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Article

Geometric Theory of Heat from Souriau Lie Groups Thermodynamics and Koszul Hessian Geometry: Applications in Information Geometry for Exponential Families

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Abstract: We introduce the Symplectic Structure of Information Geometry based on Souriau's Lie Group Thermodynamics model, with a covariant definition of Gibbs equilibrium via invariances through co-adjoint action of a group on its moment space, defining physical observables like energy, heat, and moment as pure geometrical objects. Using Geometric (Planck) Temperature of Souriau model and Symplectic cocycle notion, the Fisher metric is identified as a Souriau Geometric Heat Capacity. Souriau model is based on affine representation of Lie Group and Lie algebra that we compare with Koszul works on G/K homogeneous space and bijective correspondence between the set of G-invariant flat connections on G/K and the set of affine representations of the Lie algebra of G. In the framework of Lie Group Thermodynamics, an Euler-Poincaré equation is elaborated with respect to thermodynamic variables, and a new variational principal for thermodynamics is built through an invariant Poincaré-Cartan-Souriau integral. The Souriau-Fisher metric is linked to KKS (Kostant-Kirillov-Souriau) 2-form that associates a canonical homogeneous symplectic manifold to the co-adjoint orbits. We apply this model in the framework of Information Geometry for the action of an affine Group for exponentiel families, and provide some illustration of use cases for multivariate Gaussian densities. Information Geometry is presented in the context of seminal work of Fréchet and his Clairait-Legendre equation. Souriau model of Statistical Physics is validated as compatible with Balian gauge model of thermodynamics. We recall the precursor work of Casalis on affine group invariance for natural exponential families.

Keywords: Lie Group Thermodynamics; Moment map; Gibbs Density; Gibbs Equilibrium; Maximum Entropy; Information Geometry; Symplectic Geometry; Cartan-Poincaré Integral Invariant; Geometric Mechanics; Euler-Poincaré Equation; Fisher Metric; Gauge Theory; Affine Group

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"Lorsque le fait qu'on rencontre est en opposition avec une théorie régnante, il faut accepter le fait et abandonner la théorie, alors même que celle-ci, soutenue par de grands noms, est généralement adoptée » **- Claude Bernard**

« *Au départ, la théorie de la stabilité structurelle m'avait paru d'une telle ampleur et d'une telle généralité, qu'avec elle je pouvais espérer en quelque sorte remplacer la thermodynamique par la géométrie, géométriser en un certain sens la thermodynamique, éliminer des considérations thermodynamiques tous les aspects à caractère mesurable et stochastiques pour ne conserver que la caractérisation géométrique correspondante des attracteurs.* » **René Thom - 1982**

Lagrange works on "*Mécanique Analytique* (Analytic Mechanics)" has been interpreted by Jean-Marie Souriau in the framework of differential geometry and has initiated a new discipline called after Souriau, "*Mécanique Géométrique* (Geometric Mechanics)" [1,2, 133]. Souriau has

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observed that the collection of motions of a dynamical system is a manifold with an antisymmetric flat tensor, that is a symplectic form where the structure contains all the pertinent information of the state of the system (positions, velocities, forces, etc.). Souriau said : "*Ce que Lagrange a vu, que n'a pas vu Laplace, c'était la structure symplectique (What Lagrange saw , that has not seen Laplace was the symplectic structure)*". Using the symmetries of a symplectic manifold, Souriau introduced a mapping which he called the "***moment map***" [90, 109, 110], which takes its values in a space attached to the group of symmetries (in the dual space of its Lie algebra). He called Dynamical Groups every dimensional group of symplectomorphisms (an isomorphism between symplectic manifolds, a transformation of phase space that is volume-preserving), and introduced Galileo Group for Classical Mechanics and Poincaré Group for Relativistic Mechanics (both are sub-groups of Affine Group [80, [159]). For instance, Galileo Group could be represented in a matrix form by (with A rotation, b the boost, c space translation and **e** time translation):

$$\begin{bmatrix} x' \\ t \\ 1 \end{bmatrix} = \underbrace{\begin{bmatrix} A & b & c \\ 0 & 1 & e \\ 0 & 0 & 1 \end{bmatrix}}_{\text{GALILEO GROUP}} \begin{bmatrix} x \\ t \\ 1 \end{bmatrix} \text{ with } \begin{cases} A \in SO(3) \\ b, c \in R^3 \\ e \in R \end{cases} \text{, Lie Algebra } \begin{bmatrix} \omega & \eta & \gamma \\ 0 & 0 & \varepsilon \\ 0 & 0 & 0 \end{bmatrix} \text{ with } \begin{cases} \omega \in so(3) \\ \eta, \gamma \in R^3 \\ \varepsilon \in R^+ \end{cases} \quad (1)$$

Souriau associated to this moment map, the notion of symplectic cohomology, linked to the fact that such a moment is defined up to an additive constant that brings into play an algebraic mechanism (called cohomology). Souriau proved that the moment map is a constant of the motion, and provided geometric generalization of Emmy Noether Invariant Theorem (invariants of E. Noether theorem are the components of the moment map). For instance, Souriau gave ontological definition of mass in classical mechanics as the measure of the symplectic cohomology of the action of the Galileo group (the mass is no longer an arbitrary variable but a characteristic of the space). This is no longer true for Poincaré Group in relativistic Mechanics, where the symplectic cohomology is null, explaining the lack of conservation of mass in relativity. All the details of classical mechanics thus appear as geometric necessities, as ontological elements. Souriau has also observed that the symplectic structure has the property to be able to be reconstructed from its symmetries alone, through a 2-form (called Kirillov-Kostant-Souriau form) defined on coadjoint orbits. Souriau said that the different versions of mechanical science can be classified by the geometry that each implies for space and time ; geometry is determined by the covariance of group theory. Thus Newtonian mechanics is covariant by the group of Galileo, the Relativity by the group of Poincaré; General Relativity by the "smooth" group (the group of diffeomorphisms of space-time). But Souriau added "*However, there are some statements of mechanics whose covariance belongs to a fourth group rarely considered: the affine group, a group shown in the following diagram for inclusion. How is it possible that a unitary point of view (which would be necessarily a* ***true Thermodynamics****) , has not yet come to crown the picture? Mystery ...*".

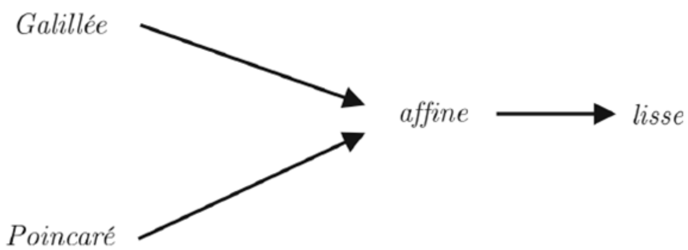

Figure 1. Souriau Scheme about mysterious "Affine Group" of a true thermodynamics between Galileo Group of Classical Mechanics, Poincaré Group of Relativistic Mechanics and Smooth Group of General Relativity.

As soon as 1966, Souriau applied his theory to Statistical Mechanics, developed it in the chapter IV of his book "Structure of Dynamical systems", and elaborated what he called a "Lie Group Thermodynamics" [172, 173, 174, 175, 176, 177, 178, 179, 180, 181, 182, 183, 184]. Using Lagrange's viewpoint, in Souriau Statistical Mechanics, a statistical state is a probability measure on the manifold of motions (and no longer in phase space [122]). Souriau observed that Gibbs equilibrium is not covariant with respect to Dynamic groups of Physics. To solve this braking of symmetry, Souriau introduced a new "Geometric Theory of Heat" where the equilibrium states are indexed by a parameter β with values in the Lie algebra of the group, generalizing the Gibbs equilibrium states, where β plays the role of a geometric (Planck) temperature. The invariance with respect to the group, and the fact that the entropy s is a convex function of this geometric temperature β, imposes very strict, universal conditions (e.g. there exist necessarily a critical temperature beyond which no equilibrium can exist). Souriau observed that the group of time translations of the classical Thermodynamics [161, 162] is not a normal subgroup of the Galilei group, proving that if a dynamical system is conservative in an inertial reference frame, it need not be conservative in another. Based on this fact, Souriau generalized the formulation of the Gibbs principle to become compatible with Galileo relativity in Classical Mechanics and with Poincaré relativity in Relativistic Mechanics. The Maximum Entropy principle [95, 96, 97, 98, 99, 100, 101, 102, 151, 196] is preserved, and the Gibbs density is given by the density of Maximum Entropy (among the equilibrium states for which the average value of the energy takes a prescribed value, the Gibbs measures are those which have the largest entropy), but with a new principle "*If a dynamical system is invariant under a Lie subgroup G' of the Galileo group, then the natural equilibria of the system forms the Gibbs ensemble of the dynamical group G'* ". The classical notion of Gibbs canonical ensemble is extended for an homogneous Symplectic Manifold on which a Lie Group (Dynamic group) has a symplectic action. When the group is not abelian (non-commutative group), the symmetry is broken, and new "cohomological" relations should be verified in Lie algebra of the group [81, 84, 85, 86]. A natural equilibrium state will thus be characterized by an element of the Lie algebra of the Lie group, determining the equilibrium temperature β. The Entropy $s(Q)$, parametrized by Q the geometric heat (mean of energy U, element of the dual Lie algebra) is defined by the Legendre transform [64, 149, 150, 154] of the Massieu Potential $\Phi(\beta)$ parametrized by β ($\Phi(\beta)$ is the minus logarithm of the partition function $\psi_\Omega(\beta)$):

$$s(Q)=\langle\beta,Q\rangle-\Phi(\beta) \quad \text{with} \quad \begin{cases} Q=\dfrac{\partial\Phi}{\partial\beta}\in \mathfrak{g}^* \\ \beta=\dfrac{\partial s}{\partial Q}\in \mathfrak{g} \end{cases} \tag{2}$$

$$p_{Gibbs}(\xi)=e^{\Phi(\beta)-\langle\beta,U(\xi)\rangle}=\frac{e^{-\langle\beta,U(\xi)\rangle}}{\int\limits_M e^{-\langle\beta,U(\xi)\rangle}d\omega}\ ,\ Q=\frac{\partial\Phi(\beta)}{\partial\beta}=\frac{\int\limits_M U(\xi)e^{-\langle\beta,U(\xi)\rangle}d\omega}{\int\limits_M e^{-\langle\beta,U(\xi)\rangle}d\omega}=\int\limits_M U(\xi)p(\xi)d\omega \tag{3}$$
$$\text{with} \quad \Phi(\beta)=-\log\int\limits_M e^{-\langle\beta,U(\xi)\rangle}d\omega$$

Souriau completed his "Geometric Heat Theory" by introducing a 2-form in the Lie algebra, that is a Riemannian metric tensor in the values of adjoint orbit of β, $[\beta,Z]$ with Z an element of the Lie algebra. This metric is given for (β,Q):

$$g_\beta([\beta,Z_1],[\beta,Z_2])=\langle\Theta(Z_1),[\beta,Z_2]\rangle+\langle Q,[Z_1,[\beta,Z_2]]\rangle \tag{4}$$

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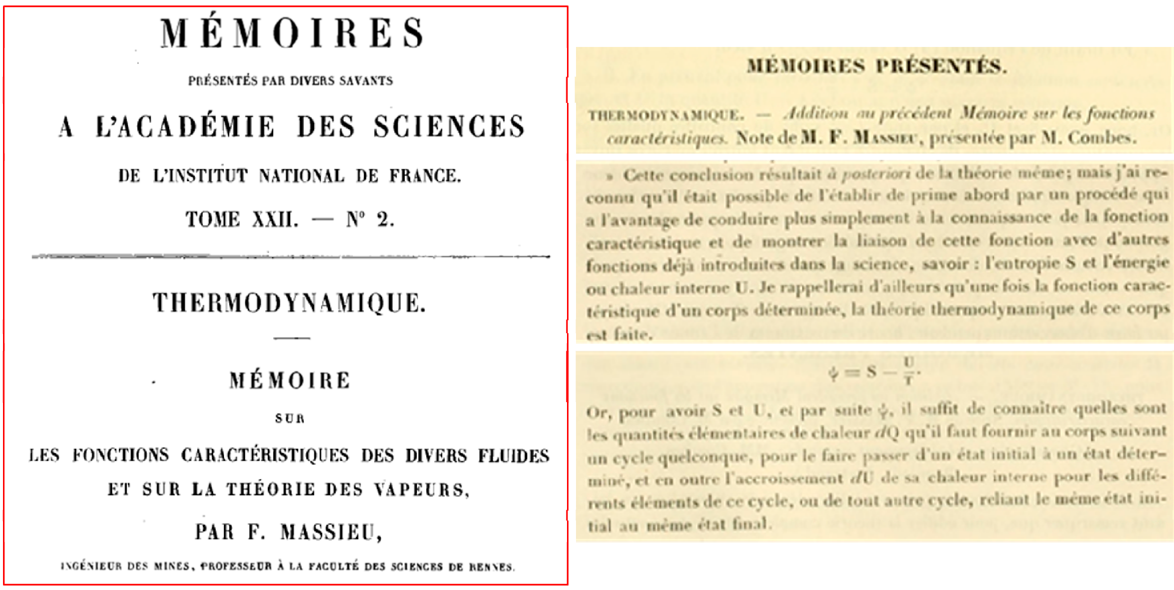
Where Θ is a cocycle of the Lie algebra, defined by $\Theta = T_e\theta$ with θ a cocycle of the Lie group defined by $\theta(M) = Q(Ad_M(\beta)) - Ad^*_M Q$. We have observed that this metric g_β is also given by the hessian of the Massieu Potential $g_\beta = -\frac{\partial^2\Phi}{\partial\beta^2} = \frac{\partial \log\psi_\Omega}{\partial\beta^2}$ as Fisher metric in classical Information Geometry theory [77], and this is a generalization of the Fisher Metric for homogeneous manifold. We call this new metric, the Souriau-Fisher metric. As $g_\beta = -\frac{\partial Q}{\partial\beta}$, Souriau compared it by analogy with classical thermodynamics to a "Geometric Specific heat" (Geometric Calorific Capacity).

The Potential theory of Thermodynamics and the introduction of "Characteristic Function" (previous function $\Phi(\beta) = -\log\psi_\Omega(\beta)$ in Souriau theory) was initiated by François Jacques Dominique Massieu [137, 138, 139, 140]. Massieu was the son of Pierre François Marie Massieu and Thérèse Claire Castel. He married in 1862 with Mlle Morand and had 2 children. Graduated from Ecole Polytechnique in 1851 and Ecole des Mines de Paris in 1956, he has integrated « Corps des Mines ». He defended his PhD in 1861 on « *Sur les intégrales algébriques des problèmes de mécanique* » and on « *Sur le mode de propagation des ondes planes et la surface de l'onde élémentaire dans les cristaux biréfringents à deux axes* » with the jury composed of Lamé, Delaunay et Puiseux. In 1870, François Massieu presented his paper to the Academy of Sciences on "*characteristic functions of the various fluids and the theory of vapors*". The design of the characteristic function is the finest scientific title of Mr. Massieu . A prominent judge , Joseph Bertrand , do not hesitate to declare , in a statement read to the Academy of Sciences July 25, 1870 , that "*the introduction of this function in formulas that summarize all the possible consequences of the two fundamental theorems seems, for the theory, a similar service almost equivalent to the Clausius has made by linking the Carnot's theorem to entropy*". The final manuscript was published by Massieu in 1873, « *Exposé des principes fondamentaux de la théorie mécanique de la chaleur (Note destinée à servir d'introduction au Mémoire de l'auteur sur les fonctions caractéristiques des divers fluides et la théorie des vapeurs)* ».


MÉMOIRES

PRÉSENTÉS PAR DIVERS SAVANTS

A L'ACADÉMIE DES SCIENCES

DE L'INSTITUT NATIONAL DE FRANCE.

TOME XXII. — N° 2.

THERMODYNAMIQUE.

MÉMOIRE

SUR

LES FONCTIONS CARACTÉRISTIQUES DES DIVERS FLUIDES

ET SUR LA THÉORIE DES VAPEURS,

PAR F. MASSIEU,

INGÉNIEUR DES MINES, PROFESSEUR A LA FACULTÉ DES SCIENCES DE RENNES.

MÉMOIRES PRÉSENTÉS.

THERMODYNAMIQUE. — *Addition au précédent Mémoire sur les fonctions caractéristiques.* Note de M. F. Massieu, présentée par M. Combes.

» Cette conclusion résultait *à posteriori* de la théorie même; mais j'ai reconnu qu'il était possible de l'établir de prime abord par un procédé qui a l'avantage de conduire plus simplement à la connaissance de la fonction caractéristique et de montrer la liaison de cette fonction avec d'autres fonctions déjà introduites dans la science, savoir : l'entropie S et l'énergie ou chaleur interne U. Je rappellerai d'ailleurs qu'une fois la fonction caractéristique d'un corps déterminée, la théorie thermodynamique de ce corps est faite.

$$\psi = S - \frac{U}{T}.$$

Or, pour avoir S et U, et par suite ψ, il suffit de connaître quelles sont les quantités élémentaires de chaleur dQ qu'il faut fournir au corps suivant un cycle quelconque, pour le faire passer d'un état initial à un état déterminé, et en outre l'accroissement dU de sa chaleur interne pour les différents éléments de ce cycle, ou de tout autre cycle, reliant le même état initial au même état final.

Figure 2. Extract from the 2nd paper of François Massieu to the French Academy of Sciences.

Massieu has introduced the following potential $\Phi(\beta)$, called "characteristic function", that is the potential used by Souriau to generalized the theory: $s(Q)=\langle\beta,Q\rangle-\Phi(\beta)\underset{\beta=\frac{1}{T}}{\Rightarrow}\Phi=\frac{Q}{T}-S$. But in his 3rd paper, Massieu was influenced by M. Bertrand to replace the variable $\beta=\frac{1}{T}$ (that he used in his two first papers) by T. We have then to wait 50 years more for the paper of Planck, who introduced again the good variable $\beta=\frac{1}{T}$, and then generalized by Souriau, giving to Planck temperature β an ontological and geometric status as element of the Lie algebra of the dynamic group. This Lie Group Thermodynamics of Souriau is able to explain astronomical phenomenon (rotation of celestial bodies: the Earth and the starts rotating about themselves). The geometric temperature β can be also interpreted as a space-time vector (generalization of the temperature vector of Planck), where the temperature vector and entropy flux are in duality unifing heat conduction and viscosity (equations of Fourier and Navier). In case of centrifuge system (e.g. used for enrichment of uranium), the Gibbs Equilibrium state [77, 78] are given by Souriau equations as the variation in concentration of the components of an inhomogeneous gas. Classical statistical mechanics corresponds to the dynamical group of time translations, for which we recover from Souriau equations the concepts and principles of dassical thermodynamics (temperature, energy, heat, work, entropy, thermodynamic potentials) and of the kinetic theory of gases (pressure, specific heats, Maxwell's velocity distribution, …).

Souriau has also studied Continuous Medium Thermodynamics, where the « Temperature Vector » is no longer constrained to be in Lie Algebra, but only contrained by phenomenologic equations (e.g. Navier equations, …). For Thermodynamic equilibrium, the « Temperature Vector » is then a Killing vector of Space-Time. For each point X, there is a « Temperature Vector » $\beta(X)$, such it is an infinitesimal conformal transform of the metric of the univers g_{ij}. Conservation equations can be then deduced for components of Impulsion-Energy tensor T^{ij} and Entropy flux S^j: $\hat{\partial}_i T^{ij}=0$ and $\partial_i S^j=0$.

$$\begin{cases}\hat{\partial}_i\beta_j+\hat{\partial}_j\beta_i=\lambda g_{ij}\\ \partial_i\beta_j+\partial_j\beta_i-2\Gamma_{ij}^k\beta_k=\lambda g_{ij}\end{cases}\quad\text{with}\quad\begin{cases}\hat{\partial}_i.:\text{covariant derivative}\\ \beta_j:\text{component of Temperature vector}\end{cases}\tag{5}$$
$$\lambda=0\Rightarrow\text{Killing Equation}$$

Before Jean-Marie Souriau, Constantin Carathéodory and Pierre Duhem [65, 66, 67, 68] initiated theoretical works to generalize Thermodynamics . The axiomatic approach of thermodynamics was published in 1909 in Mathematische Annalen [37] under the title *"Examination of the foundations of thermodynamics" [Untersuchungen überdie Grundlagen der Thermodynamik]* by Constantin Carathéodory based on Carnot works [38]. Carathéodory introduced Entropy through a mathematical approach based on the geometric behavior of a certain class of partial differential equations called Pfaffians. Carathéodory's investigations start by revisiting the first law and reformulating the second law of thermodynamics in the form of two axioms. The first axiom applies to a multiphase system change under adiabatic conditions (axiom of classical thermodynamics due to Clausius [57][61]). The second axion assumes that in the neighborhood of any equilibrium state of a system (of any number of thermodynamic coordinates), there exist states that are inaccessible by reversible adiabatic processes. In 1891, Pierre Duhem published [65] the « *On general equations of thermodynamics* » *[Sur les équations générales de la Thermodynamique]* in Annales Scientifiques de l'Ecole Normale Supérieure. Duhem writes "*We made Dynamics a special case of thermodynamics , a science that embraces common principles in all changes of state bodies , changes of places as well as changes in physical*

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qualities » [*Nous avons fait de la Dynamique un cas particulier de la Thermodynamique, une Science qui embrasse dans des principes communs tous les changements d'état des corps, aussi bien les changements de lieu que les changements de qualités physiques*]. Four scientists were credited by Duhem with having carried out "the most important researches on that subject": F. Massieu had managed to derive Thermodynamics from a "characteristic function and its partial derivatives"; J.W. Gibbs had shown that Massieu's functions "could play the role of potentials in the determination of the states of equilibrium" in a given system; H. von Helmholtz had put forward "similar ideas"; von Oettingen had given "an exposition of Thermodynamics of remarkable generality" based on general duality concept in "Die thermodynamischen Beziehungen antithetisch entwickelt" published at St. Petersburg in 1885. We have also to make reference to Henri Poincaré [121] that published the paper [155] "*Sur les tentatives d'explication mécanique des principes de la thermodynamique [On attempts of mechanical explanation forthe principles of thermodynamics]*" at the Comptes rendus de l'Académie des sciences in 1889, in which he tried to consolidate links between mechanics and thermomechanics principles. These elements were also developed in Poincaré's Lecture of 1892 [156] on "Thermodynamique" (Massieu has influenced Poincaré to introduce Massieu Characteristic function in Probability [157]). In 1906, Henri Poincaré also published a note [158] "*Reflection on The kinetic theory of gases"[Réflexions sur la théorie cinétique des gaz]*, where he said that: "*The kinetic theory of gases leaves awkward points for those who are accustomed to mathematical rigor … One of the points which embarrassed me most was the following one: it is a question of demonstrating that the entropy keeps decreasing, but the reasoning of Gibbs seems to suppose that having made vary the outside conditions we wait that the regime is established before making them vary again. Is this supposition essential, or in other words, we could arrive at opposite results to the principle of Carnot by making vary the outside conditions too fast so that the permanent regime has time to become established ?*". Leon Brillouin made the link between Boltzmann Entropy and negentropie of Information theory [27,28,29,30].

Jean-Marie Souriau has elaborated a disruptive and innovative "***théorie géométrique de la chaleur*** *(Geometric Theory of Heat)*"after the works of his predecessors: "***théorie analytique de la chaleur*** *(Analytic Theory of Heat)*" by Jean Baptiste Joseph Fourier, "***théorie mécanique de la chaleur*** *(Mechanic Theory of Heat)*" by François Clausius and François Massieu and "***théorie mathématique de la chaleur*** *(Mathematic Theory of Heat)*" by Siméon-Denis Poisson [111], as illustrated on this figure:

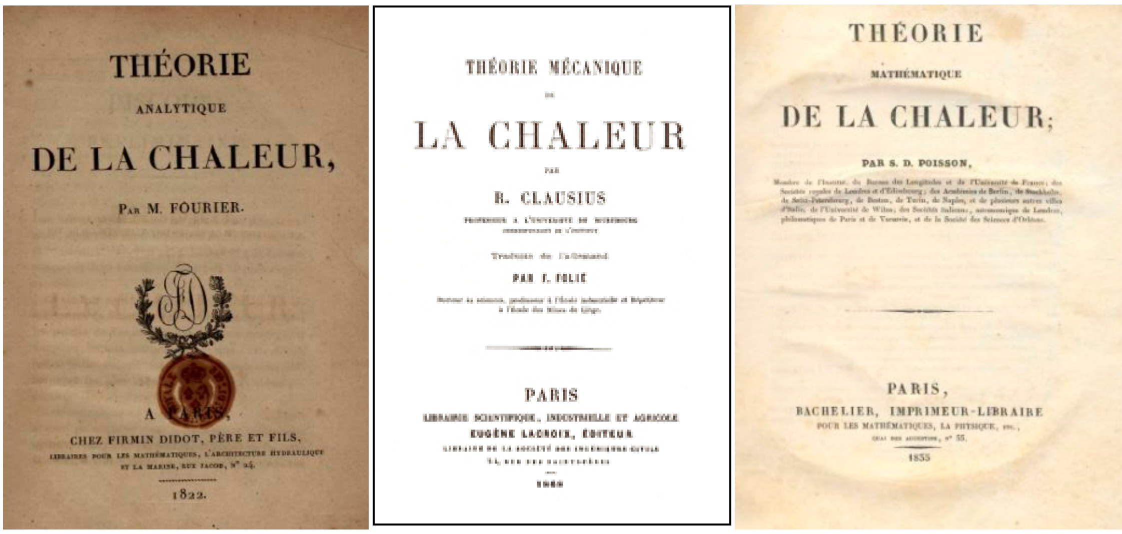

Figure 3. "théorie analytique de la chaleur (Analytic Theory of Heat)" by Jean Baptiste Joseph Fourier, "théorie mécanique de la chaleur (Mechanic Theory of Heat)" by François Clausius and "théorie mathématique de la chaleur (Mathematic Theory of Heat)" by Siméon-Denis Poisson.

Since the work of Henri Poincare and Elie Cartan, the theory of differential forms has become an essential instrument of modern differential geometry [39,40,41,42] used by Jean-Marie Souriau for identifying the space of motions as a symplectic manifold. But as said by Paulette Libermann, at the Henri Poincaré exception who wrote shortly before his death a report on the work of Elie Cartan during his application for the Sorbonne university, the French mathematicians did not see the importance of Cartan breakthroughs. Souriau followed Lectures of Elie Cartan in 1945. The 2nd student of Elie Cartan was Jean-Louis Koszul. Koszul introduced the concepts of affine spaces, affine transformations and affine representations. More especially, we are interested by Koszul definition for affine representations of Lie groups and Lie algebras. Koszul studied symmetric homogeneous spaces and defined relation between invariant flat affine connections to affine representations of Lie algebras, and characterized invariant Hessian metrics by affine representations of Lie algebras. Koszul provided correspondence between symmetric homogeneous spaces with invariant Hessian structures by using affine representations of Lie algebras, and proved that a simply connected symmetric homogeneous space with invariant Hessian structure is a direct product of a Euclidean space and a homogeneous self-dual regular convex cone. Let G be a connected Lie group and let G/K be a homogeneous space on which G acts effectively, Koszul gave a bijective correspondence between the set of G-invariant flat connections on G/K and the set of a certain class of affine representations of the Lie algebra of G. The main theorem of Koszul is that let G/K be a homogeneous space of a connected Lie group G and let g and k be the Lie algebras of G and K, assuming that G/K is endowed with a G-invariant flat connection, then g admits an affine representation (f,q) on the vector space E. Conversely, suppose that G is simply connected and that g is endowed with an affine representation, then G/K admits a G-invariant flat connection.

