Tripartite Quantum Steering Dynamics in Photonic Systems Under Non-Markovian Dynamics

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Article

Tripartite Quantum Steering Dynamics in Photonic Systems Under Non-Markovian Dynamics

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Abstract

We investigate the non-Markovian dynamics of quantum steering in a tripartite photonic system subject to dephasing noise. By developing a theoretical framework based on the single-photon dephasing model extended to three independent photons, we analyze the temporal evolution of steering measures S^{A-BC} and S^{AB-C} for two distinct classes of initial states: W-type entangled states and GHZ-type mixed entangled states. The system is studied under various environmental configurations, ranging from fully Markovian to fully non-Markovian regimes, with asymmetric distributions of memory effects across the three photons. Our results reveal that the dynamics of Gaussian steering are highly sensitive to both the number of photons coupled to non-Markovian environments and the specific partition of the system being considered. For W-states, non-Markovian effects induce oscillatory behavior with death-revival cycles, where the intervals of sudden death and revival amplitudes depend critically on the distribution of memory effects. For GHZ-states, we observe multiple death-revival cycles in some configurations and prolonged preservation of steering without complete sudden death in others. Notably, we find that non-Markovian environments can either enhance steering through information backflow or prove detrimental depending on which subsystems they are coupled to relative to the steering and steered parties. These findings demonstrate that non-Markovianity can serve as a resource for protecting specific types of quantum steering, but its effects are highly configuration-dependent, offering insights for quantum information processing tasks requiring the preservation of directional quantum correlations in photonic networks.

Keywords: quantum steering; photonic networks; fabry–pérot cavity; interference filter; open quantum systems; tripartite correlation

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1. Introduction

Quantum steering traces back to Schrödinger's 1936 reply to the EPR paradox [1,2]. It refers to the ability of one party to direct the state of a distant quantum system using shared entanglement [3–7], and provides a way to confirm entanglement without relying on the measurement devices. In 2007 Wiseman and co-workers revived the idea by giving it a precise operational definition [8]. They placed steering between ordinary entanglement and Bell nonlocality: it is strictly stronger than entanglement but weaker than the full nonlocality revealed by Bell tests. This definition underlines the one-way nature of steering, which has since been observed in both continuous-variable and discrete-variable quantum systems [9–15]. The two-party case naturally extends to the multi-party setting, where more complex forms of correlation appear. In the three-party case, Reid's monogamy relationship holds: a single party can steer both others at the same time, but two independent parties cannot steer the same third party [16,17]. Tripartite systems also display layered steering structures that include genuine three-party steering, effective two-party steering, and collective steering. These rich structures position multipartite steering as a key resource for designing and analyzing quantum networks [18].

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Consequently, the special properties of multipartite steering support new protocols for large-scale quantum communication and establish steering itself as a useful quantum resource.

Controlling the temporal evolution of quantum steering and mitigating decoherence are essential for reliable quantum information processing. In photonic systems, environmental dephasing commonly induces the sudden death of steering [19–21], thereby restricting the operational range of quantum communication and networking protocols. Photonic platforms nevertheless provide important advantages over other implementations, including extended coherence times and compatibility with high-speed, long-distance optical channels, which make them especially suitable for scalable quantum networks [22,23]. However, photon loss, phase damping, and spectral diffusion continue to challenge the long-term preservation of directional quantum correlations. Quantum steering occupies a distinctive position among quantum resources: it is strictly stronger than entanglement yet weaker than Bell nonlocality and possesses an inherent one-way character that enables device-independent entanglement verification. Although the dynamics of entanglement and nonlocality under noise have received considerable attention, the behavior of steering—particularly in the tripartite regime—requires further investigation [24]. Non-Markovian environments, where memory effects permit the backflow of information from the reservoir to the system, can counteract irreversible decay and induce periodic revivals of quantum steering [25–27].

To address these issues, we investigate the non-Markovian dynamics of quantum steering in a tripartite photonic system subject to dephasing noise. By developing a theoretical framework based on the single-photon dephasing model extended to three independent photons, we analyse the temporal evolution of the steering measures $S^{A\to BC}$ and $S^{AB\to C}$ for two distinct classes of initial states: W-type entangled states and GHZ-type mixed entangled states. The system is studied under various environmental configurations ranging from fully Markovian to fully non-Markovian regimes, with asymmetric distributions of memory effects across the three photons. Our results reveal that the dynamics of Gaussian steering are highly sensitive to both the number of photons coupled to non-Markovian environments and the specific partition of the system being considered. Non-Markovian environments can either enhance steering through information backflow or prove detrimental depending on which subsystems they are coupled to relative to the steering and steered parties. These findings demonstrate that non-Markovianity can serve as a resource for protecting specific types of quantum steering in photonic networks, but its effects are highly configuration-dependent.

The structure of this paper is as follows. In Sec. 2, we develop the theoretical framework for non-Markovian dynamics in photonic networks. We begin by introducing the single-photon dephasing model, which we then extend to describe three independent photons interacting with their respective environments. A matrix representation in the polarization basis is presented to accurately capture the dephasing mechanisms affecting each photon. In Sec. 3, we apply this framework to analyze the temporal evolution of W-type entangled states under both Markovian and non-Markovian dephasing dynamics. Sec. 4 extends this analysis to GHZ-type mixed entangled states, examining how environmental memory effects influence their decay and revival patterns. The concept of quantum steering in three-photon systems is introduced in Sec. 5, where we define the steering measures S^{A-BC} and S^{AB-C} used to quantify directional quantum correlations between different bipartitions. Our main findings are presented and discussed in Sec. 6, where we systematically compare the steering dynamics for both W and GHZ states under various environmental configurations, ranging from fully Markovian to fully non-Markovian regimes with asymmetric distributions of memory effects across the three photons. Finally, Sec. 7 summarizes our conclusions and discusses the implications of our results for quantum information processing tasks requiring the preservation of directional quantum correlations in photonic networks.

2. Non-Markovian Evolution in Structured Photonic Reservoirs

In this section we formulate a general theoretical framework describing memory-dependent dynamics in photonic systems. We begin by analyzing a single-photon polarization qubit interacting

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with a structured frequency reservoir, and then extend the description to a three-photon configuration. This progressive construction highlights how environmental memory effects become more intricate as the Hilbert-space dimension increases and provides a systematic approach for modeling decoherence in multi-photon quantum networks.

2.1. Single-Photon Dephasing Mechanism

We consider an open quantum system where the system is encoded in the polarization of a photon, while the environmental degree of freedom is associated with its frequency. Experimentally, controllable non-Markovian dynamics can be realized using a tunable Fabry–Pérot (FP) cavity combined with an interference filter and a birefringent quartz plate (see Figure 1) [28–32]. The FP cavity generates a comb-like spectral distribution, while the interference filter isolates two dominant frequency components.

