\text{Mo}} = 230$ GPa and $E_{\text{Tm}} = 45$ GPa. The q^* value can be obtained according to equation (20) and is $q^* = 7.79$ GPa. Therefore, when the pressure is over 7.79 GPa, the radii of the Mo atom and Tm atom will become equivalent. In fact, during high-energy ball milling, a pressure of 7.79 GPa is not difficult to obtain and in fact the maximum pressure value q_0 will be up to 8 GPa according to the following theoretical calculations.

Assuming that two impacting balls have the same elastic properties, and taking γ = 0.3, the calculation formula is shown as follows [30]:

$$\alpha = 1.23 \times \sqrt[3]{\frac{P^2}{E^2} \frac{R_1 + R_2}{R_1 R_2}}$$
 (21)

$$q_0 = 1.5 \times \frac{p}{\pi a^2} = 1.5\overline{q} = 0.388 \times \sqrt[3]{PE^2(\frac{R_1 + R_2}{R_1 R_2})^2}$$
 (22)

where α is the amount of elastic deformation between two colliding mill balls; P is the compression force under the colliding process; E is Young's modulus; R_1 and R_2 are the radius of the colliding balls; q_0 is the pressure at the center of the surface of contact, which is the maximum pressure of the contacting surface; and \overline{q} is the average pressure on the contact surface of two colliding mill balls.

The velocity of the ball in the present work is taken as v = 2 m/s. The weight of a mill ball is 0.51 g. The calculated impact energy is $E_k = 2 \times 0.5 \text{m}v^2 = 0.0024$ J. Therefore, the pressure q of two impacting balls can be obtained using the following integral equation (23).

$$E_k = 2\int_0^{\alpha/2} p dx \tag{23}$$

Using equation (21) and (23), E_k is: $E_k = 1.923 \times 10^9 \alpha^{5/2}$. The compression force P can then be calculated and is P = 355.623 N. According to equation (22), $q_0 = 8.369$ GPa and $\overline{q} = 5.58$ GPa.

Therefore, the value of $q_0 = 8.369$ GPa is larger than that of $q^* = 7.79$ Gpa. The Tm atoms are easily dissolved into the Mo crystal structure when the Tm atom radius is less than or equal to the Mo atom radius. Certainly, at the above discussion, the calculation is an ideal condition. Actually, the collision of two balls is random and is not always frontal collision because of random motion of grinding balls in the can. The values of collisional pressure are in a Gaussian distribution. Part pressure can over the q^* value, which will induce the radius of Tm atom to reach that of Mo atoms. Therefore, a long time is needed to make all Tm atoms dissolving into Mo crystal structure.

Simultaneously, the temperature at the colliding surface will instantaneously increase several hundred degrees Celsius resulting from the impact of the balls, which promotes the diffusion and dissolution process of Tm atoms. Davis reported that the maximum temperature rise during ball milling was approximately 300 °C in a Fe-1.2wt% C steel system [31]. Tonejc's research showed that the local rise in temperature was at least 570 °C in a Ga-Sb alloy system [32]. Meanwhile, the large defect densities increase the free energies of atoms, which also promotes Tm atom dissolution.

