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Article

Biased Random Process of Randomly Moving Particles with Fixed Mean and Curl

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Abstract: In this study, we investigate the behavior of randomly moving particles with a defined average energy, focusing on their biased motion around a specific point. These particles' velocity magnitudes follow a Maxwell distribution characterized by certain parameters. Our analysis derives the stochastic rotation rules for these particles and introduces the Itô equation conditions for biased stochastic rotation at a given curl value. We also consider the special relativistic-like effects in this context. The findings are validated through a representative example.

Keywords: randomly moving particles; biased motion; special relativistic-like effects; curl

1. Introduction

For a collection of particles with a defined average energy, their magnitudes of velocities follow a Maxwell distribution characterized by a specific parameter. When a subset of the particles exhibits biased motion in their movement direction, they can be characterized as three distinct and prevalent types of biased stochastic motion (At an extremely large particle count and scale, even scenarios with exceedingly low probabilities can manifest with notable frequency within a specific region). These biased motions encompass: 1) stochastic movement with a preferential direction while maintaining equal probabilities in other directions; 2) stochastic movement favoring a specific spatial point but equi-probable in alternative directions; and 3) stochastic movement with a predilection for revolving around a designated rotational center. Within these scenarios of biased stochastic motion, systems exhibiting characteristics of the first situation display the special relativistic-like effects [1,2]. On the other hand, biased stochastic motions with higher probabilities directed towards or revolving around a specific center incorporate these special relativistic-like effects into processes of spatial aggregation [3] or rotational phenomena [4].

For the rotational effect produced by biased randomly moving particles, the constituent particles of the rotating body inherently experience a displacement in linear velocity. Consequently, such motion should also exhibit special relativistic-like effects. Previous research [4] has investigated the rotational behavior generated by a group (radius is R) of randomly moving particles with consistent linear velocities c . The conclusion drawn is that the magnitude of the angular velocity produced by k randomly moving particles follows a Maxwell distribution with a scale parameter of $\frac{\sqrt{6}c}{3R\sqrt{k}}$. In actual physical systems, it is more common for these particles to not have perfectly consistent linear speeds. Instead, when the energy is fixed, their speeds follow a Maxwell distribution determined by a specific parameter. Furthermore, nearly every celestial body in the universe is rotating. The underlying mechanism for this rotation might be a result of the collective actions of particles whose speeds follow a Maxwell distribution. In summary, it is essential to thoroughly investigate the distribution of angular speeds generated by particles with speeds following a Maxwell distribution as they move randomly in space. Based on this premise, when the magnitude of the angular velocity is a specific value, particles undergoing pure random motion transition to biased random motion. Presenting the corresponding Itô equation expression under this scenario is the primary objective of this paper.

With the aid of mathematical packages, certain complex theoretical problems can be effortlessly addressed. Currently, among the world's most powerful symbolic computation packages are Maple (Cybernet Systems Co., Ltd., Japan) and Mathematica (Wolfram Research Inc., USA). The intricately designed algorithms within these platforms have facilitated the work of countless scientists and engineers worldwide. However, these mathematical packages are neither universal nor omnipotent. Real-world problems are often intricate, and relying solely on one specific package might not provide an optimal solution. It is essential to utilize a combination of software tools and appropriately decompose the actual problem to achieve a comprehensive resolution. In this study, Mathematica v13.2 was employed for theoretical derivations, while Maple v2023 was utilized for solving the system of partial differential equations.

This study first examines the distribution of rotational angular velocity magnitudes in a particle system with a defined average energy, where the magnitudes of the particle velocities follow a Maxwell distribution with a specific scale parameter. Subsequently, a set of Itô equations that describe the motion characteristics of biased moving particles with rotational effects is introduced. Finally, an illustrative example is provided to further elucidate the characteristics of these Itô equations. The biased moving particles discussed in this study, which exhibit special relativistic-like effects, differ from the conventional understanding where particle motion must follow the constraints of special relativity. This study assumes that these particles move at relatively low speeds, essentially unaffected by special relativistic effects, or they represent ideal particles completely unaffected by special relativistic effects. In this study, the average speed c represents the mean magnitude of particle velocities rather than the speed of light. These biased moving particles can exhibit common significant phenomena. Understanding their essence is crucial for comprehending the effects of special relativity and the inherent nature behind rotations influenced by special relativistic effects.