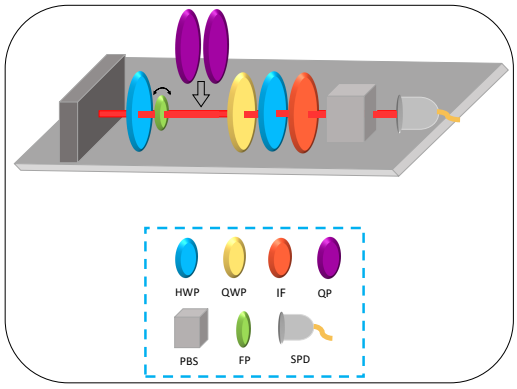

Figure 1. This figure provides a schematic illustration of the theoretical model examined in this work. It conceptually depicts the interaction between polarization and frequency in a standard optical configuration, thereby outlining the overall theoretical framework. The diagram features several labeled optical components—half-wave plate (HWP), quarter-wave plate (QWP), interference filter (IF), quartz plate (QP), polarizing beam splitter (PBS), Fabry–Pérot cavity (FP), and single-photon detector (SPD)—to demonstrate how these polarization and frequency degrees of freedom can, in principle, be controlled and manipulated.

The normalized spectral probability density of the photon is denoted by $g(\omega)$ and is modeled as a weighted sum of two Gaussian functions [33]:

$$g(\omega) = \frac{\cos^2\theta}{\sqrt{2\pi}\sigma} e^{-\frac{(\omega-\omega_1)^2}{2\sigma^2}} + \frac{\sin^2\theta}{\sqrt{2\pi}\sigma} e^{-\frac{(\omega-\omega_2)^2}{2\sigma^2}}, \tag{1}$$

where ω_1 and ω_2 denote the central frequencies, σ is the common spectral width, and $\theta \in [0, \pi/2]$ determines the relative weight of the two spectral peaks. The tilt of the FP cavity effectively controls the value of θ, allowing one to tune the spectral asymmetry.

When the photon propagates through the birefringent quartz plate, its polarization components accumulate different frequency-dependent phases, resulting in a pure dephasing process. The interaction Hamiltonian describing the coupling between polarization and frequency degrees of freedom is given by [34]

$$H_{SE} = -\int (n_V|V\rangle\langle V| + n_H|H\rangle\langle H|) \otimes \omega|\omega\rangle\langle\omega|\, d\omega, \tag{2}$$

where n_V and n_H are the refractive indices corresponding to vertical and horizontal polarizations, respectively. This interaction generates a relative phase shift between polarization components that depends on frequency.

The initial system–environment state is assumed to be separable,

$$\rho_{SE}(0) = \rho_S(0) \otimes \rho_E(0),$$

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with the environmental density operator

$$\rho_E(0) = \int d\omega\, d\omega'\, G(\omega)G^*(\omega')|\omega\rangle\langle\omega'|, \qquad |G(\omega)|^2 = g(\omega). \tag{3}$$

The reduced system dynamics is obtained through

$$\rho_S(t) = \mathrm{Tr}_E\Big[U(t)\rho_{SE}(0)U^\dagger(t)\Big], \qquad U(t) = e^{-iH_{SE}t}. \tag{4}$$

In the polarization basis $\{|V\rangle, |H\rangle\}$, the evolved density matrix takes the form [35]

$$\rho_S(t) = \begin{pmatrix} \rho_{VV}(0) & \rho_{VH}(0)\eta(t) \\ \rho_{HV}(0)\eta^*(t) & \rho_{HH}(0) \end{pmatrix}, \tag{5}$$

where the dephasing function is

$$\eta(t) = \int g(\omega)\, e^{i\omega\Delta n t}\, d\omega, \qquad \Delta n = n_V - n_H. \tag{6}$$

Substituting Eq. (1) into Eq. (6) yields

$$\eta(t) = e^{-\frac{1}{2}\sigma^2(\Delta n t)^2}\left(e^{i\omega_1\Delta n t}\cos^2\theta + e^{i\omega_2\Delta n t}\sin^2\theta\right). \tag{7}$$

2.2. Generalization to Three Independent Photons

We now consider three non-interacting polarization qubits labeled A, B, and C, each associated with an individual photon. Each subsystem interacts locally with its own structured environment E_A, E_B, and E_C.

The reduced dynamics of subsystem $S \in \{A, B, C\}$ can be expressed using a completely positive trace-preserving map,

$$\rho_S(t) = \sum_{m,n} K_{mn}^{(S)}(t)\, \rho_S(0)\, K_{mn}^{(S)\dagger}(t), \tag{8}$$

where the Kraus operators encode local decoherence characterized by the function $\eta_S(t)$ determined by the local spectral distribution [36].

Since the photons do not interact directly, the total evolution operator factorizes as

$$U_{ABC}(t) = U_A(t) \otimes U_B(t) \otimes U_C(t). \tag{9}$$

The global density matrix evolves according to

$$\rho_{ABC}(t) = \sum_{m,n,p} K_m^{(A)}(t) \otimes K_n^{(B)}(t) \otimes K_p^{(C)}(t)\, \rho_{ABC}(0)\, \left(K_m^{(A)} \otimes K_n^{(B)} \otimes K_p^{(C)}\right)^\dagger. \tag{10}$$

2.3. Matrix Representation in the Polarization Basis

Using the local basis $\{|V\rangle, |H\rangle\}$, the reduced dynamics of each subsystem can be written as

$$\rho_{\mu'\nu'}^{(S)}(t) = \sum_{\mu,\nu} \mathcal{T}_{\mu'\nu';\mu\nu}^{(S)}(t)\rho_{\mu\nu}^{(S)}(0), \tag{11}$$

where

$$\mathcal{T}_{\mu'\nu';\mu\nu}^{(S)}(t) = \sum_{m,n} \langle\mu'|K_{mn}^{(S)}|\mu\rangle\langle\nu|K_{mn}^{(S)\dagger}|\nu'\rangle. \tag{12}$$

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The tripartite Hilbert space is spanned by

$$\begin{aligned}\mathcal{B} = \{&|V_A V_B V_C\rangle, |V_A V_B H_C\rangle, |V_A H_B V_C\rangle, |V_A H_B H_C\rangle,\\ &|H_A V_B V_C\rangle, |H_A V_B H_C\rangle, |H_A H_B V_C\rangle, |H_A H_B H_C\rangle\}.\end{aligned} \tag{13}$$

Under local dephasing dynamics, the populations remain unchanged,

$$\rho_{ii}(t) = \rho_{ii}(0), \quad i = 1,\ldots,7, \qquad \rho_{88}(t) = 1 - \sum_{i=1}^{7} \rho_{ii}(0). \tag{14}$$