Koszul has proved the following. Let Ω be a convex domain in R^n containing no complete straight lines, and an associated convex cone $V(\Omega)=\left\{(\lambda x,x)\in R^n\times R/x\in\Omega,\lambda\in R^+\right\}$. Then there exists an affine embedding:

$$\ell : x\in\Omega\mapsto\begin{bmatrix}x\\1\end{bmatrix}\in V(\Omega) \tag{6}$$

If we consider η the group of homomorphism of $A(n,R)$ into $GL(n+1,R)$ given by:

$$s\in A(n,R)\mapsto\begin{bmatrix}\mathbf{f}(s) & \mathbf{q}(s)\\ 0 & 1\end{bmatrix}\in GL(n+1,R) \tag{7}$$

and associated affine representation of Lie Algebra: $\begin{bmatrix}f & q\\ 0 & 0\end{bmatrix}$ (8)

with $A(n,R)$ the group of all affine transformations of R^n. We have $\eta(G(\Omega))\subset G(V(\Omega))$ and the pair (η,ℓ) of the homomorphism $\eta : G(\Omega)\to G(V(\Omega))$ and the map $\ell:\Omega\to V(\Omega)$ is equivariant.

An Hessian structure *(D,g)* on a homogeneous space *G/K* is said to be an invariant Hessian structure if both *D* and *g* are *G*-invariant. A homogeneous space *G/K* with an invariant Hessian structure *(D,g)* is called a homogeneous Hessian manifold and is denoted by *(G/K, D, g)*. Another result of Koszul is that an homogeneous self-dual regular convex cone is characterized as a simply connected symmetric homogeneous space admitting an invariant Hessian structure that is defined by the positive definite second Koszul form (we have identified in a previous paper, that this second Koszul form is related to Fisher metric). In parallel, Vinberg [197, 198] gave a realization of a homogeneous regular convex domain as a real Siegel domain. Koszul has observed that regular convex cones admit canonical Hessian structures, improving some results of Pyateckii-Shapiro that studied realizations of homogeneous bounded domains by considering Siegel domains in connection with automorphic forms. Koszul defined a characteristic function ψ_Ω of a regular convex cone Ω, and showed that $\psi_\Omega = Dd\log\psi_\Omega$ is a Hessian metric on Ω invariant under affine automorphisms of Ω. If Ω is a homogeneous self dual cone, then the gradient mapping is a symmetry with respect to the canonical Hessian metric, and is a symmetric homogeneous

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Riemannian manifold. More information on Koszul Hessian Geometry can be found in [32,33, 141, 142, 143, 144, 145, 146, 147, 148].

In this paper, we make the link of Jean-Louis Koszul work with Souriau Model that uses an affine representations of a Lie group and of a Lie algebra in a finite-dimensional vector space, seen as special examples of actions. Souriau Model of Lie Group Thermodynamics is linked with Affine representation of Lie Group and Lie Algebra. Let G be a Lie group and E a finite-dimensional vector space. A map $A: G \to Aff(E)$ always can be written as $A(g)(x) = R(g)(x) + \theta(g)$ with $g \in G, x \in E$ where the maps $R: G \to GL(E)$ and $\theta: G \to E$ are determined by A. The map A is an affine representation of G in E. The map $\theta: G \to E$ is a one-cocycle of G with values in E, for the linear representation R; it means that θ is a smooth map which satisfies, for all $g, h \in G$: $\theta(gh) = R(g)(\theta(h)) + \theta(g)$. The linear representation R is called the linear part of the affine representation A, and θ is called the one-cocycle of G associated to the affine representation. Let $\mathfrak{g}$ be a Lie algebra and E a finite-dimensional vector space. A linear map $a: \mathfrak{g} \to aff(E)$ always can be written as $a(X)(x) = r(X)(x) + \Theta(X)$ with $X \in \mathfrak{g}, x \in E$ where the linear maps $r: \mathfrak{g} \to gl(E)$ and $\Theta: \mathfrak{g} \to E$ are determined by a. The map a is an affine representation of G in E . The linear map $\Theta: \mathfrak{g} \to E$ is a one-cocycle of G with values in E, for the linear representation r; it means that Θ satisfies, for all $X, Y \in \mathfrak{g}$: $\Theta([X,Y]) = r(X)(\Theta(Y)) - r(Y)(\Theta(X))$. $\Theta: \mathfrak{g} \to E$ is called the one-cocycle of $\mathfrak{g}$ associated to the affine representation a, and is related to the one-cocycle $\theta: G \to E$ by: $\Theta(X) = T_e\theta(X(e))$, $X \in \mathfrak{g}$.

The plan of the paper is the following. In chapter 2, we develop Souriau Symplectic Model of Statistical Physics with illustration for Multivariate Gaussian densities and some link with seminal work of Maurice Fréchet. In chapter 3, we develop Lie Group Thermodynamics model of Jean-Marie Souriau and in chaper 4, the explanation of Souriau Affine representation of Lie Group and Lie Algebra including: Affine representations and cocycles, Souriau Moment Map and Cocycles, Equivariance of Souriau Moment Map, Action of Lie Group on a Symplectic Manifold and Dual spaces of finite-dimensional Lie Algebras. In chapter 5, we identify what we call Souriau-Fisher Metric of Lie Group Thermodynamics by analogy with Fisher metric of Information geometry [82, 83] and the Souriau interpretation as geometric heat capacity. In chapter 6, we introduce new Souriau-Euler-Poincaré equations of Lie Group Thermodynamics. In chapter 7, we introduce Poincaré-Cartan Integral Invariant and Variational Principle for Souriau Lie Groups Thermodynamics. In chapter 8, we make the link of Souriau works with Koszul Affine representation of Lie Group and Lie Algebra. In chapter 9, we illustrate Koszul and Souriau Lie Group models of Information Geometry for Multivariate Gaussian laws and in chapter 10 Souriau metric for Multivariate Gaussian Densities. We give a conclusion in chapter 11 with research prospects towards affine Poisson Geometry [112]. We have 3 appendices: Appendix A develops the Clairaut(-Legendre) Equation of Maurice Fréchet associated to "distinguished functions" as fundamental equation of Information geometry; Appendix B is about Balian Gauge Model of Thermodynamics and its compliance with Souriau model; Appendix C is devoted to the link of Casalis-Letac works on Affine Group Invariance for Natural Exponential Families with Souriau works.

1. Souriau Symplectic Model of Statistical Physics

In 1970, Souriau introduced the concept of co-adjoint action of a group on its momentum space (or "*moment map*": mapping induced by symplectic manifold symmetries), based on the orbit method works, that allows to define physical observables like energy, heat and momentum or moment as pure geometrical objects (the moment map takes its values in a space determined by the group of symmetries: the dual space of its Lie algebra). The moment(um) map is a constant of the motion and is associated to symplectic cohomology (assignment of algebraic invariants to a

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topological space that arises from the algebraic dualization of the homology construction). Souriau introduced the moment map in 1965 in a lecture notes at Marseille university and published it in 1966. Souriau gave the formal definition and its name based on its physical interpretation in 1967. Souriau then studied its properties of equivariance, and formulated the coadjoint orbit theorem in his book in 1970. But in its book, Souriau also observed in chapter IV that Gibbs equilibrium states are not covariant by dynamical groups (Galileo or Poincaré groups) and then he developed a covariant model that he called "***Lie Group Thermodynamics***", where equilibriums are indexed by a "***geometric (planck) temperature***", given by a vector β that lies in the Lie algebra of the dynamical group. For Souriau, all the details of classical mechanics appear as geometric necessities (e.g., mass is the measure of the symplectic cohomology of the action of a Galileo group). Based on this new covariant model of thermodynamic Gibbs equilibrium, Souriau has formulated statistical mechanics and thermodynamics in the framework of Symplectic Geometry by use of symplectic moments and distribution-tensor concepts, giving a geometric status for temperature, heat and entropy.

There is a controversy about the name "momentum map" or "moment map". Smale referred to this map as the "angular momentum", while Souriau used the French word "moment". Cushman and Duistermaat have suggested that the proper English translation of Souriau's French word was "momentum" which fit better with standard usage in mechanics. On the other hand, Guillemin and Sternberg have validated the name given by Souriau and have used "moment" in English. In this paper, we will see that name "moment" given by Souriau was the most appropriate word. In his Chapter IV of his book, studying statistical mechanics, Souriau has geniously observed that moments of inertia in Mechanics are equivalent to moments in Probability in his new geometric model of Statistical Physics . We will see that in Souriau Lie Group Thermodynamics model, these statistical moments will be given by the Energy and the Heat defined geometrically by Souriau, and will be associated with "moment map" in dual lie algebra.

This work has been extended by Claude Vallée [192, 193] and Gery de Saxcé [163, 164, 165, 166]. More recently, M. Kapranov has also given a thermodynamical interpretation of the moment map for toric varieties [107] and Pavlov, Thermodynamics from the differential geometry standpoint [152].

The conservation of the moment of a Hamiltonian action was called by Souriau the "*Symplectic or Geometric Noether theorem*". Considering phases space as symplectic manifold, cotangent fiber of configuration space with canonical symplectic form, if Hamiltonian has Lie algebra, moment map is constant along system integral curves. Noether theorem is obtained by considering independently each component of moment map.

We will enlighten Souriau's Model with Koszul Information Geometry [113, 114, 115, 116, 117, 118, 119, 120], recently studied in [13,14,15], where we have shown that this last Geometry is founded on the notion of Koszul-Vinberg Characteristic function $\psi_\Omega(x)=\int_{\Omega^*} e^{-\langle x,\xi\rangle}d\xi,\ \forall x\in\Omega$ where Ω is a convex cone and Ω^* the dual cone with respect to Cartan-Killing inner product $\langle x,y\rangle=-B(x,\theta(y))$ invariant by automorphisms of Ω, with $B(.,.)$ the Killing form and $\theta(.)$ the Cartan involution:

$$\psi_\Omega(x+\lambda u)=\psi_\Omega(x)-\lambda\langle x^*,u\rangle+\frac{\lambda^2}{2}\langle K(x)u,u\rangle+\ldots \tag{9}$$

$$\text{with } x^*=\frac{d\Phi(x)}{dx}\ ,\Phi(x)=-\log\psi_\Omega(x)\ \text{ and }\ K(x)=\frac{d^2\Phi(x)}{dx^2} \tag{10}$$

This characteristic function is at the cornerstone of modern concept of Information Geometry, defining Koszul density by Solution of Maximum Koszul-Shannon Entropy [127]:

$$\underset{p}{Max}\left[-\int_{\Omega^*}p_{\hat{\xi}}(\xi)\log p_{\hat{\xi}}(\xi).d\xi\right]\text{such that}\int_{\Omega^*}p_{\hat{\xi}}(\xi)d\xi=1\text{ and }\int_{\Omega^*}\xi.p_{\hat{\xi}}(\xi)d\xi=\hat{\xi} \tag{11}$$

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$$p_{\hat{\xi}}(\xi)=\frac{e^{-\langle \Theta^{-1}(\hat{\xi}),\xi\rangle}}{\int\limits_{\Omega^*} e^{-\langle \Theta^{-1}(\hat{\xi}),\xi\xi\rangle}.d\xi}\quad \hat{\xi}=\Theta(\beta)=\frac{\partial\Phi(\beta)}{\partial\beta}\ \text{where}\ \Phi(\beta)=-\log\psi_{\Omega}(\beta)$$
$$\psi_{\Omega}(\beta)=\int\limits_{\Omega^*} e^{-\langle\beta,\xi\rangle}d\xi\ ,\ S(\hat{\xi})=-\int\limits_{\Omega^*} p_{\hat{\xi}}(\xi)\log p_{\hat{\xi}}(\xi).d\xi\ \text{and}\ \beta=\Theta^{-1}(\hat{\xi}) \tag{12}$$
$$S(\hat{\xi})=\left\langle \hat{\xi},\beta\right\rangle-\Phi(\beta)$$

This last relation is a Legendre transform between the logarithm of characteristic function and the Entropy:

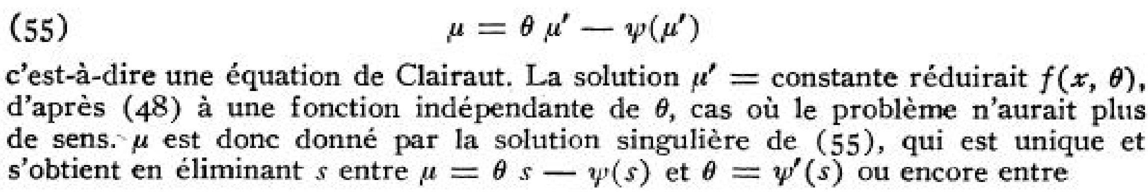
$$\log p_{\hat{\xi}}(\xi)=-\langle\xi,\beta\rangle+\Phi(\beta)$$
$$S(\hat{\xi})=-\int\limits_{\Omega^*} p_{\hat{\xi}}(\xi).\log p_{\hat{\xi}}(\xi).d\xi=-E\left[\log p_{\hat{\xi}}(\xi)\right] \tag{13}$$
$$S(\hat{\xi})=\langle E[\xi],\beta\rangle-\Phi(\beta)=\left\langle\hat{\xi},\beta\right\rangle-\Phi(\beta)$$

The inversion $\Theta^{-1}(\hat{\xi})$ is given by the Legendre transform based on the property that the Koszul-Shannon Entropy is given by the Legendre transform of minus the logarithm of the characteristic function:

$$S(\hat{\xi})=\left\langle\beta,\hat{\xi}\right\rangle-\Phi(\beta)\ \text{with}\ \Phi(\beta)=-\log\int\limits_{\Omega^*} e^{-\langle\xi,\beta\rangle}d\xi\quad\forall\beta\in\Omega\ \text{and}\ \forall\xi,\hat{\xi}\in\Omega^* \tag{14}$$

We can observe the fundamental property that $E[S(\xi)]=S(E[\xi])$, $\xi\in\Omega^*$, and also as observed by Maurice Fréchet that "distinguished functions" (densities with estimator reaching the Fréchet-Darmois bound) are solutions of the ***Alexis Clairaut Equation*** introduced by Clairaut in 1734 [74]:

$$S(\hat{\xi})=\left\langle\Theta^{-1}(\hat{\xi}),\hat{\xi}\right\rangle-\Phi\left[\Theta^{-1}(\hat{\xi})\right]\quad\forall\hat{\xi}\in\{\Theta(\beta)/\beta\in\Omega\} \tag{15}$$


(55) $\mu=\theta\,\mu'-\psi(\mu')$

c'est-à-dire une équation de Clairaut. La solution μ' = constante réduirait $f(x,\theta)$, d'après (48) à une fonction indépendante de θ, cas où le problème n'aurait plus de sens. μ est donc donné par la solution singulière de (55), qui est unique et s'obtient en éliminant s entre $\mu=\theta\,s-\psi(s)$ et $\theta=\psi'(s)$ ou encore entre

Figure 4. Clairaut-Legendre Transformation introduced by Maurice Fréchet in his 1943 paper

Details of Fréchet elaboration for this Clairaut(-Legendre) equation for "distinguished function" is given in Appendix A, and other elements are available on Fréchet's papers [73, 74, 75, 76].

In this structure, the Fisher metric $I(x)$ makes appear naturally a ***Koszul hessian geometry*** [167, 168], if we observe that

$$\log p_{\hat{\xi}}(\xi)=-\langle\xi,\beta\rangle+\Phi(\beta)$$
$$S(\hat{\xi})=-\int\limits_{\Omega^*} p_{\hat{\xi}}(\xi).\log p_{\hat{\xi}}(\xi).d\xi=-E\left[\log p_{\hat{\xi}}(\xi)\right] \tag{16}$$
$$S(\hat{\xi})=\langle E[\xi],\beta\rangle-\Phi(\beta)=\left\langle\hat{\xi},\beta\right\rangle-\Phi(\beta)$$

Then we can recover the relation with Fisher metric:

$$I(\beta)=-E\left[\frac{\partial^2\log p_\beta(\xi)}{\partial\beta^2}\right]=-E\left[\frac{\partial^2\left(-\langle\xi,\beta\rangle+\Phi(\beta)\right)}{\partial\beta^2}\right]=-\frac{\partial^2\Phi(\beta)}{\partial\beta^2}$$
$$\hat{\xi}=\frac{\partial\Phi(\beta)}{\partial\beta} \tag{17}$$
$$I(\beta)=E\left[\frac{\partial\log p_\beta(\xi)}{\partial\beta}\frac{\partial\log p_\beta(\xi)}{\partial\beta}^T\right]=E\left[\left(\xi-\hat{\xi}\right)\left(\xi-\hat{\xi}\right)^T\right]=E\left[\xi^2\right]-E[\xi]^2=Var(\xi)$$

with Crouzeix relation established in 1977 [59, 88], $\frac{\partial^2 \Phi}{\partial \beta^2} = \left[\frac{\partial^2 S}{\partial \hat{\xi}^2}\right]^{-1}$ giving the dual metric, in dual space, where Entropy S and (minus) logarithm of characteristic function, Φ , are dual potential functions.
The 1st Metric of Information Geometry [55, 56], the Fisher Metric is given by the hessian of the characteristic function logarithm:

$$I(\beta) = -E\left[\frac{\partial^2 \log p_\beta(\xi)}{\partial \beta^2}\right] = -\frac{\partial^2 \Phi(\beta)}{\partial \beta^2} = \frac{\partial^2 \log \psi_\Omega(\beta)}{\partial \beta^2} \tag{18}$$

$$ds_g^2 = d\beta^T I(\beta) d\beta = \sum_{ij} g_{ij} d\beta_i d\beta_j \quad \text{with} \quad g_{ij} = [I(\beta)]_{ij} \tag{19}$$

The 2nd Metric of Information Geometry is given by hessian of the Shannon Entropy:

$$\frac{\partial^2 S(\hat{\xi})}{\partial \hat{\xi}^2} = \left[\frac{\partial^2 \Phi(\beta)}{\partial \beta^2}\right]^{-1} \quad \text{with} \quad S(\hat{\xi}) = \langle \hat{\xi}, \beta \rangle - \Phi(\beta) \tag{20}$$

$$ds_h^2 = d\hat{\xi}^T \left[\frac{\partial^2 S(\hat{\xi})}{\partial \hat{\xi}^2}\right] d\hat{\xi} = \sum_{ij} h_{ij} d\hat{\xi}_i d\hat{\xi}_j \quad \text{with} \quad h_{ij} = \left[\frac{\partial^2 S(\hat{\xi})}{\partial \hat{\xi}^2}\right]_{ij} \tag{21}$$

Both metrics will provide the same distance:

$$ds_g^2 = ds_h^2 \tag{22}$$

This Information geometry has been intensively studied for structured matrices [20,22,23,24,25, 34, 35, 36, 53, 54, 58, 104, 105, 106, 131, 186] and in statistics [89] and is linked to seminal work of Siegel [169] on symmetric bounded domains.

We will see hereafter that Souriau has generalized this Fisher metric for Lie Group Thermodynamics, and interpreted the Fisher Metric as a Geometric Heat Capacity.

We can illustrate Information Geometry for multivariate Gaussian law [201] with computation of Fisher Metric:

$$p_{\hat{\xi}}(\xi) = \frac{1}{(2\pi)^{n/2} \det(R)^{1/2}} e^{-\frac{1}{2}(z-m)^T R^{-1}(z-m)} \tag{23}$$

If we develop:

$$\begin{aligned} \frac{1}{2}(z-m)^T R^{-1}(z-m) &= \frac{1}{2}\left[z^T R^{-1} z - m^T R^{-1} z - z^T R^{-1} m + m^T R^{-1} m\right] \\ &= \frac{1}{2} z^T R^{-1} z - m^T R^{-1} z + \frac{1}{2} m^T R^{-1} m \end{aligned} \tag{24}$$

We can write the density as a Gibbs density:

$$\begin{gathered} p_{\hat{\xi}}(\xi) = \frac{1}{(2\pi)^{n/2} \det(R)^{1/2} e^{\frac{1}{2} m^T R^{-1} m}} e^{-\left[-m^T R^{-1} z + \frac{1}{2} z^T R^{-1} z\right]} = \frac{1}{Z} e^{-\langle \xi, \beta \rangle} \\ \xi = \begin{bmatrix} z \\ zz^T \end{bmatrix} \text{ and } \beta = \begin{bmatrix} -R^{-1} m \\ \frac{1}{2} R^{-1} \end{bmatrix} = \begin{bmatrix} a \\ H \end{bmatrix} \text{ with } \langle \xi, \beta \rangle = a^T z + z^T H z = Tr\left[z a^T + H^T z z^T\right] \end{gathered} \tag{25}$$

We can then rewrite density with canonical variables:

$$\begin{gathered} p_{\hat{\xi}}(\xi) = \frac{1}{\int_{\Omega^*} e^{-\langle \xi, \beta \rangle} . d\xi} e^{-\langle \xi, \beta \rangle} = \frac{1}{Z} e^{-\langle \xi, \beta \rangle} \quad \text{with } \log(Z) = n \log(2\pi) + \frac{1}{2} \log \det(R) + \frac{1}{2} m^T R^{-1} m \\ \xi = \begin{bmatrix} z \\ zz^T \end{bmatrix}, \hat{\xi} = \begin{bmatrix} E[z] \\ E[zz^T] \end{bmatrix} = \begin{bmatrix} m \\ R + mm^T \end{bmatrix}, \beta = \begin{bmatrix} a \\ H \end{bmatrix} = \begin{bmatrix} -R^{-1} m \\ \frac{1}{2} R^{-1} \end{bmatrix} \text{ with } \langle \xi, \beta \rangle = Tr\left[z a^T + H^T z z^T\right] \\ R = E\left[(z-m)(z-m)^T\right] = E\left[zz^T - mz^T - zm^T + mm^T\right] = E\left[zz^T\right] - mm^T \end{gathered} \tag{26}$$

The 1st Potential function (Free Energy / logarithm of characteristic function) is given by:

$$\psi_\Omega(\beta) = \int_{\Omega^*} e^{-\langle \xi, \beta \rangle} . d\xi \quad \text{and} \quad \Phi(\beta) = -\log \psi_\Omega(\beta) = \frac{1}{2}\left[-Tr\left[H^{-1} a a^T\right] + \log\left[(2)^n \det H\right] - n \log(2\pi)\right] \tag{27}$$

We verify the relation between 1st Potential function and moment:

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$$\frac{\partial \Phi(\beta)}{\partial \beta} = \frac{\partial[-\log \psi_\Omega(\beta)]}{\partial \beta} = \int_{\Omega^*} \xi \frac{e^{-\langle \xi,\beta \rangle}}{\int_{\Omega^*} e^{-\langle \xi,\beta \rangle}.d\xi}.d\xi = \int_{\Omega^*} \xi.p_{\hat{\xi}}(\xi)d\xi = \hat{\xi}$$

$$\frac{\partial \Phi(\beta)}{\partial \beta} = \begin{bmatrix} \dfrac{\partial \Phi(\beta)}{\partial a} \\ \dfrac{\partial \Phi(\beta)}{\partial H} \end{bmatrix} = \begin{bmatrix} m \\ R + mm^T \end{bmatrix} = \hat{\xi} \tag{28}$$

The 2nd Potential function (Shannon Entropy) is given as Legendre Transform of 1st one:

$$S(\hat{\xi}) = \langle \hat{\xi}, \beta \rangle - \Phi(\beta) \quad \text{with} \quad \frac{\partial \Phi(\beta)}{\partial \beta} = \hat{\xi} \quad \text{and} \quad \frac{\partial S(\hat{\xi})}{\partial \hat{\xi}} = \beta$$

$$S(\hat{\xi}) = -\int_{\Omega^*} \frac{e^{-\langle \xi,\beta \rangle}}{\int_{\Omega^*} e^{-\langle \xi,\beta \rangle}.d\xi} \log \frac{e^{-\langle \xi,\beta \rangle}}{\int_{\Omega^*} e^{-\langle \xi,\beta \rangle}.d\xi}.d\xi = -\int_{\Omega^*} p_{\hat{\xi}}(\xi) \log p_{\hat{\xi}}(\xi).d\xi \tag{29}$$

$$S(\hat{\xi}) = -\int_{\Omega^*} p_{\hat{\xi}}(\xi) \log p_{\hat{\xi}}(\xi).d\xi = \frac{1}{2}\left[\log(2)^n \det\left[H^{-1}\right] + n\log(2\pi.e)\right] = \frac{1}{2}\left[\log\det[R] + n\log(2\pi.e)\right] \tag{30}$$

This remark was made by Jean-Souriau in his book as soon as 1969. He has observed tht if we take vector with tensor components $\xi = \begin{pmatrix} z \\ z \otimes z \end{pmatrix}$, components of $\hat{\xi}$ will provide moments of 1st and 2nd order of the density of probability $p_{\hat{\xi}}(\xi)$. He used this change of variable $z' = H^{1/2}z + H^{-1/2}a$, to compute the logarithm of the characteristic function $\Phi(\beta)$:

Exemple : (*loi normale*) :

Prenons le cas $V = R^n$, λ = mesure de Lebesgue, $\Psi(x) \equiv \begin{pmatrix} x \\ x \otimes x \end{pmatrix}$; un élément Z du dual de E peut se définir par la formule

$$Z(\Psi(x)) \equiv \bar{a}.x + \tfrac{1}{2}\bar{x}.H.x$$

[$a \in R^n$; H = matrice symétrique]. On vérifie que la convergence de l'intégrale I_0 a lieu si la matrice H est positive ([1]) ; dans ce cas la loi de Gibbs s'appelle *loi normale de Gauss* ; on calcule facilement I_0 en faisant le changement de variable $x^* = H^{1/2}x + H^{-1/2}a$ ([2]) ; il vient

$$\boxed{z = \tfrac{1}{2}[\bar{a}.H^{-1}.a - \log(\text{dét}(H)) + n\log(2\pi)]}$$

alors la convergence de I_1 a lieu également ; on peut donc calculer M, qui est défini par les moments du premier et du second ordre de la loi (16.196) ; le calcul montre que le moment du premier ordre est égal à $-H^{-1}.a$ et que les composantes du tenseur *variance* (16.196) sont égales aux éléments de la matrice H^{-1} ; le moment du second ordre s'en déduit immédiatement.

La formule (16.200♡) donne l'*entropie* :

$$\boxed{s = \frac{n}{2}\log(2\pi e) - \frac{1}{2}\log(\text{dét}(H))} \ ;.$$

([1]) Voir *Calcul linéaire*, tome II.
([2]) C'est-à-dire en recherchant l'*image* de la loi par l'application $x \mapsto x^*$.

Figure 5. Introduction of Potential Function for Multivariate Gaussian law

We can finally compute the metric from the matrix g_{ij}:

$$ds^2 = \sum_{ij} g_{ij} d\theta_i d\theta_j = dm^T R^{-1} dm + \frac{1}{2} Tr\left[\left(R^{-1}dR\right)^2\right] \tag{31}$$

and from classical expression of the Euler-Lagrange equation:

$$\sum_{i=1}^{n} g_{ik}\ddot{\theta}_i + \sum_{i,j=1}^{n} \Gamma_{ijk}\dot{\theta}_i\dot{\theta}_j = 0 \quad , \quad k = 1,\dots,n \quad \text{with} \quad \Gamma_{ijk} = \frac{1}{2}\left[\frac{\partial g_{jk}}{\partial \theta_i} + \frac{\partial g_{jk}}{\partial \theta_j} + \frac{\partial g_{ij}}{\partial \theta_k}\right] \tag{32}$$

That is explicitely given by:

$$\begin{cases} \ddot{R} + \dot{m}\dot{m}^T - \dot{R}R^{-1}\dot{R} = 0 \\ \ddot{m} - \dot{R}R^{-1}\dot{m} = 0 \end{cases} \tag{33}$$

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We cannot integrate this Euler-Lagrange Equation. We will see that Lie group Theory will provide new reduced equation, Euler-Poincaré equation, using Souriau theorem.

