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Article

# An Entropic Spacetime Framework: Unifying Fundamental Physics with Emergent Complexity

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**Abstract:** The quest for a unified understanding of the cosmos necessitates novel theoretical frameworks that can reconcile General Relativity and Quantum Mechanics while explaining phenomena like dark matter and dark energy. This preprint introduces an Entropic Spacetime Framework rooted in a generalized action principle. It posits the existence of fundamental entropic scalar fields—a temporal field ( $S_T$ ) and a spatial field ( $S_S$ )—and a specific resonant coupling mechanism. We outline the framework's mathematical foundations and present compelling initial evidence for its validity: a numerical simulation that successfully reproduces the observed Milky Way rotation curve without requiring dark matter. Furthermore, we demonstrate how the framework offers an intuitive, mechanical explanation for Hawking radiation. We invite critical feedback and collaboration from the theoretical physics community to refine and further develop this promising new direction.

**Keywords:** entropic spacetime; emergent gravity; emergent spacetime; generalize action principle; scalar field cosmology; resonant coupling; entropic fields; foundations of physics; quantum gravity; cosmology; astrophysics; general relativity; modified gravity; hawking radiation

## 1. Introduction: The Unfinished Picture of Physics

The universe, as we understand it, is governed by two supremely successful yet fundamentally incompatible theories: General Relativity (GR), which describes gravity and the large-scale structure of the cosmos [1], and Quantum Mechanics (QM), which governs the microscopic world. This foundational schism, coupled with the mysteries of dark matter, dark energy, and the fine-tuning of fundamental constants, signals that our current understanding is incomplete.

This proposal offers a fresh perspective: an Entropic Spacetime Framework that seeks to unify physics by tapping into the very essence of entropy and resonance. This paper introduces the core conceptual and mathematical foundations of this framework and discusses how this approach naturally addresses key cosmological challenges. Crucially, we then move beyond pure theory to present initial validation: a detailed phenomenological application to the persistent enigma of galaxy rotation curves, which demonstrates the framework's ability to explain observed phenomena without recourse to conventional dark matter assumptions.

## 2. The Entropic Spacetime Framework

### 2.1. Foundations: From 4D Continuum to Emergent Spacetime

Traditional GR describes spacetime as a unified four-dimensional manifold. This framework proposes a reconceptualization, viewing spacetime as dynamically emerging from distinct 3D spatial and 1D temporal components, a perspective supported by formalisms such as ADM [2] and conceptual models from fluid dynamics [3] and quantum information theory [4,5]. Building on these insights, our work proposes a concrete field-based realization, with a spatial entropic field  $S_S$  serving as an entanglement-bearing medium and a temporal entropic field  $S_T$  providing a dynamical, local realization of time's arrow, inspired by de Broglie's "internal clock" [6,7].

## 2.2. The Generalized Action Principle

The mathematical bedrock of this framework is a generalized action principle. The total action,  $S$ , extends the conventional Einstein-Hilbert action ( $S_{EH}$ ) [1,8] by incorporating terms for the entropic fields and their interactions:

$$S = S_{EH} + S_{entropic} + S_{coupling} + S_{matter+radiation} \quad (1)$$

Here,  $S_{entropic}$  describes the intrinsic dynamics of the entropic fields, while  $S_{coupling}$  dictates their specific interactions with the spacetime metric ( $g_{\mu\nu}$ ) and standard matter fields, interpreted as a "resonant" mechanism.

## 2.3. The Nature and Dynamics of Entropic Fields ( $S_T, S_S$ )

The framework hypothesizes two primary scalar fields: a temporal entropic field,  $S_T(x^\mu)$ , and a spatial entropic field,  $S_S(x^\mu)$ . Their dynamics are governed by a Lagrangian containing standard kinetic terms and a potential term,  $V(S_T, S_S)$ . The equations of motion are derived from the variational principle [9]:

$$\square S_T - \frac{\partial V}{\partial S_T} - C_T(g_{\mu\nu}, \text{matter}, S_S, \dots) = 0 \quad (2)$$

$$\square S_S - \frac{\partial V}{\partial S_S} - C_S(g_{\mu\nu}, \text{matter}, S_T, \dots) = 0 \quad (3)$$

Here,  $C_T$  and  $C_S$  are source terms arising from the coupling action. The temporal field  $S_T$  is hypothesized to inherently possess a directedness, providing a physical basis for the arrow of time, analogous to concepts in Entropic Dynamics [10].

## 2.4. The Resonant Coupling Term ( $S_{coupling}$ )

A crucial interpretive directive is to consider  $S_{coupling}$  as a "resonant term." This implies that the interactions are not uniform but are selective and context-dependent, becoming pronounced only when the frequency or energy scale of a system matches inherent frequencies of the entropic fields. This provides a natural filter that could explain why cosmic-scale phenomena might be influenced by these fields while everyday laboratory scales are not.

