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Article

Luminescent P₂O₅-MoO₃-Bi₂O₃-K₂O Glasses and Glass Ceramics on Their Basis: Insights from Experimental and Computational Studies

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Abstract: This paper presents results of complimentary experimental (by electron microscopy, X-ray diffraction, diffuse reflectance, photoluminescence (PL) and FTIR spectroscopy) and computational (by molecular dynamics and DFT-based electronic structure methods) studies of oxide glasses of xP2O5-yMoO3-zBi2O3-(1-x-y-z)K2O system and glass ceramics based on them (crystal@glass), where the KBi(MoO4)2 complex oxide is the crystal component (KBi(MoO4)2 @glass). The behavior of the observed PL characteristics is analyzed in synergy with results of calculations of their atomic structures and changes in the oxygen environment of bismuth atoms during the transition crystal → interphase \rightarrow glass. It has been shown that the optical absorption and PL characteristics of such systems are largely determined by content of Bi₂O₃ and MoO₃ oxides in the initial charge, and by content bismuth ions in different charge states which exist in the produced glass and glass ceramics. It was found that the blue PL (spectral range 375 - 550 nm) of both glasses and glass ceramics originates from radiative transitions ${}^{3}P_{1} \rightarrow {}^{1}S_{0}$ in bismuth ions Bi^{3+} . The yellow-red PL (range 550-850nm) should be mainly associated with the luminescence of bismuth ions in lower charge states, Bi²⁺, Bi+ and Bi0. The thickness of the interphase layers of glass ceramics is estimated to be 1.5-2.0 nm. It has been shown that the changes in the spectra of optical absorption and PL / PL excitation of the glass ceramics occur due to the decrease in the number of oxygen atoms in the nearest surrounding of bismuth ions in the interphase region and these changes can be used for spectral probing of the formation and presence of interphase layers.

Keywords: glass ceramics; phosphate-molybdate glass; atomic structure; molecular dynamics; luminescence

1. Introduction

Oxide glass ceramics are composite materials in which particles of an oxide crystal component are dispersed in an oxide glass matrix. Considerable interest in the study of such materials is associated with the wide range of their potential applications, in particular, in electronics [1], lighting [2,3], optical thermometry [4,5], biomedicine and dentistry [5,7,8], as sealing materials [9,10], etc.

Glass-ceramic materials can show a number of unique properties compared to those of their original components. These non-additive properties of glass-ceramics can be significantly determined, in particular, by existence of an interphase at the boundary between the glass and crystalline phases, which is characterized by specific atomic structure and, accordingly, unique characteristics [11–14]. For example, for electrically conductive glass-ceramics, modification of the interphase can reduce the resistivity of the material [15]. It was found that the types and strengths of chemical bonds in the interphase determine the creep resistance of the SiOC glass-ceramics [16].

Polarization of the interphase can also affect the dielectric breakdown performance of ferroelectric barium strontium titanate oxide glass ceramics [17]. As for luminescent glass-ceramics, formation of an interphase region can be manifested in a change in the luminescence spectral profile, decay time, or an increase of its intensity [13,18,19].

It is obvious that the study of peculiarities of formation of interphase regions in glass ceramics, the atomic structure and interaction processes in interphase areas is important for elaboration and application of novel composite materials. The development of new glass ceramics requires the selection of the optimal composition of the amorphous matrix, the concentration and particle size of the crystalline component, and the manufacturing and post-treatment parameters of the resulting product. Implementation of these procedures, in turn, requires a significant amount of experimental research, significant consumption of valuable reagents, electricity and time. Theoretical and computational modelling of the atomic structure and calculation of the electronic structure and properties of glass, crystal and glass ceramics can significantly reduce the number of such experiments and associated costs.

At present, theoretical and computational modelling of the electronic structure of interphases of composite materials is a top-level task of solid-state physics and materials science [20–23]. The most common approach in the computational study of the electronic structure of interfaces and interphases is the use of band-periodic methods based on the DFT formalism [1,20,22,23]. However, electronic structure calculation methods use an adiabatic approximation that does not allow effective modelling of the mutual diffusion of atoms of material components [24,25], which causes the appearance of the above-mentioned interphase layers.

The mutual diffusion of atoms of the composite components can be effectively modelled using molecular dynamics (MD) calculations [13,26,27]. It should also be noted that MD methods allow us to study bulk systems of several thousand atoms, while for more accurate DFT approximations, geometric optimization of a system of several hundred atoms usually overconsumes computer time. For this reason, in our computational studies we use molecular dynamics tools.

It should be noted that due to great complexity of atomic structure and composition glass-ceramic composites, the computational methods, and the MD method in particular, can only approximate the real atomic and electronic properties of such materials. In view of this, the correlation of the results of experiment and modelling is a necessary criterion for the adequacy of the selected calculation parameters. The luminescence experimental method used in this work makes it possible to examine the structure studied glass ceramics at the micro/nano scale. The luminescence centers (LCs) in solids are localized formations consisting of a center's core (ions of mercury-like, transition and rare earth elements, etc.) and a shell (atoms/ions/molecules/defects of the immediate environment). The luminescent characteristics of a crystal are primarily determined by the LCs core - the energy structure of its electronic and vibrational levels. However, the characteristics of the cores can be significantly modified by the interaction with the environment and the symmetry of the latter. This feature allows us to qualify the center core as a luminescent probe of the solid state structure. In this work, we use the luminescent probe, where the bismuth ion is the core.

The choice of glass ceramics for experimental studies was made in accordance with the following considerations. It is known that only certain oxides, in particular B_2O_3 , P_2O_5 , SiO_2 , GeO_2 , As_2O_3 are glass-forming, while other oxides are either only conditionally glass-forming oxides (e.g. Bi_2O_3) or modifiers (e.g. A_2O where A is an alkali metal) [28,29]. Among these, As_2O_3 is toxic, GeO_2 is quite expensive, and SiO_2 is very refractory. Glasses based on B_2O_3 and P_2O_5 are often hygroscopic, so we focused on phosphate glasses, which when the oxygen to phosphorus ratio [O]/[P] > 3 [30] are quite resistant to moisture. Here, we present results for the bismuth-containing phosphate-molybdate system xP_2O_5 -y MoO_3 -z Bi_2O_3 -(1-x-y-z) K_2O . The introduction of bismuth and molybdenum into the system provides absorption in the UV region of the spectrum and can be accompanied by luminescence of the system in a wide range of light, from the UV to red and IR [31–33]. This, along with other important useful properties of bismuth-containing glasses and glass ceramics, determines the continuous interest shown by researchers and technologists [34–36].

As for the crystalline component, KBi(MoO₄)₂, its choice was motivated by the composition of chemical elements, which is close to the composition of glass matrix. In our opinion, this approach made it possible to expect that the interphase space would have an "intermediate" chemical composition, and mechanical stresses in the glass ceramics would be minimized.

Previously, we have already studied some samples of phosphate-molybdate glass ceramics based on the amorphous matrix $31.79P_2O_5$ - $25.43MoO_3$ - $5.69Bi_2O_3$ - $37.09K_2O$ and micro/nanocrystals KBi(MoO₄)₂:RE, where RE = Eu, Pr [37,38]. The study of the luminescent properties of these glass ceramics under direct excitation of RE ions (λ_{ex} = 473 nm) showed that the emission bands of RE ions undergo a broadening compared to the case of luminescence of KBi(MoO₄)₂:RE crystals [39]. In addition, in the case of glass ceramics with KBi(MoO₄)₂:Eu phosphor, an abnormally intense broad emission band in the region of ${}^5D_0 \rightarrow {}^7F_0$ transition of Eu³⁺ ion (560-582 nm) was observed [37]. Based on these experimental results, it was assumed that RE ions partially migrated to the glass component. The rather high intensity of the red radiation of Eu³⁺/Pr³⁺ ions in these test samples indicated the prospects for further development of luminescent coatings for chip LEDs on the base of xP₂O₅-yMoO₃-zBi₂O₃-(1-x-y-z)K₂O - KBi(MoO₄)₂:RE glass-ceramics. That is why, a detailed experimental study of the optical, in particular luminescent, properties of this type of materials and their analysis using the results of theoretical modelling was found to be reasonable.