We can apply this Koszul geometry framework for cones of Symmetric Positive Definite Matrices. Let the inner product $\langle \eta, \xi \rangle = Tr(\eta^T \xi)$, $\forall \eta, \xi \in Sym(n)$ given by Cartan-Killing form, Ω be the set of symmetric positive definite matrices is an open convex cone and is self-dual $\Omega^* = \Omega$.

$$\begin{aligned} &\langle \eta, \xi \rangle = Tr(\eta^T \xi) \ , \ \forall \eta, \xi \in Sym(n) \\ &\psi_\Omega(\beta) = \int_{\Omega^*} e^{-\langle \beta, \xi \rangle} d\xi = \det(\beta)^{-\frac{n+1}{2}} \psi_\Omega(I_d) \\ &\hat{\xi} = \frac{\partial \Phi(\beta)}{\partial \beta} = \frac{\partial(-\log \psi_\Omega(\beta))}{\partial \beta} = \frac{n+1}{2} \beta^{-1} \end{aligned} \tag{34}$$

$$p_{\hat{\xi}}(\xi) = e^{-\langle \Theta^{-1}(\hat{\xi}), \xi \rangle + \Phi(\Theta^{-1}(\hat{\xi}))} = \psi_\Omega(I_d).\left[\det\left(\alpha \hat{\xi}^{-1}\right)\right] e^{-Tr\left(\alpha \hat{\xi}^{-1} \xi\right)} \quad \text{with} \quad \alpha = \frac{n+1}{2} \tag{35}$$

From the Cartan Inner Product, we can generate logarithm of the Koszul Characteristic Function, and its Legendre Transform to define Koszul Entropy, Koszul Density and Koszul Metric, as explained in the following Figure:

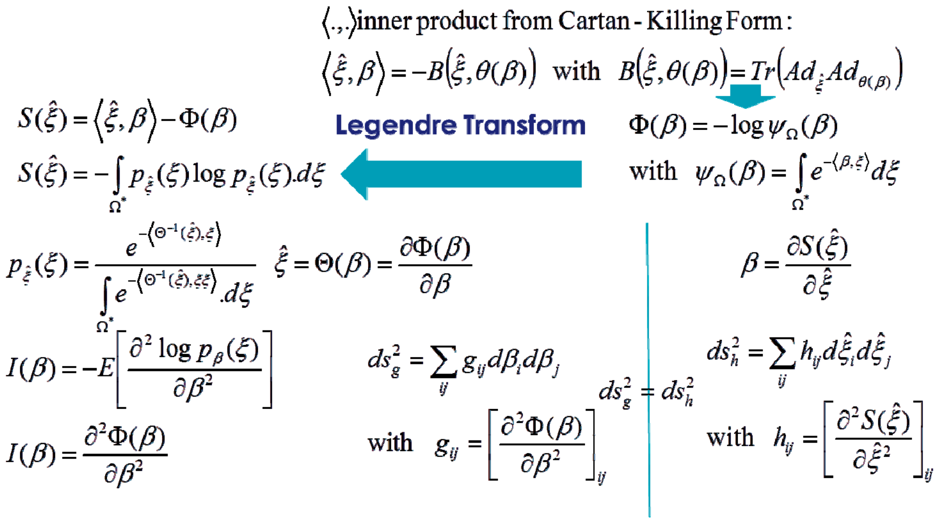

Figure 6. Generation of Koszul elements from Cartan Inner Product.

We give reference to the book of M. Deza that give a survey about distance and metric space [63].

3. Lie Group Thermodynamics model of Jean-Marie Souriau

Souriau has defined Gibbs canonical ensemble on symplectic manifold M for a Lie group action on M. In classical statistical mechanics, a state is given by the solution of Liouville equation on the phase space, the partition function. As symplectic manifolds have a completely continuous measure, invariant by diffeomorphisms, the Liouville measure λ, all statistical states will be the product of Liouville measure by the scalar function given by the generalized partition function $e^{\Phi(\beta) - \langle \beta, U(\xi) \rangle}$ defined by the energy U (defined in dual of Lie Algebra of this dynamical group) and the geometric temperature β, where Φ is a normalizing constant such the mass of probability is equal to 1, $\Phi(\beta) = -\log \int_M e^{-\langle \beta, U(\xi) \rangle} d\lambda$ [43] . Jean-Marie Souriau then generalizes the Gibbs equilibrium state to all symplectic manifolds that have a dynamical group. To ensure that all integrals, that will be defined, could converge, ***the canonical Gibbs ensemble is the largest open proper subset (in Lie algebra) where these integrals are convergent. This canonical Gibbs ensemble is***

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convex. The derivative of Φ, $Q=\frac{\partial \Phi}{\partial \beta}$ (thermodynamic heat) is equal to the mean value of the energy U. The minus derivative of this generalized heat Q, $K=-\frac{\partial Q}{\partial \beta}$ is symmetric and positive (this is a geometric heat capacity). Entropy s is then defined by Legendre transform of Φ, $s=\langle \beta, Q\rangle-\Phi$. If this approach is applied for the group of time translation, this is the classical thermodynamic theory. But ***Souriau has observed that if we apply this theory for non-commutative group (Galileo or Poincaré groups), the symmetry has been broken***. ***Classical Gibbs equilibrium states are no longer invariant by this group***. This symmetry breaking provides new equations, discovered by Souriau.

We can read in his paper this prophetical sentence "*Peut-être cette thermodynamique des groups de Lie a-t-elle un intérêt mathématique*". He explains that for dynamic Galileo group with only one axe of rotation, this thermodynamic theory is the theory of centrifuge where the temperature vector dimension is equal to 2 (sub-group of invariance of size 2), used to make "uranium 235" and "ribonucleic acid". The physical meaning of these 2 dimensions for vector-valued temperature are "thermic conduction" and "viscosity". Souriau said that the model unifies "heat conduction" and "viscosity" (Fourier and Navier equations) in the same theory of irreversible process. Souriau has applied this theory in details for relativistic ideal gas with Poincaré group for dynamical group.

Before introducing Souria Model of Lie Group Thermodynamics, we will first remind classical notation of Lie Group Theory:

- The coadjoint representation of G is the contragredient of the adjoint representation. It associates to each $g\in G$ the linear isomorphism $Ad_g^*\in GL(\mathfrak{g}^*)$, which satisfies, for each $\xi\in\mathfrak{g}^*$ and $X\in\mathfrak{g}$:
  $$\left\langle Ad^*_{g^{-1}}(\xi), X\right\rangle=\left\langle \xi, Ad_{g^{-1}}(X)\right\rangle \tag{36}$$
- The adjoint representation of the Lie algebra $\mathfrak{g}$ is the linear representation of $\mathfrak{g}$ into itself which associates, to each $X\in\mathfrak{g}$, the linear map $ad_X\in gl(\mathfrak{g})$. ad Tangent application of Ad at neutral element e of G:
  $$\begin{array}{l} ad=T_eAd:T_eG\rightarrow End(T_eG)\\ X,Y\in T_eG\mapsto ad_X(Y)=[X,Y]\end{array} \tag{37}$$
- The coadjoint representation of the Lie algebra $\mathfrak{g}$ is the contragredient of the adjoint representation. It associates, to each $X\in\mathfrak{g}$, the linear map $ad^*_X\in gl(\mathfrak{g}^*)$ which satisfies, for each $\xi\in\mathfrak{g}^*$ and $X\in\mathfrak{g}$:
  $$\left\langle ad^*_{-X}(\xi), Y\right\rangle=\left\langle \xi, Ad_{-X}(Y)\right\rangle \tag{38}$$

We can illustrate for group of matrices for $G=GL_n(K)$ with $K=R$ or C.

$$T_eG=M_n(K),\ X\in M_n(K), g\in G\quad Ad_g(X)=gXg^{-1} \tag{39}$$

$$X,Y\in M_n(K)\quad ad_X(Y)=(T_eAd)_X(Y)=XY-YX=[X,Y] \tag{40}$$

Then, the curve from $e=I_d=c(0)$ tangent to $X=c(1)$ is given by $c(t)=\exp(tX)$ and transform by Ad: $\gamma(t)=Ad\exp(tX)$

$$ad_X(Y)=(T_eAd)_X(Y)=\frac{d}{dt}\gamma(t)Y\bigg|_{t=0}=\frac{d}{dt}\exp(tX)Y\exp(tX)^{-1}\bigg|_{t=0}=XY-YX \tag{41}$$

For each temperature β, element of the Lie algebra $\mathfrak{g}$, Souriau has introduced a tensor $\tilde{\Theta}_\beta$, equal to the sum of the cocycle $\tilde{\Theta}$ and the Heat coboundary (with [.,.] Lie bracket):

$$\tilde{\Theta}_\beta(Z_1,Z_2)=\tilde{\Theta}(Z_1,Z_2)+\left\langle Q, ad_{Z_1}(Z_2)\right\rangle \text{ with } ad_{Z_1}(Z_2)=[Z_1,Z_2] \tag{42}$$

This tensor $\tilde{\Theta}_\beta$ has the following properties:

- $\widetilde{\Theta}(X,Y)=\langle \Theta(X),Y\rangle$ where the map Θ is the one-cocycle of the Lie algebra $\mathfrak{g}$ with values in $\mathfrak{g}^*$, with $\Theta(X)=T_e\theta(X(e))$ where θ the one-cocycle of the Lie group G. $\widetilde{\Theta}(X,Y)$ is constant on M and the map $\widetilde{\Theta}(X,Y):\mathfrak{g}\times\mathfrak{g}\to\Re$ is a skew-symmetric bilinear form, and is called the ***Symplectic Cocycle of Lie algebra*** $\mathfrak{g}$ associated to the ***moment map*** J, with the following properties:
  $\widetilde{\Theta}(X,Y)=J_{[X,Y]}-\{J_X,J_Y\}$ with $\{.,.\}$ Poisson Bracket and J the Moment Map (43)
  $\widetilde{\Theta}([X,Y],Z)+\widetilde{\Theta}([Y,Z],X)+\widetilde{\Theta}([Z,X],Y)=0$ (44)
  where J_X linear application from $\mathfrak{g}$ to differential function on M: $\begin{array}{l}\mathfrak{g}\to C^\infty(M,R)\\ X\to J_X\end{array}$
  and the associated differentiable application J, called moment(um) map:
  $$\begin{array}{l} J:M\to\mathfrak{g}^* \quad \text{such that } J_X(x)=\langle J(x),X\rangle,\ X\in\mathfrak{g} \\ \quad x\mapsto J(x)\end{array} \tag{45}$$
  If instead of J we take the following moment map: $J'(x)=J(x)+Q$, $x\in M$ where $Q\in\mathfrak{g}^*$ is constant, the symplectic cocycle θ is replaced by $\theta'(g)=\theta(g)+Q-Ad_g^*Q$ where $\theta'-\theta=Q-Ad_g^*Q$ is one-coboundary of G with values in $\mathfrak{g}^*$. We have also properties $\theta(g_1g_2)=Ad_{g_1}^*\theta(g_2)+\theta(g_1)$ and $\theta(e)=0$.
- $\beta\in Ker\ \widetilde{\Theta}_\beta$, such that $\widetilde{\Theta}_\beta(\beta,\beta)=0$, $\forall\beta\in\mathfrak{g}$ (46)
- The following symmetric tensor g_β, defined on all values of $ad_\beta(.)=[\beta,.]$ is positive definite:
  $g_\beta([\beta,Z_1],[\beta,Z_2])=\widetilde{\Theta}_\beta(Z_1,[\beta,Z_2])$ (47)
  $g_\beta([\beta,Z_1],Z_2)=\widetilde{\Theta}_\beta(Z_1,Z_2)$, $\forall Z_1\in\mathfrak{g},\ \forall Z_2\in \mathrm{Im}(ad_\beta(.))$ (48)
  $g_\beta(Z_1,Z_2)\geq 0$, $\forall Z_1,Z_2\in\mathrm{Im}(ad_\beta(.))$ (49)
  where the linear map $ad_X\in gl(\mathfrak{g})$ is the adjoint representation of the Lie algebra $\mathfrak{g}$ defined by $X,Y\in\mathfrak{g}(=T_eG)\mapsto ad_X(Y)=[X,Y]$, and the co-adjoint representation of the Lie algebra $\mathfrak{g}$ the linear map $ad_X^*\in gl(\mathfrak{g}^*)$ which satisfies, for each $\xi\in\mathfrak{g}^*$ and $X,Y\in\mathfrak{g}:\langle ad_X^*(\xi),Y\rangle=\langle\xi,-ad_X(Y)\rangle$

These equations are universal, because they are not dependent of the symplectic manifold but only of the dynamical group G, the symplectic cocycle Θ, the temperature β and the heat Q. Souriau called this model "***Lie Groups Thermodynamics***".

We will give the main theorem of Souriau for this "Lie Group Thermodynamics":

Theorem 1 (Souriau Theorem of Lie Group Thermodynamics). *Let Ω be the largest open proper subset of $\mathfrak{g}$, Lie algebra of G, such that $\int_M e^{-\langle\beta,U(\xi)\rangle}d\lambda$ and $\int_M \xi.e^{-\langle\beta,U(\xi)\rangle}d\lambda$ are convergent integrals, this set Ω is convex and is invariant under every transformation $Ad_g(.)$, where $g\mapsto Ad_g(.)$ is the adjoint representation of G, such that $Ad_g=T_ei_g$ with $i_g:h\mapsto ghg^{-1}$. Let $a:G\times\mathfrak{g}^*\to\mathfrak{g}^*$ a unique affine action a such that linear part is coadjoint representation of G, that is the contragradient of the adjoint representation. It associates to each $g\in G$ the linear isomorphism $Ad_g^*\in GL(\mathfrak{g}^*)$, satisfying, for each:*

$\xi\in\mathfrak{g}^*$ and $X\in\mathfrak{g}:\langle Ad_g^*(\xi),X\rangle=\langle\xi,Ad_{g^{-1}}(X)\rangle$.

Then, the fundamental equations of Lie Group Thermodynamics are:

- $\beta\to Ad_g(\beta)$ (50)
- $\Phi\to\Phi-\theta(g^{-1})\beta$ (51)
- $s\to s$ (52)
- $Q\to a(g,Q)=Ad_g^*(Q)+\theta(g)$ (53)

For Hamiltonian, actions of a Lie group on a connected symplectic manifold, the equivariance of the moment map with respect to an affine action of the group on the dual of its Lie algebra has been studied by C.M. Marle & P. Libermann [128] and Lichnerowics [129, 130]:

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Theorem 2 (Marle Theorem on cocycles). *Let G be a connected and simply connected Lie group, $R: G \to GL(E)$ be a linear representation of G in a finite-dimensional vector space E, and $r: \mathfrak{g} \to gl(E)$ be the associated linear representation of its Lie algebra $\mathfrak{g}$. For any one-cocycle $\Theta: \mathfrak{g} \to E$ of the Lie algebra $\mathfrak{g}$ for the linear representation r, there exists a unique one-cocycle $\theta: G \to E$ of the Lie group G for the linear representation R such that $\Theta(X) = T_e\theta(X(e))$, which has Θ as associated Lie algebra one-cocycle. The Lie group one-cocycle θ is a Lie group one-coboundary if and only if the Lie algebra one-cocycle Θ is a Lie algebra one-coboundary.*

Let G be a Lie group whose Lie algebra is $\mathfrak{g}$. The skew-symmetric bilinear form $\tilde{\Theta}$ on $\mathfrak{g} = T_e G$ can be extended into a closed differential two-form on G, since the identity on $\tilde{\Theta}$ means that its exterior differential $d\tilde{\Theta}$ vanishes. In other words, $\tilde{\Theta}$ is a 2-cocycle for the restriction of the de Rham cohomology of G to left invariant differential forms. In the framework of Lie Group Action on a Symplectic Manifold, equivariance of moment could be studied to prove that there is a unique action *a(.,.)* of the Lie group G on the dual $\mathfrak{g}^*$ of its Lie algebra for which the moment map J is equivariant, that means for each $x \in M$:

$$J(\Phi_g(x)) = a(g, J(x)) = Ad_g^*(J(x)) + \theta(g) \tag{54}$$

where $\Phi: G \times M \to M$ is an action of Lie Group G on differentiable manifold *M*, the fundamental field associated to an element X of Lie algebra $\mathfrak{g}$ of group *G* is the vectors field X_M on M :

$$X_M(x) = \left.\frac{d}{dt}\Phi_{\exp(-tX)}(x)\right|_{t=0} \tag{55}$$

with $\Phi_{g_1}(\Phi_{g_2}(x)) = \Phi_{g_1 g_2}(x)$ and $\Phi_e(x) = x$. Φ is hamiltonian on a Symplectic Manifold M, if Φ is symplectic and if for all $X \in \mathfrak{g}$, the fundamental field X_M is globally Hamiltonian. The cohomology class of the symplectic cocycle θ only depends on the Hamiltonian action Φ, and not on J.

In Appendix B, we observe that Souriau Lie Group Thermodynamics is compatible with Balian Gauge theory of thermodynamics [8], that is obtained by symplectization in dimension 2n+2 of contact manifold in dimension 2n+1. All elements of the Souriau geometric temperature vector are multiply by the same gauge parameter.

4. Souriau Affine representation of Lie Group and Lie Algebra

This affine representation of Lie group/algebra used by Souriau has been intensively studied by C.M. Marle [128,132, 135, 136].

4.1. Affine representations and cocycles

Souriau Model of Lie Group Thermodynamics is linked with Affine representation of Lie Group and Lie Algebra. We will give in the following main elements of this affine representation.

Let G be a Lie group and *E* a finite-dimensional vector space. A map $A: G \to Aff(E)$ always can be written as:

$$A(g)(x) = R(g)(x) + \theta(g) \text{ with } g \in G, x \in E \tag{56}$$

where the maps $R: G \to GL(E)$ and $\theta: G \to E$ are determined by *A*. The map *A* is **an affine representation of *G*** in *E*.

The map $\theta: G \to E$ is a one-cocycle of G with values in *E*, for the linear representation *R*; it means that θ is a smooth map which satisfies, for all

$$g, h \in G: \ \theta(gh) = R(g)(\theta(h)) + \theta(g) \tag{57}$$

The linear representation *R* is called the linear part of the affine representation *A*, and θ is called the one-cocycle of G associated to the affine representation *A*. *A* one-coboundary of G with values in *E*, for the linear representation *R*, is a map $\theta: G \to E$ which can be expressed as:

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$$\theta(g) = R(g)(c) - c \ , \ g \in G \text{ where } c \text{ is a fixed element in } E \tag{58}$$

and then there exist an element $c \in E$ such that, for all $g \in G$ and $x \in E$:

$$A(g)(x) = R(g)(x+c) - c \tag{59}$$

Let $\mathfrak{g}$ be a Lie algebra and E a finite-dimensional vector space. A linear map $a : \mathfrak{g} \to aff(E)$ always can be written as:

$$a(X)(x) = r(X)(x) + \Theta(X) \text{ with } X \in \mathfrak{g}, x \in E \tag{60}$$

where the linear maps $r : \mathfrak{g} \to gl(E)$ and $\Theta : \mathfrak{g} \to E$ are determined by a. The map a is an affine representation of G in E. The linear map $\Theta : \mathfrak{g} \to E$ is a one-cocycle of G with values in E, for the linear representation r; it means that Θ satisfies, for all $X, Y \in \mathfrak{g}$:

$$\Theta([X,Y]) = r(X)(\Theta(Y)) - r(Y)(\Theta(X)) \tag{61}$$

Θ is called the one-cocycle of $\mathfrak{g}$ associated to the affine representation a. A one-coboundary of $\mathfrak{g}$ with values in E, for the linear representation r, is a linear map $\Theta : \mathfrak{g} \to E$ which can be expressed as: $\Theta(X) = r(X)(c)$, $X \in \mathfrak{g}$ where c is a fixed element in E., and then there exist an element $c \in E$ such that, for all $X \in \mathfrak{g}$ and $x \in E$:

$$a(X)(x) = r(X)(x+c)$$

Let $A : G \to Aff(E)$ be an affine representation of a Lie group $\mathfrak{g}$ in a finite-dimensional vector space E, and $\mathfrak{g}$ be the Lie algebra of G. Let $R : G \to GL(E)$ and $\theta : G \to E$ be, respectively, the linear part and the associated cocycle of the affine representation A. Let $a : \mathfrak{g} \to aff(E)$ be the affine representation of the Lie algebra $\mathfrak{g}$ associated to the affine representation $A : G \to Aff(E)$ of the Lie group G. The linear part of a is the linear representation $r : \mathfrak{g} \to gl(E)$ associated to the linear representation $R : G \to GL(E)$, and the associated cocycle $\Theta : \mathfrak{g} \to E$ is related to the one-cocycle

$$\theta : G \to E \text{ by: } \Theta(X) = T_e\theta(X(e)) \ , \ X \in \mathfrak{g} \tag{62}$$

This is deduced from:

$$\left.\frac{dA(\exp(tX))(x)}{dt}\right|_{t=0} = \left.\frac{d\left(R(\exp(tX))(x) + \theta(\exp(tX))\right)}{dt}\right|_{t=0} \Rightarrow a(X)(x) = r(X)(x) + T_e\theta(X) \tag{63}$$

Let G be a connected and simply connected Lie group, $R : G \to GL(E)$ be a linear representation of G in a finite-dimensional vector space E, and $r : \mathfrak{g} \to gl(E)$ be the associated linear representation of its Lie algebra $\mathfrak{g}$. For any one-cocycle $\Theta : \mathfrak{g} \to E$ of the Lie algebra $\mathfrak{g}$ for the linear representation r, there exists a unique one-cocycle $\theta : G \to E$ of the Lie group G for the linear representation R such that:

$$\Theta(X) = T_e\theta(X(e)) \tag{64}$$

in other words, which has Θ as associated Lie algebra one-cocycle. The Lie group one-cocycle θ is a Lie group one-coboundary if and only if the Lie algrebra one-cocycle Θ is a Lie algebra one-coboundary.

$$\left.\frac{d\theta(g\exp(tX))}{dt}\right|_{t=0} = \left.\frac{d\left(\theta(g) + R(g)(\theta(\exp(tX)))\right)}{dt}\right|_{t=0} \Rightarrow T_g\theta(TL_g(X)) = R(g)(\Theta(x)) \tag{65}$$

which prove that if it exists, the Lie group one-cocycle θ such that $T_e\theta = \Theta$ is unique.

4.2. Souriau Moment Map and Cocycles

Souriau has introduced first the Moment map in his book. We will give the link with previous cocycles of affine representation.

There exist J_X linear application from $\mathfrak{g}$ to differential function on M:

$$\begin{aligned} &\mathfrak{g} \to C^\infty(M,R) \\ &X \to J_X \end{aligned} \tag{66}$$

We can then associate a differentiable application J, called moment(um) map for the Hamiltonian Lie group action Φ:

$$J : M \to \mathfrak{g}^*$$
$$x \mapsto J(x) \text{ such that } J_X(x) = \langle J(x), X \rangle, X \in \mathfrak{g} \tag{67}$$

Let J moment map, for each $(X,Y) \in \mathfrak{g} \times \mathfrak{g}$, we associate a smooth function $\tilde{\Theta}(X,Y) : M \to \Re$ defined by:

$$\tilde{\Theta}(X,Y) = J_{[X,Y]} - \{J_X, J_Y\} \quad \text{with} \quad \{.,.\}: \text{Poisson Bracket} \tag{68}$$

It is a Casimir of the Poisson algebra $C^\infty(M,\Re)$, that satisfy:

$$\tilde{\Theta}([X,Y],Z) + \tilde{\Theta}([Y,Z],X) + \tilde{\Theta}([Z,X],Y) = 0 \tag{69}$$

When the Poisson Manifold is a connected symplectic manifold, the function $\tilde{\Theta}(X,Y)$ is constant on M and the map:

$$\tilde{\Theta}(X,Y) : \mathfrak{g} \times \mathfrak{g} \to \Re \tag{70}$$

is a skew-symmetric bilinear form, and is called the Symplectic Cocycle of Lie algebra $\mathfrak{g}$ associated to the moment map J.

Let $\Theta : \mathfrak{g} \to \mathfrak{g}^*$ be the map such that for all:

$$X, Y \in \mathfrak{g} : \ \langle \Theta(X), Y \rangle = \tilde{\Theta}(X,Y) \tag{71}$$

The map Θ is therefore the one-cocycle of the Lie algebra $\mathfrak{g}$ with values in $\mathfrak{g}^*$ for the coadjoint representation $X \mapsto ad_X^*$ of $\mathfrak{g}$ associated to the affine action of $\mathfrak{g}$ on its dual:

$$a_\Theta(X)(\xi) = ad_{-X}^*(\xi) + \Theta(X) \quad , \quad X \in \mathfrak{g} \ , \xi \in \mathfrak{g}^* \tag{72}$$

Let G be a Lie group whose Lie algebra is $\mathfrak{g}$. The skew-symmetric bilinear form $\tilde{\Theta}$ on $\mathfrak{g} = T_e G$ can be extended into a closed differential two-form on G, since the identity on $\tilde{\Theta}$ means that its exterior differential $d\tilde{\Theta}$ vanishes. In other words, $\tilde{\Theta}$ is a 2-cocycle for the restriction of the de Rham cohomology of G to left (or right) invariant differential forms.

4.3. Equivariance of Souriau Moment Map

There exist a unique affine action a such that linear part is coadjoint representation:

$$a : G \times \mathfrak{g}^* \to \mathfrak{g}^*$$
$$a(g,\xi) = Ad_{g^{-1}}^* \xi + \theta(g) \tag{73}$$

with $\left\langle Ad_{g^{-1}}^* \xi, X \right\rangle = \left\langle \xi, Ad_{g^{-1}} X \right\rangle$ and that induce equivariance of moment J.

4.4. Action of Lie Group on a Symplectic Manifold

Let $\Phi : G \times M \to M$ be an action of Lie Group G on differentiable manifold M, the fundamental field associated to an element X of Lie algebra $\mathfrak{g}$ of group G is the vectors field X_M on M:

$$X_M(x) = \left. \frac{d}{dt} \Phi_{\exp(-tX)}(x) \right|_{t=0} \quad \text{With} \quad \Phi_{g_1}(\Phi_{g_2}(x)) = \Phi_{g_1 g_2}(x) \quad \text{and} \quad \Phi_e(x) = x \tag{74}$$

Φ is hamiltonian on a Symplectic Manifold M, if Φ is symplectic and if for all $X \in \mathfrak{g}$, the fundamental field X_M is globally Hamiltonian.

There is a unique action a of the Lie group G on the dual $\mathfrak{g}^*$ of its Lie algebra for which the moment map J is equivariant, that means satisfies for each $x \in M$

$$J(\Phi_g(x)) = a(g, J(x)) = Ad_{g^{-1}}^*(J(x)) + \theta(g) \tag{75}$$

$\theta : G \to \mathfrak{g}^*$ is called Cocycle associated to the differential $T_e\theta$ of 1-cocyle θ associated to J at neutral element e:

$$\langle T_e\theta(X), Y \rangle = \tilde{\Theta}(X,Y) = J_{[X,Y]} - \{J_X, J_Y\} \tag{76}$$

If instead of J we take the moment map $J'(x) = J(x) + \mu$, $x \in M$, where $\mu \in \mathfrak{g}^*$ is constant, the symplectic cocycle θ is replaced by:

$$\theta'(g) = \theta(g) + \mu - Ad_g^* \mu \tag{77}$$

where $\theta' - \theta = \mu - Ad_g^* \mu$ is one-coboundary of G with values in $\mathfrak{g}^*$.

Therefore the cohomology class of the symplectic cocycle θ only depends on the Hamiltonian action Φ, not on the choice of its moment map J. We have also:

$$\tilde{\Theta}'(X,Y) = \tilde{\Theta}(X,Y) + \langle \mu, [X,Y] \rangle \qquad (78)$$

This property is used by Jean-Marie Souriau to offer a very nice cohomological interpretation of the total mass of a classical (nonrelativistic) isolated mechanical system. He proves that the space of all possible motions of the system is a symplectic manifold on which the Galilean group acts by a Hamiltonian action. The dimension of the symplectic cohomology space of the Galilean group (the quotient of the space of symplectic one-cocycles by the space of symplectic one-coboundaries) is equal to 1. The cohomology class of the symplectic cocycle associated to a moment map of the action of the Galilean group on the space of motions of the system is interpreted as the total mass of the system.

For Hamiltonian, actions of a Lie group on a connected symplectic manifold, the equivariance of the moment map with respect to an affine action of the group on the dual of its Lie algebra has been proved by C.M. Marle. C.M. Marle has also developed the notion of symplectic cocycle and has proved that given a Lie algebra symplectic cocycle, there exists on the associated connected and simply connected Lie group a unique corresponding Lie group symplectic cocycle. C.M. Marle has also proved that there exists a two-parameter family of deformations of these actions (the Hamiltonian actions of a Lie group on its cotangent bundle obtained by lifting the actions of the group on itself by translations) into a pair of mutually symplectically orthogonal Hamiltonian actions whose moment maps are equivariant with respect to an affine action involving any given Lie group symplectic cocycle. C.M. Marle has also explained why a reduction occurs for Euler-Poncaré equation mainly when the Hamiltonian can be expressed as the moment map composed with a smooth function defined on the dual of the Lie algebra; the Euler-Poincaré equation is then equivalent to the Hamilton equation written on the dual of the Lie algebra.