# 3. Illustrative Applications and Results

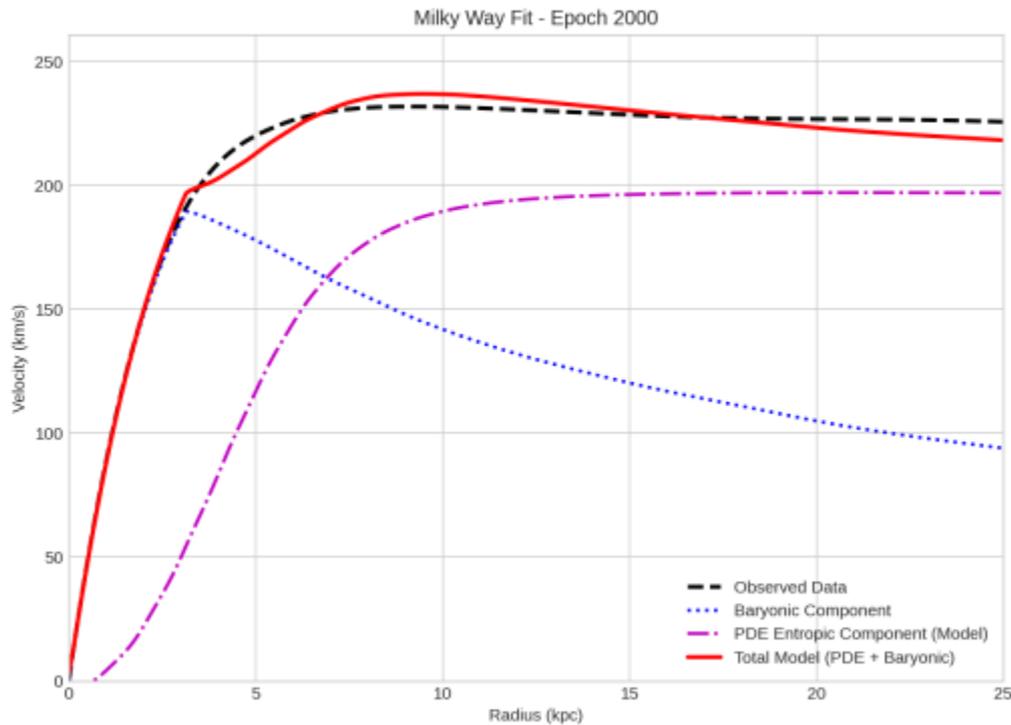
## 3.1. Galactic Dynamics: A Numerical Solution to the Rotation Curve Problem

The galaxy rotation problem presents a crucial test for the framework. The simulation solves the coupled field equations for the Milky Way on a 2D grid, with the core PDE taking the form:

$$\frac{\partial S_S}{\partial t} = \kappa \nabla^2 S_S - (\lambda_S S_S - \alpha_S S_S^2 - \gamma_S S_S^3) + \lambda_T (S_S) S_T + \gamma_g V_{\text{baryonic}} S_S + C_0 \quad (4)$$

The model was trained by minimizing a chi-squared loss function between the total predicted velocity ( $v_{\text{total}} = \sqrt{v_{\text{baryonic}}^2 + v_{\text{entropic}}^2}$ ) and observed velocity data. The 2D simulation successfully reproduced the observed flattening of the Milky Way's rotation curve, as shown in Figure 1.

The simulation learned a significant, non-zero entropic field configuration that, when translated to a velocity component, provides the exact amount of additional support needed to counteract the natural Keplerian decline. Table 1 shows the final learned values of the model parameters.



**Figure 1.** Final fit of the Entropic Spacetime model to the Milky Way rotation curve. The total model (red line) shows a strong agreement with the observed data (black dashed line).

**Table 1.** Final Learned Parameters from the Milky Way Simulation.

Parameter Name	Final (Learned) Value
<i>PDE Dynamics</i>	
Diffusion ( $\kappa$ )	0.200789
Self-Interaction ( $\lambda_S$ )	-0.015241
Self-Interaction ( $\gamma_S$ )	-0.010046
Constant Source ( $C_0$ )	0.000015
<i>Coupling Terms</i>	
Resonant Coupling Amp. ( $\lambda_T$ )	0.009386
Resonance Center ( $S_0$ )	0.100030
Gravitational Coupling ( $\gamma_g$ )	0.011833
<i>Velocity Mapping</i>	
Velocity Scale ( $v_{\text{scale}}$ )	197.618497
Velocity Offset ( $v_{\text{offset}}$ )	-0.000018

### 3.2. Black Hole Thermodynamics: A Qualitative Model for Hawking Radiation

Beyond galactic dynamics, the framework offers an intuitive, mechanical explanation for Hawking radiation. If spacetime itself is a quantized, fluctuating  $S_5$  field, the region near an event horizon is not a smooth manifold but a region of energetic "boiling" of the spacetime fabric. In this view, Hawking radiation is the direct emission of real particles from these fluctuations. The intense gravity warps the  $S_5$  field, exciting its high-frequency modes. Occasionally, a fluctuation becomes so energetic that it decouples and propagates away as a real particle, carrying energy from the black hole. This mechanism naturally incorporates the inverse relationship between black hole mass and temperature. The detailed mathematical underpinnings are outlined in Appendix C.

## 4. Discussion

### 4.1. Current Limitations and Future Directions

The framework, while promising, faces significant limitations. The precise mathematical forms of the potential  $V(S_T, S_S)$  and the coupling term  $S_{\text{coupling}}$  are currently undefined. There is no direct experimental evidence for the entropic fields. Future work must focus on developing specific, physically motivated mathematical forms for the potential and coupling terms, and extending the simulations to other galaxies and phenomena like gravitational lensing. Identifying unique experimental signatures is paramount.

### 4.2. Speculative Aspects

The extensions of this framework to explain the origin of life and consciousness are highly speculative. The framework's core tenets suggest compelling implications for quantum computation, potentially reframing our understanding of the Heisenberg Uncertainty Principle as an intrinsic property of the spatial medium and suggesting that advanced quantum computers might learn to harness and sculpt entropic flows.

