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Article

Structural, Elastic, and Electronic Properties of $\text{Y}_{1-x}\text{Sr}_x\text{TiO}_3$ ($x=0.01, 0.10, 0.50, 0.99$) by Means of Density Functional Theory

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Abstract: The perovskite crystal structure has garnered significant attention due to its unique properties and potential applications. Particularly, YTiO_3 is a ferromagnetic Mott insulator with strong electron correlation effects and orbital ordering, so it is a good candidate for electronic applications. In this research, we investigate the structural, elastic, and electronic properties of $\text{Y}_{1-x}\text{Sr}_x\text{TiO}_3$ ($x=0.01, 0.10, 0.50, 0.99$) perovskites within the $\text{Pm}\bar{3}\text{m}$ space group using density functional theory (DFT) to improve mechanical properties while preserving electronic properties. Initial geometry optimizations were performed using the GGA-PBE functional with a $6\times 6\times 6$ Monkhorst-Pack k-point mesh. The study utilized the CASTEP code with ultrasoft pseudopotentials for DFT calculations. Complete structural relaxations were achieved, and the elastic constants which confirm the mechanical stability of the materials were determined using the finite strain theory. The electronic properties were analyzed through band structure and partial density of states (PDOS) calculations. Despite variations in band structures due to Sr doping, no bandgap was observed, indicating the metallic nature of the compounds. The study highlights that p and d orbitals are predominant in electronic states. The findings suggest that Sr doping significantly influences the elastic properties, making these materials potentially valuable for various applications requiring robust mechanical properties.

Keywords: Perovskite; YTiO_3 ; SrTiO_3 ; density functional theory; elastic properties; electronic structure; mechanical properties; band structure

1. Introduction

The perovskite crystal structure has garnered significant attention due to its unique properties and potential applications. One of the key compounds within this family is Yttrium titanate (YTiO_3) which has been a subject of interest, with studies focusing on the use of molecular dynamics simulations to explore the crystal symmetry and oxygen-ion transport in perovskite materials, shedding light on their behavior at different temperatures [1]. YTiO_3 is a ferromagnetic Mott insulator with strong electron correlation effects and orbital ordering [2,3]. High-quality epitaxial YTiO_3 thin films have been synthesized using pulsed laser deposition and molecular beam epitaxy [2,4]. The electronic structure of YTiO_3 is characterized by a Mott-Hubbard band gap of ~ 1.5 eV and a lower Hubbard band at 1.1 eV below the Fermi level [2]. Both structural distortions and multiorbital electron interactions contribute to orbital polarization in YTiO_3 [3]. Doping with La or Ca significantly affects the electronic and optical properties of YTiO_3 , including spectral weight transfer and changes in the band structure and optical conductivity [3,5]. Furthermore, YTiO_3 has been studied for its potential applications in energy storage, electronic devices, and sensors [6]. One study investigated the metal-insulator transition at $\text{YTiO}_3/\text{LaTiO}_3$ interfaces grown by the soft chemical method [7]. The researchers aimed to understand the behavior of the interface between YTiO_3 and LaTiO_3 , shedding light on potential applications in electronics and optoelectronics. Additionally, a computational study explored the use of $\text{YTiO}_3\pm\delta$ as cathode materials in solid oxide fuel cells (SOFCs) [8]. This study delved into the native defects and oxygen diffusion within YTiO_3 , providing valuable insights for the development of advanced SOFC technologies. In another study, Yttrium-doped SrTiO_3 (SYT) demonstrated high electrical conductivity, structural stability, and chemical compatibility with YSZ

and LSGM, making it a promising anode material for solid oxide fuel cells, achieving a maximum power density of 58 mW/cm² at 900 °C [9]. Furthermore, investigations on functional properties of $\text{Al}_{0.8}\text{Eu}_y\text{La}_{0.2-y}\text{TiO}_3$ ($y = 0.01 - 0.04$) nanoparticles synthesized by the hydrothermal method have also been conducted [10]. The research focused on the potential applications of these nanoparticles and their unique properties for specific electronic or optical devices. The effect of high-temperature annealing on the phase transformations of the perovskite YTiO_{3-x} system was also explored, providing important insights into the structural characteristics of YTiO_3 -based materials [11]. Measurement of elastic modulus in YTiO_3 suggests the appearance of the magnetostriction effect [12].

