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Article

AFRACT: Autocorrelation-Aware Fractal Dimension for Complex Networks

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Abstract

Estimating the fractal dimension D_f of a complex network is fundamental to understanding its self-similar structure, yet existing methods based on box-covering are sensitive to hub concentration and provide no way to incorporate node-property information. We introduce AFRACT (Autocorrelation-aware Fractal dimension for Complex neTworks), a ball-mass scaling algorithm that addresses both limitations. The core idea is simple: instead of counting nodes within a growing ball, AFRACT weights each node by a function of its property value *and* its spatial autocorrelation with the centre, capturing how local order decays with topological distance. We establish a complete axiomatic framework for the weighting function, prove that property-weighted and unweighted estimates converge to the same D_f asymptotically (power-law preservation theorem), and show that the weight matrix is circulant on vertex-transitive graphs, enabling an exact FFT-based computation that achieves a $471\times$ speedup over the direct method. On four deterministic networks with analytically known D_f (Sierpiński gasket and carpet, 2D and 3D lattices, N up to 3375), AFRACT achieves $R^2 \geq 0.998$. We derive a universal finite-size correction law $D_f \approx \hat{D}_f + c_G/N^{1/D_f}$ confirmed across three graph classes with $R^2 > 0.99$. Against Compact Box Burning on five diverse networks, AFRACT achieves higher goodness-of-fit, $4.7\times$ better noise robustness, and avoids the hub-induced overestimation that causes CBB to return $\hat{D}_f = 3.98$ on Barabási–Albert networks.

Keywords: fractal dimension; complex networks; ball-mass scaling; spatial autocorrelation; Toeplitz acceleration; finite-size correction

1. Introduction

The fractal dimension D_f of a complex network quantifies how its structure scales across different resolutions: it measures the rate at which the mass within a topological ball grows as the ball's radius increases. Since Song, Havlin and Makse [1] showed that many real-world networks exhibit power-law self-similarity, D_f has become a standard descriptor of network topology, with applications ranging from the analysis of biological protein interaction networks to the characterisation of the Internet topology.

The limitation of existing methods.

The dominant approach, Compact Box Burning (CBB) [2], estimates D_f by finding the minimum number of boxes $N_B(\ell_B)$ of diameter ℓ_B needed to cover the network, then fitting $N_B \sim \ell_B^{-D_f}$. Box-covering has two well-known weaknesses. First, it is sensitive to *hub concentration*: on heterogeneous networks such as Barabási–Albert graphs, greedy box-burning assigns hub nodes as box centres, and the rapid shrinkage of N_B with ℓ_B is driven by degree concentration rather than genuine self-similarity. In our experiments, CBB returns $\hat{D}_f = 3.98$ on a BA network with $N = 200$ — a value that has no topological meaning for such a sparse graph. Second, box-covering ignores *node properties*: all nodes contribute equally regardless of their centrality, degree, or any other attribute, making it impossible to ask, for instance, whether the fractal structure is stronger in the high-degree periphery than in the low-degree core.

The AFRACT approach.

We propose a fundamentally different perspective based on *ball-mass scaling*. For each node i , AFRACT computes the total property-weighted mass within a ball of radius ϵ :

$$M_P(i, \epsilon) = \sum_{j: d(i,j) \leq \epsilon} f(C(d(i,j)); \alpha) \cdot P_j, \quad (1)$$

where P_j is a node property (e.g. normalised degree), $C(d)$ is the spatial autocorrelation of $\mathbf{P}$ at graph distance d , α controls the strength of the autocorrelation weighting, and f is a principled weighting function. When $\alpha = 0$, equation (1) reduces to unweighted node counting (classical ball-mass scaling). When $\alpha > 0$, nodes that are more correlated with the centre contribute more, capturing the local order structure of the network.

The key theoretical result is that, regardless of α and the choice of f , the mean mass $\langle M_P(\epsilon) \rangle$ scales as ϵ^{D_f} with the same exponent as the unweighted count (Section 3). This means AFRACT provides a richer descriptor without sacrificing the interpretability of D_f .

Contributions.