2. Results and Discussion

For particle assemblies with consistent magnitudes of velocity and directions arbitrarily dispersed within a three-dimensional domain, the resultant rotational effects in space have been comprehensively examined in our prior publication [4]. This study primarily investigates the rotational behavior in space exhibited by particle assemblies under identical conditions, where the norms of their velocities follow a Maxwell distribution. Analogous to the investigative approach presented in reference [4], we first examine the distribution of the norm $\|\Omega_S\|$ of the random angular velocity generated by a random vector V_S with its origin uniformly distributed on a unit sphere S , where the norm of V_S follows a Maxwell distribution with a mean of c (Figure 1). Assuming r is the radius of the unit sphere S , then

$$\omega_S = r \times v_S. \quad (1)$$

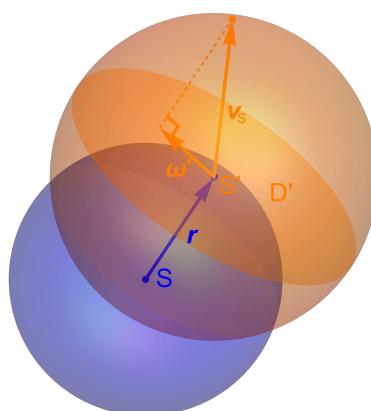


Figure 1. Schematic diagram of the generation method for the vector ω' . The radius $\|V_S\|$ of S' denotes the mean of norms of the spatial vectors that follow a Maxwell distribution with a parameter $\frac{1}{2} \sqrt{\frac{\pi}{2}} c$.

Let us establish a 3-dimensional Cartesian coordinate system for S' . Suppose that the disk D' is perpendicular to the z -axis and consider the random variables $\Theta \sim U(-1, 1)$ and $H \sim U(-\pi, \pi)$. Then, the norm of the random vector Ω' obtained by projecting the uniformly distributed points on sphere S' onto disc D' is

$$\|\Omega'\| = \|V_S\| \cdot \sqrt{1 - \Theta^2}. \quad (2)$$

Therefore, the probability of $\|\Omega'\|$ is

$$\|\omega'\|(x) = \begin{cases} \frac{8x e^{-\frac{4x^2}{\pi c^2}}}{\pi c^2}, & x \geq 0, \\ 0, & \text{otherwise.} \end{cases} \quad (3)$$

Assuming $F_X(x)$ is one coordinate X that is equivalent to the three coordinates of R . By taking a product $F_X(x) \cdot \|\omega'\|(x)$ of random variables, we can obtain the probability density of X , which represents one of the three equivalent coordinates of the angular velocity Ω_S contributed by the random vector V_S , specifically,

$$\omega_{S,X}(x) = \begin{cases} \frac{\operatorname{erfc}\left(\frac{2x}{c\sqrt{\pi}}\right)}{c}, & x \geq 0, \\ \frac{\operatorname{erfc}\left(\frac{2x}{c\sqrt{\pi}}\right) + 1}{c}, & \text{otherwise.} \end{cases} \quad (4)$$

The probability density of equivalent coordinate $\Omega_{B,X}$ of Ω_B contributed by V_B which is uniformly distributed throughout the whole unit ball enclosed by the sphere S is

$$\omega_{B,X}(x, r) = \begin{cases} \frac{r \operatorname{erfc}\left(\frac{2rx}{c\sqrt{\pi}}\right)}{c}, & x \geq 0, \\ \frac{r \left[\operatorname{erfc}\left(\frac{2rx}{c\sqrt{\pi}}\right) + 1 \right]}{c}, & \text{otherwise.} \end{cases} \quad (5)$$

Thus, the distribution function $\Omega_{B,X}(r)$ of the contribution of V_B to one of the equivalent coordinates of Ω_B is obtained. Next, $\Omega_{B,X}(r)$ is integrated over the whole unit ball and the probability density of the contribution of V_B in the whole unit ball to an equivalent coordinate X of the angular velocity Ω_B can be obtained:

$$\omega_{B,X}(x) = \begin{cases} \frac{3}{256cx^4} \left[(3\pi^2c^4 - 64x^4) \operatorname{erf}\left(\frac{2x}{c\sqrt{\pi}}\right) - 4cx(3\pi c^2 + 8x^2) e^{-\frac{4x^2}{\pi c^2} + 64x^4} \right], & x \geq 0, \\ \frac{3}{256cx^4} \left[(64x^4 - 3\pi^2c^4) \operatorname{erf}\left(\frac{2x}{c\sqrt{\pi}}\right) + 4cx(3\pi c^2 + 8x^2) e^{-\frac{4x^2}{\pi c^2} + 64x^4} \right], & \text{otherwise,} \end{cases} \quad (6)$$