Additionally, the magnetic properties of YTiO_3 have been studied extensively in the literature, revealing several intriguing phenomena. Gossling et al. [13] investigated optical excitations in Mott-Hubbard insulators YTiO_3 and SmTiO_3 , observing absorption bands and peak anomalies attributed to charge-neutral quasibound states. Akimitsu et al. [14] used the polarized neutron diffraction technique to directly observe the orbital ordering in ferromagnetic YTiO_3 . Additionally, a study by Zhou et al. [15] explored the magnetic properties and electronic structures of YTiO_3 -based superlattices, revealing unexpected magnetic orders and electronic behavior at the interfaces. In a different approach, a joint refinement model for the spin-resolved one-electron reduced density matrix of YTiO_3 was proposed by Gueddida et al. [16], which provided insights into the magnetic properties of YTiO_3 using magnetic structure factors and magnetic Compton profiles data. Chae et al. [17] investigated epitaxial growth and the magnetic properties of orthorhombic YTiO_3 thin films, revealing ferromagnetic transitions and the influence of substrate-induced atomic arrangements. Moreover, the transition from orbital liquid to the Jahn-Teller insulator in orthorhombic perovskites RTiO_3 , including YTiO_3 , was examined by Cheng et al. [18], demonstrating a spin/orbital transition at a magnetic transition temperature. Furthermore, a first-principles study by Weng et al. [19] investigated electron doping of $\text{SmNiO}_3 / \text{YTiO}_3$ via interfacial charge transfer and its effects on the electronic and magnetic properties, highlighting the complex interplay between different materials. An interesting work by An et al. [20] explored the possibility of ferrimagnetism and ferroelectricity in half-substituted rare-earth titanates, including $\text{Y}_{0.5}\text{La}_{0.5}\text{TiO}_3$.

In this study, we investigate the effect of doping Sr into YTiO_3 with different ratios (0.01, 0.10, 0.50, 0.99). Our investigation covers elastic and electronic properties. The outline of this article is as follows. In the next section, the main methods are explained. The investigation results are given in the Results and Discussion section. Lastly, the outcomes of this research are given in the conclusion.

2. Materials and Methods

$\text{Y}_{1-x}\text{Sr}_x\text{TiO}_3$ in $\text{Pm}\bar{3}\text{m}$ space group is doped by Sr by various ratios ($x=0.01, 0.10, 0.50, 0.99$). In **Figure 1**, the initial structure of YTiO_3 is given with the lattice constant $a=3.88\text{\AA}$. The geometry computation is completed with GGA-PBE, and cutoff energy is 630 eV which is defined based on used pseudopotentials and finite basis set correction [21]. The Monkhorst-Pack scheme was utilized to perform special k-points sampling integration across the Brillouin zone with the $6\times6\times6$ mesh [22].

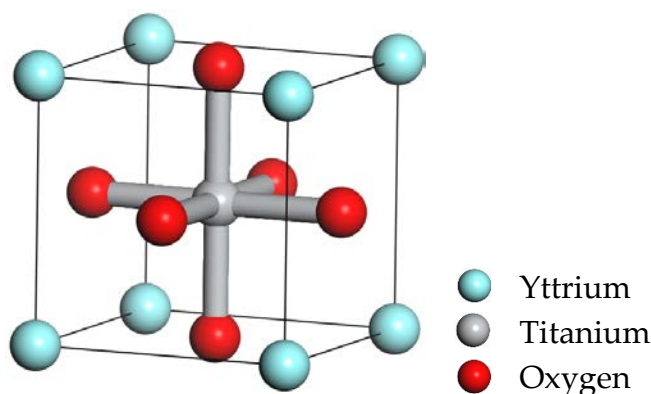


Figure 1. YTiO_3 crystal structure

Our DFT calculations are performed using the projector-augmented wave approach as implemented in the CASTEP [23,24]. We described the exchange and correlation potential using the

GGA-PW91 functional, and the electron-core interactions with the ultrasoft pseudopotential [25]. The energy convergence tolerance was set at $5.0 \times 10^{-6} \text{ eV/atom}$, the maximum force was set at $0.01 \text{ eV/\AA}$, the maximum stress was set at 0.02 GPa , and the maximum displacement was set at $5.0 \times 10^{-4} \text{ \AA}$. The Perdew-Burke-Ernzerhof generalized gradient approximation was used to solve the Kohn-Sham equations [26]. The crystal structures underwent complete relaxation using the Broyden-Fletcher-Goldfarb-Shanno (BFGS) minimization method [27].