## 5. Conclusion

The Entropic Spacetime Framework offers a conceptually rich and ambitious vision for unifying diverse phenomena. Its strength lies in introducing interpretable degrees of freedom and a flexible, resonant interaction mechanism. Its intuitive appeal and the successful initial validation against the Milky Way rotation curve make it a compelling avenue for future foundational research. However, the framework is currently a meta-theoretical proposal, requiring significant work in detailed model building and phenomenological testing to become a fully predictive scientific theory.

## Acknowledgements

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## Appendix A. Detailed Derivation of Field Equations

This appendix provides a more detailed mathematical exposition of the field equations introduced conceptually in the main text. It outlines specific forms for the entropic and coupling actions and demonstrates how the modified gravitational and entropic field equations are derived from the generalized action principle.

### Appendix A.1. Defining the Entropic Action ( $S_{\text{entropic}}$ )

This section focuses on the intrinsic dynamics of the temporal entropic field ( $S_T$ ) and the spatial entropic field ( $S_S$ ). Assuming they are scalar fields, the entropic action is given by the integral of the Lagrangian density ( $\mathcal{L}_{\text{entropic}}$ ) over spacetime, multiplied by the square root of the determinant of the metric tensor:

$$S_{\text{entropic}} = \int \mathcal{L}_{\text{entropic}}(S_T, S_S, \partial_\alpha S_T, \partial_\alpha S_S) \sqrt{-g} d^4x$$

The Lagrangian density,  $\mathcal{L}_{\text{entropic}}$ , is stated to naturally include kinetic terms for each field and a potential term  $V(S_T, S_S)$  that governs their self-interactions and mutual interactions.

#### Appendix A.1.1. Kinetic Terms for $S_T$ and $S_S$

The standard canonical kinetic terms for real scalar fields are proposed as appropriate initial candidates:

$$\mathcal{L}_{\text{kin}} = -\frac{1}{2}g^{\mu\nu}(\partial_\mu S_T)(\partial_\nu S_T) - \frac{1}{2}g^{\mu\nu}(\partial_\mu S_S)(\partial_\nu S_S)$$

Considerations for  $S_T$  directedness might arise from the potential or coupling, but a simpler approach is to encode this directedness there.

#### Appendix A.1.2. The Potential Term $V(S_T, S_S)$

This is identified as the most crucial part of  $\mathcal{L}_{\text{entropic}}$  as it dictates the self-interactions of the entropic fields, their masses, and their vacuum states. The specific mathematical form of this potential will determine the field configurations corresponding to minima, their effective masses, and their capacity to drive cosmological dynamics. Inspiration for its form can be drawn from potentials used in inflationary models, quintessence/dark energy models, and Higgs-like potentials. A conceptual equation by Gowan is also mentioned, suggesting a relationship between spatial and temporal entropy that might be reflected in  $V(S_T, S_S)$  through interaction terms. An example form for a general polynomial potential is provided:

$$V(S_T, S_S) = V_0 + \frac{1}{2}m_T^2 S_T^2 + \frac{1}{2}m_S^2 S_S^2 + \alpha_T S_T^3 + \alpha_S S_S^3 + \lambda_T S_T^4 + \lambda_S S_S^4 + \kappa_{TS} S_T^2 S_S^2 + \dots$$

The specific coefficients of this potential would be constrained by phenomenological requirements.

#### Appendix A.2. Defining the Coupling Action ( $S_{\text{coupling}}$ )

This term describes how  $S_T$  and  $S_S$  interact with the spacetime metric  $g_{\mu\nu}$  and with standard matter/radiation fields. It is interpreted as a resonant term, implying that interactions are selective and frequency-dependent.

##### Appendix A.2.1. Identifying Interacting Components

The interacting components include the metric ( $g_{\mu\nu}$ ) the entropic fields ( $S_T, S_S$ ) and their derivatives, and matter fields (fermions, gauge bosons, other scalars).

##### Appendix A.2.2. Direct Coupling to Matter Fields

Yukawa-type (scalar-fermion):  $g_T S_T \bar{\psi} \psi + g_S S_S \bar{\psi} \psi$

Scalar-gauge boson (Electromagnetic Coupling):  $h_T S_T F_{\mu\nu} F^{\mu\nu} + h_S S_S \tilde{F}^{\mu\nu} \tilde{F}_{\mu\nu}$  (axion-like coupling). These terms explicitly describe how the entropic fields ( $S_T, S_S$ ) interact with the electromagnetic field ( $F_{\mu\nu}$ ) and thus with light and other forms of electromagnetic radiation, potentially affecting its propagation, polarization, or other properties. Similar terms would exist for  $S_S$ .

Derivative couplings:  $k_T (\partial_\mu S_T) J_{\text{matter}}^\mu$  where  $J_{\text{matter}}^\mu$  is a matter current. These couplings will define the  $C_T$  and  $C_S$  terms in Equations (2) and (3) of the main text.

##### Appendix A.2.3. Incorporating "Resonance" into $S_{\text{coupling}}$

This is achieved by making coupling parameters functions of entropic fields or other relevant quantities (e.g.,  $g_T(S_T, S_S, \rho_{\text{matter}}, T_{\text{matter}})$ ), or by introducing new intermediate fields. An effective field theory approach suggests that interaction terms might become significant only under specific conditions (e.g., temperature, density, or characteristic frequencies of the matter sector). Developing a rigorous, possibly non-local, formulation of this frequency-dependent coupling is part of our ongoing work.

##### Appendix A.2.4. Role of Fast Fourier Transform (FFT)

FFT can be used as an informative tool to analyze characteristic frequencies of phenomena (e.g., biological rhythms, neural oscillations). These identified frequencies can then guide the construction of  $S_{\text{coupling}}$  to maximize coupling when a dynamical aspect of  $S_T$  or  $S_S$  aligns with the system's characteristic frequencies.