4. Conclusions

The microstructural characteristics of Mo-21%Tm₂O₃ (mass fraction,%) powder mixtures during ball milling were investigated by using transmission electron microscopy and X-ray diffraction. The thermodynamics and kinetics of phase evolution were also analyzed and discussed. The experimental results showed that the powder mixtures were first fined and nano-crystallized, and then transformed into supersaturated nanocrystalline solid solutions. The diffraction peaks of the Mo phase shifted to low angles during ball milling. The diffraction peaks of the Tm₂O₃ phase disappeared completely after 96 h of ball milling, which was attributed to ball milling induced decomposition of the Tm₂O₃ component followed by solvation into the Mo crystal structure. The elements of Mo, Tm and O were distributed uniformly in the ball-milled particles. After 96 h of ball milling, a supersaturated nanocrystalline solid solution of Mo (Tm, O) was formed. Based on the semi-experimental theory of Miedema, a thermodynamic model was developed to calculate the driving force of phase evolution. There was no chemical driving force to form a crystal solid solution of Tm atoms in the Mo crystal structure or an amorphous phase because the Gibbs free energy was higher than zero for both processes. For Mo-21%Tm₂O₃, it was mechanical work, not negative heat of mixing, that provided the driving force to form supersaturated nanocrystalline Mo (Tm, O) solid solution. The current experimental results of the ball-milled powder mixtures can provide sufficient support for subsequent bulk preparation. The relationships between mechanical properties, thermophysical properties, physical efficiency and irradiation properties of bulk materials and the characteristic microstructure of several kinds of ball-milled powders will be achieved in the future based on the current study results.

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Author Contributions

Guang Ran conceived and designed the experiments; Yong Luo performed the experiments; Yong Luo and Yaoli Zhang obtained Fig.5; Qiang Shen did TEM experiments; Nanjun Chen, Yong Luo and Guang Ran wrote and edited the paper.

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References

[1] H.S. Kim, C.Y. Joung, B.H. Lee, H.S. Kim, D.S. Sohn, Applicability of CeO_2 as a surrogate for PuO_2 in a MOX fuel development, J. Nucl. Mater 378 (2008) 98-104.

[2] R.T. Lahey, F.J. Moody, The thermal-hydraulics of a boiling water nuclear reactor, America Nuclear Society, 1993.