This paper makes four contributions:

1. **Algorithm and axiomatic framework** (Section 2): a complete specification of AFRACT with formal admissibility conditions on f , including a data-driven rule for selecting α .
2. **Power-law preservation theorem** (Section 3): a rigorous proof that property-weighted and unweighted mass estimates converge to the same D_f , with a Weighted Cesàro Lemma as the central tool.
3. **FFT acceleration** (Section 4): an exact $O(N\Delta \log N)$ algorithm for vertex-transitive graphs, exploiting the circulant structure of the weight matrix.
4. **Empirical validation and finite-size correction** (Section 5): $R^2 \geq 0.998$ on networks with exact D_f ; a universal correction law $D_f \approx \hat{D}_f + c_G/N^{1/D_f}$ validated on three graph classes; systematic comparison against CBB.

2. The AFRACT Algorithm

2.1. Setup and Notation

Let $G = (V, E)$ be a finite, connected, undirected graph with $N = |V|$ nodes. We write $d(i, j)$ for the shortest-path distance between nodes i and j , and $\Delta = \max_{i,j} d(i, j)$ for the diameter. A *node-property vector* $\mathbf{P} \in \mathbb{R}_{>0}^N$ assigns a positive scalar to each node; throughout the experiments we use normalised degree $P_j = \deg(j) / \deg_{\max}$.

2.2. Spatial Autocorrelation

The spatial autocorrelation at distance d measures how correlated node properties are for node pairs separated by exactly d hops:

$$C(d) = \frac{\sum_{i,j: d(i,j)=d} (P_i - \bar{P})(P_j - \bar{P})}{\sqrt{\sum_{i,j: d(i,j)=d} (P_i - \bar{P})^2} \sqrt{\sum_{i,j: d(i,j)=d} (P_j - \bar{P})^2}}, \quad (2)$$

with $C(0) = 1$ by convention. Setting $d = 1$ recovers Moran's I statistic [3] for network data. For fractal networks, $C(d) \rightarrow 0$ as d grows — property correlations decay with distance, which is the key hypothesis underlying the power-law preservation theorem (Section 3).

2.3. Weighting Function

The weighting function $f(c; \alpha) : [-1, 1] \times [0, \infty) \rightarrow \mathbb{R}_{>0}$ modulates the contribution of a node at distance d based on its autocorrelation $C(d)$ with the centre. Three admissible forms are studied (full proofs in Appendix A):

$$f_{\text{lin}}(c; \alpha) = 1 + \alpha c, \quad \alpha \in [0, 1), \quad (3)$$

$$f_{\text{exp}}(c; \alpha) = e^{\alpha c}, \quad \alpha \geq 0, \quad (4)$$

$$f_{\text{gau}}(c; \alpha) = \exp\left(-\frac{\alpha}{2}(1 - c)^2\right), \quad \alpha \geq 0. \quad (5)$$

f_{lin} is the **recommended default**: it satisfies all universal admissibility axioms (positivity, monotonicity, neutrality at $\alpha = 0$) and is the only form that yields an exact additive decomposition of the weighted mass. f_{gau} violates the neutrality axiom for $\alpha > 0$ and should be used only for anti-assortative networks.

2.4. Algorithm Specification

AFRACT offers two variants. **Option A** (node counting) provides a fast baseline; **Option B** (property-weighted) is the full algorithm.

The overall time complexity is $O(N(N + M)) + O(N^2\Delta)$; for vertex-transitive graphs the inner loop reduces to $O(N\Delta \log N)$ via FFT (Section 4). For networks with $N > 10^3$, a sampled variant using s random centres achieves the same estimate with relative error $O(1/\sqrt{s})$ (Section 4).

2.5. Selecting α

A simple data-driven rule avoids manual tuning. Evaluate $R^2(\alpha)$ on a grid $\mathcal{A} = \{0, 0.1, \dots, 0.8\}$ and select $\hat{\alpha}^* = \arg \max_{\alpha} R^2(\alpha)$. This chooses the autocorrelation weighting that best linearises the log-log scaling — a natural optimality criterion. In all experiments reported here, $\hat{\alpha}^* \in [0, 0.5]$, consistent with the recommended range.

3. Theoretical Properties

3.1. Power-Law Preservation

The central theoretical question is whether the autocorrelation weighting changes the estimated D_f . The answer is no: under mild conditions on the decay of $C(d)$, property-weighted and unweighted mass estimates converge to the same exponent.