This paper presents the results of comprehensive studies using experimental and computational methods of oxide glasses of the xP₂O₅-yMoO₃-zBi₂O₃-(1-x-y-z)K₂O system and glass ceramics based on them, where the KBi(MoO₄)₂ complex oxide is the crystal component (KBi(MoO₄)₂ @Glass) . The photoluminescence (PL) of such glasses and glass ceramics based on them was detected for the first time. The behavior of the observed PL characteristics is analyzed in synergy of results of calculations of their atomic structure and changes in the oxygen environment of bismuth atoms during the "crystal \rightarrow interphase \rightarrow glass" transitions. In particular, it has been found that the structure of the nearest oxygen environment of bismuth in the interphase regions of the composite differs significantly from the structure of the bismuth environment both in KBi(MoO₄)₂ crystal and in the glass, which affects the optical macro-characteristics of the studied materials.

2. Materials and Methods

2.1. Samples

Phosphate-molybdate glasses of the xP2O5-yMoO3-zBi2O3-(1-x-y-z)K2O system (x – in the range 23.01 – 40.31; y – in the range 13.55 – 46.02; z – in the range 1.00 – 5.69) were made by conventional melt quenching technique. The thoroughly grinded mixtures of K2CO3, KH2PO4, Bi2O3 and MoO3 analytically pure raw materials have been heated in platinum crucibles up to 1000 °C and kept at this temperature for 2 hours for homogenization. The melt obtained was poured onto a pre-heated (t \approx 400 °C) glass-carbon crucible and annealed in a furnace with temperature 450 °C for 2 h. Afterwards the furnace was turned off for natural cooling. The compositions of the glasses and some other experimental and calculated characteristics are accumulated in the Table S1 (Supplementary Materials). The samples are marked as G1, G1.48, ..., consistently with the Bi2O3 content increasing.

Glass-ceramics samples have been prepared in a similar manner. The calculated amount of crystalline powder has been added to the glass prepared by the given above melt quenching technique in a glass-carbon crucible at 450 °C under constant stirring. The mixture obtained has been annealed for 2 hours. Afterwards the furnace has been switched off for natural cooling. The glass-ceramic samples containing 10 wt. % of KBi(MoO₄)₂ crystalline powders in the G5.69 glass are noted as GC5.69. The details of the crystal synthesis can be found in [40].

The photoluminescence (PL) and PL excitation spectra (PLE) were recorded on as-prepared samples of glass and glass-ceramics. X-ray diffraction and diffuse light reflection studies were performed on grinding glasses and glass-ceramics.

Densities of the samples were measured three times for each glass by well-known Archimedes method. The obtained averaged values with standard deviations are collected in Table S1. Scanning electron microscope Tescan Mira 3 LMU with a 20-nm electronic beam diameter was used for SEM measurements. Powder X-ray diffraction (XRD) patterns were measured with use of Shimadzu XRD-6000 diffractometer (curved pyrolytic graphite counter monochromator, CuK α radiation with λ = 1.54184 Å) operating in Bragg–Brentano (Θ /2 Θ) geometry (2 Θ = 5 – 70°). The IR spectra were measured with user of a PerkinElmer Spectrum BX FTIR device. Diffuse reflection spectra of the samples were measured with a Perkin Elmer Lambda 950 spectrometer. The PL emission and excitation spectra were measured using Horiba/Jobin-Yvon Fluorolog-3 spectrofluorometer equipped with 450 W xenon lamp. All the experiments were performed at room temperature of the samples.

2.3. Calculation Approach

To model the atomic structure of glasses and interphase regions in glass-ceramics based on glass matrices xP₂O₅-yMoO₃-zBi₂O₃-(1-x-y-z)K₂O and KBi(MoO₄)₂ crystals, three-dimensional periodic cells were constructed using structural data processing utilities of the BIOVIA Materials Studio 2019 software package [41].

The initial filling of the periodic cells (the vacuum layer of the cells in the case of interface, see Supplementary Material, Figure S1) was performed using Amorphous Cell Packing utility [41]. The molecular fragments (molecular groups and cations) of initial filling of the cells were chosen of the base of analysis of experimental data on FTIR spectroscopy (details will be given below, in Sec. 3.2.1). The atomic structures of glasses and interphases were calculated using MD methods implemented in Forcite utility of Materials Studio package. Using additional calculations, we found that for the best compliance between the calculation results and experimental data, we should use the parameters of interatomic interaction corresponding to the set of Universal Forcefield Types [41] with options for specific elements given in Supplementary Material (Table S2). The charges of all atoms in the cells were set using the Charge equilibration utility with the set of parameters QEq_charged1.1 [41].

We used the following scheme for calculating the atomic structure of glasses MD methods. Such a scheme is a typical approach in the present-day modelling of the atomic structure of oxide glasses (see [42]), and in the case of the glass cells included the following steps:

- 1. Melting at $T_{melt} = 5000$ K (NVT, $\delta t = 1$ fs, t = 100 ps), Nose Thermostat, other calculation parameters were set default [41].
- 2. Quenching to T = 300 K (NVT, $\delta t = 1$ fs, quenching rate = 40 K/ps).
- 3. Equilibration at T = 300 K and P = 1 atm (NPT, $\delta t = 1$ fs, t = 100 ps). This step allows us to obtain the calculated densities of all modelled glass samples.
- 4. Production at T = 300 K (NVT, $\delta t = 1 \text{ fs}$, t = 100 ps). We obtained 10 different structures to provide sufficient statistical weight in the subsequent analysis of the properties of the studied systems. The use of about 10 such structures to average the calculated characteristics is usually considered sufficient to adequately describe the properties of glasses [43].
- 5. Geometry optimization (Fine optimization criteria, external pressure 10-4 GPa, full cell optimization [41]). It is believed that a necessary step in obtaining an adequate atomic structure of glass by MD methods is its additional geometric optimization in the same approximations [43,44].

In the case of the Crystal@Glass interphase cells, the same MD scheme was applied as for the glass cells, except step 3) "Equilibration at RT", since the NPT ensemble cannot be used in the presence of frozen atoms in the cell (in the case of glass, this step was needed primarily to obtain the calculated density value ϱ and compare it with the experimental data).

The structural characteristics and macro-properties of glasses and interphase layers, the atomic structure of which was calculated by MD methods, were calculated using Analysis procedures of Forcite utility [41].

Due to significant consumption of computer resources, it is practically impossible to perform DFT calculations for such large periodic cells as we used in the MD method (Figure S1). Therefore, for geometric structure optimizations by the DFT method of glasses and interphases, we used smaller cells, which (in the starting geometry) had a size of 12x12x12 Å and orthogonal lattice vectors (example of such cell is presenter in Supplementary Material, Figure S2). To obtain the starting geometry of atomic positions, the 12x12x12 Å cells were filled with atoms and molecular fragments of glass or interphase in certain proportions. Such proportions (chemical composition of elements) in the case of glass were taken from experimental data, and in the case of interphase, they were determined from the results of MD calculations of the atomic structure. The initial filling of the 12x12x12 Å cells was performed using the Amorphous Cell Packing utility with a set of atoms and molecular fragments. The parameters of the procedure for initial cell packing for DFT calculations were chosen the same as in the case of filling large cells used in MD calculations.