The coherence terms evolve according to products of local decoherence functions:

$$\begin{aligned}
&\rho_{12}(t) = \rho_{12}(0)\eta_C(t), && \rho_{13}(t) = \rho_{13}(0)\eta_B(t), && \rho_{15}(t) = \rho_{15}(0)\eta_A(t),\\
&\rho_{24}(t) = \rho_{24}(0)\eta_B(t), && \rho_{26}(t) = \rho_{26}(0)\eta_A(t), && \rho_{37}(t) = \rho_{37}(0)\eta_A(t),\\
&\rho_{34}(t) = \rho_{34}(0)\eta_C(t), && \rho_{48}(t) = \rho_{48}(0)\eta_A(t), && \rho_{56}(t) = \rho_{56}(0)\eta_C(t),\\
&\rho_{57}(t) = \rho_{57}(0)\eta_B(t), && \rho_{68}(t) = \rho_{68}(0)\eta_B(t), && \rho_{78}(t) = \rho_{78}(0)\eta_C(t),\\
&\rho_{14}(t) = \rho_{14}(0)\eta_B(t)\eta_C(t), && \rho_{16}(t) = \rho_{16}(0)\eta_A(t)\eta_C(t), &&\\
&\rho_{17}(t) = \rho_{17}(0)\eta_A(t)\eta_B(t), && \rho_{23}(t) = \rho_{23}(0)\eta_B(t)\eta_C^*(t), &&\\
&\rho_{25}(t) = \rho_{25}(0)\eta_A(t)\eta_C^*(t), && \rho_{28}(t) = \rho_{28}(0)\eta_A(t)\eta_B(t), &&\\
&\rho_{35}(t) = \rho_{35}(0)\eta_A(t)\eta_B^*(t), && \rho_{38}(t) = \rho_{38}(0)\eta_A(t)\eta_C(t), &&\\
&\rho_{46}(t) = \rho_{46}(0)\eta_A(t)\eta_B^*(t), && \rho_{47}(t) = \rho_{47}(0)\eta_A(t)\eta_C^*(t), &&\\
&\rho_{58}(t) = \rho_{58}(0)\eta_B(t)\eta_C(t), && \rho_{67}(t) = \rho_{67}(0)\eta_B(t)\eta_C^*(t), &&\\
&\rho_{18}(t) = \rho_{18}(0)\eta_A(t)\eta_B(t)\eta_C(t), && \rho_{27}(t) = \rho_{27}(0)\eta_A(t)\eta_B(t)\eta_C^*(t), &&\\
&\rho_{36}(t) = \rho_{36}(0)\eta_A(t)\eta_B^*(t)\eta_C(t), && \rho_{45}(t) = \rho_{45}(0)\eta_A(t)\eta_B^*(t)\eta_C^*(t). &&
\end{aligned} \tag{15}$$

Hermiticity is preserved automatically, ensuring $\rho_{ABC}(t) = \rho_{ABC}^{\dagger}(t)$.

Equations (14)–(15) provide a complete description of the three-qubit density matrix evolution under local non-Markovian dephasing, where collective coherence decay is governed by products of individual decoherence functions. This formalism enables systematic analysis of coherence revival, information backflow, and multipartite entanglement persistence in structured photonic reservoirs.

3. Temporal Evolution of W-Type Entangled State

This section examines how a W-type state changes over time in a system made up of three qubits. We consider the three-qubit W-state as the initial resource:

$$|\Psi_W\rangle = \frac{1}{\sqrt{3}}(|H_A V_B V_C\rangle + |V_A H_B V_C\rangle + |V_A V_B H_C\rangle), \tag{16}$$

which is a paradigmatic example of a tripartite entangled state. W-states are distinguished by their robustness: if one qubit is lost, the remaining two qubits retain some entanglement, making them particularly suitable for noisy quantum systems and distributed quantum protocols [37]. They can be generated experimentally [38] using sequences of quantum gates or optical setups, and they play a central role in quantum secret sharing, quantum networks, and error-correcting schemes.

The initial density matrix is defined as

$$\rho(0) = |\Psi_W\rangle \langle\Psi_W| . \tag{17}$$

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In the standard three-qubit basis (13), the density matrix takes the explicit 8×8 form

$$\rho(0) = \frac{1}{3}\begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 1 & 0 & 1 & 0 & 0 & 0 \\ 0 & 1 & 1 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 1 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}. \tag{18}$$

This matrix represents the W-state in the full three-qubit Hilbert space, with nonzero entries only in the block 3×3 corresponding to the states $|H_A V_B V_C\rangle$, $|V_A H_B V_C\rangle$, and $|V_A V_B H_C\rangle$. The time evolution of the density matrix under decoherence can be obtained using the general formalism of Eqs. (14) and (15). The corresponding time-dependent density matrix is

$$\rho(t) = \frac{1}{3}\begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & \eta_B(t)\eta_C^*(t) & 0 & \eta_A(t)\eta_C^*(t) & 0 & 0 & 0 \\ 0 & \eta_B^*(t)\eta_C(t) & 1 & 0 & \eta_A(t)\eta_B^*(t) & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & \eta_A^*(t)\eta_C(t) & \eta_A^*(t)\eta_B(t) & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}. \tag{19}$$

This explicit representation allows for a straightforward analysis of the temporal decay of W-type entanglement and provides a clear basis for calculating multipartite entanglement measures.

4. Temporal Evolution of GHZ-Type Mixed Entangled State

This section explores the temporal evolution of a GHZ-type state in a system composed of three qubits. We consider a GHZ–type entangled configuration as the initial resource of the system. The GHZ state represents a fundamental class of tripartite entanglement distinct from the W-type family, exhibiting maximal nonlocal correlations and perfect coherence between the logical states $|V_A V_B V_C\rangle$ and $|H_A H_B H_C\rangle$. It is defined as

$$|\mathrm{GHZ}\rangle = \frac{1}{\sqrt{2}}(|V_A V_B V_C\rangle + |H_A H_B H_C\rangle). \tag{20}$$

To account for realistic imperfections and partial mixing with environmental noise, we describe the initial state using a convex combination of a maximally mixed state and the pure GHZ state:

$$\rho(0) = \frac{1-a}{8}\,\mathbb{I}_8 + a\,|\mathrm{GHZ}\rangle\langle\mathrm{GHZ}|, \qquad 0 \le a \le 1, \tag{21}$$

where a denotes the degree of purity or coherence. This mixed GHZ-type density operator serves as a realistic model for partially decohered entangled sources in photonic or superconducting qubit platforms. The GHZ state plays a central role in foundational studies of quantum nonlocality and multipartite correlations. Unlike W-type states, GHZ-type entanglement exhibits stronger violation of local realism and is particularly useful in testing Bell-type inequalities, quantum secret sharing, distributed quantum computation, and error-resistant quantum communication protocols. Its preparation and manipulation have been demonstrated in several physical systems, including trapped ions, photons, and superconducting circuits [39–45]. In contrast to the W-state, which retains partial bipartite entanglement after the loss of one qubit, the GHZ state is highly sensitive to particle loss—its

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entanglement vanishes if any qubit is traced out. However, this fragility is compensated by its stronger correlations and nonlocal features, making it a valuable testbed for studying decoherence mechanisms and the transition from quantum to classical behavior.