4.5. Dual spaces of finite-dimensional Lie Algebras

Dual spaces of finite-dimensional Lie algebras. Let $\mathfrak{g}$ be a finite-dimensional Lie algebra, and $\mathfrak{g}^*$ its dual space. The Lie algebra $\mathfrak{g}$ can be considered as the dual of $\mathfrak{g}^*$, that means as the space of linear functions on $\mathfrak{g}^*$, and the bracket of the Lie algebra $\mathfrak{g}$ is a composition law on this space of linear functions. This composition law can be extended to the space $C^\infty(\mathfrak{g}^*,\Re)$ by setting:

$$\{f,g\}(x) = \langle x, [df(x), dg(x)] \rangle \ , \ f \text{ and } g \in C^\infty(\mathfrak{g}^*,\Re),\ x \in \mathfrak{g}^* \qquad (79)$$

If we apply this formula for Souriau Lie Group Thermodynamics, and for Entropy s(Q) depending of Geometric heat Q:

$$\{s_1,s_2\}(Q) = \langle Q, [ds_1(Q), ds_2(Q)] \rangle \ , \ s_1 \text{ and } s_2 \in C^\infty(\mathfrak{g}^*,\Re),\ Q \in \mathfrak{g}^* \qquad (80)$$

This bracket on $C^\infty(\mathfrak{g}^*,\Re)$ defines a Poisson structure on $\mathfrak{g}^*$, called its canonical Poisson structure. It implicitly appears in the works of Sophus Lie, and was rediscovered by Alexander Kirillov [108], Bertram Kostant and Jean-Marie Souriau.

The above defined canonical Poisson structure on $\mathfrak{g}^*$ can be modified by means of a symplectic cocycle $\tilde{\Theta}$ by defining the new bracket:

$$\{f,g\}_{\tilde{\Theta}}(x) = \langle x, [df(x), dg(x)] \rangle - \tilde{\Theta}(df(x), dg(x)) \qquad (81)$$

with $\tilde{\Theta}$ a symplectic cocycle of the Lie algebra $\mathfrak{g}$ is a skew-symmetric bilinear map $\tilde{\Theta} : \mathfrak{g} \times \mathfrak{g} \to \Re$ which satisfies:

$$\tilde{\Theta}([X,Y],Z) + \tilde{\Theta}([Y,Z],X) + \tilde{\Theta}([Z,X],Y) = 0 \qquad (82)$$

This Poisson structure is called the modified canonical Poisson structure by means of the symplectic cocycle $\tilde{\Theta}$. The symplectic leaves of $\mathfrak{g}^*$ equipped with this Poisson structure are the orbits of an affine action whose linear part is the coadjoint action, with an additional term determined by $\tilde{\Theta}$.

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5. Souriau-Fisher Metric of Lie Group Thermodynamics

If we differentiate this relation of Souriau theorem $Q(Ad_g(\beta))=Ad_g^*(Q)+\theta(g)$, this relation occurs:

$$\frac{\partial Q}{\partial \beta}(-[Z_1,\beta],.)=\tilde{\Theta}(Z_1,[\beta,.])+\langle Q, Ad_{.Z_1}([\beta,.])\rangle=\tilde{\Theta}_\beta(Z_1,[\beta,.]) \quad (83)$$

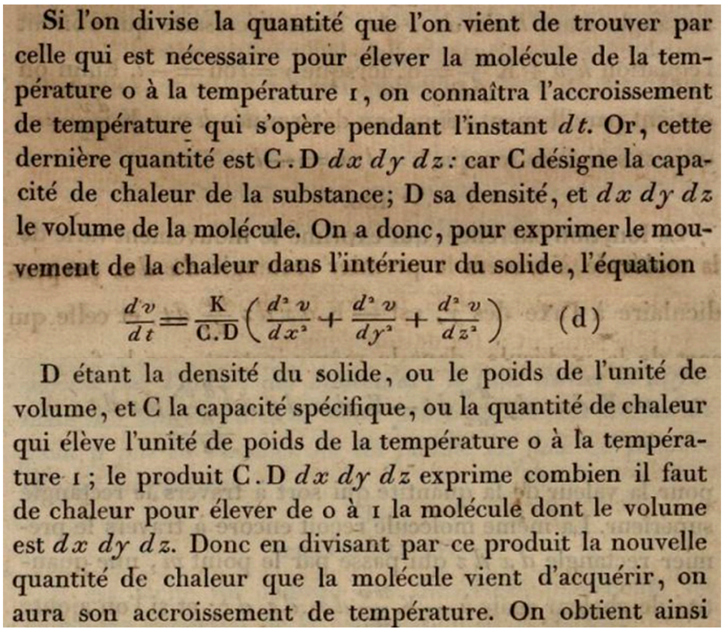
$$-\frac{\partial Q}{\partial \beta}([Z_1,\beta],Z_2.)=\tilde{\Theta}(Z_1,[\beta,Z_2])+\langle Q, Ad_{.Z_1}([\beta,Z_2])\rangle=\tilde{\Theta}_\beta(Z_1,[\beta,Z_2]) \quad (84)$$

$$\Rightarrow -\frac{\partial Q}{\partial \beta}=g_\beta([\beta,Z_1],[\beta,Z_2]) \quad (85)$$

We observe that the Fisher Metric $I(\beta)=-\frac{\partial Q}{\partial \beta}$ is exactly the Souriau Metric defined through Symplectic cocycle:

$$I(\beta)=\tilde{\Theta}_\beta(Z_1,[\beta,Z_2])=g_\beta([\beta,Z_1],[\beta,Z_2]) \quad (86)$$

The Fisher Metric $I(\beta)=-\frac{\partial^2\Phi(\beta)}{\partial\beta^2}=-\frac{\partial Q}{\partial \beta}$ has been considered by Souriau as **a *generalization of "Heat Capacity"***. Souriau called it K the **"*Geometric Capacity*"**.



Si l'on divise la quantité que l'on vient de trouver par celle qui est nécessaire pour élever la molécule de la température o à la température 1, on connaîtra l'accroissement de température qui s'opère pendant l'instant dt. Or, cette dernière quantité est C.D $dx\,dy\,dz$: car C désigne la capacité de chaleur de la substance; D sa densité, et $dx\,dy\,dz$ le volume de la molécule. On a donc, pour exprimer le mouvement de la chaleur dans l'intérieur du solide, l'équation

$$\frac{dv}{dt}=\frac{K}{C.D}\left(\frac{d^2v}{dx^2}+\frac{d^2v}{dy^2}+\frac{d^2v}{dz^2}\right) \quad (d)$$

D étant la densité du solide, ou le poids de l'unité de volume, et C la capacité spécifique, ou la quantité de chaleur qui élève l'unité de poids de la température o à la température 1; le produit C.D $dx\,dy\,dz$ exprime combien il faut de chaleur pour élever de o à 1 la molécule dont le volume est $dx\,dy\,dz$. Donc en divisant par ce produit la nouvelle quantité de chaleur que la molécule vient d'acquérir, on aura son accroissement de température. On obtient ainsi

Figure 7. Fourier heat equation in seminal manuscript of Joseph Fourier

For $\beta=\frac{1}{kT}$, $K=-\frac{\partial Q}{\partial \beta}=-\frac{\partial Q}{\partial T}\left(\frac{\partial(1/kT)}{\partial T}\right)^{-1}=kT^2\frac{\partial Q}{\partial T}$ linking the geometric capacity to calorific capacity, then Fisher metric can be introduced in Fourier heat equation:

$$\frac{\partial T}{\partial t}=\frac{\kappa}{C.D}\Delta T \text{ with } \frac{\partial Q}{\partial T}=C.D \Rightarrow \frac{\partial \beta^{-1}}{\partial t}=\kappa\left[(\beta^2/k)I_{Fisher}(\beta)\right]^{-1}\Delta\beta^{-1} \quad (87)$$

Souriau has built a thermometer (θερμός) device principle that could measure the Geometric Temperature using "Relative Ideal Gas Thermometer" based on a theory of Dynamical Group Thermometry and has also recovered the (Geometric) Laplace barometric law

We can also observe that Q is related to the mean, and K to the variance of U:

$$K=I(\beta)=-\frac{\partial Q}{\partial \beta}=\mathrm{var}(U)=\int_M U(\xi)^2.p_\beta(\xi)d\omega-\left(\int_M U(\xi).p_\beta(\xi)d\omega\right)^2 \quad (88)$$

We observe that the entropy s is unchanged, and Φ is changed but with linear dependence to β, with consequence that Fisher Souriau metric is invariant:

$$s\left[Q\left(Ad_g(\beta)\right)\right]=s(Q(\beta)) \quad \text{and} \quad I\left(Ad_g(\beta)\right)=\frac{\partial^2\left(\Phi-\theta\left(g^{-1}\right)\beta\right)}{\partial\beta^2}=\frac{\partial^2\Phi}{\partial\beta^2}=I(\beta) \tag{89}$$

General definition of Heat Capacity has also been introduced by Pierre Duhem .

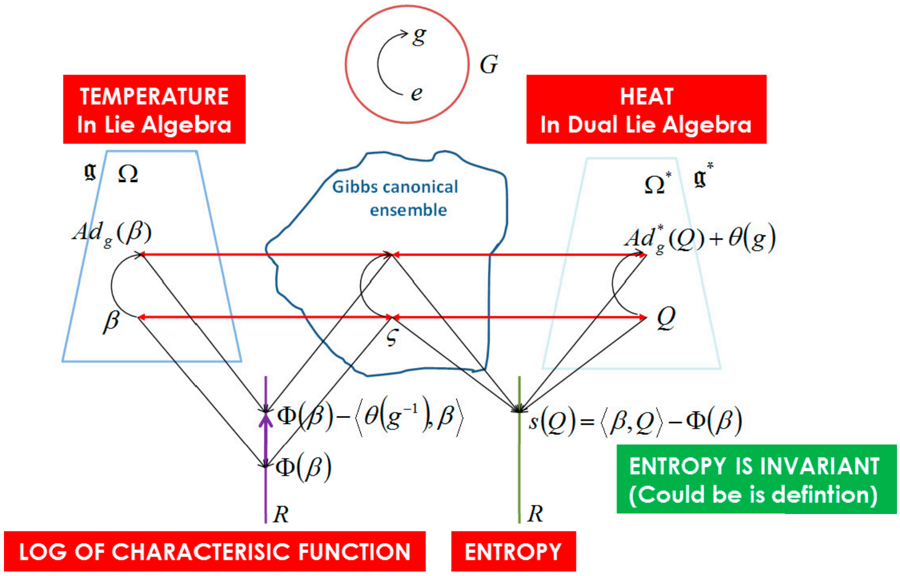


Figure 8. Global Souriau scheme of Lie Group Thermodynamics

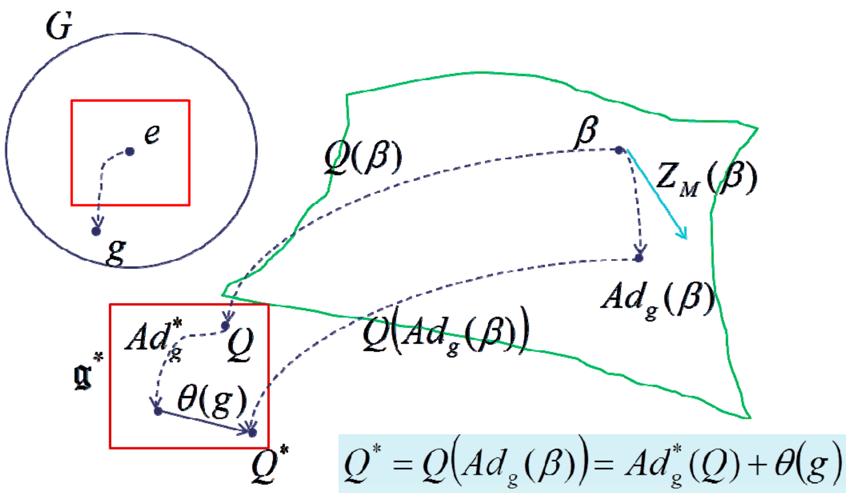


Figure 9. broken of symmetric on geometric heat Q due to adjoint action of the Group on temperature β as element of the Lie algebra

We have deduced from this Souriau Model, by reduction, the Euler-Poincaré equation describing geodesic:

$$\frac{dQ}{dt}=ad_\beta^* Q \quad \text{and} \quad \begin{cases} s(Q)=\langle\beta,Q\rangle-\Phi(\beta) \\ \beta=\dfrac{\partial s(Q)}{\partial Q}\in\mathfrak{g},\ Q=\dfrac{\partial\Phi(\beta)}{\partial\beta}\in\mathfrak{g}^* \end{cases} \tag{90}$$

Back to Koszul model of Information Geometry, we can then deduce Euler-Poincaré equation for statistical models

$$\frac{dx^*}{dt} = ad_x^* x^* \quad \text{and} \quad \begin{cases} \Phi^*(x^*) = \langle x, x^* \rangle - \Phi(x) \\ x = \dfrac{\partial \Phi^*(x^*)}{\partial x} \in \Omega, \; x^* = \dfrac{\partial \Phi(x)}{\partial x} \in \Omega^* \end{cases} \tag{91}$$

6. Souriau-Euler-Poincaré equations of Lie Group Thermodynamics

When a Lie algebra acts locally transitively on the configuration space of a Lagrangian mechanical system, Henri Poincaré proved that the Euler-Lagrange equations are equivalent to a new system of differential equations defined on the product of the configuration space with the Lie algebra. C.M. Marle has written the Euler-Poincaré equations [134], under an intrinsic form, without any reference to a particular system of local coordinates, proving that they can be conveniently expressed in terms of the Legendre and moment maps of the lift to the cotangent bundle of the Lie algebra action on the configuration space. The Lagrangian is a smooth real valued function L defined on the tangent bundle TM. To each parameterized continuous, piecewise smooth curve $\gamma : [t_0, t_1] \to M$, defined on a closed interval $[t_0, t_1]$, with values in M, one associates the value at γ of the action integral: $I(\gamma) = \int_{t_0}^{t_1} L\left(\frac{d\gamma(t)}{dt} \right) dt$ (92)

The partial differential of the function $L : M \times \mathfrak{g} \to \Re$ with respect to its second variable $d_2 \overline{L}$, which plays an important part in the Euler-Poincaré equation, can be expressed in terms of the moment and Legendre maps: $d_2 \overline{L} = p_{\mathfrak{g}^*} \circ \varphi^t \circ \mathrm{L} \circ \varphi$ with $J = p_{\mathfrak{g}^*} \circ \varphi^t (\Rightarrow d_2 \overline{L} = J \circ \mathrm{L} \circ \varphi)$ the moment map, $p_{\mathfrak{g}^*} : M \times \mathfrak{g}^* \to \mathfrak{g}^*$ the canonical projection on the second factor, $\mathrm{L} : TM \to T^* M$ the Legendre transform, with:

$$\varphi : M \times \mathfrak{g} \to TM \;/\; \varphi(x, X) = X_M(x) \quad \text{and} \quad \varphi^t : T^* M \to M \times \mathfrak{g}^* \;/\; \varphi^t(\xi) = (\pi_M(\xi), J(\xi)) \tag{93}$$

The Euler-Poincaré equation can therefore be written under the form:

$$\left(\frac{d}{dt} - ad_{V(t)}^* \right) (J \circ \mathrm{L} \circ \varphi(\gamma(t), V(t))) = J \circ d_1 \overline{L}(\gamma(t), V(t)) \quad \text{with} \quad \frac{d\gamma(t)}{dt} = \varphi(\gamma(t), V(t)) \tag{94}$$

With $H(\xi) = \langle \xi, \mathrm{L}^{-1}(\xi) \rangle - L(\mathrm{L}^{-1}(\xi))$, $\xi \in T^* M$, $\mathrm{L} : TM \to T^* M$, $H : T^* M \to R$ (95)

Following the remark made by Poincaré at the end of his note, the most interesting case is when the map $\overline{L} : M \times \mathfrak{g} \to R$ only depends on its second variable $X \in \mathfrak{g}$. The Euler-Poincaré equation becomes:

$$\left(\frac{d}{dt} - ad_{V(t)}^* \right) (d\overline{L}(V(t))) = 0 \tag{96}$$

We can use analogy of structure when the convex Gibbs ensemble is homogeneous [185]. We can then apply Euler-Poincaré equation for Lie Group Thermodynamics. Considering Clairaut equation: $s(Q) = \langle \beta, Q \rangle - \Phi(\beta) = \langle \Theta^{-1}(Q), Q \rangle - \Phi(\Theta^{-1}(Q))$ (97)

with $Q = \Theta(\beta) = \frac{\partial \Phi}{\partial \beta} \in \mathfrak{g}^*$, $\beta = \Theta^{-1}(Q) \in \mathfrak{g}$, a ***Souriau-Euler-Poincaré equation*** can be elaborated for Souriau Lie Group Thermodynamics:

$$\frac{dQ}{dt} = ad_\beta^* Q \tag{98}$$

or

$$\frac{d}{dt}\left(Ad_g^* Q \right) = 0 \tag{99}$$

An associated equation on Entropy is: $\frac{ds}{dt} = \left\langle \frac{d\beta}{dt}, Q \right\rangle + \langle \beta, ad_\beta^* Q \rangle - \frac{d\Phi}{dt}$ that reduces to

$$\frac{ds}{dt} = \left\langle \frac{d\beta}{dt}, Q \right\rangle - \frac{d\Phi}{dt} \tag{100}$$

due to $\langle \xi, ad_V X \rangle = -\langle ad_V^* \xi, X \rangle \Rightarrow \langle \beta, ad_\beta^* Q \rangle = \langle Q, ad_\beta \beta \rangle = 0$.

With these new equation of thermodynamics $\frac{dQ}{dt} = ad_{\beta}^{*}Q$ and $\frac{d}{dt}\left(Ad_{g}^{*}Q\right) = 0$, we can observe that the new important notion is related to co-adjoint orbits, that are associated to a Symplectic manifold by Souriau or KKS 2-form.

7. Poincaré-Cartan Integral Invariant and Variational Principle of Souriau Lie Groups Thermodynamics

We will define the Poincaré-Cartan Integral Invariant for Lie Group Thermodynamics. Classically in mechanics, the Pfaffian form $\omega = p.dq - H.dt$ is related to Poincaré-Cartan integral invariant [26]. P. Dedecker has observed, based on the relation $\omega = \partial_{\dot{q}}L.dq - \left(\partial_{\dot{q}}L.\dot{q} - L\right).dt = L.dt + \partial_{\dot{q}}L\varpi$ with $\varpi = dq - \dot{q}.dt$, that the property that among all forms $\chi \equiv L.dt \bmod \varpi$ the form $\omega = p.dq - H.dt$ is the only one satisfying $d\chi \equiv 0 \bmod \varpi$, is a particular case of more general T. Lepage congruence.

Analogies between Geometric Mechanics & Geometric Lie Group Thermodynamics, provides the following similarities of structures:

$$\begin{cases} \dot{q} \leftrightarrow \beta \\ p \leftrightarrow Q \end{cases} , \quad \begin{cases} L(\dot{q}) \leftrightarrow \Phi(\beta) \\ H(p) \leftrightarrow s(Q) \\ H = p.\dot{q} - L \leftrightarrow s = \langle Q, \beta \rangle - \Phi \end{cases} \quad \text{and} \quad \begin{cases} \dot{q} = \frac{dq}{dt} = \frac{\partial H}{\partial p} \leftrightarrow \beta = \frac{\partial s}{\partial Q} \\ p = \frac{\partial L}{\partial \dot{q}} \leftrightarrow Q = \frac{\partial \Phi}{\partial \beta} \end{cases} \tag{101}$$

We can then consider a similar ***Poincaré-Cartan-Souriau Pfaffian form***:

$$\omega = p.dq - H.dt \leftrightarrow \omega = \langle Q, (\beta.dt) \rangle - s.dt = \left(\langle Q, \beta \rangle - s\right).dt = \Phi(\beta).dt \tag{102}$$

This analogy provides an associated ***Poincaré-Cartan-Souriau Integral Invariant***:

$$\int_{C_a} p.dq - H.dt = \int_{C_b} p.dq - H.dt \quad \text{is transformed in} \quad \int_{C_a} \Phi(\beta).dt = \int_{C_b} \Phi(\beta).dt \tag{103}$$

We can then deduce an ***Euler-Poincaré-Souriau Variational Principle*** for Thermodynamics: *The Variational Principle holds on* $\mathfrak{g}$, *for variations* $\delta\beta = \dot{\eta} + [\beta, \eta]$, *where* $\eta(t)$ *is an arbitrary path that vanishes at the endpoints,* $\eta(a) = \eta(b) = 0$:

$$\delta \int_{t_0}^{t_1} \Phi(\beta(t)).dt = 0 \tag{104}$$

8. Koszul Affine representation of Lie Group and Lie Algebra

Previously, we have developed Souriau works on affine representation of Lie group used to elaborate the Lie Group Thermodynamics. We will study here some extension of affine representation of Lie group and Lie algebra given by Jean-Louis Koszul.

Koszul has proved that on a complex homogeneous space , an invariant volume defines with the complex structure, an invariant Hermitian form. If this space is a bounded domain, then this hermitian form is positive definite and coincides with the classical Bergman metric of this domain. During his stay at Institute for Advanced Study in Princeton, Koszul has also demonstrated the reciprocal for a class of complex homogeneous spaces, defined by open orbits of complex affine transformation groups.

Let G a connex Lie Group and E a real or complex vector space of finite dimension, Koszul has introduced an affine representation of G in E such that:

$$\begin{array}{l} E \to E \\ a \mapsto sa \quad \forall s \in G \end{array} \tag{105}$$

is an affine transformation. We set $A(E)$ the set of all affine transformations of a vector space E, a Lie Group called affine transformation group of E. The set $GL(E)$ of all regular linear transformations of E, a subgroup of $A(E)$.

We define a linear representation from G to $GL(E)$:

$$\begin{aligned}&\mathbf{f}: G \to GL(E)\\ &\quad s \mapsto \mathbf{f}(s)a = sa - so \quad \forall a \in E\end{aligned} \tag{106}$$

and an application from G to E:

$$\begin{aligned}&\mathbf{q}: G \to E\\ &\quad s \mapsto \mathbf{q}(s) = so \quad \forall s \in G\end{aligned} \tag{107}$$

Then we have $\forall s,t \in G$:

$$\mathbf{f}(s)\mathbf{q}(t) + \mathbf{q}(s) = \mathbf{q}(st) \tag{108}$$

deduced from $\mathbf{f}(s)\mathbf{q}(t) + \mathbf{q}(s) = s\mathbf{q}(t) - so + so = s\mathbf{q}(t) = sto = \mathbf{q}(st)$.

On the contrary, if an application $\mathbf{q}$ from G to E and a linear representation $\mathbf{f}$ from G to $GL(E)$ verify previous equation, then we can define an affine representation of G in E, written $(\mathbf{f},\mathbf{q})$:

$$\mathrm{Aff}(s): a \mapsto sa = \mathbf{f}(s)a + \mathbf{q}(s) \quad \forall s \in G, \forall a \in E \tag{109}$$

The condition $\mathbf{f}(s)\mathbf{q}(t) + \mathbf{q}(s) = \mathbf{q}(st)$ is equivalent to requiring the following mapping to be an homomorphism:

$$\mathrm{Aff}: s \in G \mapsto \mathrm{Aff}(s) \in A(E) \tag{110}$$

We write f the linear representation of Lie algebra $\mathfrak{g}$ of G, defined by $\mathbf{f}$ and q the restriction to $\mathfrak{g}$ of the differential to $\mathbf{q}$ (f and q the differential of $\mathbf{f}$ and $\mathbf{q}$ respectively), Koszul has proved that:

$$\begin{aligned}&f(X)q(Y) - f(Y)q(X) = q([X,Y]) \quad \forall X,Y \in \mathfrak{g}\\ &\text{with } f:\mathfrak{g} \to gl(E) \text{ and } q:\mathfrak{g} \mapsto E\end{aligned} \tag{111}$$

where $gl(E)$ the set of all linear endomorphisms of E, the Lie algebra of $GL(E)$.

Using the computation,

$$q(Ad_s Y) = \left.\frac{d\mathbf{q}(s.e^{tY}.s^{-1})}{dt}\right|_{t=0} = \mathbf{f}(s)f(Y)\mathbf{q}(s^{-1}) + \mathbf{f}(s)q(Y) \tag{112}$$

We can obtain:

$$q([X,Y]) = \left.\frac{d\mathbf{q}(Ad_{e^{tX}}Y)}{dt}\right|_{t=0} = f(X)q(Y)\mathbf{q}(e) + \mathbf{f}(e)f(Y)(-q(X)) + f(X)q(Y) \tag{113}$$

where e is the unit element in G. Since $\mathbf{f}(e)$ is the identity mapping and $\mathbf{q}(e) = 0$, we have the equality: $f(X)q(Y) - f(Y)q(X) = q([X,Y])$.

A pair (f,q) of a linear representation f of a Lie algebra $\mathfrak{g}$ on E and a linear mapping q from $\mathfrak{g}$ to E is an affine representation of $\mathfrak{g}$ on E, if it satisfies $f(X)q(Y) - f(Y)q(X) = q([X,Y])$.

Conversely, if we assume that $\mathfrak{g}$ admits an affine representation (f,q) on E, using an affine coordinate system $\{x^1,...,x^n\}$ on E, we can express an affine mapping $v \mapsto f(X)v + q(Y)$ by an $(n+1)\times(n+1)$ matrix representation:

$$aff(X) = \begin{bmatrix} f(X) & q(X) \\ 0 & 0 \end{bmatrix} \tag{114}$$

where $f(X)$ is a $n \times n$ matrix and $q(X)$ is a n row vector.

$X \mapsto aff(X)$ is an injective Lie algebra homomorphism from $\mathfrak{g}$ in the Lie algebra of all $(n+1)\times(n+1)$ matrices, $gl(n+1,R)$:

$$\left|\begin{aligned}&\mathfrak{g} \to gl(n+1,R)\\ &X \mapsto aff(X)\end{aligned}\right. \tag{115}$$

If we denote $\mathfrak{g}_{aff} = aff(\mathfrak{g})$, we write G_{aff} the linear Lie subgroup of $GL(n+1,R)$ generated by $\mathfrak{g}_{aff}$.

An element of $s \in G_{aff}$ is expressed by:

$$Aff(s) = \begin{bmatrix} \mathbf{f}(s) & \mathbf{q}(s) \\ 0 & 1 \end{bmatrix} \tag{116}$$

Let M_{aff} be the orbit of G_{aff} through the origin o, then $M_{aff} = \mathbf{q}(G_{aff}) = G_{aff} / K_{aff}$ where $K_{aff} = \{s \in G_{aff} / \mathbf{q}(s) = 0\} = Ker(\mathbf{q})$.