In order to create a sufficient statistical base for analysis, 8 cells of 12x12x12 Å of glass composition # 7 (see Supplementary Material, Table S1) with the same chemical composition and density but different atomic structures were generated using Amorphous Cell Packing. Similarly, 8 cells were generated with the composition and density corresponding to the interphase of KBi(MoO₄)₂ crystal with glass # 7. A full geometric optimization procedure (both atomic coordinates and cell parameters were optimized) in the DFT approximation was applied to each of the 16 structures. In subsequent analysis, the results of calculated atomic structures were averaged over 8 glass and 8 interphase DFT-optimized cells.

For geometry optimizations of the glass and interphase cells, we used the DFT-based band-periodic pseudopotential electronic structure method implemented in CASTEP, which is part of the Materials Studio package [41]. In these calculations, ion-electron interactions were modelled by nonlocal pseudopotentials of the Vanderbilt type [45]. Exchange-correlation effects were taken into account in the GGA approximation using the PBE potential approximation [46]. The energy threshold defining the size of the plane wave basis was 340.0 eV. The geometry of the nuclei was optimised using the Broyden-Fletcher-Goldfarb-Shenno (BFGS) minimisation technique [47] with the simultaneous convergence criterion for energy and force on the atoms, which were $10^{-5}\,\mathrm{eV}$ per atom and $0.03\,\mathrm{eV/Å}$ per atom, respectively. The geometric optimisation calculations were performed only for the point G of the Brillouin zone.

3. Results and Discussion

3.1. Part 1. Characterization of the Glasses and Glass-Ceramics

Characteristics of the representative glass samples of xP_2O_5 - $yMoO_3$ - zBi_2O_3 - $(1-x-y-z)K_2O$ system were experimentally measured for different combinations (x, y, z) values. Due to significant number of experimental measurements required to study the whole set of synthesized glass compositions 1-7 (see SM, Table S1), the experimental studies in this paper were limited to only several representative glass compositions (glass compositions # 1, 2, 4 and 7 in Table S1 correspond respectively to the samples denoted as G1, G1.48, G4.12 and G5.69, in Table 1 below) and GC5.69 glass ceramics.

Table 1. Optical band gaps, Bi₂O₃ and MoO₃ content, and normalized intensities of the PL bands

Glass name	Eg, eV	Oxide content,		MoO ₃ /Bi ₂ O ₃	Normalized PL intensity, r.u.		
		m Bi ₂ O ₃	ol % MoO3	ratio	I(Tot)	I(Blue)	I(YR)
G1	3.62	1.00	11.17	11.17	1.00	1.0	1.0
G1.48	3.61	1.48	16.53	11.17	1.23	1.5	1.1
G4.12	3.44	4.12	46.02	11.17	4.24	8.4	2.9
G5.69	3.60	5.69	25.43	4.47	4.89	7.5	4.1

Characteristic SEM images for the glass samples, crystal, and the glass ceramics are shown in Figure 1.

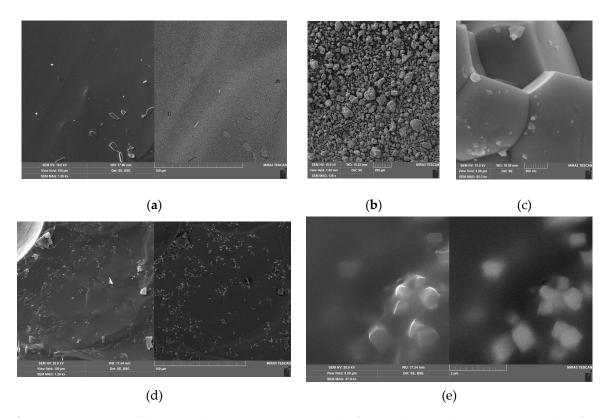


Figure 1. SEM images of the G5.69 glass (a), KBi(MoO4)2 powder (b,c), and GC5.69 glass-ceramics samples (d,e).

It can be seen that the glass surface is quite smooth. There are also small fragments of glass on the surface (the left part of Figure 1a corresponds to the image of the formed SE, and the right part to the BSE), which were formed as a result of cleavage during the preparation of the sample for measurements. The polycrystalline powder contains agglomerates up to 100 µm in size (Figure 1b). These agglomerates consist of prismatic micro- and nano-sized crystallites (Figure 1c). The studied glass ceramics contain a significant number of crystallites, which lie both near the surface (SE image, left part in Figure 1d) and in the sample bulk (BSE, right part in Figure 1d). A detailed view of the area of the glass-ceramic sample with crystalline inclusions (Figure 1e) shows similar crystal morphology in Figures 1c and Figure 1e. That is, under the above conditions of obtaining glass-ceramics, no significant transformations of the crystal component occur. This result is quite expected, since the melting point of KBi(MoO₄)₂ crystals is about 750° C [39].

3.1.2. Diffuse Reflection Data

Diffuse reflection spectra for the listed in Table 1 samples are shown in Figure 2a. As the Figure shows, the reflection coefficient for the studied glasses is low (R < 20 %) in the UV spectral range (λ < 370 nm), while in the visible range, at λ > 400 nm, the reflection coefficient increases to 50 – 80 %.

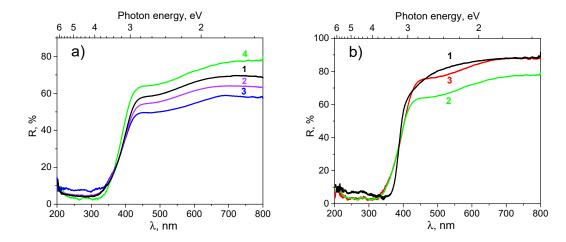


Figure 2. Diffuse reflection spectra of the: (a) G1(1), G1.48 (2), G4.12 (3), and G5.69 glasses (4); (b) KBi(MoO₄)₂ crystalline powder (1), G5.69 glass (2), and GC5.69 glass-ceramics (3)

The region of sharp change of reflection (350-400 nm) obviously corresponds to the absorption edge of the glasses. Analysis of the shape of the spectra at the absorption edge, according to the widely used Kubelka-Munk theory and Tauc plot technique [48,49], allowed us to estimate the optical band gaps E_g for these glasses. The estimated values of E_g lie in 3.44-3.62 eV range (Table 1).

It is known that absorption in 200-400 nm range can be caused by molecular phosphate and molybdate groups, molybdenum and bismuth ions (see details in [33–35,50,51] and the text below). The wide absorption bands in 450 - 650 nm range (transitions in the forbidden zone) occur due to defects in the glass matrix. Some increase in absorption at the long-wavelength limit of the measurement range (700 - 800 nm) should be of the same nature.

As Figure 2a shows, the G5.69 glass has the lowest absorption in the visible light range (370 - 800 nm) and the highest absorption in the UV range (200 - 370 nm) among all glass samples studied (if we assume that light scattering for all samples is the same). This fact makes the G5.69 glass the most interesting for further elaboration of luminescent glass ceramics based on it.