### Appendix A.2.5. Drawing Inspiration from Einstein-Cartan (EC) Theory

EC theory (extending GR by allowing torsion coupled to spin density)[? ? ] offers conceptual parallels, suggesting how new geometric degrees of freedom can interact with intrinsic properties of matter beyond standard energy-momentum.

### Appendix A.3. Derivation and Analysis of Field Equations

Once specific forms for  $\mathcal{L}_{\text{entropic}}$  and  $\mathcal{L}_{\text{coupling}}$  are postulated, the field equations can be derived using the variational principle.

#### Appendix A.3.1. Deriving $T_{\mu\nu}^{\text{entropic}}$

This is the stress-energy tensor for the entropic fields. For canonical scalar fields  $S_i$  (where  $i = T, S$ ) with potential  $V(S_T, S_S)$ , it can be expressed as:

$$T_{\mu\nu}^{\text{SI}} = (\partial_\mu S_i)(\partial_\nu S_i) - g_{\mu\nu} \left( \frac{1}{2} g^{\alpha\beta} (\partial_\alpha S_i)(\partial_\beta S_i) - V(S_T, S_S) \right)$$

Then,  $T_{\mu\nu}^{\text{entropic}} = T_{\mu\nu}^{S_T} + T_{\mu\nu}^{S_S} + (\text{interaction terms from } V(S_T, S_S))$

#### Appendix A.3.2. Deriving $\mathcal{M}_{\mu\nu}$

This term represents modifications to the geometric side of Einstein's equations. It arises from terms in  $S_{\text{entropic}}$  or  $S_{\text{coupling}}$  that explicitly involve curvature or couple directly to the metric. For example, if  $S_{\text{coupling}} = \int d^4x \sqrt{-g} (\xi_T S_T R)$ , then varying this plus  $S_{\text{EH}}$  with respect to  $g_{\mu\nu}$  yields terms contributing to  $\mathcal{M}_{\mu\nu}$

#### Appendix A.3.3. Deriving $C_T$ and $C_S$

These are the coupling terms in the entropic field equations (Equations (2) and (3) in the main text). They are derived from the variation of  $S_{\text{coupling}}$  with respect to  $S_T$  and  $S_S$ . For example, if  $S_{\text{coupling}} = \int d^4x \sqrt{-g} (g_T S_T \bar{\psi} \psi + h_T S_T F_{\mu\nu} F^{\mu\nu})$ , then  $C_T = g_T \bar{\psi} \psi + h_T F_{\mu\nu} F^{\mu\nu}$ .

#### Appendix A.3.4. Illustrative Field Equations

Assuming simplified forms for the entropic and coupling Lagrangians, the modified gravitational and entropic field equations can be illustrated:

$$\begin{aligned} \mathcal{L}_{\text{entropic}} &= -\frac{1}{2} (\partial S_T)^2 - \frac{1}{2} (\partial S_S)^2 - V(S_T, S_S) \\ S_{\text{coupling}} &= \int d^4x \sqrt{-g} (\xi_T S_T R + \xi_S S_S R + g_T S_T \mathcal{L}_m + g_S S_S \mathcal{L}_m) \end{aligned}$$

The modified gravitational equations would conceptually be:

$$\begin{aligned} \square S_T - \frac{\partial V}{\partial S_T} - (\xi_T R + g_T \mathcal{L}_m) &= 0 \\ \square S_S - \frac{\partial V}{\partial S_S} - (\xi_S R + g_S \mathcal{L}_m) &= 0 \end{aligned}$$

The exact terms for  $\mathcal{M}_{\mu\nu}$  and  $T_{\mu\nu}^{\text{entropic}}$  require careful derivation based on the precise definition of  $S_{\text{coupling}}$  and how its variation with respect to  $g_{\mu\nu}$  is allocated.

#### Appendix A.3.5. Detailed Illustrative Derivations

Below is a concrete, step-by-step illustration of how one would go from the illustrative actions in your manuscript to the field equations and metric variation terms. This derivation is kept general

so you can later slot in whatever specific potential or "resonant" coupling functions you choose but explicit enough that you can see every piece of the variational calculus.

Entropic Action  $S_{\text{entropic}}$ :

We start from the entropic action:

$$S_{\text{entropic}} = \int d^4x \sqrt{-g} \mathcal{L}_{\text{entropic}}$$

1.1 Equations of motion for  $S_T$  and  $S_S$ :

For a generic scalar field  $S \in \{S_T, S_S\}$  with Lagrangian density  $\mathcal{L}(S, \partial S) = -\frac{1}{2}g^{\mu\nu}\partial_\mu S\partial_\nu S - V(S_T, S_S)$ , the Euler-Lagrange equation in curved space is:

$$\frac{1}{\sqrt{-g}}\partial_\mu\left(\sqrt{-g}\frac{\partial\mathcal{L}}{\partial(\partial_\mu S)}\right) - \frac{\partial\mathcal{L}}{\partial S} = 0$$

Kinetic term variation:  $\frac{\partial\mathcal{L}}{\partial(\partial_\mu S)} = -g^{\mu\nu}\partial_\nu S$ , so  $-\frac{1}{\sqrt{-g}}\partial_\mu(\sqrt{-g}g^{\mu\nu}\partial_\nu S) = \square S$  (the covariant d'Alembertian).