In the standard computational polarization basis (13), the initial state (21), which represents a convex combination of the maximally mixed state and the pure GHZ state. The parameter a quantifies the degree of purity, ranging from a completely mixed state $(a = 0)$ to the pure GHZ state $(a = 1)$. The pure component $|\mathrm{GHZ}\rangle\langle\mathrm{GHZ}|$ occupies a two-dimensional subspace spanned by the first and last computational basis vectors, while all other elements of the density matrix vanish. Consequently, in this basis, the initial density matrix takes the following form

$$\rho(0) = \begin{pmatrix} c & 0 & 0 & 0 & 0 & 0 & 0 & \frac{a}{2} \\ 0 & d & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & d & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & d & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & d & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & d & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & d & 0 \\ \frac{a}{2} & 0 & 0 & 0 & 0 & 0 & 0 & c \end{pmatrix}, \tag{22}$$

where $c = \frac{1+3a}{8}$ and $d = \frac{1-a}{8}$. This form clearly identifies the coherent GHZ correlations through the off-diagonal terms $\rho_{18} = \rho_{81} = a/2$, which encode the superposition between the all-$|V\rangle$ and all-$|H\rangle$ components. As the system evolves under local decoherence channels, these terms decay, providing a quantitative signature of entanglement degradation and non-Markovian memory effects. By employing the relations established in Eqs. (14) and (15), the time-dependent density operator of the tripartite system, denoted by $\rho(t)$, can be explicitly constructed. Since dephasing processes preserve the population terms, the diagonal elements of $\rho(t)$ remain constant in time, while the coherence elements decay according to the product of the corresponding local decoherence functions $\eta_S(t)$ associated with each qubit subsystem $S = A, B, C$. Accordingly, the temporal form of the density matrix is expressed as

$$\rho(t) = \begin{pmatrix} c & 0 & 0 & 0 & 0 & 0 & 0 & \frac{a}{2}\eta_A(t)\eta_B(t)\eta_C(t) \\ 0 & d & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & d & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & d & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & d & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & d & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & d & 0 \\ \frac{a}{2}\eta_A^*(t)\eta_B^*(t)\eta_C^*(t) & 0 & 0 & 0 & 0 & 0 & 0 & c \end{pmatrix}. \tag{23}$$

This explicit representation allows for a straightforward analysis of the temporal decay of GHZ-type entanglement and provides a clear basis for calculating multipartite entanglement measures in the photonic system.

5. Quantum Steering in Three-Photon System

The concept of tripartite quantum steering in the photonic system is illustrated in Figure 2. Three independent photons A, B, and C are prepared in either W-type or GHZ-type entangled states. Each photon couples locally to its own dephasing environment E_A, E_B, or E_C (Markovian or non-Markovian). Steering is witnessed by the negativity of the matrices τ^1 ($S^{A\to BC}$) and τ^2 ($S^{AB\to C}$). The blue arrow represents one-party steering $S^{A\to BC}$ (Alice steers BC), while the red arrow represents two-

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party steering $S^{AB\to C}$ (AB steers C). This schematic clarifies how local non-Markovian environments influence directional quantum correlations across different partitions of the tripartite network.

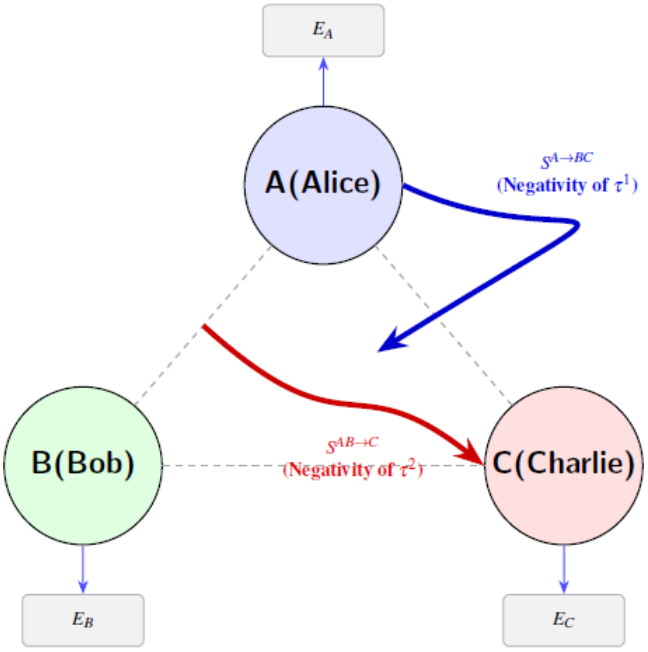

Figure 2. Schematic illustration of tripartite quantum steering in the photonic system. Three independent photons A, B, C in W-type or GHZ-type entangled states. Each photon couples to its local dephasing environment E_A, E_B, E_C (Markovian or non-Markovian). Steering is witnessed by negativity of τ^1 ($S^{A\to BC}$) and τ^2 ($S^{AB\to C}$). Blue arrow: one-party steering $S^{A\to BC}$ (Alice steers BC, negativity of τ^1). Red arrow: two-party steering $S^{AB\to C}$ (AB steers C, negativity of τ^2). The blue arrow represents the decay of each subsystem to its environment.

Negativity measures entanglement by quantifying the extent to which a system's partially transposed density matrix has negative eigenvalues. In a tripartite system, the bipartite entanglement between one subsystem and the other two is termed one-tangle, defined as

$$N_{A(BC)} = \|\rho_{ABC}^{T_A}\| - 1, \tag{24}$$

where T_A denotes the partial transpose of ρ_{ABC} with respect to subsystem A [93]. Since $\|A\| - 1$ equals twice the sum of the absolute values of the negative eigenvalues, the one-tangle can also be expressed as

$$N_{A(BC)} = 2\sum_{i=1}^{n} |\lambda_{A(BC)}^{(-)}|_i, \tag{25}$$

where $|\lambda_{A(BC)}^{(-)}|_i$ are the negative eigenvalues of the partially transposed matrix. For a general tripartite state ρ_{ABC} shared by Alice, Bob, and Charlie in the three-photon system, the steering from Alice to Bob and Charlie can be witnessed if the density matrix $\tau_{ABC}^{1(W)}$, defined as

$$\tau_{ABC}^{1(W)} = \frac{1}{\sqrt{3}}\rho_{ABC}^{(W)} + \left(\frac{\sqrt{3}-1}{2\sqrt{3}}\right) I_2 \otimes \rho_{BC}^{(W)}, \tag{26}$$

is entangled, where $\rho_{BC}^{(W)} = \mathrm{Tr}_A(\rho_{ABC}^{(W)})$ and I_2 is the 2×2 identity matrix [94]. By substituting the evolved density matrix for the W state from Eq. (18), the matrix $\tau_{ABC}^{1(W)}$ is obtained as