The elastic properties are determined by the CASTEP through the application of the finite strain theory [28]. This theory allows for the calculation of elastic constants, which are the coefficients that relate the applied strain to the resulting stress, $\sigma_i = c_{ij}\epsilon_j$. The polycrystalline bulk modulus B and shear modulus G are then evaluated using the Voigt-Reuss-Hill approximation based on the obtained c_{ij} values [29].

3. Results and Discussion

In this study, we investigated the structural, elastic, and electronic properties of $\text{Y}_{1-x}\text{Sr}_x\text{TiO}_3$ ($x=0, 0.01, 0.1, 0.5, 0.99$) perovskites using density functional theory (DFT). The materials exhibit a cubic structure (Pm $\bar{3}$ m) with optimized lattice constants. Elastic properties, including bulk modulus, shear modulus, and Young's modulus, confirm mechanical stability in accordance with the Born-Huang criteria. Electronic properties, examined through band structure and partial density of states (PDOS), reveal significant alterations in band structure due to Sr doping; however, no bandgap is observed, and the metallic nature of the materials is retained, with notable contributions from p and d orbitals.

3.1. Structural Properties

The structural properties and geometries of $\text{Y}_{1-x}\text{Sr}_x\text{TiO}_3$ ($x=0, 0.01, 0.1, 0.5, \text{ and } 0.99$) perovskites were examined using density functional theory (DFT) (See Figure 1). The materials were found to have cubic structures with the space group Pm $\bar{3}$ m (no. 221). The states of complete relaxation in the structures are achieved through the optimization of its geometry, taking into account the lattice constants and internal atomic positions. This process involves adjusting the arrangement of atoms within the structure to minimize any residual stress or strain. By optimizing these parameters, the structure can achieve a more stable and energetically favorable configuration. Wyckoff spaces of Y, Ti, and O are respectively (0, 0, 0), (1/2, 1/2, 1/2), and (1/2, 1/2, 0). The optimized lattice constants, obtained from the unit cell optimization for the ground state stable atomic configuration, are presented in Table 1. To ascertain the stability of perovskite materials, we employ a computational approach to calculate and present the elastic constants in the subsequent sections, providing comprehensive details.

Table 1. Computed crystal parameters

Compounds	Lattice Parameter (A)	Density (amu/A**3)	Final Energy (eV)
YTiO ₃ [This Work]	3.888	3.144	-3979.753
YTiO ₃ [30]	3.89	3.132	---
Y _{0.99} Sr _{0.01} TiO ₃	3.889	4.633	-3977.715
Y _{0.5} Sr _{0.5} TiO ₃	3.920	4.522	-3893.352
Y _{0.1} Sr _{0.9} TiO ₃	3.932	4.483	-3818.073
Y _{0.01} Sr _{0.99} TiO ₃	3.924	4.509	-3698.007

3.2. Elastic Properties

The study of the elastic properties of materials is essential prior to their practical use, as many device features depend on the elastic properties of the materials used in their manufacturing. Therefore, materials with superior elastic properties are highly valuable for incorporation into various devices. The elastic constant serves as a mathematical link between the mechanical and dynamical characteristics of materials, playing a crucial role in determining their properties and how they respond to external forces. Based on the principles of cubic symmetry, the elastic tensor consists of three distinct elastic components, namely c_{11} , c_{12} , and c_{44} . The bulk modulus (BM) calculates the stress needed to change a material's volume, whereas the shear modulus (SM) describes how resistant

a material is to shear stress-induced deformation. The forces per unit cross sectional area needed for an elongation or shortening dimensional change are measured by the Young modulus (YM).