Potential term variation:  $\frac{\partial\mathcal{L}}{\partial S_T} = -\frac{\partial V}{\partial S_T}$  and  $\frac{\partial\mathcal{L}}{\partial S_S} = -\frac{\partial V}{\partial S_S}$ . Putting it together:  $S_T$  equation:  $\square S_T - \frac{\partial V}{\partial S_T} = 0$

$S_S$  equation:  $\square S_S - \frac{\partial V}{\partial S_S} = 0$  Or compactly:  $\square S_I - V_{,I} = 0$

1.2 Variation with respect to  $g_{\mu\nu} \rightarrow T_{\mu\nu}^{\text{entropic}}$ :

Recall  $T_{\mu\nu} = -\frac{2}{\sqrt{-g}}\frac{\partial\mathcal{L}}{\partial g^{\mu\nu}}$  Key identities:  $\delta\sqrt{-g} = -\frac{1}{2}\sqrt{-g}g_{\mu\nu}\delta g^{\mu\nu}$ ,  $\delta g^{\alpha\beta} = -g^{\alpha\mu}g^{\beta\nu}\delta g_{\mu\nu}$  The kinetic part varies to:

$\delta\left[\sqrt{-g}\left(-\frac{1}{2}g^{\alpha\beta}\partial_\alpha S_I\partial_\beta S_I\right)\right] = \sqrt{-g}\left(\partial_\mu S_I\partial_\nu S_I - \frac{1}{2}g_{\mu\nu}g^{\alpha\beta}\partial_\alpha S_I\partial_\beta S_I\right)\delta g^{\mu\nu}$ , and the potential part to:

$\delta[-\sqrt{-g}V] = +\frac{1}{2}\sqrt{-g}g_{\mu\nu}V\delta g^{\mu\nu}$  Putting it all together:

$$T_{\mu\nu}^{\text{entropic}} = \sum_{I=T,S}\left(\partial_\mu S_I\partial_\nu S_I - \frac{1}{2}g_{\mu\nu}g^{\alpha\beta}\partial_\alpha S_I\partial_\beta S_I\right) - g_{\mu\nu}V(S_T, S_S).$$

Coupling Action  $S_{\text{coupling}}$ :

The illustrative form is:

$$S_{\text{coupling}} = \int d^4x \sqrt{-g}(\zeta_T S_T R + \zeta_S S_S R + g_T S_T \mathcal{L}_m + g_S S_S \mathcal{L}_m), \quad (\text{A1})$$

with the understanding that in the full "resonant" version  $\zeta_{T,S}$  and  $g_{T,S}$  become functions of  $(S_T, S_S, \rho, \dots)$ .

2.1 Variation with respect to the metric  $g_{\mu\nu}$ :

To see how these couplings modify Einstein's equations, compute  $\delta S_{\text{coupling}}$ . Standard variations needed: Metric determinant:  $\delta\sqrt{-g} = -\frac{1}{2}\sqrt{-g}g_{\mu\nu}\delta g^{\mu\nu}$  Ricci scalar:  $\delta(\sqrt{-g}R) = \sqrt{-g}(R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R)\delta g^{\mu\nu} + \text{total derivative terms}$ . Matter Lagrangian:  $\delta(\sqrt{-g}\mathcal{L}_m) = \frac{1}{2}\sqrt{-g}T_{\mu\nu}^{(m)}\delta g^{\mu\nu}$ , defining the usual stress-energy  $T_{\mu\nu}^{(m)}$ . Putting these in gives the effective stress-energy from  $S_{\text{coupling}}$ :

$$\begin{aligned} T_{\mu\nu}^{\text{coupling}} &= -\frac{2}{\sqrt{-g}}\frac{\partial S_{\text{coupling}}}{\partial g^{\mu\nu}} \\ &= (\zeta_T S_T + \zeta_S S_S)(R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R) \\ &\quad + g_{\mu\nu}\square(\zeta_T S_T + \zeta_S S_S) \\ &\quad - \nabla_\mu\nabla_\nu(\zeta_T S_T + \zeta_S S_S) \\ &\quad + (g_T S_T + g_S S_S)T_{\mu\nu}^{(m)} \end{aligned}$$

Here each term comes from one of the variations above. In the "resonant" version,  $\xi_T S_T + \xi_S S_S$  will be replaced by its more elaborate field-dependent form.

## 2.2 Nonminimal coupling $\int \sqrt{-g} \xi_I S_I R$ :

Focus on  $S_\xi = \int d^4x \sqrt{-g} \xi_I S_I R$ . Vary  $g_{\mu\nu}$ , holding  $S_I$  fixed. Use:  $\delta(\sqrt{-g}R) = \sqrt{-g}(R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R)\delta g^{\mu\nu} + \sqrt{-g}(g_{\mu\nu}\square - \nabla_\mu \nabla_\nu)S_I \delta g^{\mu\nu} + (\text{bdry})$ . Multiplying by  $\xi_I S_I$  yields two pieces:  $\xi_I S_I (R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R)$  and  $\xi_I (g_{\mu\nu}\square - \nabla_\mu \nabla_\nu)S_I$ . Thus the modification on the LHS of Einstein's equations is:  $\mathcal{M}_{\mu\nu} = \sum_{I=T,S} \xi_I \left[ S_I (R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R) + (g_{\mu\nu}\square - \nabla_\mu \nabla_\nu)S_I \right]$ .