Example:

Let Ω be a convex domain in R^n containing no complete straight lines, we define a convex cone $V(\Omega)$ in $R^{n+1} = R^n \times R$ by $V(\Omega) = \{(\lambda x, x) \in R^n \times R / x \in \Omega, \lambda \in R^+\}$. Then there exists an affine embedding:

$$\ell : x \in \Omega \mapsto \begin{bmatrix} x \\ 1 \end{bmatrix} \in V(\Omega) \tag{117}$$

If we consider η the group of homomorphism of $A(n,R)$ into $GL(n+1,R)$ given by:

$$s \in A(n,R) \mapsto \begin{bmatrix} \mathbf{f}(s) & \mathbf{q}(s) \\ 0 & 1 \end{bmatrix} \in GL(n+1,R) \tag{118}$$

with $A(n,R)$ the group of all affine transformations of R^n. We have $\eta(G(\Omega)) \subset G(V(\Omega))$ and the pair (η, ℓ) of the homomorphism $\eta : G(\Omega) \to G(V(\Omega))$ and the map $\ell : \Omega \to V(\Omega)$ is equivariant:

$$\ell \circ s = \eta(s) \circ \ell \quad \text{and} \quad d\ell \circ s = \eta(s) \circ d\ell \tag{119}$$

Let $\{x^1, x^2, ..., x^n\}$ be a local coordinate system on *M*, the Christoffel's symbols Γ_{ij}^k of the connection D are defined by:

$$D_{\frac{\partial}{\partial x^i}} \frac{\partial}{\partial x^j} = \sum_{k=1}^{n} \Gamma_{ij}^k \frac{\partial}{\partial x^k} \tag{120}$$

The torsion tensor *T* of *D* is given by:

$$T(X,Y) = D_X Y - D_Y X - [X,Y] \tag{121}$$

$$T\left(\frac{\partial}{\partial x^i}, \frac{\partial}{\partial x^j}\right) = \sum_{k=1}^{n} T_{ij}^k \frac{\partial}{\partial x^k} \quad \text{with} \quad T_{ij}^k = \Gamma_{ij}^k - \Gamma_{ji}^k \tag{122}$$

The curvature tensor *R* of *D* is given by:

$$R(X,Y)Z = D_X D_Y Z - D_Y D_X Z - D_{[X,Y]} Z \tag{123}$$

$$R\left(\frac{\partial}{\partial x^k}, \frac{\partial}{\partial x^l}\right)\frac{\partial}{\partial x^j} = \sum_i R_{jkl}^i \frac{\partial}{\partial x^i} \quad \text{with} \quad R_{jkl}^i = \frac{\partial \Gamma_{lj}^i}{\partial x^k} - \frac{\partial \Gamma_{kj}^i}{\partial x^l} + \sum_m \left(\Gamma_{lj}^m \Gamma_{km}^i - \Gamma_{kj}^m \Gamma_{lm}^i\right) \tag{124}$$

The Ricci tensor *Ric* of *D* is given by:

$$Ric(Y,Z) = Tr\{X \to R(X,Y)Z\} \tag{125}$$

$$R_{jk} = Ric\left(\frac{\partial}{\partial x^j}, \frac{\partial}{\partial x^k}\right) = \sum_i R_{kij}^i \tag{126}$$

Let *G/K* be a homogeneous space (on which *G*, a connected Lie group, acts transitively), Koszul has proved a bijective correspondence between the set of *G*-invariant flat connections on *G/K* and the set of affine representations of the Lie algebra of *G*. We consider a homogeneous space *G/K* endowed with a *G*-invariant flat connection *D* (homogeneous flat manifold) written (*G/K,D*).

Let (G,K) be the pair of connected Lie group G and its closed subgroup K. Let $\mathfrak{g}$ the Lie algebra of G and $\mathfrak{k}$ be the Lie subalgebra of $\mathfrak{g}$ coresponding to K. X^* is defined as the vector field on $M = G/K$ induced by the 1-parameter group of transformation e^{-tX}. We denote $A_{X^*} = L_{X^*} - D_{X^*}$, with L_{X^*} the Lie derivative.

Let V be the tangent space of G/K at $o = \{K\}$ and let consider, the following values at o:

$$f(X) = A_{X^*,o} \tag{127}$$

$$q(X) = X_o^* \tag{128}$$

where $A_{X^*} Y^* = -D_{Y^*} X^*$ (where D is a locally flat linear connection: its torsion and curvature tensors vanish identically), then:

$$f([X,Y]) = [f(X), f(Y)] \tag{129}$$

$$f(X)q(Y) - f(Y)q(X) = q([X,Y]) \tag{130}$$

where $\ker(\mathfrak{k}) = q$, and (f,q) an affine representation of the Lie algebra $\mathfrak{g}$:

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$$\forall X \in g,\quad X_a = \sum_i \left(\sum_j f(X)_i^j x^i + q(X)^i \right) \frac{\partial}{\partial x^i} \tag{131}$$

The 1-parameter transformation group generated by X_a is an affine transformation group of V, with linear parts given by $e^{-t.f(X)}$ and translation vector parts :

$$\sum_{n=1}^{\infty} \frac{(-t)^n}{n!} f(X)^{n-1} q(X) \tag{132}$$

These relations are proved by using:

$$\begin{cases} A_{X^*}Y^* - A_{Y^*}X^* = [X^*, Y^*] \\ [A_{X^*}, A_{Y^*}] = A_{[X^*,Y]^*} \end{cases} \quad \text{with} \quad A_{X^*}Y^* = -D_{Y^*}X^* \tag{133}$$

based on the property that the connection D is locally flat and there is local coordinate systems on M such that $D_{\frac{\partial}{\partial x^i}} \frac{\partial}{\partial x^j} = 0$ with a vanishing torsion and curvature:

$$T(X,Y) = 0 \Rightarrow D_X Y - D_Y X = [X,Y] \tag{134}$$

$$R(X,Y)Z = 0 \Rightarrow D_X D_Y Z - D_Y D_X Z = D_{[X,Y]}Z \tag{135}$$

deduced from the fact the a locally flat linear connection (vanishing of torsion and curvature).

Let ω be an invariant volume element on G/K in an affine local coordinate system $\{x^1, x^2, ..., x^n\}$ in a neighborhood of o:

$$\omega = \Phi . dx^1 \wedge ... \wedge dx^n \tag{136}$$

We can write $X^* = \sum_i \chi^i \frac{\partial}{\partial x^i}$ and develop the Lie derivative of the volume element ω:

$$L_{X^*}\omega = (L_{X^*}\Phi) dx^1 \wedge ... \wedge dx^n + \sum_j \Phi . dx^1 \wedge \cdots \wedge L_{X^*}dx^j \wedge \cdots \wedge dx^n = \left(X^*\Phi + \left(\sum_j \frac{\partial \chi^j}{\partial x^j} \right) \Phi \right) dx^1 \wedge ... \wedge dx^n \tag{137}$$

Since the volume element ω is invariant by G:

$$L_{X^*}\omega = 0 \Rightarrow X^*\Phi + \left(\sum_j \frac{\partial \chi^j}{\partial x^j} \right) \Phi = 0 \Rightarrow X^* \log \Phi = -\sum_j \frac{\partial \chi^j}{\partial x^j} \tag{138}$$

By using $A_{X^*}Y^* = -D_{Y^*}X^*$, we have:

$$\left(D_{\frac{\partial}{\partial x^i}} (A_{X^*}) \right) \left(\frac{\partial}{\partial x^j} \right) = D_{\frac{\partial}{\partial x^i}} \left(A_{X^*} \left(\frac{\partial}{\partial x^j} \right) \right) - A_{X^*} \left(D_{\frac{\partial}{\partial x^i}} \frac{\partial}{\partial x^j} \right) = -D_{\frac{\partial}{\partial x^i}} D_{\frac{\partial}{\partial x^j}} \left(\sum_k \chi^k \frac{\partial}{\partial x^k} \right) = -\sum_k \frac{\partial^2 \chi^k}{\partial x^i \partial x^j} \frac{\partial}{\partial x^k} \tag{139}$$

But as D is locally flat and X^* is an infinitesimal affine transformation with respect to D:

$$D_{\frac{\partial}{\partial x^i}} (A_{X^*}) = 0 \Rightarrow \frac{\partial^2 \chi^k}{\partial x^i \partial x^j} = 0 \tag{140}$$

The Koszul form and canonical bilinear form are given by:

$$\alpha = \sum_i \frac{\partial \log \Phi}{\partial x^i} dx^i = D \log \Phi \tag{141}$$

$$D\alpha = \sum_{i,j} \frac{\partial^2 \log \Phi}{\partial x^i \partial x^j} dx^i dx^j = Dd \log \Phi \tag{142}$$

$$L_{X^*}\alpha = L_{X^*} D \log \Phi = D L_{X^*} \log \Phi = D X^* \log \Phi = -D \left(\sum_j \frac{\partial \chi^j}{\partial x^j} \right) = -\sum_{,j} \frac{\partial^2 \chi^j}{\partial x^i \partial x^j} dx^i = 0 \tag{143}$$

Then, $L_{X^*}\alpha = 0 \quad \forall X \in g$.

By using $X^* \log \Phi = -\sum_j \frac{\partial \chi^j}{\partial x^j}$, we can obtain:

$$\alpha(X^*) = (D \log \Phi)(X^*) \underset{L_{X^*}\alpha = 0}{\Rightarrow} D_{X^*} \log \Phi = -\sum_j \frac{\partial \chi^j}{\partial x^j} \tag{144}$$

By using $A_{X^*}Y^* = -D_{Y^*}X^*$, we can develop:

$$A_{X^*}\left(\frac{\partial}{\partial x^j} \right) = -D_{\frac{\partial}{\partial x^j}} X^* = -\sum_i \frac{\partial \chi^i}{\partial x^j} \frac{\partial}{\partial x^i} \tag{145}$$

As $f(X)=A_{X^*,o}$ and $q(X)=X_o^*$:

$$Tr(f(X))=Tr\left(A_{X^*,o}\right)=-\sum_i \frac{\partial \chi^i}{\partial x^i}(o)=\alpha\left(X_0^*\right)=\alpha_0(q(X)) \tag{146}$$

If we use that $L_{X^*}\alpha=0 \quad \forall X\in \mathrm{g}$, then we obtain:

$$(D\alpha)\left(X^*,Y^*\right)=\left(D_{Y^*}\alpha\right)\left(X^*\right)=-\left(A_{Y^*}\alpha\right)\left(X^*\right)=-A_{Y^*}\left(\alpha(X^*)\right)+\alpha\left(A_{Y^*}X^*\right)=\alpha\left(A_{Y^*}X^*\right) \tag{147}$$

$$D\alpha_0(q(X),q(Y))=\alpha_0(f(Y)q(X)) \tag{148}$$

To synthetize the result proved by Jean-Louis Koszul, if α_o and $D\alpha_o$ are the values of α and $D\alpha$ at o, then:

$$\alpha_o(q(X))=Tr(f(X)) \quad \forall X\in \mathrm{g} \tag{149}$$

$$D\alpha_o(q(X),q(Y))=\langle q(X),q(Y)\rangle_o=\alpha_0(f(X)q(Y)) \quad \forall X,Y\in \mathrm{g} \tag{150}$$

Jean-Louis Koszul has also proved that the inner product $\langle .,.\rangle$ on V, given by the Riemannian metric g_{ij}, satisfies the following conditions:

$$\langle f(X)q(Y),q(Z)\rangle+\langle q(Y),f(X)q(Z)\rangle=\langle f(Y)q(X),q(Z)\rangle+\langle q(X),f(Y)q(Z)\rangle \tag{151}$$

Koszul and Vey [194, 195] have also developed extended results with the following theorem for connected hessian manifolds:

Theorem (Koszul-Vey Theorem). *Let M be a connected hessian manifold with hessian metric g. Suppose that M admits a closed 1-form α such that $D\alpha=g$ and there exists a group G of affine automorphisms of M preserving α:*

- *If M/G is quasi-compact, then the universal covering manifold of M is affinely isomorphic to a convex domain Ω of an affine space not containing any full straight line.*
- *If M/G is compact, then Ω is a sharp convex cone.*

On this basis, Koszul has given a Lie Group construction of a homogeneous cone that has been developed and applied in Information Geometry by Shima and Boyom in the framework of Hessian Geometry.

To make the link with Souriau model of thermodynamics, 1st Koszul form $\alpha=D\log\Phi=Tr(f(X))$ will play the role of the geometric heat Q and the 2nd koszul form $D\alpha=Dd\log\Phi=\langle q(X),q(Y)\rangle_o$ will be the equivalent of Souriau-Fisher metric, that is G-invariant.

9. Illustration of Koszul and Souriau Lie Group models of Information Geometry for Multivariate Gaussian laws

The case of Natural Exponential families invariant by affine group has been studied by Casalis (in 1999 paper and in her PhD thesis) [44, 45, 46, 47, 48, 49, 50] and by Letac [124, 125, 126]. We give the details of Casalis development in Appendix 3. Barndorff-Nielsen has also studied transformation models for exponential families [16,17,18,19, 103].In this chapter, we will only consider the case of Multivariate Gaussian densities.

To more deeply understand Koszul and Souriau Lie Group models of Information Geometry, we will illustrate their tools for multivariate Gaussian laws.

Consider the General Linear Group $GL(n)$ consisting of the invertible nxn matrices, that is a topological group acting linearly on R^n by:

$$\begin{aligned} &GL(n)\times R^n \to R^n \\ &(A,x)\mapsto Ax \end{aligned} \tag{152}$$

The Group $GL(n)$ is a Lie group, is a subgroup of the General Affine Group $GA(n)$, composed of all pairs (A,υ) where $A\in GL(n)$ and $\upsilon\in R^n$, the group operation given by:

$$(A_1,\upsilon_1)(A_2,\upsilon_2)=(A_1A_2,A_1\upsilon_2+\upsilon_1) \tag{153}$$

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$GL(n)$ is an open subset of R^{n^2}, and may be considered as n²-dimensional differential manifold with the same differentiable structure than R^{n^2}. Multiplication and inversion are infinitely often differentiable mappings. Consider the vector space $gl(n)$ of real $n\times n$ matrices and the commutator product:

$$\begin{aligned}&gl(n)\times gl(n)\to gl(n)\\&(A,B)\mapsto AB-BA=[A,B]\end{aligned}\tag{154}$$

This is a Lie product making $gl(n)$ into a Lie Algebra. The exponential map is then the mapping defined by:

$$\begin{aligned}&\exp : gl(n)\to GL(n)\\&A\mapsto \exp(A)=\sum_{n=0}^{\infty}\frac{A^n}{n!}\end{aligned}\tag{155}$$

Restricting A to have positive determinant one obtains the Positive General Affine Group $GA_+(n)$ that acts transitively on R^n by:

$$((A,\upsilon),x)\mapsto Ax+\upsilon\tag{156}$$

In case of symmetric Positive definite matrices $Sym^+(n)$, we can use the Cholesky decomposition:

$$R=LL^T\tag{157}$$

where L is a lower triangular matrix with real and positive diagonal entries, and L^T denotes the transpose of L, to define the square root of R.

Given a positive semidefinite matrix R, according to the spectral theorem, the continuous functional calculus can be applied to obtain a matrix $R^{1/2}$ such that $R^{1/2}$ is itself positive and $R^{1/2}R^{1/2}=R$. The operator $R^{1/2}$ is the unique non-negative square root of R.

$N_n=\{\aleph(\mu,\Sigma)/\mu\in R^n,\Sigma\in Sym^+{}_n\}$ the class of regular multivariate normal distributions, where μ is the mean vector and Σ is the (symmetric positive definite) covariance matrix ,is invariant under the transitive action of $GA(n)$. The induced action of $GA(n)$ on $R^n\times Sym^+{}_n$ is then given by:

$$\begin{aligned}&GA(n)\times\left(R^n\times Sym^+n\right)\to R^n\times Sym^+n\\&((A,\upsilon),(\mu,\Sigma))\mapsto\left(A\mu+\upsilon,A\Sigma A^T\right)\end{aligned}\tag{158}$$

and

$$\begin{aligned}&GA(n)\times R^n\to R^n\\&((A,\upsilon),x)\mapsto Ax+\upsilon\end{aligned}\tag{159}$$

As the isotropy group of $(0,I_n)$ is eqal to $O(n)$, wa can observe that:

$$N_n=GA(n)/O(n)\tag{160}$$

N_n is an open subset of the vectorspace $T_n=\{(\eta,\Omega)/\eta\in R^n,\Omega\in Sym_n\}$ and is a differentiable manifold, where the tangent space at any point may be identified with T_n.

The Fisher information defines a metric given to N_n a Riemannian manifold structure. The inner product of two tangent vectors $(\eta_1,\Omega_1)\in T_n$, $(\eta_2,\Omega_2)\in T_n$ at the point $(\mu,\Sigma)\in N_n$ is given by:

$$g_{(\mu,\Sigma))}((\eta_1,\Omega_1),(\eta_1,\Omega_1))=\eta_1^T\Sigma^{-1}\eta_2+\frac{1}{2}Tr\left(\Sigma^{-1}\Omega_1\Sigma^{-1}\Omega_2\right)\tag{161}$$

Niels Christian Bang Jesperson has proved that the transformation model on R^n with parameter set $R^n\times Sym^+{}_n$ are exacltly those of the form $p_{\mu,\Sigma}=f_{\mu,\Sigma}\lambda$ where λ is the Lebesque measure, where $f_{\mu,\Sigma}(x)=h\left((x-\mu)^T\Sigma^{-1}(x-\mu)\right)/\det(\Sigma)^{1/2}$ and $h:[0,+\infty[\to R^+$ is a continuous function with $\int_0^{+\infty}h(s)s^{\frac{n}{2}-1}ds<+\infty$. Distributions with densities of this form are called elliptic distributions.

To improve understanding of tools, we will consider $GA(n)$ as a sub-group of affine group, that could be defined by a Matrix Lie group G_{aff}, that acts for Multivariate gaussians laws:

$$\begin{bmatrix} Y \\ 1 \end{bmatrix} = \begin{bmatrix} R^{1/2} & m \\ 0 & 1 \end{bmatrix} \begin{bmatrix} X \\ 1 \end{bmatrix} = \begin{bmatrix} R^{1/2}X + m \\ 1 \end{bmatrix} \quad , \quad \begin{cases} (m,R) \in R^n \times Sym^+(n) \\ M = \begin{bmatrix} R^{1/2} & m \\ 0 & 1 \end{bmatrix} \in G_{aff} \end{cases} \tag{162}$$

$X \approx \aleph(0,I) \rightarrow Y \approx \aleph(m,R)$

We can verify that M is a Lie group with classical properties, that product of M preserve the structure, the associativity, the non commutativity, and the existence of neutral element:

$$\left. \begin{aligned} M_1.M_2 &= \begin{bmatrix} R_1^{1/2} & m_1 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} R_2^{1/2} & m_2 \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} R_1^{1/2} R_2^{1/2} & R_1^{1/2} m_2 + m_1 \\ 0 & 1 \end{bmatrix} \\ M_2.M_1 &= \begin{bmatrix} R_2^{1/2} & m_2 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} R_1^{1/2} & m_1 \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} R_2^{1/2} R_1^{1/2} & R_2^{1/2} m_1 + m_2 \\ 0 & 1 \end{bmatrix} \end{aligned} \right\} \Rightarrow \begin{cases} M_1.M_2 \in G_{aff} \\ M_2.M_1 \in G_{aff} \\ M_1.M_2 \neq M_2.M_1 \\ M_1.(M_2.M_3) = (M_1.M_2).M_3 \\ M_1.I = M_1 \end{cases} \tag{163}$$

We can also observe that the inverse preserves the structure:

$$M = \begin{bmatrix} R^{1/2} & m \\ 0 & 1 \end{bmatrix} \Rightarrow M_R^{-1} = M_L^{-1} = M^{-1} = \begin{bmatrix} R^{-1/2} & -R^{-1/2} m \\ 0 & 1 \end{bmatrix} \in G_{aff} \tag{164}$$

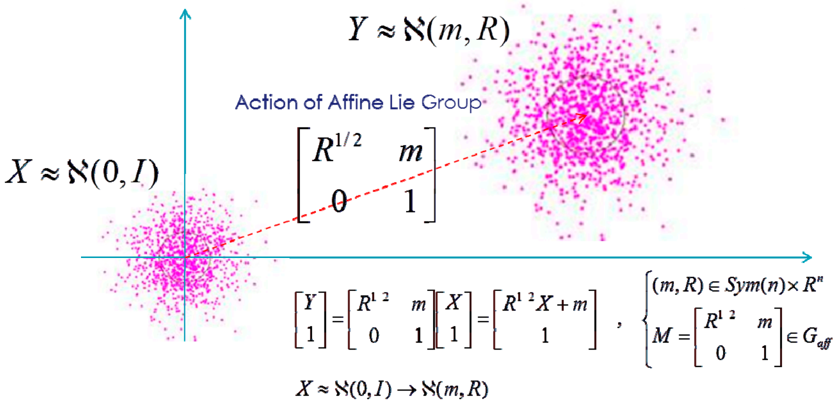

Figure 10. Affine Lie Group action for Multivariate Gaussian Law

To this Lie group we can associate a Lie algebra whose underlying vector space is the tangent space of the Lie group at the identity element and which completely captures the local structure of the group. This Lie group acts smoothly on the manifold, and acts on the vector fields. Any tangent vector at the identity of a Lie group can be extended to a left (respectively right) invariant vector field by left (respectively right) translating the tangent vector to other points of the manifold. This identifies the tangent space at the identity $\mathfrak{g} = T_I(G)$ with the space of left invariant vector fields, and therefore makes the tangent space at the identity into a Lie algebra, called the Lie algebra of G.

$$L_G : \begin{cases} G_{aff} \rightarrow G_{aff} \\ M \mapsto L_M N = M.N \end{cases} \quad \text{and} \quad R_G : \begin{cases} G_{aff} \rightarrow G_{aff} \\ M \mapsto R_M N = N.M \end{cases} \tag{165}$$

Considering the curve $\gamma(t)$ and its derivative $\dot{\gamma}(t)$:

$$\gamma(t) = \begin{bmatrix} R^{1/2}(t) & m(t) \\ 0 & 1 \end{bmatrix} \quad \text{and} \quad \dot{\gamma}(t) = \begin{bmatrix} \dot{R}^{1/2}(t) & \dot{m}(t) \\ 0 & 0 \end{bmatrix} \tag{166}$$

We can consider the curve with the point $\gamma(0)$ moved at the the identy element on the left or on the right. Then, the tangent plan at identity element provides the Lie Algebra:

$$\Gamma_L(t) = L_{M^{-1}}(\gamma(t)) = \begin{bmatrix} R^{-1/2} R^{1/2}(t) & R^{-1/2}(m(t) - m) \\ 0 & 1 \end{bmatrix} \tag{167}$$

$$\dot{\Gamma}_L(t)\Big|_{t=0} = \begin{bmatrix} R^{-1/2} \dot{R}^{1/2}(0) & R^{-1/2} \dot{m}(0) \\ 0 & 1 \end{bmatrix} = \frac{d}{dt}\left(L_{M^{-1}}(\gamma(t))\right)\Big|_{t=0} = dL_{M^{-1}} \dot{\gamma}(0) = dL_{M^{-1}} \dot{M} \tag{168}$$

Lie Algebra on the right and on the left is the defined by:

$$\begin{aligned} & dL_{M^{-1}} : T_M(G) \to \mathrm{g}_L \\ & \qquad \dot{M} \mapsto \Omega_L = dL_{M^{-1}}\dot{M} = M^{-1}\dot{M} = \begin{bmatrix} R^{-1/2}\dot{R}^{1/2} & R^{-1/2}\dot{m} \\ 0 & 0 \end{bmatrix} \end{aligned} \tag{169}$$

$$\begin{aligned} & dR_{M^{-1}} : T_M(G) \to \mathrm{g}_R \\ & \qquad \dot{M} \mapsto \Omega_R = dR_{M^{-1}}\dot{M} = \dot{M}M^{-1} = \begin{bmatrix} R^{-1/2}\dot{R}^{1/2} & \dot{m} - R^{-1/2}\dot{R}^{1/2}\dot{m} \\ 0 & 0 \end{bmatrix} \end{aligned} \tag{170}$$

We can then observe the velocities in two different ways , either by placing in a fixed outside frame, either by putting in place of the element in the process of moving by placing in the reference frame of the element.

$$\begin{bmatrix} X(t) \\ 1 \end{bmatrix} = M \begin{bmatrix} x \\ 1 \end{bmatrix} \Rightarrow \begin{bmatrix} \dot{X}(t) \\ 0 \end{bmatrix} = \Omega_R \begin{bmatrix} X(t) \\ 1 \end{bmatrix} \quad \text{with } x \text{ fixed} \tag{171}$$

$$\begin{bmatrix} x(t) \\ 1 \end{bmatrix} = M^{-1} \begin{bmatrix} X \\ 1 \end{bmatrix} \Rightarrow \begin{bmatrix} \dot{x}(t) \\ 0 \end{bmatrix} = -\Omega_L \begin{bmatrix} X \\ 1 \end{bmatrix} \quad \text{with } X \text{ fixed} \tag{172}$$

In the following, we will complete the global view by the operators which will allow to link algebra (from the left or the right) between them and also connect to their dual. We will first consider the automorphisms, the action by conjugation of the Lie group on itself, that allows this operator to carry a member of the group.

$$\begin{aligned} & AD : G \times G \to G \\ & \qquad M,N \mapsto AD_M N = M.N.M^{-1} \end{aligned} \tag{173}$$

$$\begin{cases} M_1 = \begin{bmatrix} R_1^{1/2} & m_1 \\ 0 & 1 \end{bmatrix} \ , \ M_2 = \begin{bmatrix} R_2^{1/2} & m_2 \\ 0 & 1 \end{bmatrix} \\ AD_{M_1} M_2 = \begin{bmatrix} R_2^{1/2} & -R_2^{1/2} m_1 + R_1^{1/2} m_2 + m_1 \\ 0 & 1 \end{bmatrix} \end{cases} \tag{174}$$

If now we consider a curve $N(t)$ curve on the manifold via the identity at t = 0. Its image by the previous operator will be then curve $\gamma = M.N(t).M^{-1}$ passing through Identity element at $t = 0$. As $\dot{N}(0)$ is an element of the Lie algebra and its image by previous conjugation operator is called Adjoint operator:

$$\begin{aligned} & Ad : G \times \mathrm{g} \to \mathrm{g} \\ & \qquad M,n \mapsto Ad_M n = M.n.M^{-1} = \left.\frac{d}{dt}\right|_{t=0} \left(AD_M N(t) \right) \ \text{ with } \begin{cases} N(0) = I \\ \dot{N}(0) = n \in \mathrm{g} \end{cases} \end{aligned} \tag{175}$$

We can then compute the Adjoint operator for previous Lie group:

$$\begin{cases} n_{2L} = \begin{bmatrix} R_2^{-1/2}\dot{R}_2^{1/2} & R_2^{-1/2}\dot{m}_2 \\ 0 & 0 \end{bmatrix}, n_{2R} = \begin{bmatrix} R_2^{-1/2}\dot{R}_2^{1/2} & -R_2^{-1/2}\dot{R}_2^{1/2} m_2 + \dot{m}_2 \\ 0 & 0 \end{bmatrix} \\ Ad_{M_1} n_{2L} = n_{2R} \ \text{ and } \ Ad_{M_1} n_{2R} = \begin{bmatrix} R_2^{-1/2}\dot{R}_2^{1/2} & -R_2^{-1/2}\dot{R}_2^{1/2} m_2 + \dot{R}_2^{1/2} m_2 + R_2^{1/2}\dot{m}_2 \\ 0 & 0 \end{bmatrix}, Ad_{M_1^{-1}} n_{2R} = n_{2L} \end{cases} \tag{176}$$

We recall that the Lie algebra has been defined as the tangent space at the identity of a Lie group. We will then introduced a Lie bracket $[.,.]$, the expression of the operator associated with the combined action of the Lie algebra on itself , called adjoint operator. The adjoint operator represents the action by conjugation of the Lie algebra on itself and is defined by:

$$\begin{aligned} & ad : \mathrm{g} \times \mathrm{g} \to \mathrm{g} \\ & \qquad n,m \mapsto ad_m n = m.n - n.m = \left.\frac{d}{dt}\right|_{t=0} \left(Ad_M n(t) \right) = [m,n] \ \text{ with } \begin{cases} \dot{N}(0) = n \in \mathrm{g} \\ \dot{M}(0) = m \in \mathrm{g} \end{cases} \end{aligned} \tag{177}$$

We can then compute this operator for our use case:

$$n_{1L} = \begin{bmatrix} R_1^{-1/2}\dot{R}_1^{1/2} & R_1^{-1/2}\dot{m}_1 \\ 0 & 0 \end{bmatrix}, n_{2L} = \begin{bmatrix} R_2^{-1/2}\dot{R}_2^{1/2} & R_2^{-1/2}\dot{m}_2 \\ 0 & 0 \end{bmatrix} \tag{178}$$

$$ad_{n_{1L}} n_{2L} = [n_{1L}, n_{2L}] = \begin{bmatrix} 0 & R_1^{-1/2}\left(\dot{R}_1^{1/2}\dot{m}_2 - \dot{R}_2^{1/2}\dot{m}_1\right)R_2^{-1/2} \\ 0 & 0 \end{bmatrix} \tag{179}$$

$$ad_{n_{1R}} n_{2R} = [n_{1R}, n_{2R}] = \begin{bmatrix} 0 & R_1^{-1/2}\dot{R}_1^{1/2}\left(-R_2^{-1/2}\dot{R}_2^{1/2}m_2 + \dot{m}_2\right) - R_2^{-1/2}\dot{R}_2^{1/2}\left(-R_1^{-1/2}\dot{R}_1^{1/2}m_1 + \dot{m}_1\right) \\ 0 & 0 \end{bmatrix} \quad (180)$$

To study the geodesic trajectories of the group, we consider the Lagrangian from the total kinetic energy (a quadratic form on speeds). It may therefore in particular be written in the left algebra "left", with the scalar product associated with the metric .

$$E_L = \frac{1}{2}\langle n_L, n_L\rangle = \frac{1}{2}Tr\left[n_L^T n_L\right] \quad (181)$$

If we consider as scalar product:

$$\begin{aligned} &\langle .,.\rangle : \mathfrak{g}^* \times \mathfrak{g} \to R \\ &\quad k,n \mapsto \langle k, n\rangle = Tr\left(k^T n\right) \end{aligned} \quad (182)$$

and left algebra:

$$n_L = \begin{bmatrix} R^{-1/2}\dot{R}^{1/2} & R^{-1/2}\dot{m} \\ 0 & 0 \end{bmatrix} \quad (183)$$

we obtain for the total kinetic energy

$$E_L = \frac{1}{2}\left(Tr\left(R^{-1}\dot{R}\right) + \dot{m}^T R^{-1}\dot{m}\right) \quad (184)$$

We will then introduce the coadjoint operator that will enable to work on the elements of the dual algebra of the Lie algebra defined above. Like algebra, which is physically the space of instantaneous speeds, the dual algebra is the space of moments. For dual of left algebra, the moment is given by:

$$\Pi_L = \frac{\partial E_L}{\partial n_L} = n_L \quad (185)$$

Where E_L is the kinetic energy of the system and is currently associated with Π_L is an element of the left algebra. The moment space is the dual algebra , denoted $\mathfrak{g}^*$, associated with the Lie algebra $\mathfrak{g}$. This value is deduced from the computation:

$$\begin{aligned} &\left\langle \frac{\partial E_L}{\partial n_L}, \delta U\right\rangle = \lim_{\varepsilon\to 0} \frac{E_L(n_L + \varepsilon.\delta U) - E_L(n_L)}{\varepsilon} \\ &\text{with } E_L(n_L + \varepsilon.\delta U) = \frac{1}{2}\langle n_L + \varepsilon.\delta U, n_L + \varepsilon.\delta U\rangle = \frac{1}{2}(n_L + \varepsilon.\delta U)^T (n_L + \varepsilon.\delta U) \\ &\left\langle \frac{\partial E_L}{\partial n_L}, \delta U\right\rangle = 2.\frac{1}{2}tr\left(\eta_L^T \delta U\right) = \langle n_L, \delta U\rangle \Rightarrow \frac{\partial E_L}{\partial n_L} = n_L \end{aligned} \quad (186)$$

Then the moment map is given by:

$$\begin{aligned} &\alpha_M : \mathfrak{g} \to \mathfrak{g}^* \\ &\quad n_L \mapsto \Pi_L = \eta_L \end{aligned} \quad (187)$$

We can obseve that the application that turns left algebra in its dual algebra is the identity application but physically, the first are moment and the seconds instantaneous speeds.