Next, we extend the analysis to the case in which the radius of the ball has an arbitrary value R . The probability density of the contribution of the random vector V to the single equivalent coordinate X of angular velocity Ω is

$$\omega_X(x) = \begin{cases} \frac{3}{256} \left[\left(\frac{3\pi^2 c^3}{R^3 x^4} - \frac{64R}{c} \right) \operatorname{erf} \left(\frac{2Rx}{c\sqrt{\pi}} \right) - \frac{4(3\pi c^2 + 8R^2 X^2)}{R^2 x^3} e^{-\frac{4R^2 x^2}{c^2 \pi}} + \frac{64R}{c} \right], & x \geq 0, \\ \frac{3}{256} \left[\left(\frac{64R}{c} - \frac{3\pi^2 c^3}{R^3 x^4} \right) \operatorname{erf} \left(\frac{2Rx}{c\sqrt{\pi}} \right) + \frac{4(3\pi c^2 + 8R^2 X^2)}{R^2 x^3} e^{-\frac{4R^2 x^2}{c^2 \pi}} + \frac{64R}{c} \right], & \text{otherwise,} \end{cases} \quad (7)$$

the standard deviation of which is $\frac{c\sqrt{\pi}}{2R}$. Therefore, when k independent and identically distributed random vectors V move randomly in space, according to the central limit theorem (when they are grouped together), the norm $\|\Omega\|$ of the average angular velocity generated by all of their components relative to their total centroid follows a Maxwell distribution with scale parameter $\frac{c\sqrt{\pi}}{2R\sqrt{k}}$.

Building upon this foundation, we delve into the representation of this stochastic rotation behavior using the Itô equation, specifically discussing scenarios where the curl assumes a defined value of (a, b, c) , where $a, b, c \in \mathbb{R}$. For particle assemblies with a defined energy, the magnitudes of their velocities follow a Maxwell distribution characterized by specific parameters. Previous research has provided the form of the Itô equation for particles exhibiting biased random motion in a particular direction, incorporating the special relativistic-like effects [2]. In this study, we further incorporate the special relativistic-like effects into the stochastic rotation.

For any given particle, the biased motion around a specific point fundamentally remains a movement with a linear velocity. Consequently, it still adheres to the velocity reduction inherent of special relativistic-like. Assuming that at a specific time t , the linear velocity of a particle at position (x, y, z) is $u(t)$, its projections onto the three coordinate axes are $u_1(t), u_2(t)$ and $u_3(t)$. As per the description method in the literature [2], the Itô equation at time t can be expressed as:

$$\begin{cases} dx_1(t) = u_1(t)dt + \frac{\sqrt{c^2 - u(t)^2}}{c} \sigma dw_1(t), \\ dx_2(t) = u_2(t)dt + \frac{\sqrt{c^2 - u(t)^2}}{c} \sigma dw_2(t), \\ dx_3(t) = u_3(t)dt + \frac{\sqrt{c^2 - u(t)^2}}{c} \sigma dw_3(t), \end{cases} \quad (8)$$

where $u(t) = \sqrt{u_1(t)^2 + u_2(t)^2 + u_3(t)^2}$. Differing from the literature, here we replace u_1, u_2 and u_3 with $u_1(t), u_2(t)$ and $u_3(t)$. From a macroscopic perspective, if the system exhibits rotational effects, then at a given time t , the velocity u should also be a function of the position (x, y, z) . Then, the sum of integrals of the curl of the velocity over all possible closed loops (Γ s) is not zero, that is,

$$\sum \oint_{\Gamma} \nabla \times \mathbf{u} \neq 0 \quad (9)$$

We will not delve extensively into this issue here; instead, we will examine the scenarios in the microscopic domain when particles show indications of rotation. At a given time t , the velocity u should also be a function of the position (x, y, z) . Then, the velocity components along the three axes

should satisfy the condition that a non-zero curl, meaning $u_1(x_1, x_2, x_3)$, $u_2(x_1, x_2, x_3)$ and $u_3(x_1, x_2, x_3)$ are subject to the following constraints:

$$\left(\frac{\partial u_3}{\partial x_2} - \frac{\partial u_2}{\partial x_3} \right) \mathbf{i} + \left(\frac{\partial u_1}{\partial x_3} - \frac{\partial u_3}{\partial x_1} \right) \mathbf{j} + \left(\frac{\partial u_2}{\partial x_1} - \frac{\partial u_1}{\partial x_2} \right) \mathbf{k} \neq 0 \quad (10)$$

We will focus on the form of the Itô equation when the curl is specified.