In Figure 2b, the diffuse reflection spectra of G5.69 glass are compared with corresponding spectra of GC5.69 glass ceramics and KBi(MoO₄)₂ polycrystalline powder. These data confirm the above-mentioned assumptions about the nature of absorption of the studied systems. At λ < 400 nm, the spectra of glass and glass ceramics practically coincide, while at λ > 400 nm, the reflection of the glass is much lower, although the shape of the spectra is preserved (Figure 2b, curves 2 and 3). Therefore, it can be assumed that in this region, the increase in diffuse reflection of the glass ceramics sample is due to light scattering on the KBi(MoO₄)₂ micro/nanocrystals incorporated into the glass. The increase of absorption of crystals in 275 – 375 nm range may be due to an increase in the relative amount of bismuth ions and MoO₄²- groups compared to their amount in the glass and glass ceramics. At the same time, it can be seen that the crystals practically do not show a defect-related absorption band in 450 – 650 nm range. This result evidently follows from the expected lower defectiveness of the lattice structure of crystals compared to that of glass and glass ceramics.

3.1.3. Photoluminescence Spectroscopy

The studied glasses and glass ceramics, when excited in their absorption region, exhibit photoluminescence with intensity sufficient for analysing its spectral behaviour. The most intensive PL band with a maximum near 435 nm is located in the blue luminescence range 375 - 525 nm (below we refer this band as the "blue" PL). The spectral profile of this band is asymmetric, and therefore, it consists of at least two components. In addition to this band, the spectra also reveal a wide band extending from 550 to 850 nm (Figure 3), referred below as "yellow-red" PL. For G1 glass, four components with maxima λ_{max} in the vicinity of 580, 635, 750, and 830 nm can be distinguished in this band (Figure 3a, curve 1). When moving from the glasses G1, G1.48 to glass G4.12, the total PL

intensity increases by a factor of 4.24. However, the rate of such increase is different for the "blue" and "yellow-red" bands. While the "blue" band increases by ~ 8.4 times, the "yellow-red" band increases only by 2.9 times (see Table 1).

At the same time, the components of the "yellow-red" band are almost indistinguishable for the case of G4.12 glass sample (Figure 3a, curve 3). The absence of clear correlations in changes of intensity of the "blue" and "yellow-red" bands is evidence of their different origin. The same conclusion should be drawn regarding the origin of the components of "yellow-red" band, whose intensity in transition $G1 \rightarrow G1.48 \rightarrow G4.12$ also changes uncorrelated: here, the intensity of the longwave components with λ_{max} at 750 and 830 nm increases more strongly (Figure 3a). Thus, the origins of these components are different.

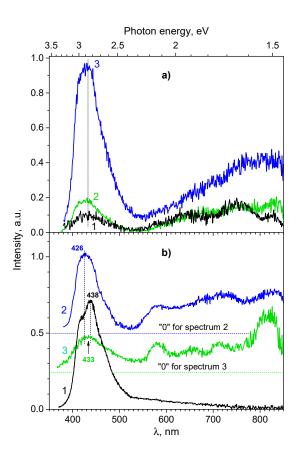


Figure 3. The PL spectra of: (a) glass samples G1 (1), G1.48 (2), and G4.12 (3); (b) KBi(MoO₄)₂ crystalline powder (1), G5.69 glass (2), and GC5.69 glass-ceramics (3); $\lambda_{\text{ex}} = 365 \text{ nm}$.

The "blue" PL is excited by short-wave UV light in a band with maximum around 265 nm and in the long-wave UV region, in a band with maximum around 367 nm (Figure 4a, curve 1). Both excitation bands likely are complex. The first one is a composition of at least two, and the second one (367 nm) of at least three components, of which the middle one has the highest intensity, and the other two look like side wings. It is easy to see that the excitation bands correspond to the features of diffuse reflection/absorption spectra described above.

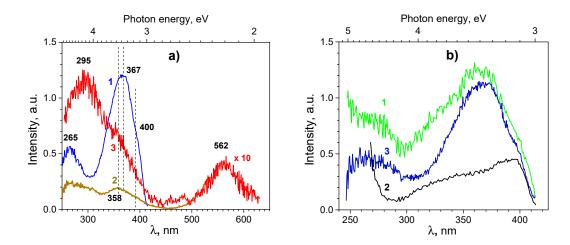


Figure 4. (a) PL excitation spectra of G4.12 glass measured at λ_{em} = 435 (1), 570 (2), and 650 nm (3). (b) PL excitation spectra of the KBi(MoO₄)₂ polycrystalline powder (1), G5.69 glass (2), and GC5.69 glass-ceramics (3); λ_{em} = 435 nm.

The spectra of "blue" PL and its excitation for the systems studied here are consistent with the PL spectra of Bi³⁺ ions in various types of crystals and glasses [40,52,53].

Comparison of the "blue" PL spectra of glass with the recorded PL spectra of KBi(MoO4)2 crystal confirms that the "blue" PL band should be associated with the luminescence of Bi3+ ions. It is well known that Bi atoms have the electronic configuration (Xe)4f¹⁴5d¹⁰6s²6p³, whose outer (s, p) electrons can combine in various chemical compounds to form charge states 0, +1, +2, +3, and +5. In particular, Bi³⁺ ions have a 6s² electronic configuration, which is quite sensitive to influence of the nearest surrounding. The ground state in this configuration is ¹S₀, and the excited states are ³P₀, ³P₁, ³P₂, and ¹P₁ (in order of the energy increasing). According to the selection rules for electro-dipole transitions with light absorption, the ${}^{1}S_{0} \rightarrow {}^{3}P_{0}$ and ${}^{1}S_{0} \rightarrow {}^{3}P_{2}$ transitions are strictly prohibited. Therefore, the photoexcitation of Bi³⁺ ions usually occurs due to the ${}^{1}S_{0} \rightarrow {}^{3}P_{1}$, ${}^{1}P_{1}$ transitions. For example, for the glasses of composition $xMoO_3-30B_2O_3-(70-x)Bi_2O_3$, transition ${}^1S_0 \rightarrow {}^1P_1$ was associated with excitation band at 225 nm, whereas the bands at 300 and 365 nm were assigned to ${}^{1}S_{0} \rightarrow {}^{3}P_{1}$ transition. The doubling of the latter transition was explained by the location of Bi3+ ions in positions with different structures of the nearest surrounding [57]. Accordingly, the luminescence of Bi³⁺ ions occurs due to ${}^{3}P_{1} \rightarrow {}^{1}S_{0}$ radiative transition, which causes the noted spectral band of "blue" PL. The peak position of this band can lie between 425 and 570 nm, which is determined by the interaction of Bi3+ ions with the environment, namely the structure and symmetry of the latter.

The spectral characteristics of observed "blue" PL bands of glasses and KBi(MoO₄)₂ crystal are quite similar, so it should be assumed that the nearest surrounding of Bi³⁺ ions in the studied glasses and KBi(MoO₄)₂ crystals is also similar. In KBi(MoO₄)₂ crystals, bismuth ions are surrounded by 8 oxygen ions located at a distance 2.55 - 2.63 Å (average < d > ≈ 2.59 Å). These oxygen ions form the first coordination sphere of the bismuth ions, while the second sphere of the environment (at distances \sim 3.80 – 4.10 Å) contains 8 molybdenum ions, 4 oxygen ions and 4 positions that can be equally occupied by both bismuth and potassium ions [39,58]. Changes in the structure and composition of Bi³⁺ environment lead to changes in the spectral characteristics of their optical absorption and luminescence. Thus, the asymmetry of the "blue" PL band and the appearance of several bands in its excitation spectra can be attributed to the presence of Bi³⁺ ions in crystals and glasses, which are in a slightly different environment.

In the PL spectrum of KBi(MoO₄)₂ crystal, one can also see a long-wavelength "tail" of low intensity extending from 500 to 800 nm (Figure 3b, curve 1). Its manifestation is associated with radiative transitions in defective molybdate groups and other defects of the crystal lattice, primarily with vacancies in the oxygen sublattice [39,40].