The obtained values for these elastic constants are presented in **Table 2** and are evaluated against the Born-Huang generalized stability criteria [31] as defined in equation (1) for the $Pm\bar{3}m$ structure. Therefore, it can be concluded that compounds exhibit mechanical stability.

$$c_{44} > 0, c_{11} > |c_{12}|, c_{11} + 2c_{12} > 0 \quad (1)$$

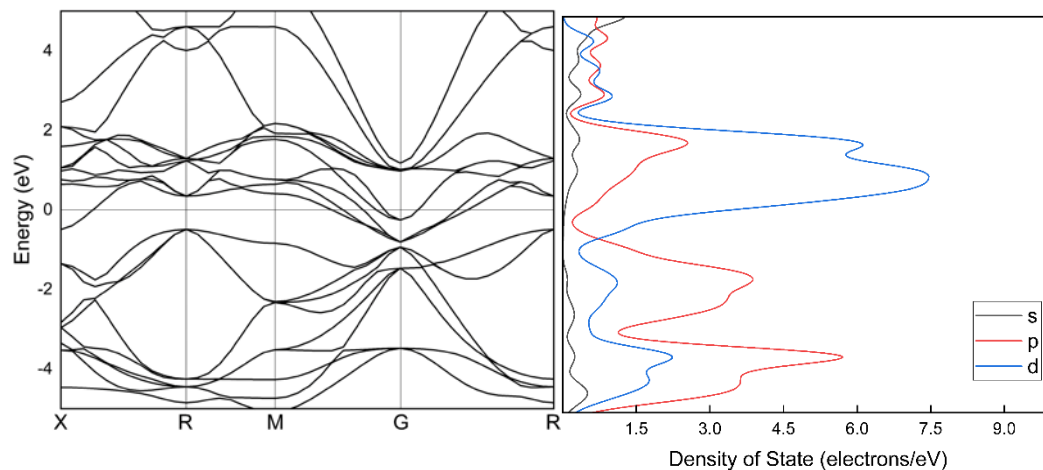
Table 2 displays the findings of our study on the elastic properties under no pressure conditions. Based on the computation results, elastic stability is confirmed.

Table 2. Elastic constants of $Y_{1-x}Sr_xTiO_3$ ($x=0.01, 0.10, 0.50, 0.99$)

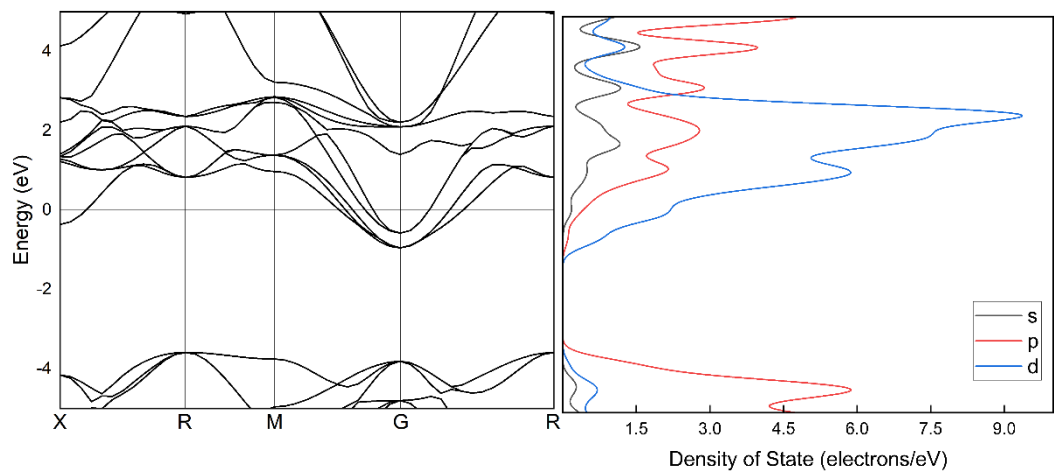
Compounds	C_{11} (GPa)	C_{12} (GPa)	C_{44} (GPa)	BM (GPa)	SM (GPa)	YM (GPa)
YTiO ₃	288.12775	15.46215	131.04480	183.40578	156.27653	113.67433
Y _{0.99} Sr _{0.01} TiO ₃	290.92433	132.91674	16.17490	185.58594	41.30646	115.36065
Y _{0.5} Sr _{0.5} TiO ₃	304.87325	73.32170	77.23745	150.50555	92.65278	230.63188
Y _{0.1} Sr _{0.9} TiO ₃	326.90365	104.50730	90.79475	178.63942	98.95612	250.59631
Y _{0.01} Sr _{0.99} TiO ₃	311.60835	98.84900	110.50245	169.76878	108.85334	269.05522