To elucidate this problem further, we present a concrete example. Assuming these particles rotate within the plane $(x_1, x_2, 0)$ around the axis $(0, 0, x_3)$, resulting in a curl value of $(0, 0, \omega)$, and the magnitude of the particle's linear velocity remains constant at u . We proceed to solve the following equation:

$$\begin{cases} \frac{\partial u_3}{\partial x_2} - \frac{\partial u_2}{\partial x_3} = 0, \\ \frac{\partial u_1}{\partial x_3} - \frac{\partial u_3}{\partial x_1} = 0, \\ \frac{\partial u_2}{\partial x_1} - \frac{\partial u_1}{\partial x_2} = \omega, \\ u_1^2 + u_2^2 + u_3^2 = u. \end{cases} \quad (11)$$

Subsequently, we select a particular solution from the aforementioned equation to illustrate the issue:

$$\begin{cases} u_1 = -\sqrt{u^2 - \omega^2 x_1^2}, \\ u_2 = \omega x_1, \\ u_3 = 0. \end{cases} \quad (12)$$

Substituting Equation (12) into Equation (8) and omitting the time term, we obtain

$$\begin{cases} dx_1(t) = -\sqrt{u^2 - \omega^2 x_1^2} dt + \frac{\sqrt{c^2 - u^2}}{c} \sigma dw_1(t), \\ dx_2(t) = \omega x_1 dt + \frac{\sqrt{c^2 - u^2}}{c} \sigma dw_2(t), \\ dx_3(t) = \frac{\sqrt{c^2 - u^2}}{c} \sigma dw_3(t). \end{cases} \quad (13)$$

We substitute the values $\omega = 0.1$, $c = 100$, $u = 60$, and $\sigma = 10$ into Equation (13) and simulate with a sample of 100 particles. The results are depicted in Figure 2:

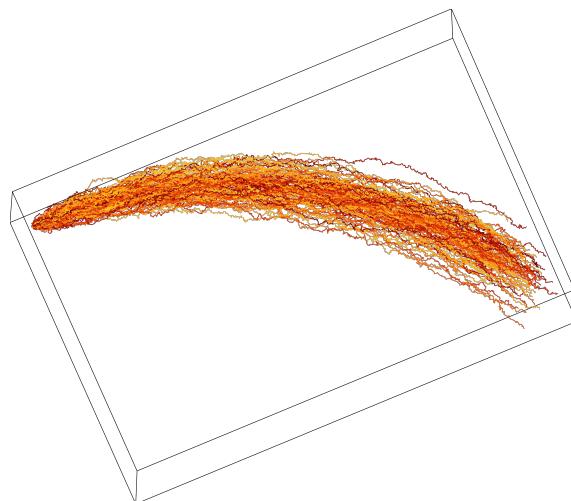


Figure 2. Simulation results of the trajectories of 100 particles randomly moving for 10 s ($\omega = 0.1$, $c = 100$, $u = 60$ and $\sigma = 10$).

In Figure 2, although only the motion of 100 particles starting from the same point is depicted, it is not difficult to infer the motion scenarios at various points when the starting positions are distributed across the entire plane. We will not elaborate further with additional examples here.

3. Conclusions

For a set of k particles with a defined average energy, when their velocity magnitudes follow a Maxwell distribution characterized by the parameter $\frac{1}{2}\sqrt{\frac{\pi}{2}}c$ (average speed is c), their random motion can lead to angular velocity magnitudes that follow a Maxwell distribution with the parameter $\frac{c\sqrt{\pi}}{2R\sqrt{k}}$. Specifically, when the curl of the particles maintains a certain value, their motion must satisfy the biased conditions defined by Equations (8) and (10). Notably, akin to the scenario where the particle group moves linearly, these rotation effects are also constrained by the special relativistic-like effects.

Supplementary Materials: The following supporting information can be downloaded at the website of this paper posted on [Preprints.org](#). Mathematica code for necessary calculation process and graphics.

Data Availability Statement: All data generated or analysed during this study are included in this published article and its Supplementary Information files.

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