Taking into account the literature and our own data, we can confidently state that the "blue" PL band at 435 nm (Figure 3a) and the associated excitation bands at 265 and 367 nm (Figure 4a, curve 1) are caused by the excitation and emission processes in Bi³⁺ ions, which, similar to the case of KBi(MoO₄)₂ crystals, are mainly surrounded by 8 oxygen O²⁻ ions. It is obvious that the structure of the nearest surrounding of Bi³⁺ ions can change when the glass composition changes. As Figure 3 and Figure 4 show, the peak position of the "blue" PL band insignificantly depends on the glass composition, however, the side wings of the band differ, resulting in a change in the band shape. This observation may be an indication of the fact that the distance between the Bi³⁺ and the nearest oxygen ions and the structure in the first coordination sphere do not change, while the second and third spheres undergo changes.

The spectra of studied glasses also show additional, if compared to the case of KBi(MoO₄)2 crystals, PL (Figure 3) and PL excitation bands (Figure 4a, curves 2 and 3). The spectrum of additional PL is located in 550 – 850 nm region, longer than the spectrum of the "blue" PL of Bi³⁺, and therefore here and below this additional PL will be referenced as "yellow-red". The spectrum of the "yellow-red" luminescence is formed by several components, the peak positions of which can be roughly defined as 580 (a weakly expressed band), 635, 750, and 830 nm. In the excitation spectra of the "yellow-red" PL, in addition to the bands inherent in the excitation of Bi³⁺ ions, bands with maxima at approximately 295, 460, 480, and 560 nm are observed. The relative intensity of these bands depends on registration wavelength λ_{em} . When registering at λ_{em} = 570 nm (Figure 4a, curve 2), a band with a maximum in the vicinity of 300 nm is clearly visible in the spectrum. At such registration wavelength, the excitation band of the "blue" PL (367 nm) is "deformed" compared to curve 1 in Figure 4a, and its maximum is shifted to 358 nm, i.e., towards the above-mentioned short-wavelengh wing. At λ_{em} = 650 nm, the band at 295 nm already dominates in the excitation spectrum, and in addition to it, the peaks at 458 and 483 nm and a relatively wide band at 562 nm are observed (Figure 4a, curve 3).

Similar spectral and luminescent manifestations were observed earlier for bismuth-containing glasses and were associated with the absorption and luminescence of bismuth ions in different charges existing in an oxygen environment of different structure and composition. Earlier, it was also emphasized that the formation of several bismuth ions of different charge states could be the result of reversible auto-thermal reduction of Bi^{3+} ions down to the Bi^0 state, which can occur during melting, or one of the transformations of intermediate valence states: $Bi^{3+} \leftrightarrow Bi^{2+} \leftrightarrow Bi^{+} \leftrightarrow Bi^{0}$ [57,59–61]. Therefore, the set of luminesce bands was associated with electronic transitions in ions/atoms/groups such as Bi^{5+} , Bi^{2+} , Bi^{+} , Bi^{0} , Bi^{2-}/Bi^{2-} , $Bi^$

Table 2. Assignments of photoluminescence (PL) and PL excitation (PLE) bands in bismuth-containing glasses to specific transitions (literature data).

Glass composition	Band location: region (λ_{max}), nm		Band assignments and comments Re	efs.	
	PL	PLE			
	(560)	200 – 230	1) Transitions in Bi3+ ions: PLE bands at		
		(218)	218, 344, and 355 nm;		
		300 - 400	2) ${}^2P_{3/2} \leftrightarrow {}^2P_{1/2}$ transitions in Bi ²⁺ ions:		
		(344)	PLE band at 480 nm and PL band at 690		
33SrO-67B ₂ O ₃ -1Bi ₂ O ₃		400 - 550	nm. [59-	-62]	
		(480)	Relative contribution of these bands		
	(600)	(350)	into overall PL or PLE spectra depends		
	(690)	(490)	on the ratio between glass-forming		
		(480)	B ₂ O ₃ and network modifier SrO oxides.		
MO ₃ -B ₂ O ₃ -CeO ₂ -Bi ₂ O ₃	(600)	(200)	PL and PLE bands ascribed to ² P _{3/2} ↔	← [25]	
(M= Mo or W)	(610)	(300)	${}^{2}P_{1/2}$ transitions in Bi ²⁺ ions.	[35]	

xMoO3-30B2O3-(70-x)Bi2O3	(600)				
(x = 0, 2.5, 5, 7.5, and 10 in mol%)	(620)	(300)		[57]	
(95–x) SiO ₂ ·xSrO·5Al ₂ O ₃ ·2Bi ₂ O ₃ (x = 30, 35, 40, 45, 50, in mol%)	(640)	(532)		[52]	
	(660)	(470)	1) ${}^{2}P_{3/2} \leftrightarrow {}^{2}P_{1/2}$ transitions in Bi ²⁺ ions: PLE band at 470 nm and PL band at 660		
MO-B ₂ O ₃ -Bi ₂ O ₃ (M=Ca, Sr, Ba)	(1300)	(808)	nm; 2) The infrared emission peak possibly comes from Bi ions in low valence state.	[59,62]	
23B ₂ O ₃ –5ZnO–72Bi ₂ O ₃ –xCuO	(804)	(530)	$^3P_2 \rightarrow ^3P_0$ transitions in Bi+ ions: PL band; $^3P_0 \rightarrow ^1S_0$ transitions in Bi+ ions: PLE band	[60]	

According to the data collected in Table 2, it can be stated that decreasing of the charge state of bismuth ions (Bi³⁺ \rightarrow Bi²⁺, Bi⁺) leads to red-shifting of the absorption and luminescence bands of Birelated centers. Oxygen coordination of bismuth is another factor that probably influences the positions of these bands. In particular, it is assumed that IR luminescence (1000 to 1600 nm) in bismuth-containing glasses can be attributed to the $X_2^2P_{3/2} \rightarrow X_1^2P_{1/2}$ radiation transition in the BiO groups [52], as Bi₂O₃ can be partially converted into the noted molecular species at high temperature.

The picture of luminescence processes described above becomes more complicated for glasses containing molybdate molecular groups and molybdenum ions. As a rule, molybdate groups in oxide glasses, if they form a network, exist in the form of MoO_4^{2-} tetrahedra. The Mo^{5+} ions can also exist in such systems in small quantities (up to 1-2 %) [64]. The weak intensity of PL band at 850 nm was attributed to emission of the Mo^{3+} ions in the lead borate glasses (P_2O_5)-(Li_2O)-(CaO)-(Al_2O_3)-(MoO_3) [65], although an earlier work [51] stated that the peak of the PL band of Mo^{3+} ions should be in the vicinity of 970 nm.

In bismuth-containing glasses of the MoO_3 – B_2O_3 – B_2O_3 – B_1O_3 type, the absorption peaks near 360, 450 and 460 nm can be related to Mo^{3+} ions, while the absorption peaks in 360 – 380 and 710 – 780 nm ranges can be associated with the Mo^{5+} ions, and the Mo^{4+} ions can cause an absorption band near 550 nm [53]). At 300 nm excitation, these glasses show a "blue" PL with a maximum at 460 nm, and glasses with MoO_3 content above 7.5 M show PL bands with maxima at 600 and 620 nm, the intensity of which increased with increasing molybdenum content. Due to the overlap of emission and absorption bands of Mo and Bi ions in such systems, excitation energy transfer between Mo and Bi ions can occur. This phenomenon also affects and complicates luminescence processes in the glasses under study.