We can also define the moment Π_R associated to the right algebra η_R by:

$$\langle \Pi_L, n_L\rangle = \langle \Pi_L, M^{-1} n_R M\rangle = \langle \Pi_R, n_R\rangle \quad (188)$$

But as $\Pi_L = n_L$, we can deduce that:

$$\begin{aligned} &\langle n_L, M^{-1} n_R M\rangle = \langle \Pi_R, n_R\rangle \\ &\text{with } M = \begin{bmatrix} R^{1/2} & m \\ 0 & 1 \end{bmatrix}, n_L = \begin{bmatrix} R^{-1/2}\dot{R}^{1/2} & R^{-1/2}\dot{m} \\ 0 & 0 \end{bmatrix} \text{ and } \eta_R = \begin{bmatrix} R^{-1/2}\dot{R}^{1/2} & \dot{m} - R^{-1/2}\dot{R}^{1/2}\dot{m} \\ 0 & 0 \end{bmatrix} \\ &\Rightarrow \Pi_R = \begin{bmatrix} R^{-1/2}\dot{R}^{1/2} + R^{-1}\dot{m}m^T & R^{-1}\dot{m} \\ 0 & 0 \end{bmatrix} \end{aligned} \quad (189)$$

Then, the operator that transform the right algebra to its dual algebra is given by:

$$\begin{aligned} &\beta_M : \mathfrak{g} \to \mathfrak{g}^* \\ &\quad n_R = \begin{bmatrix} \eta_{R1} & \eta_{R2} \\ 0 & 0 \end{bmatrix} \mapsto \Pi_R = \begin{bmatrix} \eta_{R1}\left(1 + m^T R^{-1} m\right) + \eta_{R2} m^T R^{-1} & \eta_{R1} R^{-1} m + R^{-1}\eta_{R2} \\ 0 & 0 \end{bmatrix} \end{aligned}$$

(190)

As there is an operator to change the view of algebra, there is one that did the same on the dual algebra, the co-adjoint operator that is the conjugate action of Lie group on its dual algebra:

$$\begin{cases} Ad^* : G \times \mathfrak{g}^* \to \mathfrak{g} \\ \quad M, \eta \mapsto Ad_M^* \eta \end{cases} \text{ with } \left\langle Ad_M^* \eta, n \right\rangle = \left\langle \eta, Ad_M n \right\rangle \text{ where } n \in \mathfrak{g} \tag{191}$$

We can then develop this expression for our use case of affine sup-group, we find:

$$\begin{cases} M = \begin{bmatrix} A & b \\ 0 & 1 \end{bmatrix} \in G \\ \eta = \begin{bmatrix} \eta_1 & \eta_2 \\ 0 & 0 \end{bmatrix} \in \mathfrak{g}^* \\ n = \begin{bmatrix} n_1 & n_2 \\ 0 & 0 \end{bmatrix} \in \mathfrak{g} \end{cases} \Rightarrow \begin{cases} \left\langle Ad_M^* \eta, n \right\rangle = \left\langle \eta, Ad_M n \right\rangle = \left\langle \eta, MnM^{-1} \right\rangle \\ \left\langle Ad_M^* \eta, n \right\rangle = \left\langle \begin{bmatrix} \eta_1 - \eta_2 b^T & A\eta_2 \\ 0 & 0 \end{bmatrix}, n \right\rangle \end{cases} \Rightarrow Ad_M^* \eta = \begin{bmatrix} \eta_1 - \eta_2 b^T & A\eta_2 \\ 0 & 0 \end{bmatrix} \tag{192}$$

and we can also observed that:

$$Ad_{M^{-1}}^* \eta = \begin{bmatrix} \eta_1 + A\eta_2 b^T & A\eta_2 \\ 0 & 0 \end{bmatrix} \tag{193}$$

And the following relation between the left and the right algebras:

$$Ad_M^* \Pi_R = \Pi_L \quad \text{and} \quad Ad_{M^{-1}}^* \Pi_L = \Pi_R \tag{194}$$

As we have define a commutateur on the Lie algebra , it is possible to define one on its dual algebra. This commutator on the dual algebra can also be defined using operator expressing the combined action of the algebra of its dual. This operator is called the co-adjoint operator:

$$\begin{cases} ad^* : \mathfrak{g} \times \mathfrak{g}^* \to \mathfrak{g}^* \\ \quad n, \eta \mapsto ad_n^* \eta \end{cases} \text{ with } \left\langle ad_n^* \eta, \kappa \right\rangle = \left\langle \eta, ad_n \kappa \right\rangle \text{ where } \kappa \in \mathfrak{g} \tag{195}$$

We can develop this co-adjoint operator on its dual algebra for our use-case:

$$\begin{cases} \kappa = \begin{bmatrix} \kappa_1 & \kappa_2 \\ 0 & 0 \end{bmatrix} \in G \\ \eta = \begin{bmatrix} \eta_1 & \eta_2 \\ 0 & 0 \end{bmatrix} \in \mathfrak{g}^* \\ n = \begin{bmatrix} n_1 & n_2 \\ 0 & 0 \end{bmatrix} \in \mathfrak{g} \end{cases} \Rightarrow \begin{cases} \left\langle ad_n^* \eta, \kappa \right\rangle = \left\langle \eta, ad_n \kappa \right\rangle = \left\langle \eta, n\kappa - \kappa n \right\rangle \\ \left\langle ad_n^* \eta, \kappa \right\rangle = \left\langle \begin{bmatrix} -\eta_2 n_2^T & n_1 \eta_2 \\ 0 & 0 \end{bmatrix}, \kappa \right\rangle \end{cases} \Rightarrow \begin{cases} ad_n^* \eta = \begin{bmatrix} -\eta_2 n_2^T & n_1 \eta_2 \\ 0 & 0 \end{bmatrix} \\ ad_n^* \eta = \{n, \eta\} \end{cases} \tag{196}$$

This co-adjoint operator will give the equation of Euler-Poincaré equation. While the Euler-Lagrange equation is defined on the tangent bundle (union of the tangent spaces at each point) of the manifold and give the geodesics, the equation of Euler-Poincaré equation gives a differential system on the dual Lie algebra of the group associated with the manifold.

We can also comple these maps by an additional ones. First, $p \in T_M^* G$ the moment associated with $\dot{M} \in T_M G$ in tangent space of G at M, and also two others that map the element of the dual algebra in dual tangent space, respectively on the left and on the right:

$$\begin{cases} \left\langle \Pi_L, n_L \right\rangle = \left\langle dL_{M^{-1}}^* \Pi_L, \dot{M} \right\rangle \\ \left\langle \Pi_L, dL_{M^{-1}} \dot{M} \right\rangle = \left\langle \Pi_L, M^{-1} \dot{M} \right\rangle \end{cases} \Rightarrow p = \left(M^{-1}\right)^T \Pi_L \tag{197}$$

Where $\begin{matrix} dL_{M^{-1}}^* : \mathfrak{g}_L^* \to T_M^* G \\ \Pi_L \mapsto p = \left(M^{-1}\right)^T \Pi_L \end{matrix}$ and $\begin{matrix} dR_{M^{-1}}^* : \mathfrak{g}_R^* \to T_M^* G \\ \Pi_R \mapsto p = \Pi_R \left(M^{-1}\right)^T \end{matrix}$

(198)

From these relation, we can also observe that:

$$\Pi_L = n_L = M^{-1} \dot{M}$$
$$\Rightarrow \begin{cases} p = \left(M^{-1}\right)^T M^{-1} \dot{M} \\ p = \Xi_M . \dot{M} \quad \text{with} \quad \Xi_M = \left(M^{-1}\right)^T M^{-1} \end{cases} \tag{199}$$

All theses maps could be summarized in the following figure:

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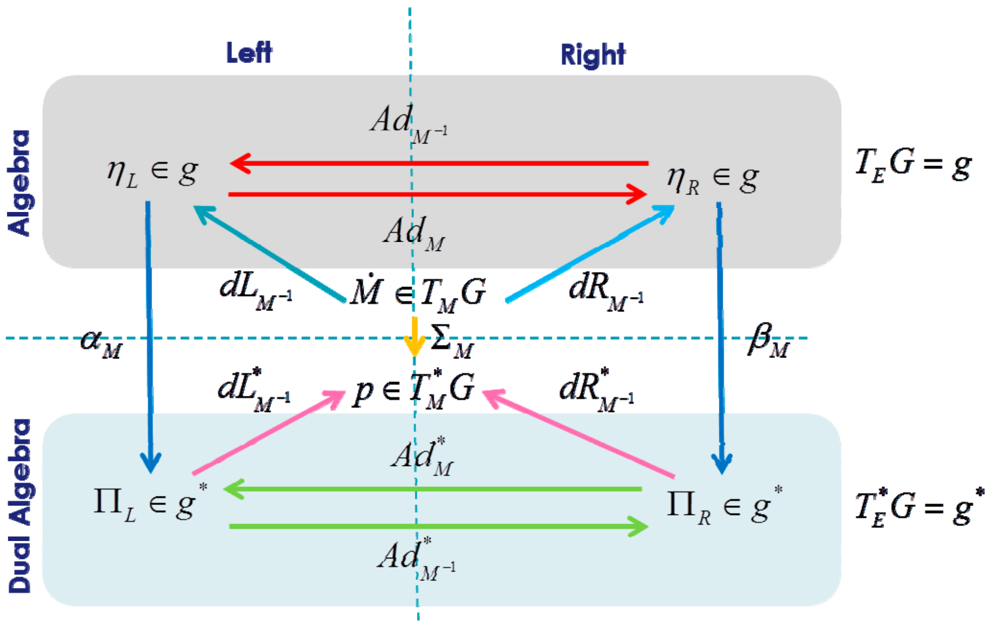

Figure 11. Maps between algebras

Heni Poincaré proved that when a Lie algebra acts locally transitively on the configuration space of a Lagrangian mechanical system, the Euler-Lagrange equations are equivalent to a new system of differential equations defined on the product of the configuration space with the Lie algebra

If we consider that the following functional is stationary for a Lagragian $l(.)$ invariant with respect to the action of the a group on the left:

$$S(\eta_L) = \int_a^b l(\eta_L)dt \quad \text{with} \quad \delta S(\eta_L) = 0 \quad \text{and} \quad l : \mathfrak{g} \to R \tag{200}$$

Solution is given by Euler-Poincaré equation:

$$\frac{d}{dt}\frac{\delta l}{\delta \eta_L} = ad^*_{\eta_L}\frac{\delta l}{\delta \eta_L} \tag{201}$$
$$\delta\eta_L = \dot{\Gamma} + ad_{\eta_L}\Gamma \quad \text{where} \quad \Gamma(t) \in \mathfrak{g}$$

If we take for the function $l(.)$, the total kinetic energy E_L, using that $\Pi_L = M^{-1}\dot{M} = \frac{\partial E_L}{\partial n_L} \in \mathfrak{g}_L$, the Euler-Poincaré equation is given by:

$$\frac{d\Pi_L}{dt} = ad^*_{n_L}\Pi_L \quad \text{with} \quad \frac{\delta l}{\delta \eta_L} = \frac{\partial E_L}{\partial n_L} = \Pi_L \in \mathfrak{g}_L \tag{202}$$

The following quantities are conserved :

$$\frac{d\Pi_R}{dt} = 0 \tag{203}$$

With this second theorem, it is possible to write the geodesic not from its coordinate system but from the quantity of motion, and in addition to determine explicitly what are the conserved quantities along the geodesic (conservations are related to the symmetries of the variety and hence the invariance of the Lagrangian under the action of the group) .

For our use-case, the Euler-Poincaré equation is given by:

$$\begin{cases}\dot{\eta}_{L1} = -\eta_{L2}\eta_{L2}^T \\ \dot{\eta}_{L2} = \eta_{L2}\eta_{L1}\end{cases} \text{ with } \begin{cases}\eta_{L1} = R^{-1/2}\dot{R}^{1/2} \\ \eta_{L2} = R^{-1/2}\dot{m}\end{cases} \Rightarrow \begin{cases}\left(R^{-1/2}\dot{R}^{1/2}\right)^{\bullet} = -R^{-1/2}\dot{m}\dot{m}^T R^{-1/2} \\ \left(R^{-1/2}\dot{m}\right)^{\bullet} = R^{-1/2}\dot{R}^{1/2}R^{-1/2}\dot{m}\end{cases} \tag{204}$$

If we remark that we have $R^{-1/2}\dot{R}^{1/2} = R^{-1/2}\left(R^{-1/2}\dot{R}\right) = R^{-1}\dot{R}$, then the conserved Souriau moment could be given by:

$$\Pi_R = \begin{bmatrix} R^{-1/2}\dot{R}^{1/2} + R^{-1}\dot{m}m^T & R^{-1}\dot{m} \\ 0 & 0 \end{bmatrix} = \begin{bmatrix} R^{-1}\dot{R} + R^{-1}\dot{m}m^T & R^{-1}\dot{m} \\ 0 & 0 \end{bmatrix} \tag{205}$$

Componants of the Souriau moment gives the conserved quantities that are the classical elements given by Emmy Noether Theorem (Souriau moment is a geometrization of Emmy Noether Theorem):

$$\frac{d\Pi_R}{dt}=\begin{bmatrix}\frac{d\left(R^{-1}\dot{R}+R^{-1}\dot{m}m^T\right)}{dt} & \frac{d\left(R^{-1}\dot{m}\right)}{dt}\\ 0 & 0\end{bmatrix}=0\Rightarrow\begin{cases}R^{-1}\dot{R}+R^{-1}\dot{m}m^T=B=cste\\ R^{-1}\dot{m}=b=cste\end{cases} \quad (206)$$

From this constant, we can obtain a reduced equation of geodesic:

$$\begin{cases}\dot{m}=Rb\\ \dot{R}=R\left(B-bm^T\right)\end{cases} \quad (207)$$

This is the Euler-Poincaré equation of geodesic. We can observe that we have obtained a reduction of the following Euler-Lagrange equation [171, 172, 34]:

$\begin{cases}\ddot{R}+\dot{m}\dot{m}^T-\dot{R}R^{-1}\dot{R}=0\\ \ddot{m}-\dot{R}R^{-1}\dot{m}=0\end{cases}$ associated to the Information Geometry metric $ds^2=dm^TR^{-1}dm+\frac{1}{2}Tr\left(\left(R^{-1}dR\right)^2\right)$

The Fisher information defines a metric turning $N_n=\{(m,R)\in R^n\times Sym^+(n)\}$ into a Riemannian Manifold. The inner product of two tangent vectors $(m_1,R_1)\in T_n$ and $(m_2,R_2)\in T_n$ at the point $(\mu,\Sigma)\in N_n$ is given by:

$$g_{(\mu,\Sigma)}((m_1,R_1),(m_2,R_2))=m_1^T\Sigma^{-1}m_2+\frac{1}{2}tr\left(\Sigma^{-1}R_1\Sigma^{-1}R_2\right) \quad (208)$$

And the geodesic is given by:

$$l(\chi)=\int_{t_0}^{t_1}\sqrt{g_{\chi(t)}(\dot{\chi}(t),\dot{\chi}(t))}dt \quad (209)$$

We can also observe that the manifold of Multivariate Gaussian is homogeneous with respect to positive affine group $GA^+(n)$:

$$ds_Y^2=ds_X^2 \quad \text{for } Y=\Sigma^{1/2}X+\mu \quad \text{with } \mathrm{GA}^+(n)=\{(\mu,\Sigma)\in R\times GL(R)/\det(\Sigma)>0\} \quad (210)$$

caracterized by the action of the group $(m,R)\mapsto\rho.(m,R)=\left(\Sigma^{1/2}m+\mu,\Sigma^{1/2}R\Sigma^{1/2T}\right),\ \rho\in GA^+(n)$

$$\text{with } \begin{bmatrix}Y\\1\end{bmatrix}=\begin{bmatrix}\Sigma^{1/2} & \mu\\ 0 & 1\end{bmatrix}\begin{bmatrix}X\\1\end{bmatrix} \quad (211)$$

$$\begin{aligned}ds_Y^2&=d\left(\Sigma^{1/2}m+\mu\right)^T\left(\Sigma^{1/2}R\Sigma^{1/2T}\right)^{-1}d\left(\Sigma^{1/2}m+\mu\right)+\frac{1}{2}Tr\left(\left(\left(\Sigma^{1/2}R\Sigma^{1/2T}\right)^{-1}d\left(\Sigma^{1/2}R\Sigma^{1/2T}\right)\right)^2\right)\\ ds_Y^2&=dm^TR^{-1}dm+\frac{1}{2}Tr\left(\left(R^{-1}dR\right)^2\right)=ds_X^2\end{aligned} \quad (212)$$

Since the special orthogonal group $SO(n)=\{\delta\in GL(R)/\det(\delta)=1\}$ is the stabilizer subgroup of $(0,I_n)$, we have the following isomorphism:

$$\begin{aligned}&GA^+(n)/SO(n)\rightarrow N_n=\{(m,R)\in R^n\times Sym^+(n)\}\\ &\rho=(\mu,\Sigma)\mapsto\rho.(0,I_n)=\left(\mu,\Sigma^{1/2}\Sigma^{1/2T}\right)=(\mu,\Sigma)\end{aligned} \quad (213)$$

We can then restrict the computation of the geodesic from $(0,I_n)$ and then we can partially integrate the system of equations:

$$\begin{cases}\dot{m}=Rb\\ \dot{R}=R\left(B-bm^T\right)\end{cases} \quad (214)$$

where $\left(R^{-1}(0)\dot{m}(0),R^{-1}(0)\left(\dot{R}(0)+\dot{m}(0)m(0)^T\right)\right)=(b,B)\in R^n\times Sym_n(R)$ are the integration constants.

From this Euler-Poincaré equation, we can compute geodesics by geodesic shooting [87, 91, 94, 153] using classical Eriksen equations [69, 70, 71, 72], by the following change of parameters:

$$\begin{cases}\Delta(t)=R^{-1}(t)\\ \delta(t)=R^{-1}(t)m(t)\end{cases}\Rightarrow\begin{cases}\dot{\Delta}=-B\Delta+bm^T\\ \dot{\delta}=-B\delta+\left(1+\delta^T\Delta^{-1}\delta\right)b\\ \Delta(0)=I_p,\delta(0)=0\end{cases}\text{ with }\begin{cases}\dot{\Delta}(0)=-B\\ \dot{\delta}(0)=b\end{cases} \quad (215)$$

The initial speed of the geodesic is given by $\left(\dot{\delta}(0),\dot{\Delta}(0)\right)$. The geodesic shooting is given by the exponential map:

$$\Lambda(t) = \exp(tA) = \sum_{n=0}^{\infty} \frac{(tA)^n}{n!} = \begin{pmatrix} \Delta & \delta & \Phi \\ \delta^T & \varepsilon & \gamma^T \\ \Phi^T & \gamma & \Gamma \end{pmatrix} \text{ with } A = \begin{pmatrix} -B & b & 0 \\ b^T & 0 & -b^T \\ 0 & -b & B \end{pmatrix} \tag{216}$$

This equation can be interpreted by Group Theory. A could be considered as an element of Lie algebra $so(n+1,n)$ of special Lorentz group $SO_O(n+1,n)$ and more specifically as the element $\mathbf{p}$ of Cartan Decomposition $\mathbf{l}+\mathbf{p}$ where $\mathbf{l}$ is the Lie algebra of a maximal compact sub-group $K = S(O(n+1)\times O(n))$ of the Group $G = SO_O(n+1,n)$. We know that its exponential map defines a geodesic on Riemannian Symetric space G/K.

This equation can be established by following developments:

$$\dot{\Lambda}(t) = \Lambda(t).A \Rightarrow \begin{pmatrix} \dot{\Delta} & \dot{\delta} & \dot{\Phi} \\ \dot{\delta}^T & \dot{\varepsilon} & \dot{\gamma}^T \\ \dot{\Phi}^T & \dot{\gamma} & \dot{\Gamma} \end{pmatrix} = \begin{pmatrix} -B & b & 0 \\ b^T & 0 & -b^T \\ 0 & -b & B \end{pmatrix} . \begin{pmatrix} \Delta & \delta & \Phi \\ \delta^T & \varepsilon & \gamma^T \\ \Phi^T & \gamma & \Gamma \end{pmatrix} \tag{217}$$

We can the deduce that:

$$\begin{cases} \dot{\Delta} = -B\Delta + b\delta^T \\ \dot{\delta} = -B\delta + \varepsilon b \end{cases} \tag{218}$$

If $\varepsilon = 1 + \delta^T \Delta^{-1}\delta$, then (Δ, δ) is solution to the geodesic equation presiously defined. Since $\varepsilon(0) = 1$, it suffices to demonstrates that $\dot{\varepsilon} = \dot{\tau}$ where $\tau = \delta^T \Delta^{-1}\delta$.

From $\dot{\Lambda}(t) = \Lambda(t).A$, using that $\dot{\delta}^T = b^T\Delta - b^T\Phi^T$, we can deduce:

$$\begin{cases} \dot{\varepsilon} = b^T\delta - b^T\gamma \\ \dot{\tau} = b^T\delta - b^T\left((\tau - \varepsilon)\Delta^{-1}\delta + \Phi^T\Delta^{-1}\delta\right) \end{cases} \tag{219}$$

Then $\dot{\varepsilon} = \dot{\tau}$, if $\gamma = (\tau - \varepsilon)\Delta^{-1}\delta + \Phi\Delta^{-1}\delta$, that could be verified using relation $\Lambda.\Lambda^{-1} = I$, by observing that:

$$\Lambda^{-1} = \exp(-tA) = \Lambda(-t) = \begin{bmatrix} \Gamma & \gamma & \Phi^T \\ \gamma^T & \varepsilon & \delta^T \\ \Phi & \delta & \Delta \end{bmatrix} \tag{220}$$

$$\Lambda.\Lambda^{-1} = I \Rightarrow \begin{cases} \Delta\gamma + \varepsilon\delta + \Phi\delta = 0 \\ \Delta\Phi^T + \delta\delta^T + \Phi\Delta = 0 \end{cases} \Rightarrow \begin{cases} \gamma = -\varepsilon\Delta^{-1}\delta - \Delta^{-1}\Phi\delta \\ \Phi^T\Delta^{-1} + \Delta^{-1}\delta\delta^T\Delta^{-1} + \Delta^{-1}\Phi = 0 \end{cases} \Rightarrow \begin{cases} \gamma = -\varepsilon\Delta^{-1}\delta - \Delta^{-1}\Phi\delta \\ \Phi^T\Delta^{-1}\delta + \tau\Delta^{-1}\delta + \Delta^{-1}\Phi\delta = 0 \end{cases} \tag{221}$$

We can then compute γ from two last equations:

$$\gamma = (\tau - \varepsilon)\Delta^{-1}\delta + \Phi^T\Delta^{-1}\delta \tag{222}$$

As $\dot{\tau} = b^T\delta - b^T\left((\tau - \varepsilon)\Delta^{-1}\delta + \Phi^T\Delta^{-1}\delta\right)$ then we can deduce that $\dot{\tau} = b^T\delta - b^T\gamma$ and then $\dot{\tau} = \dot{\varepsilon}$.

To interprete elements of Λ, $(\Gamma(t), \gamma(t)) = (\Delta(-t), \delta(-t))$, opposite points to $(\Delta(t), \delta(t))$, and $\varepsilon = 1 + \delta^T\Delta^{-1}\delta = 1 + \gamma^T\Gamma^{-1}\gamma$.

Then the geodesic that goes through the origin $(0, I_n)$ with initial tangent vector $(b, -B)$ is the curve given by $(\delta(t), \Delta(t))$. Then the distance computation is reduced to estimate the initial tangent vector space related by $\left(R^{-1}(0)\dot{m}(0), R^{-1}(0)\left(\dot{R}(0) + \dot{m}(0)m(0)^T\right)\right) = (b, B) \in R^n \times Sym_n(R)$

The distance will be then given by the initial tangent vector:

$$d = \sqrt{\dot{m}(0)^T R^{-1}(0)\dot{m}(0) + \frac{1}{2}Tr\left[\left(R^{-1}(0)\dot{R}(0)\right)^2\right]} \tag{223}$$

This initial tanget vector vector will be identified by "Geodesic Shooting". Let $V = \log_A B$:

$$\begin{cases} \dfrac{dV_m}{dt} = \dfrac{1}{2}\left(\dfrac{dR}{dt}\right)R^{-1}V_m + \dfrac{1}{2}V_R R^{-1}\left(\dfrac{dm}{dt}\right) \\ \dfrac{dV_R}{dt} = \dfrac{1}{2}\left(\left(\dfrac{dR}{dt}\right)R^{-1}V_m + V_R R^{-1}\left(\dfrac{dR}{dt}\right)\right) - \dfrac{1}{2}\left(\left(\dfrac{dm}{dt}\right)V_m^T + V_m^T\left(\dfrac{dm}{dt}\right)\right) \end{cases} \tag{224}$$

Geodesic Shooting is corrected by using Jacobi Field J and parallel transport:

$J(t) = \left.\dfrac{\partial \chi_\alpha(t)}{\partial \alpha}\right|_{t=0}$ solution to $\dfrac{d^2 J(t)}{dt^2} + R(J(t), \dot{\chi}(t))\dot{\chi}(t) = 0$ with R the Riemann Curvature tensor.