3.3. Electronic Properties

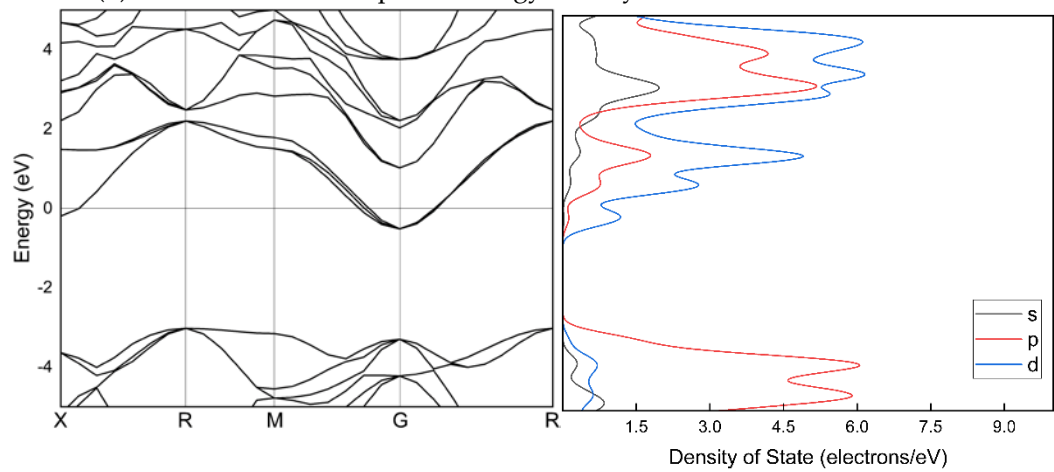
Figure 2 illustrates the band arrangement and PDOS of $Y_{1-x}Sr_xTiO_3$ ($x=0, 0.01, 0.1, 0.5$, and 0.99), where the lattice parameters have been optimized along the high symmetry directions within the first Brillouin zone. The electronic band structure helps to identify the valence band and the area empty of electrons. The possible energies of electrons are explained by the conduction band. The conduction and valence bands are distinguished by the E_F . The valence bands are seen behind the E_F , while the conduction band is located above it. The findings indicate that, despite a considerable alteration in band structures due to Sr doping, there is no bandgap. The valence and conduction band spacing is expanding in doped crystals, but the metal nature of the crystals is preserved. In the band structure, p and d orbitals are particularly stronger than s orbitals.



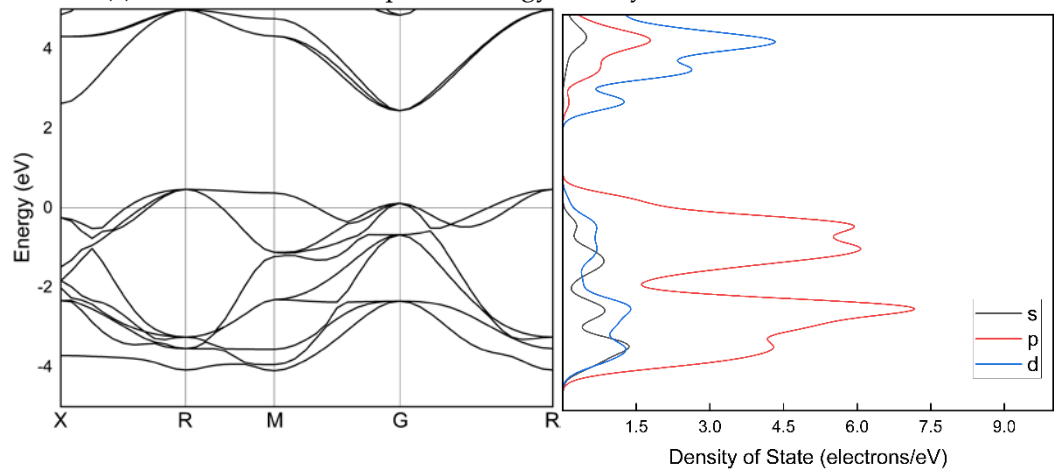
(a) Band structure and partial energy density of states of YTiO₃