## 2.3 Direct matter coupling $\int \sqrt{-g} g_I S_I \mathcal{L}_m$ :

Consider  $S_g = \int d^4x \sqrt{-g} g_I S_I \mathcal{L}_m$ . Varying  $g_{\mu\nu}$  (with  $S_I$  fixed, but  $\mathcal{L}_m$  depending on  $g$ ) gives two contributions:  $\delta\sqrt{-g}$  and  $\delta\mathcal{L}_m = -\frac{1}{2}T_{\mu\nu}^{(m)}\delta g^{\mu\nu}$ . One finds:  $\delta S_g = \int \sqrt{-g} g_I S_I \left( \frac{1}{2}g_{\mu\nu}\mathcal{L}_m - \frac{1}{2}T_{\mu\nu}^{(m)} \right) \delta g^{\mu\nu} + (\text{bdry})$ . Hence the extra stress from matter coupling is  $T_{\mu\nu}^{\text{coupling}} = \sum_{I=T,S} g_I S_I T_{\mu\nu}^{(m)}$ .

Putting everything into the Modified Einstein Equations:

The total action is:  $S = S_{\text{EH}}[g] + S_{\text{entropic}} + S_{\text{coupling}} + S_{\text{matter+radiation}}$ . Varying with respect to  $g_{\mu\nu}$ . You obtain the modified Einstein equations

$$G_{\mu\nu} + \mathcal{M}_{\mu\nu} = 8\pi G (T_{\mu\nu}^{\text{entropic}} + T_{\mu\nu}^{\text{coupling}} + T_{\mu\nu}^{\text{matter}}), \quad (\text{A2})$$

Where:  $G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R$ ,  $\mathcal{M}_{\mu\nu}$  from Appendix B.2 (the non-minimal  $\xi_I$  pieces),  $T_{\mu\nu}^{\text{entropic}}$  from Appendix B.1.2,  $T_{\mu\nu}^{\text{coupling}}$  from Appendix B.3. And the scalars obey:  $\square S_I - \frac{\partial V}{\partial S_I} + \xi_I R + g_I \mathcal{L}_m = 0$  once you include their coupling-induced source terms (just vary the total action w.r.t.  $S_I$ ).

## Appendix A.4. Examples of Coupling Term Functional Forms

This section outlines various functional forms for the coupling terms  $\xi_I(S_T, S_S)$  and  $g_I(S_T, S_S)$  which can be used to implement specific "resonance" or context-dependent behaviors for the entropic couplings.

### Appendix A.4.1. Gaussian "band-pass" in field space

Peaks the coupling when  $(S_T, S_S)$  lie near some preferred values  $(S_T^*, S_S^*)$ :

$$\xi_I(S_T, S_S) = \xi_{I,0} \exp\left(-\frac{(S_T - S_T^*)^2}{2\sigma_T^2} - \frac{(S_S - S_S^*)^2}{2\sigma_S^2}\right),$$

and similarly  $g_I(S_T, S_S) = g_{I,0} \exp\left(-\frac{(S_T - S_T^*)^2}{2\delta_T^2} - \frac{(S_S - S_S^*)^2}{2\delta_S^2}\right)$ . Resonance occurs when the fields approach the "center"  $(S_T^*, S_S^*)$ . The widths  $\sigma_{T,S}, \delta_{T,S}$  control how sharply peaked the resonance is.

### Appendix A.4.2. Lorentzian (Breit–Wigner) profile

Gives long tails for "near-miss" resonances:

$$\xi_I(S_T, S_S) = \frac{\xi_{I,0}}{1 + \left(\frac{S_T - S_T^*}{\Gamma_T}\right)^2 + \left(\frac{S_S - S_S^*}{\Gamma_S}\right)^2},$$

(and analogously for  $g_I$ ). Here  $\Gamma_{T,S}$  set the half-width at half-maximum.

#### Appendix A.4.3. Logistic (step-like) gating

Ideal if you want almost zero coupling below a threshold and nearly constant above:

$$\zeta_I(S_T, S_S) = \frac{\zeta_{I,0}}{1 + \exp\{(-\alpha_T(S_T - S_T^*) - \alpha_S(S_S - S_S^*))\}}$$

$\alpha_{T,S}$  large  $\Rightarrow$  sharp switch-on at the “resonant” field values. Can be used alone or multiplied by one of the peaked forms above for combined gating and tuning.

#### Appendix A.4.4. Frequency-domain resonance

If in your simulation you can track the local time-series  $S_I(t)$ , you can Fourier-analyze it and let the coupling depend on the spectral amplitude at some frequency  $\omega_0$ . E.g.

$$\tilde{S}_I(\omega) = \int dt e^{i\omega t} S_I(t), \quad \zeta_I = \zeta_{I,0} \exp\left(-\frac{(\omega - \omega_0)^2}{2(\Delta\omega)^2}\right),$$

where  $\omega_{\text{peak}}$  is the frequency at which  $|\tilde{S}_I(\omega)|$  is maximal in your local patch. This realizes a truly dynamical resonance in time.

#### Appendix A.4.5. Combined forms

$$\zeta_I = \zeta_{I,0} \exp\left(-\frac{(S_T - S_T^*)^2}{2\sigma^2}\right) \times \frac{1}{1 + \exp\{(-\Lambda(\nabla_\mu S_T \nabla^\mu S_T - C))\}}.$$

Here the coupling only “turns on” both when  $S_I \approx S_I^*$  and its local gradient exceeds some threshold  $\Lambda$ .

#### Appendix A.4.6. How to choose parameters

- Centers  $S_I^*, \omega_0$ : pick based on where/when you want coupling to peak in your simulation.
- Widths  $\sigma, \Gamma, \Delta\omega$ : tune so the resonance is broad enough to capture the phenomenon but narrow enough to remain “selective.”
- Amplitudes  $\zeta_{I,0}, g_{I,0}$ : set by the overall strength of the entropic effects you want to explore.