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We consider a geodesic χ between θ_0 and θ_1 with an initial tangent vector V, and we suppose that V is perturbated by W, to $V+W$. The variation of the final point θ_1 can be determined thanks to the Jacobi field with $J(0)=0$ and $\dot{J}(0)=W$. In term of the exponential map, this could be written:

$$J(t)=\frac{d}{d\alpha}\exp_{\theta_0}\left(t(V+\alpha W)\right)\Big|_{\alpha=0} \tag{225}$$

This could be illustrated in these figures:

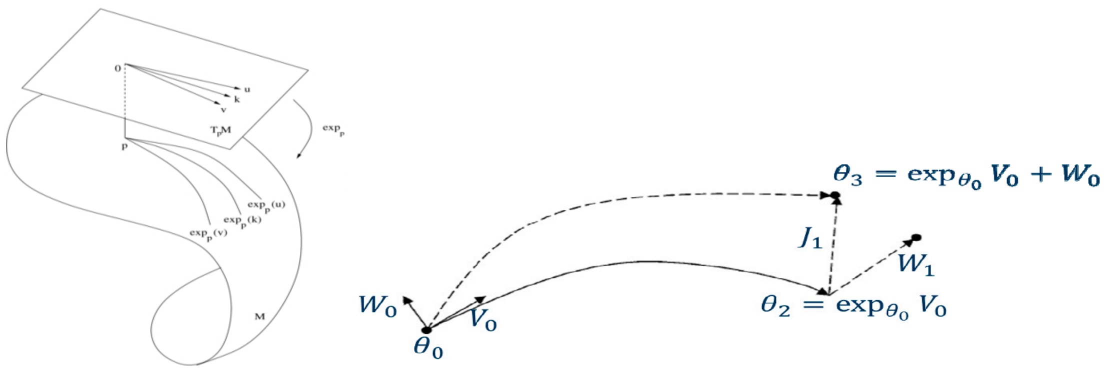

Figure 12. Geodesic Shooting Principle

We give some illustration of geodesic shooting to compute distance between multivariate Gaussian density for the case n=2:

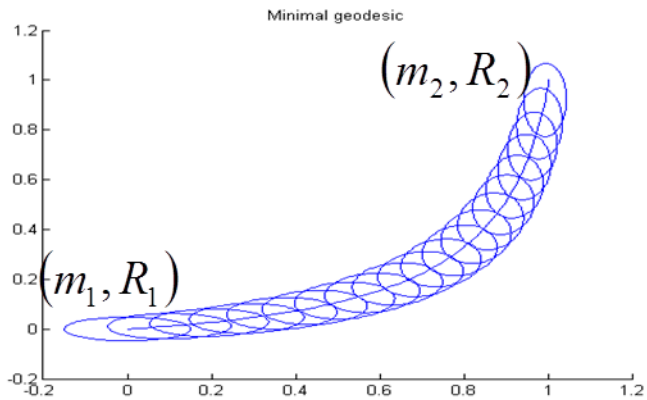

Figure 13. Geodesic Shooting between two multivariate Gaussian in case n=2

10. Souriau metric for Multivariate Gaussian Densities

To illustrate the Souriau-Fisher metric, we will consider the family of Multivariate Gaussian densities and will develop some elements that we have previously developed purely theoretically.

For the families of Multivariate Gaussian densities, that we have identified as homogeneous manifold with the associated sub-group of the affine group $\begin{bmatrix} R^{1/2} & m \\ 0 & 1 \end{bmatrix}$, we have seen that if we consider them as elements of exponential families, we can write $\hat{\xi}$ (element of the dual Lie Algebra) that play the role of geometric heat Q in Souriau Lie Group Thermodynamics, and β the geometric (planck) temperature.

$$\hat{\xi}=\begin{bmatrix} E[z] \\ E\left[zz^T\right] \end{bmatrix}=\begin{bmatrix} m \\ R+mm^T \end{bmatrix},\ \beta=\begin{bmatrix} -R^{-1}m \\ \frac{1}{2}R^{-1} \end{bmatrix} \tag{226}$$

These elements are homeomorph to the matrix elements in Matrix Lie Algebra and Dual Lie Algebra:

$$\hat{\xi} = \begin{bmatrix} R + mm^T & m \\ 0 & 0 \end{bmatrix} \in \mathrm{g}^* \ , \beta = \begin{bmatrix} \frac{1}{2}R^{-1} & -R^{-1}m \\ 0 & 0 \end{bmatrix} \in \mathrm{g} \tag{227}$$

If we consider $M = \begin{bmatrix} R'^{1/2} & m' \\ 0 & 1 \end{bmatrix}$, then we can compute the co-adjoint operator:

$$Ad_M^* \hat{\xi} = \begin{bmatrix} R + mm^T - mm'^T & R'^{1/2}m \\ 0 & 0 \end{bmatrix} \tag{228}$$

We can also compute the adjoint operator:

$$Ad_M \beta = M.\beta.M^{-1} = \begin{bmatrix} R'^{1/2} & m' \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \frac{1}{2}R^{-1} & -R^{-1}m \\ 0 & 0 \end{bmatrix} \begin{bmatrix} R'^{-1/2} & -R'^{-1/2}m' \\ 0 & 1 \end{bmatrix}$$

$$Ad_M \beta = \begin{bmatrix} \frac{1}{2}R'^{1/2}R^{-1}R'^{-1/2} & -\frac{1}{2}R'^{1/2}R^{-1}R'^{-1/2}m' - R'^{1/2}R^{-1}m \\ 0 & 0 \end{bmatrix} \tag{229}$$

We can rewrite $Ad_M \beta$ with the following identification:

$$Ad_M \beta = \begin{bmatrix} \frac{1}{2}\Omega^{-1} & -\Omega^{-1}n \\ 0 & 0 \end{bmatrix} \tag{230}$$

with $\Omega = R'^{1/2} R R'^{-1/2}$ and $n = \left(\frac{1}{2}m' + R'^{1/2}m \right)$

We have then to develop $\hat{\xi}(Ad_M(\beta))$, that is to say $\hat{\xi}(\beta)$ after action of the group on the Lie Algebra for β, given by $Ad_M(\beta)$. By analogy of structure between $\hat{\xi}(\beta)$ and β, we can write :

$$\left. \begin{array}{l} \beta = \begin{bmatrix} \frac{1}{2}R^{-1} & -R^{-1}m \\ 0 & 0 \end{bmatrix} \\ \hat{\xi}(\beta) = \begin{bmatrix} R + mm^T & m \\ 0 & 0 \end{bmatrix} \end{array} \right\} \Rightarrow \left\{ \begin{array}{l} Ad_M \beta = \begin{bmatrix} \frac{1}{2}\Omega^{-1} & -\Omega^{-1}n \\ 0 & 0 \end{bmatrix} \\ \hat{\xi}(Ad_M(\beta)) = \begin{bmatrix} \Omega + nn^T & n \\ 0 & 0 \end{bmatrix} \end{array} \right. \tag{231}$$

We have then to identify the cohomology cycle $\theta(M)$ from $\hat{\xi}(Ad_M(\beta)) = Ad_M^*(\hat{\xi}) + \theta(M)$ $\Rightarrow \theta(M) = \hat{\xi}(Ad_M(\beta)) - Ad_M^* \hat{\xi}$ where :

$$Ad_M^* \hat{\xi} = \begin{bmatrix} R + mm^T - mm'^T & R'^{1/2}m \\ 0 & 0 \end{bmatrix} \tag{232}$$

$$\hat{\xi}(Ad_M(\beta)) = \begin{bmatrix} R'^{1/2} R R'^{-1/2} + \left(\frac{1}{2}m' + R'^{1/2}m \right)\left(\frac{1}{2}m' + R'^{1/2}m \right)^T & \left(\frac{1}{2}m' + R'^{1/2}m \right) \\ 0 & 0 \end{bmatrix} \tag{233}$$

The cocycle is then given by:

$$\theta(M) = \begin{bmatrix} R'^{1/2} R R'^{-1/2} + \left(\frac{1}{2}m' + R'^{1/2}m \right)\left(\frac{1}{2}m' + R'^{1/2}m \right)^T & \left(\frac{1}{2}m' + R'^{1/2}m \right) \\ 0 & 0 \end{bmatrix} - \begin{bmatrix} R + mm^T - mm'^T & R'^{1/2}m \\ 0 & 0 \end{bmatrix}$$

$$\theta(M) = \begin{bmatrix} \left(R'^{1/2} R R'^{-1/2} - R\right) + \left(R'^{1/2} mm^T R'^{1/2T} - mm^T\right) + \left(\frac{1}{2}m'm^T R'^{1/2T} + \frac{1}{2}R'^{1/2} mm'^T - mm'^T \right) & \frac{1}{2}m' \\ 0 & 0 \end{bmatrix} \tag{234}$$

From $\theta(M) = \hat{\xi}(Ad_M(\beta)) - Ad_M^* \hat{\xi}$, we can compute cocycle in Lie Algebra

$$\Theta = T_e \theta \tag{235}$$

used to define the tensor:

$$\begin{array}{l} \tilde{\Theta}(X,Y) : \mathrm{g} \times \mathrm{g} \to \Re \\ \qquad X, Y \mapsto \langle \Theta(X), Y \rangle \end{array} \tag{236}$$

The Souriau-Fisher Metric is then given by:

$$g_\beta([\beta,Z_1],[\beta,Z_2])=\tilde{\Theta}_\beta(Z_1,[\beta,Z_2]) \tag{237}$$

with

$$\tilde{\Theta}_\beta(Z_1,Z_2)=\tilde{\Theta}(Z_1,Z_2)+\left\langle \hat{\xi},ad_{Z_1}Z_2\right\rangle=\left\langle \Theta(Z_1),Z_2\right\rangle+\left\langle \hat{\xi},[Z_1,Z_2]\right\rangle \tag{238}$$

$$\begin{aligned} g_\beta([\beta,Z_1],[\beta,Z_2])&=\tilde{\Theta}_\beta(Z_1,[\beta,Z_2])=\tilde{\Theta}(Z_1,[\beta,Z_2])+\left\langle \hat{\xi},[Z_1,[\beta,Z_2]]\right\rangle \\ &=\left\langle \Theta(Z_1),[\beta,Z_2]\right\rangle+\left\langle \hat{\xi},[Z_1,[\beta,Z_2]]\right\rangle \end{aligned} \tag{239}$$

where $\beta=\begin{bmatrix}\frac{1}{2}R^{-1} & -R^{-1}m \\ 0 & 0\end{bmatrix}$ and $\hat{\xi}=\begin{bmatrix}R+mm^T & m \\ 0 & 0\end{bmatrix}$ (240)

If we set $Z_1=\begin{bmatrix}\frac{1}{2}\Omega_1^{-1} & -\Omega_1^{-1}n_1 \\ 0 & 0\end{bmatrix}$ and $Z_2=\begin{bmatrix}\frac{1}{2}\Omega_2^{-1} & -\Omega_2^{-1}n_2 \\ 0 & 0\end{bmatrix}$ (241)

With $\langle\ldots,\ldots\rangle$ given by $\xi=\begin{bmatrix}L & b \\ 0 & 0\end{bmatrix}$, $\beta=\begin{bmatrix}H & a \\ 0 & 0\end{bmatrix}$ with $\langle\xi,\beta\rangle=Tr\left[ba^T+H^TL\right]$ (242)

$$\begin{aligned} [\beta,Z_2]&=\beta Z_2-Z_2\beta=\begin{bmatrix}\frac{1}{2}R^{-1} & -R^{-1}m \\ 0 & 0\end{bmatrix}\begin{bmatrix}\frac{1}{2}\Omega_2^{-1} & -\Omega_2^{-1}n_2 \\ 0 & 0\end{bmatrix}-\begin{bmatrix}\frac{1}{2}\Omega_2^{-1} & -\Omega_2^{-1}n_2 \\ 0 & 0\end{bmatrix}\begin{bmatrix}\frac{1}{2}R^{-1} & -R^{-1}m \\ 0 & 0\end{bmatrix} \\ [\beta,Z_2]&=\begin{bmatrix}\frac{1}{4}\left(R^{-1}\Omega_2^{-1}-\Omega_2^{-1}R^{-1}\right) & -\frac{1}{2}\left(R^{-1}\Omega_2^{-1}n_2-\Omega_2^{-1}R^{-1}m\right) \\ 0 & 0\end{bmatrix}\end{aligned} \tag{243}$$

$$\begin{aligned} [Z_1,[\beta,Z_2]]&=\begin{bmatrix}\frac{1}{2}\Omega_1^{-1} & -\Omega_1^{-1}n_1 \\ 0 & 0\end{bmatrix}\begin{bmatrix}\frac{1}{4}\left(R^{-1}\Omega_2^{-1}-\Omega_2^{-1}R^{-1}\right) & -\frac{1}{2}\left(R^{-1}\Omega_2^{-1}n_2-\Omega_2^{-1}R^{-1}m\right) \\ 0 & 0\end{bmatrix} \\ &\quad-\begin{bmatrix}\frac{1}{4}\left(R^{-1}\Omega_2^{-1}-\Omega_2^{-1}R^{-1}\right) & -\frac{1}{2}\left(R^{-1}\Omega_2^{-1}n_2-\Omega_2^{-1}R^{-1}m\right) \\ 0 & 0\end{bmatrix}\begin{bmatrix}\frac{1}{2}\Omega_1^{-1} & -\Omega_1^{-1}n_1 \\ 0 & 0\end{bmatrix} \\ &=\begin{bmatrix}\frac{1}{8}\left(\Omega_1^{-1}\left(R^{-1}\Omega_2^{-1}-\Omega_2^{-1}R^{-1}\right)-\left(R^{-1}\Omega_2^{-1}-\Omega_2^{-1}R^{-1}\right)\Omega_1^{-1}\right) & -\frac{1}{4}\left(\Omega_1^{-1}\left(R^{-1}\Omega_2^{-1}n_2-\Omega_2^{-1}R^{-1}m\right)-\left(R^{-1}\Omega_2^{-1}-\Omega_2^{-1}R^{-1}\right)\Omega_1^{-1}n_1\right) \\ 0 & 0\end{bmatrix}\end{aligned} \tag{244}$$

We can then compute:

$$\begin{aligned}\left\langle \hat{\xi},[Z_1,[\beta,Z_2]]\right\rangle&=Tr\left[\frac{1}{4}m\left(\left(R^{-1}\Omega_2^{-1}-\Omega_2^{-1}R^{-1}\right)\Omega_1^{-1}n_1-\Omega_1^{-1}\left(R^{-1}\Omega_2^{-1}n_2-\Omega_2^{-1}R^{-1}m\right)\right)^T\right] \\ &+Tr\left[\left(\frac{1}{8}\left(\Omega_1^{-1}\left(R^{-1}\Omega_2^{-1}-\Omega_2^{-1}R^{-1}\right)-\left(R^{-1}\Omega_2^{-1}-\Omega_2^{-1}R^{-1}\right)\Omega_1^{-1}\right)\right)\left(R+mm^T\right)\right]\end{aligned} \tag{245}$$

The Souriau-Fisher metric is defined in Lie Algebra $g_\beta([\beta,Z_1],[\beta,Z_2])$ where:

$$[\beta,Z_1]=\begin{bmatrix}\frac{1}{4}\left(R^{-1}\Omega_1^{-1}-\Omega_1^{-1}R^{-1}\right) & -\frac{1}{2}\left(R^{-1}\Omega_1^{-1}n_1-\Omega_1^{-1}R^{-1}m\right) \\ 0 & 0\end{bmatrix}=\begin{bmatrix}\frac{1}{2}G_1^{-1} & -G_1^{-1}g_1 \\ 0 & 0\end{bmatrix}$$

with $G_1=2\left(\Omega_1R-R\Omega_1\right)$ and $g_1=(I-R\Omega_1R^{-1}\Omega_1^{-1})n_1+(\Omega_1R\Omega_1^{-1}R^{-1}-I)m$ (246)

$$[\beta,Z_2]=\begin{bmatrix}\frac{1}{4}\left(R^{-1}\Omega_2^{-1}-\Omega_2^{-1}R^{-1}\right) & -\frac{1}{2}\left(R^{-1}\Omega_2^{-1}n_2-\Omega_2^{-1}R^{-1}m\right) \\ 0 & 0\end{bmatrix}=\begin{bmatrix}\frac{1}{2}G_2^{-1} & -G_2^{-1}g_2 \\ 0 & 0\end{bmatrix}$$

with $G_2=2\left(\Omega_2R-R\Omega_2\right)$ and $g_2=(I-R\Omega_2R^{-1}\Omega_2^{-1})n_2+(\Omega_2R\Omega_2^{-1}R^{-1}-I)m$

and $\beta=\begin{bmatrix}\frac{1}{2}R^{-1} & -R^{-1}m \\ 0 & 0\end{bmatrix}$ (247)

Another approach to develop the Souriau-Fisher Metric $g_\beta([\beta,Z_1],[\beta,Z_2])$ is to compute the tensor $\tilde{\Theta}(X,Y)$ from the moment map J:

$\tilde{\Theta}(X,Y)=J_{[X,Y]}-\{J_X,J_Y\}$ with $\{.,.\}$ Poisson Bracket and J the Moment Map (248)

$$\widetilde{\Theta}(X,Y):\mathfrak{g}\times\mathfrak{g}\rightarrow\Re \tag{249}$$

We can then write the Souriau-Fisher metric as:

$$\widetilde{\Theta}_\beta(Z_1,Z_2)=J_{[Z_1,Z_2]}-\{J_{Z_1},J_{Z_2}\}+\langle\hat{\xi},[Z_1,Z_2]\rangle \tag{250}$$

Where the associated differentiable application J, called moment map is:

$$\begin{aligned} J : M &\rightarrow \mathfrak{g}^* \quad \text{such that} \quad J_X(x)=\langle J(x),X\rangle,\ X\in\mathfrak{g} \\ x &\mapsto J(x) \end{aligned} \tag{251}$$

This moment map could be identified with the operator that transform the right algebra to an element of its dual algebra given by:

$$\beta_M : \mathfrak{g}\rightarrow\mathfrak{g}^*$$
$$Z=\begin{bmatrix} \mathrm{N} & \eta \\ 0 & 0 \end{bmatrix}\mapsto J=\begin{bmatrix} \mathrm{N}(1+m^T R^{-1} m)+\eta m^T R^{-1} & \mathrm{N}R^{-1}m+R^{-1}\eta \\ 0 & 0 \end{bmatrix} \tag{252}$$

11. Conclusion

In this paper, we have developed Souriau's model of Lie Group Thermodynamics that recovers the symmetry broken by lack of covariance of Gibbs density in classical statistical mechanics with respect to dynamic groups action in physics (Galileo and Poincaré groups, sub-group of Affine group). Ontological model of Souriau gives geometric status to (Planck) temperature (element of Lie alebra), heat (element of dual Lie algebra) and Entropy. Souriau said in one of his paper on this new "Lie Group Thermodynamics" that "*these formulas are universal , in that they do not involve the symplectic manifold, but only Group G , the symplectic cocycle. Perhaps this Lie group thermodynamics could be of interest for mathematics*".

We have observed that Souriau has introduced a generalization of Fisher Metric, that we call Souriau-Fisher metric, that preserves the property to be defined as hessian of partition function logarithm $g_\beta=-\frac{\partial^2\Phi}{\partial\beta^2}=\frac{\partial^2\log\psi_\Omega}{\partial\beta^2}$ as in classical Information Geometry, but when the partition function (Massieu Characteristic function) has been replaced by Souriau model.

$$g_\beta([\beta,Z_1],[\beta,Z_2])=\langle\Theta(Z_1),[\beta,Z_2]\rangle+\langle Q,[Z_1,[\beta,Z_2]]\rangle \tag{253}$$

This Souriau-Fisher metric, as observed by Souriau, is equal to minus the first derivative of the heat $g_\beta=-\frac{\partial Q}{\partial\beta}$, and then could be compared by analogy to "specific heat" or "calorific Capacity". This equivalence between Fisher metric and "Heat Capacity" is of major importance and should be related to Pierre Duhem theory of thermodynamics where notion of "Capacities" is at the heart of general equations of thermodynamics.

Based on the Poincaré's idea exposed in his paper of 1889 « Sur les tentatives d'explication mécanique des principes de la thermodynamique » , we have proposed, based on Souriau's Lie group model and on analogy with mechanical variables, a variational principle of Thermodynamics deduced from Poincaré-Cartan integral invariant:

The Variational Principle holds on $\mathfrak{g}$, *for variations* $\delta\beta=\dot{\eta}+[\beta,\eta]$, *where* $\eta(t)$ *is an arbitrary path that vanishes at the endpoints,* $\eta(a)=\eta(b)=0$:

$$\delta\int_{t_0}^{t_1}\Phi(\beta(t)).dt=0 \tag{254}$$

where the Poincaré-Cartan invariant $\int_{C_a} \Phi(\beta).dt = \int_{C_b} \Phi(\beta).dt$ is defined by $\Phi(\beta)$, the Massieu characteristic function, with the analogy of $\omega = \Phi(\beta).dt = (\langle Q,\beta\rangle - s).dt = \langle Q,(\beta.dt)\rangle - s.dt$ where

$$\omega = p.dq - H.dt \quad \text{where} \quad \begin{cases} \dot{q} = \dfrac{dq}{dt} = \dfrac{\partial H}{\partial p} \leftrightarrow \beta = \dfrac{\partial s}{\partial Q} \\ p = \dfrac{\partial L}{\partial \dot{q}} \leftrightarrow Q = \dfrac{\partial \Phi}{\partial \beta} \end{cases} \quad \text{and} \quad \begin{cases} L(\dot{q}) \leftrightarrow \Phi(\beta) \\ H(p) \leftrightarrow s(Q) \\ H = p.\dot{q} - L \leftrightarrow s = \langle Q,\beta\rangle - \Phi \end{cases} \tag{255}$$

We have also defined an Euler-Poincaré Equations for Souriau model:

$$\frac{dQ}{dt} = ad^*_\beta Q \quad \text{and} \quad \frac{d}{dt}\left(Ad^*_g Q\right) = 0 \tag{256}$$

For this new covariant Thermodynamics, the fundamental notion is the coadjoint orbit that is linked to positive definite KKS (Kostant-Kirillov-Souriau) 2 form [] :

$$\omega_w(X,Y) = \langle w,[U,V]\rangle \quad \text{with} \quad X = ad_w U \in T_w\mathrm{M} \quad \text{and} \quad Y = ad_w V \in T_w\mathrm{M} \tag{257}$$

that is the Kähler-form of a *G*-invariant kähler structure compatible with the canonical complex structure of *M*, and determines a canonical Symplectic structure on *M*. When the cocycle is equal to zero, the KKS and Souriau-Fisher metric are equal. This 2-form introduced by Jean-Marie Souriau is linked to the coadjoint action and the coadjoint orbits of the group on its moment space. Souriau provided a classification of the homogeneous symplectic manifolds with this moment map. The coadjoint representation of a Lie group G is the dual of the adjoint representation. If $\mathfrak{g}$ denotes the Lie algebra of *G*, the corresponding action of *G* on $\mathfrak{g}^*$, the dual space to $\mathfrak{g}$, is called the coadjoint action. Souriau proved based on the moment map that a symplectic manifold is always a coadjoint orbit, affine of its group of Hamiltonian transformations, deducing that coadjoint orbits are the universal models of symplectic manifolds: a symplectic manifold homogeneous under the action of a Lie group, is isomorphic, up to a covering, to a coadjoint orbit. So the link between Souriau-Fisher metric and KKS 2-form will provide symplectic structure and foundation to Information Manifolds. For Souriau Thermodynamics, the Souriau-Fisher metric is the canonical structure linked to KKS 2-form, modified by the cocycle (its symplectic leaves are the orbits of the affine action that makes equivariant the moment map). This last property allows to determine all homogeneous spaces of a Lie group admitting an invariant symplectic structure by the action of this group: there are the orbits of the coadjoint representation of this group or of a central extension of this group (the central extension allowing to suppress the cocycle). For affine coadjoint orbits, we give reference to Alice Tumpach PhD [189, 190, 191] that has developed previous works of K.H. Neeb, O. Biquard and P. Gauduchon.

other promising domains of research are theory of Generating maps [51, 52, 199, 200] and the link with Poisson geometry through affine Poisson group. As observed by Pierre Dazord [62] in his paper “Groupe de Poisson Affines”, extension of Poisson Group to affine Poisson group due to Drinfel’d, includes affine structures of Souriau on dual Lie algebra. Let an affine Poisson group, its universal covering could be identified to a vector space with an associated affine structure. In case that this vector space is an abelian affine Poisson group, we find affine structure of Souriau. For abelian group $(R^3,+)$, affine Poisson groups are the affine structures of Souriau.

This Souriau’s model of Lie Group Thermodynamics could be the promising way to achieve René Thom dream to replace Thermodynamics by Geometry [187, 188], and could be extended to the Second Order Extension of the Gibbs State [92,93].

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"*Si on ajoute que la critique qui accoutume l'esprit, surtout en matière de faits, à recevoir de simples probabilités pour des preuves, est, par cet endroit, moins propre à le former, que ne le doit être la géométrie qui lui fait contracter l'habitude de n'acquiescer qu'à l'évidence; nous répliquerons qu'à la rigueur on pourrait conclure de cette différence même, que la critique donne, au contraire, plus d'exercice à l'esprit que la géométrie: parce que l'évidence, qui est une et absolue, le fixe au premier aspect sans lui laisser ni la liberté de douter, ni le mérite de choisir; au lieu que les probabilités étant susceptibles du plus et du moins, il faut, pour se mettre en état de prendre un parti, les comparer ensemble, les discuter et les peser. Un genre d'étude qui rompt, pour ainsi dire, l'esprit à cette opération, est certainement d'un usage plus étendu que celui où tout est soumis à l'évidence; parce que les occasions de se déterminer sur des vraisemblances ou probabilités, sont plus fréquentes que celles qui exigent qu'on procède par démonstrations: pourquoi ne dirions –nous pas que souvent elles tiennent aussi à des objets beaucoup plus importants ?*" **- Joseph de Maistre**

« *Le cadavre qui s'acoutre se méconnait et imaginant l'éternité s'en approrie l'illusion … C'est pourquoi j'abandonnerai ces frusques et jetant le masque de mes jours, je fuirai le temps où, de concert avec les autres, je m'éreinte à me trahir* ». **Emile Cioran – Précis de décomposition**

Appendix A: Clairaut(-Legendre) Equation of Maurice Fréchet associated to "distinguished functions" as fundamental equation of Information geometry

Before Rao [160, 31], in 1943, Maurice Fréchet [74] wrote a seminal paper introducing what was then called the Cramer-Rao bound. This paper contains in fact much more that this important discovery. In particular, Maurice Fréchet introduces more general notions relative to "distinguished functions", densities with estimator reaching the bound, defined with a function, solution of Clairaut's equation. The solutions "envelope of the Clairaut's equation" are equivalents to standard Legendre transform without convexity constraints but only smoothness assumption. This Fréchet's analysis can be revisited on the basis of Jean-Louis Koszul works as seminal foundation of "Information Geometry".