In addition to the noted above, in molybdate glasses containing bismuth, mutual redox reactions between Bi and Mo can occur. These reactions result in the partial oxidation of Mo³+ and Mo⁵+ into Mo⁶+ ions (the latter do not exhibit noticeable luminescence) and the simultaneous partial reduction of Bi³+ into Bi²+ ions [35 and elsewhere]. These transformations should quench the intensity of the "blue" fluorescence of Mo³+ and Bi³+ ions and, conversely, promote the growth of the intensity of the "yellow-red" luminescence of bismuth ions with increasing Bi₂O₃ content. It is easy to see that this is exactly the behavior of the spectral and luminescent characteristics of the glasses studied here. Indeed, Table 1 shows that for the first three samples, the total PL intensity (ITot), the intensity of the "blue" (IBlue) and the "yellow-red" PL (IYR) increase by 4.2, 8.4 and 2.9 times, respectively, with a 4.12-fold increase of Bi₂O₃ content. As can be seen, the increase in PL (ITot) corresponds to an increase in the content of the bismuth-oxide component of the glass, which confirms association of PL with bismuth ions. With a further increase in Bi₂O₃ content by a factor of 1.38, up to 5.69 M, ITot increases by a factor of 1.15, IYR by a factor of 1.39, and the intensity of the blue PL even decreases by a factor of 0.9. The latter changes are manifestations of the above-mentioned transformations of bismuth ions

towards a decrease in their charge and, accordingly, transformation of the "blue" PL spectra into the bands of "red" and near-IR emission.

3.2. Part 2. Insight from Interphase Layers

Results of experiments presented below were used in computational studies of the interphase layers.

3.2.1. FTIR Spectroscopy Data

The FTIR spectroscopy of studied glasses (Figure 5) provided information on active vibrational system of the investigated materials, which clarified the sets of molecular groups (fragments) to be taken into account for adequate computational modelling.

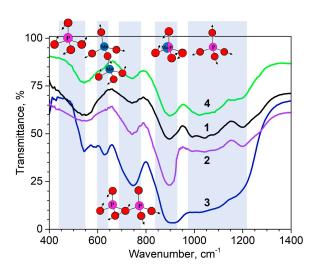


Figure 5. FTIR spectra of G1 (1), G1.48 (2), G4.12 (3) and G5.69 (4) glasses.

The FTIR spectra show a set of characteristic bands which have been expected for phosphate-molybdate glasses (Figure 5). A wide band with a maximum near 540 – 550 cm⁻¹ corresponds to O–P–O bending vibration mode [66,67]. The absorption band at 640 cm⁻¹ found for glasses with the highest content of MoO₃ (36 – 46 %) has been attributed to MoO₄²⁻ and Mo₂O₇²⁻ bending vibration modes [68]. The bands observed near 745 cm⁻¹ were attributed to stretching modes of the bridging oxygen atoms (P–O–P) in P₂O₇, P₃O₉, P₃O₁₀, P₄O₁₃ chains. This band is shifted to higher frequencies, from 739 to 752 cm⁻¹ with the MoO₃ content increase. The complex bands located around 900 – 1204 cm⁻¹ are attributed to symmetric stretching of non-bridging oxygen in PO₄³⁻ tetrahedra [67], and asymmetric stretching vibration of P–O–P in the PO₂ groups [69]. The band observed at 1222 cm⁻¹ is assigned to symmetric stretching vibrations of PO₃ groups. The band with low intensity at 1245 cm⁻¹ is assigned to P=O asymmetric stretching vibration of phosphate tetrahedra located in the chains.

It should be noted that intensity of various vibration bands of phosphate structures is increased with increase of MoO₃ content. Consequently, molybdenum oxide as a glass modifier increases the number of non-bridging oxygens in the phosphate network and acts as ionic cross-linker among phosphate groups lowering the polymerization level of phosphate nets.

In accordance with the results of IR spectroscopy and literature data, the unit cells used in calculations of glasses and crystal/glass interfaces were initially filled (see details in Sec. 2.3) with molecular groups PO₄, P₃O₉, P₃O₁₀, P₄O₁₃, MoO₄, Mo₂O₇, as well as individual K and Bi atoms. The ratio between the numbers of molecular groups was chosen so that the XRD spectra and density of the modelled structures, after MD optimization were as close as possible to the experimental data. The ratios of the number of molecular groups used for initial filling of the cells are given in Table S1.

Comparison of the calculated and experimentally measured XRD patterns (Figure 6) and density (Figure 7) shows a fairly good compliance between the calculation results and experimental data. Indeed, Figure 6 shows that the calculated XRD patterns reproduce well such features of the experimental diffractograms as the main peak in the 25-30 $^{\circ}$ region of 2 Θ , an additional shoulder in 10-22 $^{\circ}$ region, and the general profile of the curve above 35 $^{\circ}$.

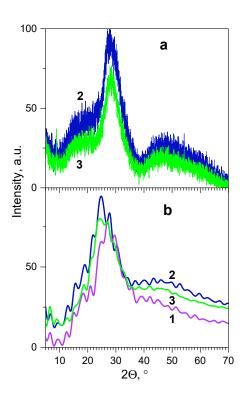


Figure 6. Experimental (**a**) and calculated (**b**) XRD patterns for the studied glasses: G1.48 (1), G4.12 (2), and G5.69 (3).

The data on the density of the samples (Figure 7) show that, despite a certain error, the calculations satisfactorily reproduce the trends in the density of studied glasses. This compliance of the calculation results with experimental data evidences the adequacy of chosen approximations, parameters and scheme of MD calculations.

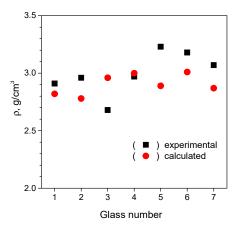


Figure 7. Experimental and calculated density versus glass composition number (see Table S1).

3.2.3. Determination of Thickness and Chemical Composition of Interphase Region

As a result of applying the MD procedure to the interphase cell, part of which contained crystal and the other part glass, a significant mixing (mutual diffusion) of glass and crystal atoms occurred

(an example of such mixed structure is presented in SM, Figure S1b). The degree of mixing can be quantitatively characterized by the calculated concentration profiles of chemical elements. Note that the part of the cell corresponding to the z-coordinate values from 0 to ~18 Å was initially filled with KBi(MoO₄)₂ crystal atoms, while the rest of the cell was filled with the glass of composition # 4 (G4.12). Atomic coordinates of the crystal layer located below the initial crystal/glass boundary (z in range from 0 to ~4 Å) were "frozen" at all steps of MD calculation scheme (see Sec. 2.3).

Figure 8 shows the calculated concentration profiles of chemical elements calculated along the normal to the "KBi(MoO₄)₂/glass" interface, i.e. along the z-coordinate. Concentration profiles were calculated for two glass compositions, #4 (G4.12) and #7 (G5.69), after initial filling (black curves in Figure 8) and after applying the MD calculation procedure (red curves). In order to obtain a sufficient statistical base, the curves shown in the figure were obtained by averaging the corresponding data over 10 different structures. The vertical dashed line indicates the projection of the conditional crystal/glass boundary (interface plane) that existed at the initial filling stage (the plane corresponds to $z\approx 17\pm0.5$ Å).