These functional forms comprise a flexible toolkit for building in the “resonant,” context-dependent behavior of entropic couplings.

#### Appendix A.4.7. Next steps

- Choose a specific potential  $V(S_T, S_S)$  and explicit functional forms  $\zeta_I(S_T, S_S), g_I(S_T, S_S)$  to implement your “resonance.”
- Plug those into the above general formulas.
- Use symbolic algebra software (e.g. xAct in Mathematica) to handle the tensor algebra and covariant derivatives.

## Appendix B. Variational Analysis

This appendix provides a fully detailed walk-through of every variational step needed to go from your ansatz actions to the field equations and stress–energy tensors. It is broken into three parts:

- Varying  $S_{\text{entropic}}$  to get the scalar EoMs and  $T_{\mu\nu}^{\text{entropic}}$
- Varying the non-minimal curvature couplings  $\int \sqrt{-g} \zeta_I S_I R$  to extract their contribution  $\mathcal{M}_{\mu\nu}$  on the LHS of Einstein’s equations.
- Varying the direct matter couplings  $\int \sqrt{-g} g_I S_I \mathcal{L}_m$  to get the additional stress tensor  $T_{\mu\nu}^{\text{coupling}}$ .

### Appendix B.1. $S_{\text{entropic}} \rightarrow$ scalar EoMs and entropic stress tensor

We start with

$$S_{\text{entropic}} = \int d^4x \sqrt{-g} \mathcal{L}_{\text{entropic}},$$

$$\mathcal{L}_{\text{entropic}} = -\frac{1}{2} g^{\mu\nu} (\partial_\mu S_T \partial_\nu S_T + \partial_\mu S_S \partial_\nu S_S) - V(S_T, S_S).$$

#### Appendix B.1.1. Variation w.r.t. $S_T$ (same for $S_S$ )

Compute  $\delta \mathcal{L}$ .  $\delta \mathcal{L} = -g^{\mu\nu} \partial_\mu S_T \partial_\nu (\delta S_T) - \frac{\partial V}{\partial S_T} \delta S_T$ . Integrate by parts the kinetic term:  $\int \sqrt{-g} (-g^{\mu\nu} \partial_\mu S_T \partial_\nu (\delta S_T)) d^4x = \int \sqrt{-g} \square S_T \delta S_T d^4x + (\text{boundary})$ . Here we used:  $\nabla_\mu (\sqrt{-g} g^{\mu\nu} \partial_\nu S_T) = 0$  and dropped total derivatives. Euler–Lagrange equation:  $0 = \square S_T - \frac{\partial V}{\partial S_T}$ ,  $0 = \square S_S - \frac{\partial V}{\partial S_S}$ . Or compactly  $\square S_I - V_{,I} = 0$ .

#### Appendix B.1.2. Variation w.r.t. $g_{\mu\nu} \rightarrow T_{\mu\nu}^{\text{entropic}}$

Recall  $T_{\mu\nu} = -\frac{2}{\sqrt{-g}} \frac{\delta \mathcal{L}}{\delta g^{\mu\nu}}$ . We need variations of  $\sqrt{-g} g^{\alpha\beta} \partial_\alpha S_I \partial_\beta S_I V(S_T, S_S)$  Key identities:  $\delta \sqrt{-g} = -\frac{1}{2} \sqrt{-g} g_{\mu\nu} \delta g^{\mu\nu}$ ,  $\delta g^{\alpha\beta} = -g^{\alpha\mu} g^{\beta\nu} \delta g_{\mu\nu}$ . Kinetic piece:  $\delta \left[ \sqrt{-g} \left( -\frac{1}{2} g^{\alpha\beta} \partial_\alpha S_I \partial_\beta S_I \right) \right] = \sqrt{-g} \left( \partial_\mu S_I \partial_\nu S_I - \frac{1}{2} g_{\mu\nu} g^{\alpha\beta} \partial_\alpha S_I \partial_\beta S_I \right) \delta g^{\mu\nu}$ . Carefully collecting signs yields the standard scalar–field stress tensor. Potential piece:  $\delta [-\sqrt{-g} V] = -\sqrt{-g} \left( -\frac{1}{2} g_{\mu\nu} \delta g^{\mu\nu} \right) V = +\frac{1}{2} \sqrt{-g} g_{\mu\nu} V \delta g^{\mu\nu}$ . Putting it all together,  $T_{\mu\nu}^{\text{entropic}} = \partial_\mu S_T \partial_\nu S_T + \partial_\mu S_S \partial_\nu S_S - \frac{1}{2} g_{\mu\nu} ((\partial S_T)^2 + (\partial S_S)^2) - g_{\mu\nu} V(S_T, S_S)$ .

#### Appendix B.2. Non-minimal coupling $\int \sqrt{-g} \zeta_I S_I R$

Focus on  $S_\zeta = \int d^4x \sqrt{-g} \zeta_I S_I R$ . We vary w.r.t.  $g_{\mu\nu}$ , holding  $S_I$  fixed. Use the well-known identity  $\delta(\sqrt{-g} R) = \sqrt{-g} \left( R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R \right) \delta g^{\mu\nu} + \sqrt{-g} (\square g_{\mu\nu} - \nabla_\mu \nabla_\nu) \delta g^{\mu\nu} + (\text{bdry})$ . Multiplying by  $\zeta_I S_I$ , we get two pieces: Einstein–tensor shift  $\zeta_I S_I (R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R)$ . Derivative terms  $\zeta_I (g_{\mu\nu} \square - \nabla_\mu \nabla_\nu) S_I$ . These together define the  $\mathcal{M}_{\mu\nu}$  modification on the LHS of Einstein’s equations:  $\mathcal{M}_{\mu\nu} = \sum_{I=T,S} \zeta_I \left[ S_I (R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R) + (g_{\mu\nu} \square - \nabla_\mu \nabla_\nu) S_I \right]$ .