We will use Maurice Fréchet notations, to consider the estimator:

$$T = H(X_1,...,X_n) \tag{258}$$

and the random variable $A(X) = \dfrac{\partial \log p_\theta(X)}{\partial \theta}$ (259)

that are associated to : $U = \sum_i A(X_i)$ (260)

The normalizing constraint $\int_{-\infty}^{+\infty} p_\theta(x)dx = 1$ implies that : $\int_{-\infty}^{+\infty}...\int_{-\infty}^{+\infty} \prod_i p_\theta(x_i)dx_i = 1$

If we consider the derivative if this last expression with respect to θ, then $\int_{-\infty}^{+\infty}...\int_{-\infty}^{+\infty}\left[\sum_i A(x_i)\right]\prod_i p_\theta(x_i)dx_i = 0$ gives : $E_\theta[U] = 0$ (261)

Similarly, if we assume that $E_\theta[T] = \theta$, then $\int_{-\infty}^{+\infty}...\int_{-\infty}^{+\infty} H(x_1,...,x_n)\prod_i p_\theta(x_i)dx_i = \theta$, and we obtain by derivation with respect to θ :

$$E[(T-\theta)U] = 1 \tag{262}$$

But as $E[T]=\theta$ and $E[U]=0$, we immedialty deduce that :

$$E[(T-E[T])(U-E[U])]=1 \tag{263}$$

From Schwarz inequality, we can develop the following relations :

$$[E(ZT)]^2 \leq E[Z^2]E[T^2]$$

$$1 \leq E[(T-E[T])^2]E[(U-E[U])^2]=(\sigma_T\sigma_U)^2 \tag{264}$$

U being the summation of independant variables, Bienaymé equality could be applied :

$$(\sigma_U)^2=\sum_i[\sigma_{A(X_i)}]^2=n(\sigma_A)^2 \tag{265}$$

From which, Fréchet deduced the bound, rediscoved by Cramer and Rao 2 years later :

$$(\sigma_T)^2 \geq \frac{1}{n(\sigma_A)^2} \tag{266}$$

Fréchet observed that it is a remarkable inequality where the second member is independent of the choice of the function H defining the " emperical value " T, where the first member can be taken to any empirical value $T=H(X_1,...,X_n)$ subject to the unique condition $E_\theta[T]=\theta$ regardless is θ.

The classic condition that the Schwarz inequality becomes an equality helps us to determine when σ_T reaches its lower bound $\frac{1}{\sqrt{n}\sigma_n}$.

The previous inequality becomes an equality if there are two numbers α and β (not random and not both zero) such that $\alpha(H'-\theta)+\beta U=0$, with H' particular function among eligible H as we have the equality . This equality is rewritten $H'=\theta+\lambda' U$ with λ' a non-random number.

If we use the previous equation, then :

$$E[(T-E[T])(U-E[U])]=1 \Rightarrow E[(H'-\theta)U]=\lambda' E_\theta[U^2]=1 \tag{267}$$

We obtain : $U=\sum_i A(X_i) \Rightarrow \lambda' n E_\theta[A^2]=1$ (268)

From which we obrain λ' and the form of the associated estimator H':

$$\lambda'=\frac{1}{nE[A^2]} \Rightarrow H'=\theta+\frac{1}{nE[A^2]}\sum_i\frac{\partial \log p_\theta(X_i)}{\partial\theta} \tag{269}$$

It is therefore deduced that the estimator that reaches the terminal is of the form:

$$H'=\theta+\frac{\sum_i \frac{\partial \log p_\theta(X_i)}{\partial\theta}}{n\int_{-\infty}^{+\infty}\left[\frac{\partial p_\theta(x)}{\partial\theta}\right]^2\frac{dx}{p_\theta(x)}} \tag{270}$$

with $E[H']=\theta+\lambda' E[U]=\theta$.

H' would be one of the eligible functions, if H' would be independent of θ. Indeed, if we consider $E_{\theta_0}[H']=\theta_0$, $E[(H'-\theta_0)^2]\leq E_{\theta_0}[(H-\theta_0)^2]$ $\forall H$ such that $E_{\theta_0}[H]=\theta_0$.

$H=\theta_0$ satisfies the equation and inequality shows that it is almost certainly equal to θ_0.

So to look for θ_0, we should know beforehand θ_0.

At this stage, Fréchet looked for "distinguished functions" ("densités distinguées" in French), as any probability density $p_\theta(x)$ such that the function :

$$h(x)=\theta+\frac{\frac{\partial \log p_\theta(x)}{\partial\theta}}{\int_{-\infty}^{+\infty}\left[\frac{\partial p_\theta(x)}{\partial\theta}\right]^2\frac{dx}{p_\theta(x)}} \tag{271}$$

is independant of θ. The objective of Fréchet is then to determine the minimizing function $T=H'(X_1,...,X_n)$ that reaches the bound. We can deduce from previous relations that:

$$\lambda(\theta)\frac{\partial \log p_\theta(x)}{\partial\theta}=h(x)-\theta \tag{272}$$

But as $\lambda(\theta)>0$, *we can consider* $\frac{1}{\lambda(\theta)}$ *as the second derivative of a function* $\Phi(\theta)$ such that :

$$\frac{\partial \log p_\theta(x)}{\partial \theta}=\frac{\partial^2\Phi(\theta)}{\partial\theta^2}[h(x)-\theta] \quad (273)$$

Wich we deduce that :

$$\ell(x)=\log p_\theta(x)-\frac{\partial\Phi(\theta)}{\partial\theta}[h(x)-\theta]-\Phi(\theta) \quad (274)$$

Is an independant quantity of θ. ***A distinguished function*** will be then given by :

$$p_\theta(x)=e^{\frac{\partial\Phi(\theta)}{\partial\theta}[h(x)-\theta]+\Phi(\theta)+\ell(x)} \quad (275)$$

With the normalizing constraint $\int_{-\infty}^{+\infty}p_\theta(x)dx=1$.

These two conditions are sufficient. Indeed, reciprocally, let three functions $\Phi(\theta)$, $h(x)$ et $\ell(x)$ that we have, for any θ : $\int_{-\infty}^{+\infty}e^{\frac{\partial\Phi(\theta)}{\partial\theta}[h(x)-\theta]+\Phi(\theta)+\ell(x)}dx=1$ (276)

Then the function is distinguished :

$$\theta+\frac{\frac{\partial \log p_\theta(x)}{\partial\theta}}{\int_{-\infty}^{+\infty}\left[\frac{\partial p_\theta(x)}{\partial\theta}\right]^2\frac{dx}{p_\theta(x)}}=\theta+\lambda(x)\frac{\partial^2\Phi(\theta)}{\partial\theta^2}[h(x)-\theta] \quad (277)$$

If $\lambda(x)\frac{\partial^2\Phi(\theta)}{\partial\theta^2}=1$, when $\frac{1}{\lambda(x)}=\int_{-\infty}^{+\infty}\left[\frac{\partial\log p_\theta(x)}{\partial\theta}\right]^2p_\theta(x)dx=(\sigma_A)^2$ (278)

The function is reduced to $h(x)$ and then is not dependant of θ.

We have then the following relation:

$$\frac{1}{\lambda(x)}=\int_{-\infty}^{+\infty}\left(\frac{\partial^2\Phi(\theta)}{\partial\theta^2}\right)^2[h(x)-\theta]^2e^{\frac{\partial\Phi(\theta)}{\partial\theta}(h(x)-\theta)+\Phi(\theta)+\ell(x)}dx \quad (279)$$

The relation is valid for any θ, on peut dériver l'expression précédente par rapport à θ :

$$\int_{-\infty}^{+\infty}e^{\frac{\partial\Phi(\theta)}{\partial\theta}(h(x)-\theta)+\Phi(\theta)+\ell(x)}\left(\frac{\partial^2\Phi(\theta)}{\partial\theta^2}\right)[h(x)-\theta]dx=0 \quad (280)$$

We can divide by $\frac{\partial^2\Phi(\theta)}{\partial\theta^2}$ because it doesn't depend on x.

If we derive again with respect to θ, we will have :

$$\int_{-\infty}^{+\infty}e^{\frac{\partial\Phi(\theta)}{\partial\theta}(h(x)-\theta)+\Phi(\theta)+\ell(x)}\left(\frac{\partial^2\Phi(\theta)}{\partial\theta^2}\right)[h(x)-\theta]^2dx=\int_{-\infty}^{+\infty}e^{\frac{\partial\Phi(\theta)}{\partial\theta}(h(x)-\theta)+\Phi(\theta)+\ell(x)}dx=1 \quad (281)$$

Combining this relation with that of $\frac{1}{\lambda(x)}$, we can deduce that $\lambda(x)\frac{\partial^2\Phi(\theta)}{\partial\theta^2}=1$ and as $\lambda(x)>0$ then $\frac{\partial^2\Phi(\theta)}{\partial\theta^2}>0$.

Fréchet emphasizes at this step, another way to approach the problem. We can select arbitrarily $h(x)$ and $l(x)$ and then $\Phi(\theta)$ is determined by:

$$\int_{-\infty}^{+\infty}e^{\frac{\partial\Phi(\theta)}{\partial\theta}[h(x)-\theta]+\Phi(\theta)+\ell(x)}dx=1 \quad (282)$$

That could be rewritten : $e^{\theta.\frac{\partial\Phi(\theta)}{\partial\theta}-\Phi(\theta)}=\int_{-\infty}^{+\infty}e^{\frac{\partial\Phi(\theta)}{\partial\theta}h(x)+\ell(x)}dx$ (283)

If we then fixed arbitrarily $h(x)$ and $l(x)$ and let s an arbitrary variable, the following function will be an explicit positive function given by $e^{\Psi(s)}$:

$$\int_{-\infty}^{+\infty}e^{s.h(x)+\ell(x)}dx=e^{\Psi(s)} \quad (284)$$

Fréchet obtained fanally the function $\Phi(\theta)$ as solution of the equation :

$$\Phi(\theta)=\theta.\frac{\partial\Phi(\theta)}{\partial\theta}-\Psi\left(\frac{\partial\Phi(\theta)}{\partial\theta}\right) \tag{285}$$

Fréchet noted that this is the Alexis Clairaut Equation.

The case $\frac{\partial\Phi(\theta)}{\partial\theta}=cste$ would reduce the density to a function that would be independant of θ, and so $\Phi(\theta)$ is given by a singular solution of this Clairaut equation, that is unique and could be computed by eliminating the variable s between :

$$\Phi=\theta.s-\Psi(s) \text{ and } \theta=\frac{\partial\Psi(s)}{\partial s} \tag{286}$$

Or between :

$$e^{\theta.s-\Phi(\theta)}=\int_{-\infty}^{+\infty}e^{s.h(x)+\ell(x)}dx \text{ and } \int_{-\infty}^{+\infty}e^{s.h(x)+\ell(x)}\left[h(x)-\theta\right]dx=0 \tag{287}$$

$\Phi(\theta)=-\log\int_{-\infty}^{+\infty}e^{s.h(x)+\ell(x)}dx+\theta.s$ where s is given implicitely by $\int_{-\infty}^{+\infty}e^{s.h(x)+\ell(x)}\left[h(x)-\theta\right]dx=0$.

What is then, when we known the distinguished function, H' among functions $H(X_1,...,X_n)$ verifying $E_\theta[H]=\theta$ and such that σ_H reaches for each value of θ, an absolute minimum, equal to $\frac{1}{\sqrt{n}\sigma_A}$.

For the previous equation:

$$h(x)=\theta+\frac{\frac{\partial\log p_\theta(x)}{\partial\theta}}{\int_{-\infty}^{+\infty}\left[\frac{\partial p_\theta(x)}{\partial\theta}\right]^2\frac{dx}{p_\theta(x)}} \tag{288}$$

We can rewrite the estimator as :

$$H'(X_1,...,X_n)=\frac{1}{n}\left[h(X_1)+...+h(X_n)\right] \tag{289}$$

And compute the associated empirical value :

$$t=H'(x_1,...,x_n)=\frac{1}{n}\sum_i h(x_i)=\theta+\lambda(\theta)\sum_i\frac{\partial\log p_\theta(x_i)}{\partial\theta}$$

And if we take $\theta=t$, we have as $\lambda(\theta)>0$:

$$\sum_i\frac{\partial\log p_t(x_i)}{\partial t}=0 \tag{290}$$

When $p_\theta(x)$ is a distinguished function, the emperical value t of θ corresponding to a sample $x_1,...,x_n$ is a root of previous equation in t. This equation has a root and only one when X is a distinguished variable. Indeed, as we have:

$$p_\theta(x)=e^{\frac{\partial\Phi(\theta)}{\partial\theta}\left[h(x)-\theta\right]+\Phi(\theta)+\ell(x)} \tag{291}$$

$$\sum_i\frac{\partial\log p_t(x_i)}{\partial t}=\frac{\partial^2\Phi(t)}{\partial t^2}\left[\frac{\sum_i h(x_i)}{n}-t\right] \text{ with } \frac{\partial^2\Phi(t)}{\partial t^2}>0 \tag{292}$$

We can then recover the unique root: $t=\frac{\sum_i h(x_i)}{n}$.

This function $T\equiv H'(X_1,...,X_n)=\frac{1}{n}\sum_i h(X_i)$ can have an arbitrary form, that is a sum of functions of each only one of the quantities and it is even the arithmetic average of N values of a same auxiliary random variable $Y=h(X)$. The dispersion is given by:

$$\left(\sigma_{T_n}\right)^2 = \frac{1}{n(\sigma_A)^2} = \frac{1}{n\int\limits_{-\infty}^{+\infty}\left[\frac{\partial p_\theta(x)}{\partial\theta}\right]^2\frac{dx}{p_\theta(x)}} = \frac{1}{n\frac{\partial^2\Phi(\theta)}{\partial\theta^2}} \tag{293}$$

and T_n follows the probability density:

$$p_\theta(t) = \sqrt{n}\frac{1}{\sigma_A\sqrt{2\pi}}e^{-\frac{n(t-\theta)^2}{2.\sigma_A^2}} \quad \text{with} \quad (\sigma_A)^2 = \frac{\partial^2\Phi(\theta)}{\partial\theta^2} \tag{294}$$

- **Clairaut Equation and Legendre Transform**

We have just observed that Fréchet shows that distinguished functions depend on a function $\Phi(\theta)$, solution of the Clairaut equation:

$$\Phi(\theta) = \theta.\frac{\partial\Phi(\theta)}{\partial\theta} - \Psi\left(\frac{\partial\Phi(\theta)}{\partial\theta}\right) \tag{295}$$

Or given by the Legendre Transform:

$$\Phi = \theta.s - \Psi(s) \quad \text{and} \quad \theta = \frac{\partial\Psi(s)}{\partial s} \tag{296}$$

Fréchet also observed that this function $\Phi(\theta)$ could be rewritten:

$\Phi(\theta) = -\log\int\limits_{-\infty}^{+\infty} e^{s.h(x)+\ell(x)}dx + \theta.s$ where s is given implicitely by $\int\limits_{-\infty}^{+\infty} e^{s.h(x)+\ell(x)}\left[h(x)-\theta\right]dx = 0$.

This equation is the fundamental equation of Information Geometry.

The "Legendre" transform was introduced by Adrien-Marie Legendre in 1787 to solve a minimal surface problem Gaspard Monge in 1784. Using a result of Jean Baptiste Meusnier, a student of Monge, it solves the problem by a change of variable corresponding to the transform which now entitled with his name. Legendre wrote: "*I have just arrived by a change of variables that can be useful in other occasions.*" About this transformation, Darboux [60] in his book gives an interpretation of Chasles: "*This comes after a comment by Mr. Chasles, to substitute its polar reciprocal on the surface compared to a paraboloïd.*" The equation of Clairaut was introduced 40 years earlier in 1734 by Alexis Clairaut [123]. Solutions "envelope of the Clairaut equation" are equivalent to the Legendre transform with unconditional convexity, but only under differentiability constraint. Indeed, for a non-convex function, Legendre transformation is not defined where the Hessian of the function is canceled, so that the equation of Clairaut only make the hypothesis of differentiability. The portion of the strictly convex function *g* in Clairaut equation $y = px - g(p)$ to the function *f* giving the envelope solutions by the formula $y = f(x)$ is precisely the Legendre transformation. The approach of Fréchet may be reconsidered in a more general context on the basis of the work of Jean-Louis Koszul.

Appendix B: Balian Gauge Model of Thermodynamics and its compliance with Souriau model

Supported by TOTAL group, Roger Balian has introduced in a Gauge Theory of Thermodynamics [8] and has also developed Information Geometry in Statistical Physics and Quantum Physics [3,4,5,6,7,8,9,10,11,12]. Balian has observed that the Entropy S (we use Balian notation, contrary with previous chapter where we use $-S$ as neg-Entropy) can be regarded as an extensive variable $q^0 = S(q^1,...,q^n)$, with q^i $(i=1,...,n)$, n independent quantities, usually extensive and conservative, characterizing the system. The n intensive variables γ_i are defined as the partial derivatives:

$$\gamma_i = \frac{\partial S(q^1,...,q^n)}{\partial q^i} \tag{297}$$

Balian has introduced a non-vanishing gauge variable p_0, without physical relevance, which multiplies all the intensive variables, defining a new set of variables:

$$p_i = -p_0.\gamma_i \quad , \quad i=1,...,n \tag{298}$$

The $2n+1$-dimensional space is thereby extended into a $2n+2$-dimensional thermodynamic space T spanned by the variables p_i, q^i with $i=0,1,...,n$, where the physical system is associated with a n+1-dimensional manifold M in T, parameterized for instance by the coordinates $q^1,...,q^n$ and

p_0. A gauge transformation which changes the extra variable p_0 while keeping the ratios $p_i / p_0 = -\gamma_i$ invariant ***is not observable***, so that a state of the system is represented by any point of a one-dimensional ray lying in M, along which the physical variables $q^0,...,q^n,\gamma_1,...,\gamma_n$ are fixed. Then, the relation between contact and canonical transformations is a direct outcome of this gauge invariance: the contact structure $\tilde{\omega} = dq^0 - \sum_{i=1}^{n} \gamma_i . dq^i$ in *2n+1* dimension can be embedded into a symplectic structure in $2n$+2 dimension, with 1-form:

$$\omega = \sum_{i=0}^{n} p_i . dq^i \tag{299}$$

as ***symplectization***, with geometric interpretation in the theory of fibre bundles.

The n *+1*-dimensional thermodynamic manifolds M are characterized by the vanishing of this form $\omega = 0$. The 1-form induces then ***a symplectic structure on*** T:

$$d\omega = \sum_{i=0}^{n} dp_i \wedge dq^i \tag{300}$$

Any thermodynamic manifold M belongs to the set of the so-called Lagrangian manifolds in T, which are the integral submanifolds of $d\omega$ with maximum dimension (n +1). Moreover, M ***is gauge invariant***, which is implied by $\omega = 0$. The extensivity of the entropy function $S(q^1,...,q^n)$ is expressed by the Gibbs-Duhem relation $S = \sum_{i=1}^{n} q^i \frac{\partial S}{\partial q^i}$, rewritten with previous relation $\sum_{i=0}^{n} p_i q^i = 0$, defining a *2n+1*-dimensional extensivity sheet in T, where the thermodynamic manifolds M should lie. Considering an infinitesimal canonical transformation, generated by the Hamiltonian $h(q^0,q^1,...,q^n,p_0,p_1,...,p_n)$, $\dot{q}_i = \frac{\partial h}{\partial p_i}$ and $\dot{p}_i = \frac{\partial h}{\partial q^i}$, the Hamilton's equations are given by Poisson bracket:

$$\dot{g} = \{g,h\} = \sum_{i=0}^{n} \frac{\partial g}{\partial q^i}\frac{\partial h}{\partial p_i} - \frac{\partial h}{\partial q_i}\frac{\partial g}{\partial p_i} \tag{301}$$

The concavity of the entropy $S(q^1,...,q^n)$, as function of the extensive variables, expresses the stability of equilibrium states. This property produces constraints on the physical manifolds M in the $2n$+2-dimensional space. It entails the existence of a metric structure in the n-dimensional space q_i relying on the quadratic form:

$$ds^2 = -d^2 S = -\sum_{i,j=1}^{n} \frac{\partial^2 S}{\partial q^i \partial q^j} dq^i dq^j \tag{302}$$

which defines a distance between two neighboring thermodynamic states.

As $d\gamma_i = \sum_{j=1}^{n} \frac{\partial^2 S}{\partial q^i \partial q^j} dq^j$, then: $$ds^2 = -\sum_{i=1}^{n} d\gamma_i dq_i = \frac{1}{p_0} \sum_{i=0}^{n} dp_i dq^i \tag{303}$$

The factor $1/p_0$ ***ensures gauge invariance***. In a continuous transformation generated by h, the metric evolves according to:

$$\frac{d}{d\tau}(ds^2) = \frac{1}{p_0}\frac{\partial h}{\partial q^0} ds^2 + \frac{1}{p_0}\sum_{i,j=0}^{n}\left(\frac{\partial^2 h}{\partial q^i \partial p_j} dp_i dp_j - \frac{\partial^2 h}{\partial q^i \partial q^j} dq^i dq^j\right) \tag{304}$$

We can observe that ***this Gauge Theory of Thermodynamics is compatible with Souriau Lie Group Thermodynamics***, where we have to consider the Souriau vector $\beta = \begin{bmatrix} \gamma_1 \\ \vdots \\ \gamma_n \end{bmatrix}$, transformed in a new vector:

$$p_i = -p_0 . \gamma_i , \quad p = \begin{bmatrix} -p_0\gamma_1 \\ \vdots \\ -p_0\gamma_n \end{bmatrix} = -p_0 . \beta \tag{305}$$

Appendix C: Casalis-Letac Affine Group Invariance for Natural Exponential Families

The characterization of the natural exponential families of R^d which are preserved by a group of affine transformationsis has been examined by Muriel Casalis in her PhD and her different papers. Her method has consisted in translating the invariance property of the family into a property concerning the measures which generate it, and to characterize such measures.

Let E a vector space of finite size, E^* its dual. $\langle \theta, x\rangle$ duality braket with $(\theta, x) \in E^* \times E$. μ Positive Radon measure on E, Laplace transform is :

$$L_\mu : E^* \to [0,\infty] \text{ with } \theta \mapsto L_\mu(\theta) = \int_E e^{\langle \theta, x\rangle} \mu(dx) \tag{306}$$

Let transformation $k_\mu(\theta)$ defined on $\Theta(u)$ interior of $D_\mu = \{\theta \in E^*, L_\mu < \infty\}$:

$$k_\mu(\theta) = \log L_\mu(\theta) \tag{307}$$

natural exponential families are given by:

$$F(\mu) = \{P(\theta,\mu)(dx) = e^{\langle \theta, x\rangle - k_\mu(\theta)} \mu(dx), \theta \in \Theta(\mu)\} \tag{308}$$

with injective function (domian of means):

$$k'_\mu(\theta) = \int_E x P(\theta,\mu)\mu(dx) \tag{309}$$

the inverse function:

$$\psi_\mu : M_F \to \Theta(\mu) \text{ with } M_F = \mathrm{Im}(k'_\mu(\Theta(\mu))) \tag{310}$$

and the Covariance operator:

$$V_F(m) = k''_\mu(\psi_\mu(m)) = (\psi'_\mu(m))^{-1}, \; m \in M_F \tag{311}$$

Measure generetad by a family F is then given by:

$$F(\mu) = F(\mu') \Leftrightarrow \exists (a,b) \in E^* \times R, \text{such that } \mu'(dx) = e^{\langle a, x\rangle + b} \mu(dx) \tag{312}$$

Let F an exponential family of E generated by μ and $\varphi : x \mapsto g_\varphi x + v_\varphi$ with $g_\varphi \in GL(E)$ automorphisms of E and $v_\varphi \in E$, then the family $\varphi(F) = \{\varphi(P(\theta,\mu)), \theta \in \Theta(\mu)\}$ is an exponential familly of E generated by $\varphi(\mu)$

Definition:

An exponential family F is invariant by a group G (affine group of E), if $\forall \varphi \in G, \varphi(F) = F$:

$$\forall \mu, F(\varphi(\mu)) = F(\mu) \tag{313}$$

(the contrary could be false)

Then Muriel Casalis has established the following theorem:

Theorem (Casalis):

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Let $F = F(\mu)$ an exponential familly of E and G affine group of E, then F is invariant by G if and only:

$$\exists a : G \to E^*, \exists b : G \to R, \text{such that}:$$
$$\forall(\varphi,\varphi') \in G^2, \begin{cases} a(\varphi\varphi') = {}^t g_\varphi^{-1} a(\varphi') + a(\varphi) \\ b(\varphi\varphi') = b(\varphi) + b(\varphi') - \langle a(\varphi'), g_\varphi^{-1} v_\varphi \rangle \end{cases} \tag{314}$$
$$\forall \varphi \in G, \varphi(\mu)(dx) = e^{\langle a(\varphi),x\rangle + b(\varphi)} \mu(dx)$$

When G is a linear subgroup, b is a character of G and a could be obtained by the help of Cohomology of Lie groups.

If we define action of G on E^* by:

$$g.x = {}^t g^{-1} x, g \in G, x \in E^* \tag{315}$$

It can be verified that:

$$a(g_1 g_2) = g_1.a(g_2) + a(g_1) \tag{316}$$

the action a is an inhomogeneous 1-cocycle:

$\forall n > 0$, let the set of all functions from G^n to E^*, $\Im(G^n, E^*)$ called inhomogenesous n-cochains,

then we can define the operators $d^n : \Im(G^n, E^*) \to \Im(G^{n+1}, E^*)$ by:

$$\begin{aligned} d^n F(g_1,\cdots,g_{n+1}) = g_1.F(g_2,\cdots,g_{n+1}) + \sum_{i=1}^{n} (-1)^i F(g_1,g_2,\cdots,g_i g_{i+1},\cdots,g_n) \\ + (-1)^{n+1} F(g_1,g_2,\cdots,g_n) \end{aligned} \tag{317}$$

Let $Z^n(G,E^*) = Ker(d^n), B(G,E^*) = \mathrm{Im}(d^{n-1})$, with Z^n inhomogneous n-cocycles , the quotient:

$$H^n(G,E^*) = Z^n(G,E^*) / B^n(G,E^*) \tag{318}$$

is the Cohomology Group of G with value in E^*. We have:

$$\begin{aligned} d^0 : E^* \to \Im(G,E^*) \\ x \mapsto (g \mapsto g.x - x) \end{aligned} \tag{319}$$

$$Z^0 = \{x \in E^*; g.x = x, \forall g \in G\} \tag{320}$$

$$\begin{aligned} d^1 : \Im(G,E^*) \to \Im(G^2,E^*) \\ F \mapsto d^1 F \ , \ d^1 F(g_1,g_2) = g_1.F(g_2) - F(g_1 g_2) + F(g_1) \end{aligned} \tag{321}$$

$$Z^1 = \{F \in \Im(G,E^*); F(g_1 g_2) = g_1.F(g_2) + F(g_1), \forall(g_1,g_2) \in G^2\} \tag{322}$$

$$B^1 = \{F \in \Im(G,E^*); \exists x \in E^*, F(g) = g.x - x\} \tag{323}$$

When the Cohomology Group $H^1(G,E^*) = 0$ then:

$$Z^1(G,E^*) = B^1(G,E^*) \tag{324}$$

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$$\Rightarrow \exists c \in E^*, \text{such that } \forall g \in G, a(g) = \left(I_d - {}^t g^{-1}\right)c \tag{325}$$

Then if $F = F(\mu)$ is an exponential familly invariant by G, μ verifies:

$$\forall g \in G, g(\mu)(dx) = e^{\langle c,x\rangle - \langle c, g^{-1}x\rangle + b(g)}\mu(dx) \tag{326}$$

$$\forall g \in G, g\left(e^{\langle c,x\rangle}\mu(dx)\right) = e^{b(g)}e^{\langle c,x\rangle}\mu(dx) \text{ with } \mu_0(dx) = e^{\langle c,x\rangle}\mu(dx) \tag{327}$$

For all compact Group, $H^1\left(G, E^*\right) = 0$ and we can express a:

$$\begin{array}{l} A : G \rightarrow GA(E) \\ \quad g \mapsto A_g \ , \ A_g(\theta) = {}^t g^{-1}\theta + a(g) \end{array} \tag{328}$$

$$\begin{array}{l} \forall (g, g') \in G^2, A_{gg'} = A_g A_{g'} \\ A(G) \text{ compact sub-group of } GA(E) \end{array} \tag{329}$$

$$\exists \text{fixed point} \Rightarrow \forall g \in G, A_g(c) = {}^t g^{-1}c + a(g) = c \Rightarrow a(g) = \left(I_d - {}^t g^{-1}\right)c \tag{330}$$

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