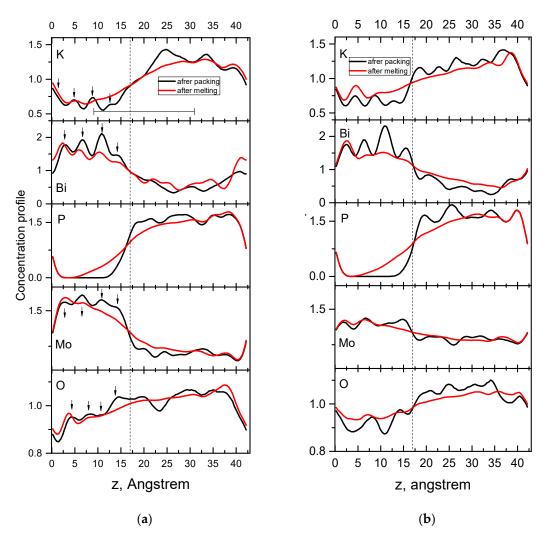


Figure 8. Concentration profiles of constituent elements of KBi(MoO₄)2@glass interphase cell calculated along the normal to initial boundary for glasses G5.69 (a) and G4.12 (b): after initial filling (black curves) and after the MD calculations (red curves). The initial crystal/glass boundary is marked with a vertical dashed line.

As Figure 8 shows, in the case of G5.69 glass (left part of the figure), the initial concentration profiles in the crystal region (black curves) are characterized by a set of peaks and dips with a clear periodicity along the z coordinate. This behavior of the profiles is an obvious consequence of the crystal atomic structure periodicity, namely, the existence of layers of potassium and bismuth ions

and molybdate groups in it. The local concentration maxima (indicated by vertical arrows in the figure) obviously correspond to position of the layer of a particular element kind.

After the MD calculation procedure, the crystal and glass atoms mix to form an interphase layer. The width (thickness) of the latter can be estimated as the range z, where the change in the concentration profiles "before" and "after" the MD procedure is significant. This assessment is most indicative of the phosphorus profile, since its concentration in the KBi(MoO₄)₂ crystal region is zero. As a result, according to the phosphorus profile, the interphase lies within the range of z from 7.5 to 22.5 Å. It can be seen that significant changes in the concentration of all other elements also occur in this range (the range is indicated by a horizontal line in the left top part of Figure 8). Thus, according to results of the calculations, the thickness of the interphase layer in the KBi(MoO₄)₂@G5.69 glass ceramics should be estimated at ~15 Å.

As the left part of Figure 8 shows, the concentration curves for K and Mo reveal significant mixing of these elements, and the peaks corresponding to the outermost layers of the crystal practically disappear after the MD procedure. In contrast, the situation for Bi is somewhat different: the Bi layers in the crystal region remain clearly visible in the concentration curve even after the MD procedure. This fact indicates that bismuth ions move relatively weakly (compared to other elements in the crystal) towards the glass.

Phosphorus atoms, on the contrary, diffuse from the glass to the crystal (corresponding concentration curve decreases monotonically in the glass-to-crystal direction). The change in the oxygen concentration curves, although also noticeable, lies within the statistical error of the averaging, since the oxygen content of the crystal and glass G5.69 was approximately the same at the stage of initial filling.

As the data in the right part of Figure 8 show, the concentration profiles of KBi(MoO₄)₂@G4.12 interphase are generally similar to those described above and lead to similar conclusions, in particular, regarding the thickness of the interfacial layer. However, the concentration profile of bismuth is an exception. As can be seen from the figure, in the case of the G4.12 glass, the three peaks of the Bi concentration curve in the KBi(MoO₄)₂ crystal, closest to the initial boundary, disappear after the MD procedure. This means that during formation of KBi(MoO₄)₂@G4.12 interphase, deeper layers of bismuth atoms of the crystal are disordered if compared with KBi(MoO₄)₂@G5.69 case. This is probably due to the higher (almost twice) content of molybdate groups in G4.12 glass compared to G5.69 one (see Table S1). Thus, the increase in the content of molybdate groups in the glass enhances the diffusion of bismuth cations from the KBi(MoO₄)₂ crystal to phosphate-molybdate glass.

This conclusion is consistent with the change in the composition of the interphase layer compared to the composition of the glass phase (see below). Concentration profiles of the KBi(MoO₄)₂@G4.12 interphase cell (Figure 8, right) show that in the *z-coordinate* region between ~10 and ~30 Å there is a significant redistribution of chemical element concentrations relative to the corresponding concentrations in the composite components - KBi(MoO₄)₂ crystal and G4.12 glass. Right this region with a thickness of ~20 Å can be considered the interphase layer of the composite system KBi(MoO₄)₂@G4.12.

Similar results regarding the thickness of the interphase layer were obtained for composite glass-ceramic systems with other compositions. Thus, according to the calculation results, the thickness of the interphase layers in the "KBi(MoO₄)2@(xP₂O₅-yMoO₃-zBi₂O₃-(1-x-y-z)K₂O) glass-ceramics can be estimated as ~1.5-2.0 nm.

By counting the number of atoms in the noted part of the interphase cell (between 10 and 30 Å along the z-axis), we found that the chemical composition of the interphase region approximates the formula K₁₂P₆Mo₁₀Bi₅O₆₁. If we compare this formula with composition of the glass component (K₁₈P₁₃Mo₅Bi₃O₆₁, see SM, Table S1), significant changes become evident. In particular, a twofold increase in the molybdenum content in the interphase is accompanied by a 1.67-fold increase in the bismuth content; the number of oxygen atoms does not change, and the potassium and phosphorus content decreases by 1.5 and 2.17 times, respectively. Thus, it can be argued that in the interphase region, the number of oxygen atoms near molybdenum and bismuth ions should decrease compared

to the case of the glass medium. In the case of interphase of KBi(MoO₄)2@G5.69 system, the DFT-based geometry optimizations were carried out with use of rather small cells (12x12x12 Å). These cells were filled with the same fragments as in the G5.69 glass case (see SM, Table S1), however with different amounts: (P_3O_9) - 2 units, (MoO_4) - 6 units, (Mo_2O_7) - 2 units, K - 12 units, and Bi - 5 units. The chosen ratio of fragments corresponds to the chemical composition of the interphase determined above - $(K_{12}P_6Mo_{10}Bi_5O_{61})$. The results of geometric optimization of structures are analyzed below.

3.2.4. Analysis of Structure of the Nearest Surrounding of Bismuth Cations

Using results of the structural geometry optimizations by the DFT method, we analyzed the composition and structure of the nearest oxygen surrounding of bismuth cations located in different spatial regions of KBi(MoO₄)2@G5.69 system. For this purpose, we estimated content of different types of BiO_n polyhedra (with coordination number n = 2, 3, 4, ..., 8) in the crystal, glass, and interphase of this glass-ceramic system using SHAPE utility program [70]. The analysis was applied to 8 geometry-optimized structures of G5.69 glass and KBi(MoO₄)2@G5.69 interphase with the subsequent statistical averaging of the results. For the case of the KBi(MoO₄)2 crystal, the analysis was performed using its crystal structure taken from literature [71]. The analysis took into account all oxygen atoms at a distance of less than 3.2 Å from bismuth atoms.

Results of analysis by SHAPE utility are summarized in Table 3. Among other structural characteristics, the Table contains the calculated parameters S, which characterize degree of deviation of the polyhedron from the ideal one. The value of S is calculated as a measure of deviation from the ideal polyhedron of a certain dimension, which is expressed through the standard deviation of the lengths of each of the B-O bonds. The parameter is normalized to 100% so that a value of S = 0 corresponds to a complete coincidence of the structure with an ideal polyhedron of a certain dimension and symmetry, and a value of S = 100% is not achieved in principle [72]. Thus, the smaller is the value of S, the closer the structure is to the ideal polyhedron of the corresponding dimension and symmetry.

Table 3. Characteristics of the oxygen surrounding of bismuth atoms in the crystal, glass and interphase of KBi(MoO₄)₂@G5.69 system.