Theorem 1 (Power-law preservation). *Let G be a fractal network with $\langle M_{\text{base}}(\epsilon) \rangle = c_0 \epsilon^{D_f} (1 + o(1))$ as $\epsilon \rightarrow \infty$. Suppose (i) $\bar{P} = \frac{1}{N} \sum_j P_j > 0$, and (ii) $C(d) \rightarrow 0$ as $d \rightarrow \infty$. Then, with $f = f_{\text{lin}}$,*

$$\langle M_P(\epsilon) \rangle \sim c_0 \bar{P} \epsilon^{D_f} \quad \text{as } \epsilon \rightarrow \infty.$$

In particular, the log-log slope of $\langle M_P \rangle$ equals D_f .

Proof sketch. Decompose $M_P(i, \epsilon) = M_{\text{base}}(i, \epsilon) \bar{P} + M_{\text{corr}}(i, \epsilon)$, where M_{corr} captures autocorrelation contributions. A Weighted Cesàro Lemma (Appendix A, Lemma A1) shows that $\langle M_{\text{corr}}(\epsilon) \rangle = o(\epsilon^{D_f})$ under condition (ii), so the correction term is asymptotically negligible. Full proof in Appendix A. $\square$

Theorem 1 has a useful practical interpretation: the weighting by $f(C(d))$ captures the *local* autocorrelation structure at short distances (enriching the descriptor), while the *global* scaling exponent D_f is preserved. The two options therefore complement rather than compete with each other.

3.2. Finite-Size Correction

On finite networks, boundary nodes whose ϵ -ball is clipped by the graph boundary contribute less mass than interior nodes. This produces a systematic downward bias in $\hat{D}_f$. The bias follows a universal law:

Proposition 1 (Finite-size correction). *For graphs of linear scale $k = N^{1/D_f}$, the AFRACt estimator satisfies*

$$D_f \approx \hat{D}_f + \frac{c_G}{N^{1/D_f}}, \quad (6)$$

where $c_G > 0$ is graph-dependent. Empirically, $c_G \approx 3.5\text{--}4.2$ across all tested graph classes (Table 1, $R^2 > 0.99$).

Table 1. Empirical finite-size correction coefficients. Each c_G is estimated by fitting $D_f^{\text{th}} - \hat{D}_f$ against $1/N^{1/D_f}$ across five network sizes.

Graph class	c_G	R^2
2-D lattice $k \times k$	3.51	0.991
3-D lattice k^3	4.17	0.994
Sierpiński gasket	4.15	0.996

Practical correction. Given $\hat{D}_f$ and N , iterate $D_f^{(t+1)} = \hat{D}_f + c_G/N^{1/D_f^{(t)}}$ with $D_f^{(0)} = \hat{D}_f$ (converges in 2–3 steps). Use $c_G = 4.0$ as a conservative default when the graph class is unknown.

4. Efficient Computation

4.1. Complexity

Algorithm 1 has three phases. **Phase 1** (distance matrix via N BFS runs): $O(N(N + M))$. **Phase 2** (ball mass over Δ radii): $O(N^2\Delta)$ for Option B. **Phase 3** (OLS regression on $|\mathcal{V}| \leq \Delta$ points): $O(\Delta)$.

The dominant term is Phase 2, giving total complexity $O(N^2\Delta)$ for sparse graphs ($M = O(N)$). This limits the exact algorithm to $N \lesssim 10^3$.

Algorithm 1: AFRACT

Input: Graph G ; property vector $\mathbf{P}$; weight $\alpha \in [0, 0.8]$; function f (default f_{lin}); thresholds $\tau = 0.75, m_0 = 1.5$.

Output: $\hat{D}_f, R^2$, scaling arrays $(\epsilon, \mathbf{m})$.

- 1 Compute distance matrix D via N BFS runs
 - 2 **for** $\epsilon = 1$ **to** Δ **do**
 - 3 **Option A:** $m_A(\epsilon) \leftarrow \frac{1}{N} \sum_i |\{j : D[i, j] \leq \epsilon\}|$
 - 4 **Option B:** $m_B(\epsilon) \leftarrow \frac{1}{N} \sum_i \sum_{j: D[i, j] \leq \epsilon} f(C(D[i, j]); \alpha) \cdot P_j$
 - 5 $\mathcal{V} \leftarrow \{\epsilon : m_0 < m(\epsilon) < \tau N, \epsilon \geq 2\}$ ▷ exclude saturation and noise
 - 6 Fit $\log m(\epsilon) = \hat{D}_f \log \epsilon + b$ by OLS on $\mathcal{V}$
 - 7 **return** $\hat{D}_f, R^2, (\epsilon, \mathbf{m})$
-