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$$\tau_{ABC}^{1(W)} = \frac{1}{3}\begin{pmatrix} \alpha & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & \frac{\sqrt{3}}{3}+\alpha & \left(\frac{\sqrt{3}}{3}+\alpha\right)\eta_B\eta_C^* & 0 & \frac{\sqrt{3}}{3}\eta_A\eta_C^* & 0 & 0 & 0 \\ 0 & \left(\frac{\sqrt{3}}{3}+\alpha\right)\eta_B^*\eta_C & \frac{\sqrt{3}}{3}+\alpha & 0 & \frac{\sqrt{3}}{3}\eta_A\eta_B^* & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & \frac{\sqrt{3}}{3}\eta_A^*\eta_C & \frac{\sqrt{3}}{3}\eta_A^*\eta_B & 0 & \frac{\sqrt{3}}{3}+\alpha & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & \alpha & \alpha\,\eta_B\eta_C^* & 0 \\ 0 & 0 & 0 & 0 & 0 & \alpha\,\eta_B^*\eta_C & \alpha & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}, \tag{27}$$

where $\alpha = \frac{\sqrt{3}-1}{2\sqrt{3}}$. In this representation, the global prefactor $1/3$ ensures the proper normalization of the density matrix and reflects the statistical weight inherited from the underlying W-state construction. The matrix elements inside the brackets therefore represent dimensionless population and coherence weights. These weights naturally separate into distinct physically meaningful sectors: The coefficient α characterizes the weakest population sector and is associated with mixing-induced or noise-like contributions introduced by the state construction. The constant contribution $\sqrt{3}/3$ represents the intrinsic population and coherence scale originating from the ideal W-state component. The combined coefficient $(\sqrt{3}/3+\alpha)$ governs the dominant population sector and also appears in several coherence terms, indicating that both the intrinsic W-state structure and the mixing parameter α jointly determine the strength of quantum correlations between basis states.

Similarly, the steering from Alice and Bob to Charlie can be witnessed if the following density matrix is followed.

$$\tau_{ABC}^{2(W)} = \frac{1}{3}\rho_{ABC}^{(W)} + \frac{1}{6}I_4 \otimes \rho_C^{(W)}, \tag{28}$$

is entangled, where $\rho_C^{(W)} = \mathrm{Tr}_{AB}(\rho_{ABC}^{(W)})$ and I_4 is the identity matrix 4×4 [94]. The explicit form of $\tau_{ABC}^{2(W)}$ is

$$\tau_{ABC}^{2(W)} = \frac{1}{18}\begin{pmatrix} 2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 3 & 2\,\eta_B\eta_C^* & 0 & 2\,\eta_A\eta_C^* & 0 & 0 & 0 \\ 0 & 2\,\eta_B^*\eta_C & 4 & 0 & 2\,\eta_A\eta_B^* & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 2\,\eta_A^*\eta_C & 2\,\eta_A^*\eta_B & 0 & 4 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 2 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix} \tag{29}$$

This expression is derived for the W state under local non-Markovian dephasing in the photonic system, where the off-diagonal elements are damped by products of the dephasing functions $\eta_A(t)$, $\eta_B(t)$, and $\eta_C(t)$ corresponding to the qubits that differ between basis states. Due to the mathematical equivalence of the effective density operators derived in the non-Markovian dephasing model to those arising in photonic systems where one photon experiences asymmetric decoherence, the steering quantifiers $S^{A\to BC}$ and $S^{AB\to C}$ exhibit behaviors modulated by the environmental memory effects.

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The explicit form of $\tau_{ABC}^{1(G)}$ for the GHZ state is

$$\tau_{ABC}^{1(G)} = \frac{1}{3}\begin{pmatrix} K_c & 0 & 0 & 0 & 0 & 0 & 0 & \frac{\sqrt{3}}{2}a\,\eta_A\eta_B\eta_C \\ 0 & K_d^{(1)} & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & K_d^{(1)} & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & K_d^{(2)} & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & K_d^{(2)} & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & K_d^{(1)} & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & K_d^{(1)} & 0 \\ \frac{\sqrt{3}}{2}a\,\eta_A^*\eta_B^*\eta_C^* & 0 & 0 & 0 & 0 & 0 & 0 & K_c \end{pmatrix}, \tag{30}$$

where $K_c \equiv \sqrt{3}\,c + 3\alpha\,(c+d)$, $K_d^{(1)} \equiv \sqrt{3}\,d + 6\alpha\,d$, $K_d^{(2)} \equiv \sqrt{3}\,d + 3\alpha\,(c+d)$., and the shorthand $\eta_A = \eta_A(t)$, $\eta_B = \eta_B(t)$, $\eta_C = \eta_C(t)$ (and similarly for complex conjugates) has been used, and $\rho_{BC}(t) = \mathrm{Tr}_A[\rho(t)]$. The explicit form of $\tau_{ABC}^{2(G)}$ for the GHZ state is

$$\tau_{ABC}^{2(G)} = \frac{1}{6}\begin{pmatrix} L_c & 0 & 0 & 0 & 0 & 0 & 0 & a\,\eta_A\eta_B\eta_C \\ 0 & L_d & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & L_d & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & L_d & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & L_d & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & L_d & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & L_d & 0 \\ a\,\eta_A^*\eta_B^*\eta_C^* & 0 & 0 & 0 & 0 & 0 & 0 & L_c \end{pmatrix}, \tag{31}$$

where $L_c = 3(c+d)$, $L_d = 5d + c$, and $\rho_C(t) = \mathrm{Tr}_{AB}[\rho(t)] = \mathrm{diag}(c+6d, c+6d)$ (again up to normalization). The matrices are Hermitian, with off-diagonal coherence terms preserved only in the $|000\rangle \leftrightarrow |111\rangle$ sector and scaled by the respective coefficients in the steering witnesses.

6. Results and Discussion

The dynamics of Gaussian steering for a three-mode photonic system, initially prepared in a W-state and subject to dephasing, are presented in Figure 3. The figure illustrates the time evolution of two distinct steering measures: S^{A-BC} (panel a), which quantifies the ability of party A to steer the composite subsystem BC, and S^{AB-C} (panel b), which quantifies the ability of the composite system AB to steer party C. The system evolves under various configurations of Markovian and non-Markovian dephasing environments, characterized by the angle parameter θ, with the spectral properties of the environments defined by the detuning $\Delta n = 0.01$, central frequencies $\nu_1 = 2.676$ PHz and $\nu_2 = 2.692$ PHz, and bandwidth $\sigma = 1.8$ THz [29].

Panel (a) depicts the steering from mode A to the combined modes B and C. A key observation is the pronounced difference in the longevity of steering under different environmental conditions. In the fully Markovian symmetric case (dashed blue line, $\theta_A = \theta_B = \theta_C = 0$), the steering decays monotonically and slowly, vanishing entirely by approximately $t \approx 60$ ps. This behavior is characteristic of memoryless environments, where information is lost irreversibly. In stark contrast, the fully non-Markovian symmetric case (solid green line, $\theta_A = \theta_B = \theta_C = \pi/4$) exhibits a rapid initial decay leading to a sudden death of steering. After a finite interval, however, the steering experiences a