Types of	KBi(MoO ₄) ₂ crystal		Glass G5.69		Interphase	
polyhedra	Content, %.	Δ S ¹	Content, %.	ΔS	Content, %.	•
BiO ₂ + BiO ₃	-	-	-	-	5.3	-
BiO ₄	-	-	4.2	0.697	15.8	0.674-7.591
BiO ₅	-	-	50.0	0.433-14.867	50.0	0.339-12.064
BiO_6	-	-	41.6	1.263-11.995	28.9	1.503-24.535
BiO ₇	-	-	4.2	2.974	-	-
BiO_8	100	0.481	-	-	-	-

 $^{^{1}}$ ΔS is the numerical range in which fall the S values calculated for all polyhedra of corresponding system.

For convenience, in calculations of the S parameter the following approximations were used. If Bi cation is surrounded by 5 oxygen atoms, then it is considered to be an octahedron with an oxygen vacancy - the symmetry of the position in such an ideal "Vacant octahedron" is C_{4v}. If Bi has 6 oxygens in its surrounding, then it is treated as "Octahedron" of symmetry O_h. Deviations from the octahedron were considered only for both 5- and 6-coordinated bismuth atoms.

Results of analysis show that in KBi(MoO₄)₂ crystal, all bismuth atoms are surrounded by 8 oxygen atoms, which form distorted triangular dodecahedrons (point group of symmetry of bismuth ion is D_{2d}). In contrast, for glass, there is a significant decrease in the coordination number of the oxygen environment of bismuth atoms. The BiO₈ polyhedra are absent there at all, and the most common form of oxygen surroundings are BiO₅ and BiO₆ polyhedra (see Table 3). In the case of interphase, the coordination of bismuth is even simpler: the number of polyhedra of BiO₆ decreases by about half compared to glass, the number of BiO₄ increases fourfold, and the simplest structures

of bismuth surrounding appear - BiO₃ and BiO₂. Degree of deviation from an ideal form of BiO₆ in the interphase (see the corresponding values of ΔS in the Table 3) is obviously higher than for such polyhedra in the glass.

Thus, according to the results of our analysis, the structure of the nearest oxygen surrounding of bismuth in the KBi(MoO₄)₂ crystal (only BiO₈ polyhedra) is fundamentally different from the corresponding structure of bismuth environment in glass (mainly BiO₅ and BiO₆ polyhedra), and in the interphase regions of glass ceramics. In the latter, BiO₂ and BiO₃ structures appear and the number of BiO₄ increases.

4. Concluding Remarks

It is worth to compare results of experimental studies of glass systems presented in the first part of this paper with results of computational studies of glass-ceramic systems presented in the second part. From this comparison, it can be seen that an increase in the bismuth oxide content in the glass, as well as formation of an interphase layer in glass ceramics, result in an increase in the number of bismuth ions coordinated by (bonded with) fewer oxygen atoms. Consequently, some peculiarities of spectral/luminescent properties should be common for both studied glasses and glass ceramics.

For real glass-ceramic systems with a sufficient total volume of interphase regions, changes in the quantitative composition of different types of bismuth polyhedra can be observed experimentally. Indeed, Figure 2 shows that in the 450-650 nm region, the reflection spectrum of GC5.69 glass-ceramics differs significantly from the reflection spectra of both G5.69 glass and the KBi(MoO₄)₂ crystal. This difference can be attributed to the influence of interphase regions.

The luminescence manifestations of interphase layers are more noticeable. So, in the excitation spectra of the "blue" PL of glass (Figure 4, curve 2), in contrast to the excitation spectrum of such PL in crystals (Figure 4, curve 1), the main band is observed at 367 nm (see Figure 3). The same band dominates the excitation spectrum of glass-ceramics (Figure 4, curve 3), but here the contribution of bands in the 300 - 350 and 350 - 425 nm regions, which are characteristic of crystals, increases significantly. We also attributed these differences to the appearance of interphase layers in the composition of GC5.69 glass ceramics.

Regarding the manifestation of these layers in the PL spectra of the glass ceramics, it should be emphasized that according to data in Table 1, the "yellow-red" PL band increases by a factor of 1.55 compared to the "blue" PL band in G4.12 \rightarrow G5.69 transition. The relative intensity of Iyr/IBlue is 1.73. Comparison of the PL spectra characteristics of G5.69 glass and GC5.69 glass ceramics shows that G5.69 \rightarrow GC5.69 transition is accompanied by an increase in the relative intensity of the "yellow-red" PL: Iyr/IBlue increases from 1.73 to 1.96. This increase in Iyr/IBlue, as can be seen from Figure 2, is mainly due to the longest-wavelength components of the "yellow-red" PL in the region of 700 – 850 nm. It has already been shown above (Part 1) that such changes in the spectra of glass are due to an increase in the contribution of bismuth ions with low charge and low coordination with oxygen atoms. Thus, the described luminescence manifestations confirm the influence of interphase layers on the properties of the studied glass-ceramic samples.

5. Conclusions

The properties of a set of glasses of the xP_2O_5 -yMoO₃-zBi₂O₃-(1-x-y-z)K₂O system (where, x – in 23.01 – 40.31; y - in 13.55 – 46.02; z - in 1.00 – 5.69 range) and based on them crystal@glass ceramics where the crystal component is luminescent crystalline particles KBi(MoO₄)₂ have been produced and studied by experimental and theoretical methods.

It has been shown that the optical absorption and photoluminescence characteristics of such systems are largely determined by content of Bi₂O₃ and MoO₃ oxides in the initial charge, and by content bismuth ions in different charge states which exist in the produced glass and glass ceramics.

In particular, the "blue" photoluminescence (spectral range 375 – 550 nm) of both glasses and glass ceramics originates from radiative transitions ${}^{3}P_{1} \rightarrow {}^{1}S_{0}$ in bismuth ions Bi³⁺. The "Yellow-red"

PL (spectral range 550 - 850 nm) should be mainly associated with the luminescence of bismuth ions in lower charge states – Bi^{2+} , Bi^{+} and Bi^{0} , and therefore it should be assumed that the increasing contribution of "yellow-red" PL to the total luminescence spectrum occurs due to an increase in the number of such lower-charge ions.

In glass-ceramics based on the xP_2O_5 -yMoO₃-zBi₂O₃-(1-x-y-z)K₂O glass matrix and crystals KBi(MoO₄)₂, interphase layers are formed between the glass and crystal phases, the thickness of which is estimated to be 1.5-2.0 nm. According to theoretical results, the structure and composition of the interphase layers differ from those of both the glass and crystal components. In particular, if compared to the case of glass, content of bismuth and molybdenum in the interphase layer increases by 1.7 and 2 times respectively, and the content of potassium and phosphorus decreases by 1.5 and 2.17 times respectively, whereas content of oxygen is almost unchanged.

This decrease in the number of oxygen atoms in the nearest surrounding of bismuth ions in the interphase region is accompanied by changes in the spectra of optical absorption and PL / PL excitation, which allows to use these changes for spectral probing of the formation and presence of interphase layers.

Supplementary Materials: The following supporting information can be downloaded at: Preprints.org, Procedure of atomic structure modeling, Figure S1: Periodic cell constructed for interphase of KBi(MoO₄)2@G5.69 system: initially filled and after MD calculations; Figure S2: An example 12x12x12 Å periodic cell of G5.69 glass, used in the DFT-based geometry optimizations; Table S1: Chemical composition, experimental and calculated densities of synthesized and modeled samples of the P₂O₅ - MoO₃-Bi₂O₃-K₂O glass system; Table S2: Parameters of interaction used in MD calculations.

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