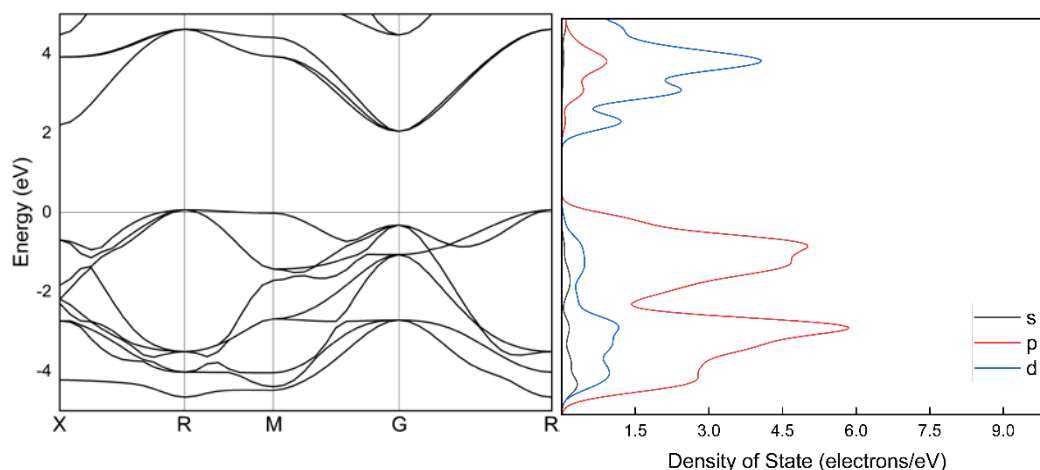
(b) Band structure and partial energy density of states of $\text{Y}_{0.99}\text{Sr}_{0.01}\text{TiO}_3$



(c) Band structure and partial energy density of states of $\text{Y}_{0.5}\text{Sr}_{0.5}\text{TiO}_3$



(d) Band structure and partial energy density of states of $\text{Y}_{0.1}\text{Sr}_{0.9}\text{TiO}_3$



(e) Band structure and partial energy density of states of $\text{Y}_{0.01}\text{Sr}_{0.99}\text{TiO}_3$

Figure 2. The electronic structures of $\text{Y}_{1-x}\text{Sr}_x\text{TiO}_3$. The band structures are shown at the equilibrium lattice parameters, highlighting the high symmetry wave-vectors, k . The partial energy densities of states are presented and the Fermi level is set as 0 eV.

4. Conclusion

In conclusion, the structural, elastic, and electronic properties of $\text{Y}_{1-x}\text{Sr}_x\text{TiO}_3$ ($x=0.01, 0.10, 0.50, 0.99$) perovskites within the $\text{Pm}\bar{3}\text{m}$ space group were thoroughly investigated using density functional theory (DFT). The initial geometry optimizations, performed with the GGA-PBE functional, revealed stable configurations with varying Sr doping levels. Detailed elastic property calculations confirmed the mechanical stability of these compounds, with significant variations in elastic constants attributed to Sr doping. The electronic structure analysis, including band structure and partial density of states (PDOS), indicated that despite changes in the band structures, the materials retained their metallic nature, dominated by p and d orbitals. These findings underscore the potential of Sr-doped YTiO_3 perovskites for applications that demand superior mechanical properties and highlight the influence of Sr doping on the materials' structural and electronic characteristics.

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Conflicts of Interest: The authors declare no conflicts of interest.

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