- [3] L. Chen, Q. X. Ding, G. W. Bi, et al, inventor; Nuclear Industry Patent Center, assignee. A gray control rods and absorber. The People's Republic of China patent 201210127376.X. 2012 Apr.27.
- [4] V.D. Risovany, E.E. Varlashova, D.N. Suslov, Dysprosium titanate as an absorber material for control rods, J. Nucl. Mater. 281(2000) 84-89.
- [5] V.D. Risovany, E.P. Klochkov, E.E. Varlashova, Hafnium and dysprosium titanate based control rods for thermal water-cooled reactors, At. Energ. 81(1996) 764-769.
- [6] J. Bourgoin, F. Couvreur, D. Gosset, F. Defoort, M. Monchanin, X. Thibault, The behavior of control rod absorber under irradiation, Journal of Nuclear Materials. 275 (1999) 296-304.
- [7] J. Q. Lu, C. T. Tang, H. Li, B. Yang, J. W. Li, Q. X. Ding, L. B. Zhu, J. Z. Liu, inventor; Nuclear Industry Patent Center, assignee. An advanced gray control rods and absorber. The People's Republic of China patent 201210127037.1. 2012 Apr.27.
- [8] V.D. Risovany, E.E. Varlashova, D.N. Suslov, Dysprosium titanate as an absorber material for control rods, Journal of Nuclear Materials. 281 (2000) 84-89(6).
- [9]E.M. Baum, M.C. Ernesti, H.D. Knox, T.R. Miller, A.M. Waston, S.D. Travis, Nuclides and Isotopes, 7th edn, Knolls Atomic Power Laboratory, 2009, p 66-67.
- [10] A.O. Brunfelt, E. Steinnes, Determination of Lutetium, Ytterbium and terbium in rocks by neutron activation and mixed solvent anion-exchange chromatography, Analyst. 94 (1969) 979-984.
- [11] C. Suryanarayana, Mechanical alloying and milling, Progress in Materials Science. 46 (2001) 1-184.
- [12]S.Q. Xi, K.S. Zuo, X.G. Lin, G. Ran, J.E. Zhou, Study on the solid solubility extension of Mo in Cu by mechanical alloying Cu with amorphous Cr(Mo), Acta Materialia.56 (2008) 6050-6060.
- [13]G. Ran, J.E. Zhou, Formation of nanocrystalline and amorphous phase of Al-Pb-Si-Sn-Cu powder during mechanical alloying, Materials Science and Engineering A. 416 (2006) 45–50.
- [14] A.O. Aning., Z. Wang, T.H. Courtney., Tungsten solution kinetic and amorphization of nickel in mechanically alloyed Ni-W alloys, Acta Metall Mater. 41 (1993) 165-174.
- [15]P. Bowen, C. Carry, From powders to sintered pieces: forming, transformations and sintering of nanostructured ceramic oxides, Powder Technol. 128 (2002) 248–255.
- [16]H.Ferkel, R.J. Hellmig, Effect of nanopowder deagglomeration on the densities of nanocrystalline ceramic green bodies and their sintering behavior, Nanostruct. Mater. 11 (1999) 617–622.
- [17] C. Suryanarayana, M. Grant Norton, X-ray Diffraction A Practical Approach, Plenum Press, 1998 (in New York).
- [18] L. Zhang, S. Ukai, T. Hoshino, S. Hayashi, X. H. Qu, Y₂O₃ evolution and dispersion refinement in Co-base ODS alloys, Acta Materialia. 57 (2009) 3671–3682.
- [19] K.G. Raghavendra, Arup Dasgupta, Pragna Bhaskar, K. Jayasankar, C.N. Athreya, Padmalochan Panda, S. Saroja, V. Subramanya Sarma, R. Ramaseshan, Synthesis and characterization of Fe-15 wt.% ZrO₂ nanocomposite powders by mechanical milling, Powder Technology. 287 (2016) 190–200.
- [20] L. Toualbi, M. Ratti, G. André, F. Onimus, Y. de Carlan, Use of neutron and X-ray diffraction to study the precipitation mechanisms of oxides in ODS materials, Journal of Nuclear Materials. 417 (2011) 225–228.
- [21] J.S. Benjamin, Mechanical Alloying-A Perspective, Metal Powder Rep. 45 (1990) 122-127.
- [22] A. K. Niessen, A. R. Miedema, Enthalpies of formation of liquid and solid binary alloys based on 3d metals, Physica. 151B (1988) 401-432.
- [23] A. R. Miedema, F. R. De Boer, R. Boom, Predicting heat effects in alloys, Physica. 103B (1981) 67-81.
- [24] J. M. López, J. A. Alonso, L. J. Gallego, Determination of the glass-forming concentration range in binary alloys from a semiempirical theory: Application to Zr based alloys, Phys. Rev. B. 36 (1987) 3716-3722.

- [25] A. K. Niessen, A. R. Miedema, The enthalpy of solution for solid binary alloys of two 4d-trasition metals, Phys. Chem. 7 (1983) 27-36.
- [26] A. Gabrlel, H. L. Lukas, C. H. Allibet., I. Ansara, Experimental and calculated phase diagrams of the Ni-W, Co-W and Co-Ni-W system, Z. Metallkde. 76 (1985) 589-595.
- [27] C. K. Gupta, N. Krishnamurthy, Extractive metallurgy of rare earths, CRC Press, 2005.
- [28] E. A. Brands, Smithells Metals Reference Book, Elsevier Science & Technilogy Books, London, 1983, pp. 151-153.
- [29] The Photographic Periodic Table of the Elements. http://periodictable.com/index.html, Accessed September 3,2013.
- [30] S.P. Timoshenkoo, J.N. Goodier, Theory of Elasticity, third ed., McGraw-Hill Book Company, US, 1970, pp. 410–413.
- [31] R.M. Davis, B. McDermott, C.C. Koch, Mechanical alloying of brittle materials, Metall Trans, A19 (1988) 2 867-2874.
- [32] A. Tonejc, A.M. Tonejc, D. Duzl evicÂ, Estimation of peak temperature reached by particles trapped among colliding balls in the ball-milling process using excessive oxidation of antimony, Scripta Metall Mater.25 (1991) 1 111-1113.



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