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revival, displaying oscillatory behavior: it passes through a maximum before decaying once more to another sudden death. This oscillatory behavior is a direct consequence of the memory effects arising from the non-Markovian environments, which facilitate a temporary backflow of information between the tripartite system and its reservoirs, leading to the recurrent revival and decay of quantum correlations. The intermediate cases reveal a nuanced dependence on how many photons are coupled to a non-Markovian reservoir. The scenario with a single photon in a non-Markovian environment (case 1, dash-dotted red line, $\theta_A = \pi/4, \theta_B = \theta_C = 0$) exhibits a similar oscillatory behavior to the fully non-Markovian symmetric case, characterized by an initial rapid decay to sudden death followed by a revival. However, in this configuration, both the intervals of sudden death and the subsequent revival are noticeably smaller, and the overall amplitude of the decay is also increased compared to the symmetric case. This suggests that while the Markovian baths on modes B and C drive the initial loss, the single non-Markovian bath on mode A introduces memory effects that are sufficient to induce revivals, though with diminished strength and duration. The configuration with two photons in a non-Markovian environment (case 2, long-dashed black line, $\theta_A = \theta_B = \pi/4, \theta_C = 0$) displays the same qualitative oscillatory pattern observed in the previous cases. However, compared to case 1, the intervals of sudden death are further reduced, and the revival occurs more rapidly, with the steering maintaining a higher overall value (~ 0.07 at $t = 40$ ps). This indicates a cumulative effect; coupling more modes to non-Markovian environments enhances the memory-driven recovery process, leading to greater overall robustness of the steering shared by A over BC. An intriguing anomaly is observed for the configuration $\theta_A = 0, \theta_B = \pi/4, \theta_C = 0$ (spaced dashed orange line). Despite only one photon being in a non-Markovian environment (mode B), the steering S^{A-BC} exhibits an oscillatory decay that is more rapid than in the fully Markovian case, notably characterized by the absence of complete sudden death over short intervals, unlike the distinct death-revival cycles seen in previous non-Markovian configurations. This counter-intuitive result suggests that the distribution of non-Markovian resources is critical. If the steering party (A) is in a Markovian environment, placing a non-Markovian bath on a steered party (B) does not aid A's ability to steer and may, through back-action effects mediated by the global tripartite correlations, even accelerate the loss of its steering capability while simultaneously preventing the complete suppression of steerability.

Panel (b) shows the steering from the composite system AB to party C. The dynamics here are qualitatively different from panel (a), with the initial values being notably smaller across all configurations compared to the S^{AB-C} case. The fully Markovian symmetric case (dashed blue line, $\theta_A = \theta_B = \theta_C = 0$) exhibits a monotonic decay, with steering slowly decreasing and undergoing sudden death by approximately $t = 40$ ps, after which it remains zero with no further revival. In stark contrast, the fully non-Markovian symmetric case (solid green line, $\theta_A = \theta_B = \theta_C = \pi/4$) displays a pronounced oscillatory behavior. After an initial rapid decay leading to a sudden death around $t \approx 9$ ps, the steering experiences a strong revival, passing through a maximum value of approximately 0.01 at $t \approx 38$ ps. This oscillatory behavior is a clear signature of strong system-environment memory effects, where information lost to the environment is subsequently fed back into the system, temporarily restoring quantum correlations. The behavior of the mixed cases is also highly distinct. The configuration with one non-Markovian photon on the steering side (case 1, dash-dotted red line, $\theta_A = \pi/4, \theta_B = \theta_C = 0$) exhibits the same qualitative oscillatory behavior observed in the fully non-Markovian symmetric case, characterized by an initial decay to sudden death followed by a revival. However, in this configuration, the interval of sudden death is notably smaller (occurring around $t = 10$ ps), and the amplitude of the subsequent revival is sufficiently greater than in the symmetric case. This enhanced revival amplitude, despite the shorter death interval, suggests that the asymmetric distribution of non-Markovian resources—where only the steering party A benefits from memory effects—can, under certain conditions, amplify the backflow of information and temporarily strengthen the steering capability of AB over C, even though the steered party C itself is in a Markovian environment. Most notably, the configuration with two non-Markovian photons on the steering side (case 2, long-dashed black line, $\theta_A = \theta_B = \pi/4, \theta_C = 0$) exhibits a dramatically different dynamic

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compared to the previous cases. While it follows the same qualitative pattern of oscillatory behavior, the initial decay is extremely rapid, leading to sudden death after a very short time ($t \approx 5$ ps). However, unlike case 1 where a revival with greater amplitude followed, here the steering shows no evidence of recovery; its value drops below the resolution of the plot and remains effectively zero. This suggests that while non-Markovian environments on the steering parties (A and B) preserve their individual steering over C (as seen in panel a), they completely and irreversibly disrupt the collective steering capability of the pair AB over C, effectively suppressing the oscillatory recovery observed when only one steering party was non-Markovian. The anomalous case from panel (b) ($\theta_A = 0, \theta_B = \pi/4, \theta_C = 0$, spaced dashed orange line) here exhibits an oscillatory decay, but notably, it does not undergo sudden death within the simulated time frame, maintaining small oscillations without ever reaching zero. This reinforces that a non-Markovian environment on a non-steering party (B) introduces memory effects that induce oscillatory behavior in the steering dynamics of $AB \to C$, yet without causing the complete death of steering observed in other configurations. In summary, the dynamics of Gaussian steering in a dephasing W-state are highly sensitive to both the number of photons coupled to non-Markovian environments and the specific partition of the system being considered. For the S^{A-BC} steering measure, non-Markovian effects can induce oscillatory behavior characterized by death-revival cycles, with the intervals of sudden death and revival amplitudes depending critically on how many photons share the memory effects. However, when the steering party itself is Markovian, placing a non-Markovian environment on a steered party leads to oscillatory decay without complete sudden death, yet accelerates the overall loss of steering capability. For the S^{AB-C} steering measure, non-Markovian environments can produce revivals with significantly enhanced amplitudes when asymmetrically distributed, but can also completely suppress oscillatory recovery when both steering parties are non-Markovian, leading to irreversible sudden death. These results highlight the complex interplay between multipartite quantum correlations and environmental memory, demonstrating that non-Markovianity can be a resource for protecting specific types of quantum steering, but its effects are highly configuration-dependent, sometimes enhancing correlations and other times proving detrimental depending on the correlation structure and the distribution of memory effects across the system.