#### Appendix B.3. Direct matter coupling $\int \sqrt{-g} g_I S_I \mathcal{L}_m$

Consider  $S_g = \int d^4x \sqrt{-g} g_I S_I \mathcal{L}_m$ . Again vary  $g_{\mu\nu}$ , treating  $S_I$  fixed but allowing  $\mathcal{L}_m$  to depend on  $g$ . One finds:

Two sources of variation  $\delta \sqrt{-g} \delta \mathcal{L}_m = -\frac{1}{2} T_{\mu\nu}^{(m)} \delta g^{\mu\nu}$ . Result  $\delta S_g = \int \sqrt{-g} g_I S_I \left( \frac{1}{2} g_{\mu\nu} \mathcal{L}_m - \frac{1}{2} T_{\mu\nu}^{(m)} \right) \delta g^{\mu\nu} + (\text{bdry})$ . Hence the extra stress from matter coupling is  $T_{\mu\nu}^{\text{coupling}} = \sum_{I=T,S} g_I S_I T_{\mu\nu}^{(m)}$ .

#### Appendix B.4. Putting it all together

Vary the full action  $S = S_{\text{EH}} + S_{\text{entropic}} + S_\zeta + S_g + S_{\text{matter+radiation}}$  w.r.t.  $g_{\mu\nu}$ . You obtain the modified Einstein equations  $G_{\mu\nu} + \mathcal{M}_{\mu\nu} = 8\pi G (T_{\mu\nu}^{\text{entropic}} + T_{\mu\nu}^{\text{coupling}} + T_{\mu\nu}^{\text{matter}})$ , Where:  $G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R$   $\mathcal{M}_{\mu\nu}$  from Appendix B.2 (the non-minimal  $\zeta_I$  pieces)  $T_{\mu\nu}^{\text{entropic}}$  from Appendix B.1.2  $T_{\mu\nu}^{\text{coupling}}$  from Appendix B.3 And the scalars satisfy  $\square S_I - \frac{\partial V}{\partial S_I} + \zeta_I R + g_I \mathcal{L}_m = 0$  once you include their coupling-induced source terms (just vary the total action w.r.t.  $S_I$ ).

#### Appendix B.5. Next steps

- Choose a specific potential  $V(S_T, S_S)$  and explicit functional forms  $\zeta_I(S_T, S_S)$ ,  $g_I(S_T, S_S)$  to implement your “resonance.”
- Plug those into the above general formulas.

- Use symbolic algebra software (e.g. xAct in Mathematica) to handle the tensor algebra and covariant derivatives.

content for Appendix B goes here, detailing the variational calculus.)

## Appendix C. Quantization and Advanced Applications

### Appendix C.1. Application to Black Hole Thermodynamics and Hawking Radiation

The conceptual model of Hawking radiation described in Section 3.2 can be formalized. We consider the dynamics of the entropic fields in a fixed Schwarzschild background metric. The core of the new physics comes from how we model the "boiling" of spacetime. We propose that the quantum fluctuations of the  $S_5$  field become amplified by the gravitational potential. A simplified, effective model for a fluctuation field  $\phi$  near the horizon  $r_s = 2GM$  can be described by a stochastic differential equation:

$$\square\phi - m_{\text{eff}}^2(r)\phi = \sigma(r)\dot{W}(t) \quad (\text{A3})$$

where  $\sigma(r)\dot{W}(t)$  is a stochastic noise term whose amplitude is amplified by the local gravitational field strength, posited as  $\sigma(r) \propto 1/(1 - 2GM/r)$ . Particle emission occurs when a fluctuation's energy exceeds a critical threshold. The rate of emission,  $\Gamma$ , is proportional to the probability of this occurrence:

$$\Gamma \propto P(\phi > \phi_{\text{crit}}) \propto \exp\left(-\frac{\phi_{\text{crit}}^2}{2\langle\phi^2\rangle}\right) \propto \exp(-k \cdot M^2) \quad (\text{A4})$$

This last step connects to the standard Bekenstein-Hawking temperature result,  $T_H \propto 1/M$ , providing a qualitative path to deriving black hole thermodynamics from the framework.

## Appendix D. The Human-AI Collaboration: A New Paradigm for Theoretical Physics

This Entropic Spacetime Framework was developed over approximately two months through a close collaboration between human intuition and advanced Artificial Intelligence. This appendix details the nature and dynamics of this human-AI partnership.

### Appendix D.1. Roles and Interplay in the Development Process

The collaboration evolved through distinct roles. The human contributor served as the initiator and intuitive guide, defining the core problem, providing high-level validation, and setting the strategic research path. The AI (Gemini and ChatGPT) acted as a knowledge synthesizer and ideation partner, performing rapid literature reviews, identifying relevant mathematical formalisms, suggesting standard Lagrangian forms, and assisting in drafting and refining the manuscript.

### Appendix D.2. Implications for Future Scientific Discovery

This project exemplifies a burgeoning new paradigm for theoretical science: a symbiotic partnership where human creativity and strategic direction are amplified by AI's capacity for information synthesis and rapid ideation. It suggests that future breakthroughs may increasingly arise from such augmented intelligence, accelerating the pace at which grand challenges can be tackled.

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