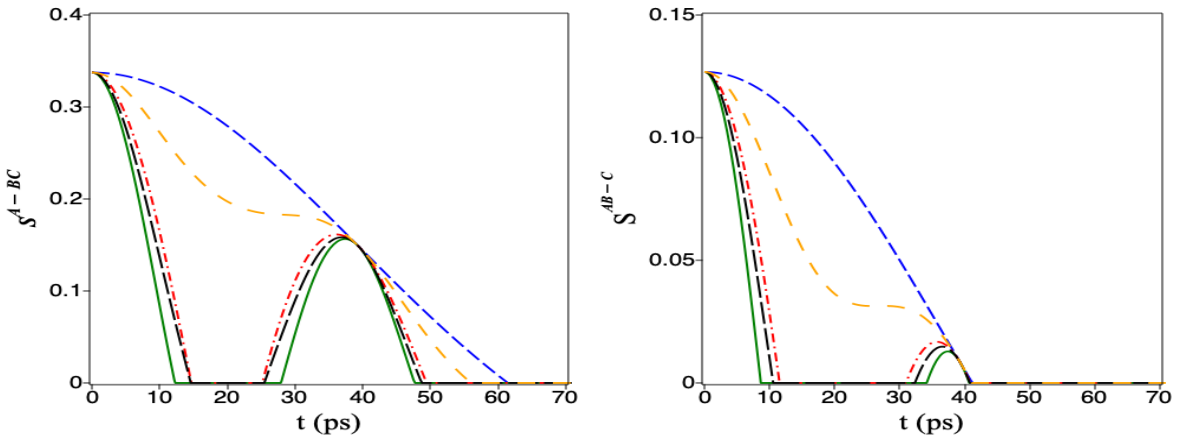

Figure 3. Time evolution of the steering S^{A-BC} (panel (a)) and S^{AB-C} (panel (b)), respectively; for a tripartite photonic system initially prepared in W-state and for different angle values $\theta = 0, \pi/4$, under both Markovian and non-Markovian dephasing dynamics. The system parameters are $\Delta n = 0.01$, $\nu_1 = 2.676$ PHz (~ 704.5 nm), $\nu_2 = 2.692$ PHz (~ 700.3 nm), and $\sigma = 1.8$ THz [29]. Subfigures (**a**) and (**b**) correspond W state. Line styles indicate different environmental configurations: dashed blue — fully Markovian symmetric case ($\theta_A = \theta_B = \theta_C = 0$); solid green — fully non-Markovian symmetric case ($\theta_A = \theta_B = \theta_C = \pi/4$); dash dotted red — case 1 with one photon in non-Markovian environment ($\theta_B = 0, \theta_C = 0, \theta_A = \pi/4$); long doshed black — case 2 with two photons in a non-Markovian environment ($\theta_A = \theta_B = \pi/4, \theta_C = 0$); space dashed orange: $\theta_A = 0, \theta_B = \pi/4, \theta_C = 0$. Units used are: petahertz (PHz $= 10^{15}$ Hz) and terahertz (THz $= 10^{12}$ Hz).

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The dynamics of Gaussian steering for a three-mode photonic system, initially prepared in a GHZ-state with parameter $a = 1$ and subject to dephasing, are presented in Figure 4. The figure illustrates the time evolution of two distinct steering measures: S^{A-BC} (panel a), which quantifies the ability of party A to steer the composite subsystem BC, and S^{AB-C} (panel b), which quantifies the ability of the composite system AB to steer party C. The system evolves under various configurations of Markovian and non-Markovian dephasing environments, characterized by the angle parameter θ, with the spectral properties defined by the detuning $\Delta n = 0.01$, central frequencies $\nu_1 = 2.676$ PHz (~ 704.5 nm), $\nu_2 = 2.692$ PHz (~ 700.3 nm), and bandwidth $\sigma = 1.8$ THz [29].

Panel (a) depicts the steering from mode A to the combined modes B and C for the GHZ-state. The dynamics reveal a strong sensitivity to the environmental configuration, with notable differences compared to the W-state case, particularly in the initial values and decay patterns. The fully Markovian symmetric case (dashed blue line, $\theta_A = \theta_B = \theta_C = 0$) exhibits a rapid monotonic decay, with steering undergoing sudden death around $t \approx 45$ ps and remaining zero thereafter. This behavior is characteristic of memoryless environments where information is irreversibly lost. In stark contrast, the fully non-Markovian symmetric case (solid green line, $\theta_A = \theta_B = \theta_C = \pi/4$) displays a distinct oscillatory behavior. After an initial decay, the steering experiences a revival, passing through a small maximum around $t \approx 38$ ps, followed by another decay. Notably, after $t \approx 42$ ps, which the steering ultimately undergoes complete sudden death. This oscillatory behavior, characterized by multiple death-revival cycles, is a direct consequence of the environment's memory, which allows for a recurrent backflow of information to the system, temporarily restoring quantum correlations before their eventual irreversible loss. The configuration with one non-Markovian photon (case 1, dash-dotted red line, $\theta_A = \pi/4, \theta_B = \theta_C = 0$) exhibits a behavior qualitatively similar to the fully non-Markovian symmetric case, with oscillatory decay and revival. However, in this configuration, the amplitude of the revival is notably larger, and the steering maintains a higher value throughout the dynamics compared to the symmetric non-Markovian case. This suggests that when only the steering party A benefits from memory effects, the backflow of information is more efficiently channeled into preserving its steering capability over BC, despite modes B and C being in Markovian environments. The configuration with two non-Markovian photons on the non-steering sides (case 2, long-dashed black line, $\theta_A = \theta_B = \pi/4, \theta_C = 0$) displays yet another distinct dynamic. It follows the same oscillatory pattern but with a much longer-lived revival. After an initial decay to near-zero around $t \approx 30$ ps, the steering experiences a narrow revival that persists with small oscillations up to $t \approx 44$ ps before eventually decaying. This indicates a synergistic effect; coupling more modes to non-Markovian environments enhances the memory-driven recovery process, leading to shorted preservation of the steering shared by A over BC.

Panel (b) shows the steering from the composite system AB to party C for the GHZ-state. The dynamics here are qualitatively different from panel (a), with overall smaller initial values and distinct decay patterns across configurations. The fully Markovian symmetric case (dashed blue line, $\theta_A = \theta_B = \theta_C = 0$) exhibits a monotonic decay, with steering undergoing sudden death around $t \approx 38$ ps and remaining zero thereafter with no revival. The fully non-Markovian symmetric case (solid green line, $\theta_A = \theta_B = \theta_C = \pi/4$) displays a markedly different behavior. After an initial rapid decay, the steering undergoes sudden death around $t \approx 8$ ps and remains zero thereafter, exhibiting no revival. This absence of recovery indicates that, despite the non-Markovian character of the environments, the collective steering capability of AB over C is irreversibly lost. The configuration with one non-Markovian photon on the steering side (case 1, dash-dotted red line, $\theta_A = \pi/4, \theta_B = \theta_C = 0$) exhibits a similar oscillatory pattern to the fully non-Markovian symmetric case, but with notably different characteristics. The initial decay is slower, and the revival amplitude is smaller, with steering maintaining a smaller value throughout the dynamics. However, like the symmetric case, it undergoes sudden death around $t \approx 38$ ps with no further revival. This suggests that while a single non-Markovian environment on the steering party A enhances the steering capability, it cannot prevent the eventual sudden death when the steered party C is in a Markovian environment. Most notably, the

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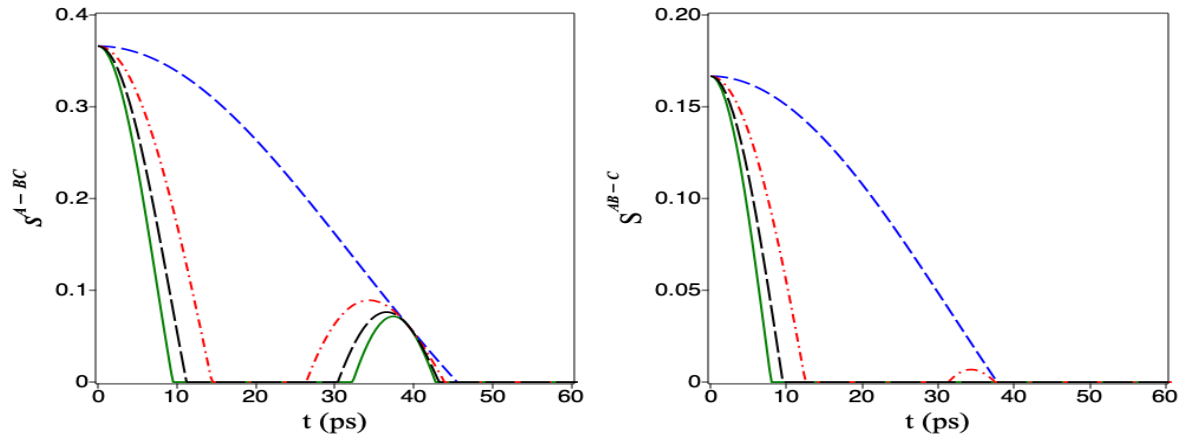
configuration with two non-Markovian photons on the steering side (case 2, long-dashed black line, $\theta_A = \theta_B = \pi/4, \theta_C = 0$) exhibits a dramatically different dynamic. It decays monotonically without any oscillatory behavior, undergoing sudden death around $t \approx 10$ ps and remaining zero thereafter. This suggests that while non-Markovian environments on the steering parties (A and B) preserve their individual steering over C (as seen in panel a), they completely disrupt the oscillatory recovery of the collective steering capability of the pair AB over C, leading to irreversible sudden death.

In summary, the dynamics of Gaussian steering in a dephasing GHZ-state are highly sensitive to both the number of photons coupled to non-Markovian environments and the specific partition of the system being considered. For the S^{A-BC} steering measure, non-Markovian effects can induce oscillatory behavior with multiple death-revival cycles, and in some configurations, lead to prolonged preservation of steering without complete sudden death. In contrast, for the S^{AB-C} steering measure, non-Markovian environments on the steering parties can either induce a single revival before irreversible sudden death or, in the case of two non-Markovian steering parties, result in rapid decay with no revival whatsoever. These results highlight that the efficacy of non-Markovian environments depends critically on which subsystems they are coupled to, relative to the steering and steered parties. The complex interplay between multipartite quantum correlations and environmental memory demonstrates that non-Markovianity can serve as a resource for protecting specific types of quantum steering in GHZ-states, but its effects are highly configuration-dependent, sometimes leading to enhanced preservation and other times resulting in accelerated irreversible loss.



(a) (b)

Figure 4. Time evolution of the steering S^{A-BC} (panel (a)) and S^{AB-C} (panel (b)), respectively; for a tripartite photonic system initially prepared in GHZ-state with ($a = 1$) and for different angle values $\theta = 0, \pi/4$, under both Markovian and non-Markovian dephasing dynamics. The system parameters are $\Delta n = 0.01$, $\nu_1 = 2.676$ PHz (~ 704.5 nm), $\nu_2 = 2.692$ PHz (~ 700.3 nm), and $\sigma = 1.8$ THz [29]. Subfigures (**a**) and (**b**) correspond GHZ state. Line styles indicate different environmental configurations: dashed blue — fully Markovian symmetric case ($\theta_A = \theta_B = \theta_C = 0$); solid green — fully non-Markovian symmetric case ($\theta_A = \theta_B = \theta_C = \pi/4$); dash dotted red — case 1 with one photon in non-Markovian environment ($\theta_B = 0, \theta_C = 0, \theta_A = \pi/4$); long doshed black — case 2 with two photons in a non-Markovian environment ($\theta_A = \theta_B = \pi/4, \theta_C = 0$). Units used are: petahertz (PHz $= 10^{15}$ Hz) and terahertz (THz $= 10^{12}$ Hz

7. Conclusions

In this work, we have systematically investigated the non-Markovian dynamics of quantum steering in a tripartite photonic system, with a particular focus on how environmental memory effects influence the temporal evolution of quantum correlations. We developed a comprehensive theoretical framework beginning with the single-photon dephasing model, which we extended to describe three independent photons interacting with their respective environments. By employing a matrix representation in the polarization basis, we were able to accurately model the dephasing mechanisms affecting each photon and subsequently derive the dynamics for composite three-photon

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systems. We applied this framework to two distinct classes of initial states: the W-type entangled state and the GHZ-type mixed entangled state, examining their evolution under both Markovian and non-Markovian dephasing dynamics. The steering measures S^{A-BC} and S^{AB-C} were employed to quantify the directional quantum correlations between different bipartitions of the tripartite system, providing insight into how steering is distributed and preserved across the network. Our results reveal that the dynamics of Gaussian steering are highly sensitive to both the number of photons coupled to non-Markovian environments and the specific partition of the system being considered. For the W-state, the S^{A-BC} steering measure exhibited oscillatory behavior characterized by death-revival cycles, with the intervals of sudden death and revival amplitudes depending critically on how many photons shared the memory effects. Notably, when the steering party itself was Markovian, placing a non-Markovian environment on a steered party led to oscillatory decay without complete sudden death, yet paradoxically accelerated the overall loss of steering capability. For the S^{AB-C} steering measure in W-states, non-Markovian environments produced revivals with significantly enhanced amplitudes when asymmetrically distributed, but completely suppressed oscillatory recovery when both steering parties were non-Markovian, leading to irreversible sudden death. For the GHZ-state, the dynamics displayed even richer behavior. In the S^{A-BC} case, non-Markovian effects induced multiple death-revival cycles, with some configurations leading to prolonged preservation of steering without complete sudden death. The asymmetric distribution of memory effects proved particularly beneficial when only the steering party benefited from non-Markovian environment, resulting in enhanced revival amplitudes. In contrast, the S^{AB-C} steering measure for GHZ-states showed that non-Markovian environments on the steering parties could either induce a single revival before irreversible sudden death or, in the case of two non-Markovian steering parties, result in rapid decay with no revival whatsoever. These findings demonstrate that the efficacy of non-Markovian environments depends critically on which subsystems they are coupled to, relative to the steering and steered parties. The complex interplay between multipartite quantum correlations and environmental memory reveals that non-Markovianity can serve as a resource for protecting specific types of quantum steering, but its effects are highly configuration-dependent. In some instances, memory effects enhance correlations through information backflow; in others, they prove detrimental depending on the correlation structure and the distribution of memory effects across the system. Our results contribute to the fundamental understanding of open quantum systems and provide practical insights for quantum information processing tasks that rely on the preservation of steering-based quantum correlations. The ability to control and harness non-Markovian effects through engineered environments offers a promising avenue for protecting entanglement and steering in photonic networks, with potential applications in quantum communication, quantum cryptography, and distributed quantum computation. Future work may explore the extension of these results to larger photonic networks, the influence of different spectral densities, and the optimization of environmental parameters to maximize the preservation of quantum correlations in realistic experimental settings.

Author Contributions: Writing—original draft preparation, K.B. and S.B. All authors have read and agreed to the published version of the manuscript.

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