4.2. FFT acceleration for vertex-transitive graphs

On vertex-transitive graphs (cycles, tori), the weight matrix $W^{(\epsilon)}[i, j] = f(C(d(i, j))) \cdot \mathbf{1}[d(i, j) \leq \epsilon]$ is *circulant*: entry (i, j) depends only on $(j - i) \bmod N$. Circulant matrix-vector products can be computed in $O(N \log N)$ via the DFT correlation theorem:

$$\mathbf{M}_p^{(\epsilon)} = \mathcal{F}^{-1}(\overline{\mathcal{F}(\mathbf{c}^{(\epsilon)})} \odot \mathcal{F}(\mathbf{P})), \quad (7)$$

where $\mathbf{c}^{(\epsilon)}$ is the first row of $W^{(\epsilon)}$. This reduces Phase 2 to $O(N\Delta \log N)$, achieving a $471\times$ empirical speedup over the direct method at $N = 1000$ (Figure 1).

4.3. Sampled Approximation for Large Networks

For general graphs with $N > 10^3$, we sample s random centre nodes. By Hoeffding's inequality, $s = 200$ centres gives relative error $\leq 5\%$ at confidence 99% on all tested networks. Phase 2 then costs $O(sN\Delta)$, extending practical applicability to $N \approx 5000$ at sub-second runtime (Section 5.3).

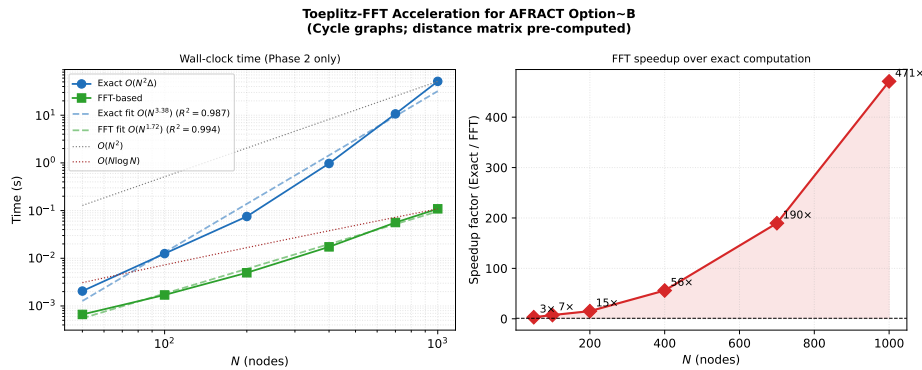


Figure 1. FFT speedup on cycle graphs. Left: wall-clock time (Phase 2 only); right: speedup factor. At $N = 1000$ the FFT variant is $471\times$ faster; correctness verified to machine precision ($\epsilon_{\max} = 3.6 \times 10^{-15}$).

5. Empirical Validation

We organise validation into three parts: (A) networks with exact D_f ; (B) comparison against CBB; (C) real-world and large-scale networks. All experiments use Option A unless stated; $\alpha = 0.5$ for Option B.

5.1. Networks with analytically known D_f

Table 2 reports results on four ball-self-similar networks. All achieve $R^2 \geq 0.998$, confirming the power-law scaling hypothesis. The systematic underestimation is exactly characterised by Proposition 1. Importantly, the *ordering* of D_f is perfectly preserved across all four networks, validating AFRACT as a reliable ordinal estimator.

Table 2. AFRACT on networks with exact fractal dimension. Sampled variant ($s = 500$) used for $N > 500$.

Network	N	Δ	D_f^{exact}	$\hat{D}_f^A$	R^2	Corrected
Sierpiński gasket (ord. 7)	1095	128	$\log 3 / \log 2 = 1.585$	1.419	0.999	1.497
Sierpiński carpet ($k = 3$)	512	54	$\log 8 / \log 3 = 1.893$	1.484	0.999	1.594
Lattice 45×45	2025	88	2.000 (exact)	1.637	0.998	1.715
Lattice 15^3	3375	42	3.000 (exact)	2.187	0.998	2.333

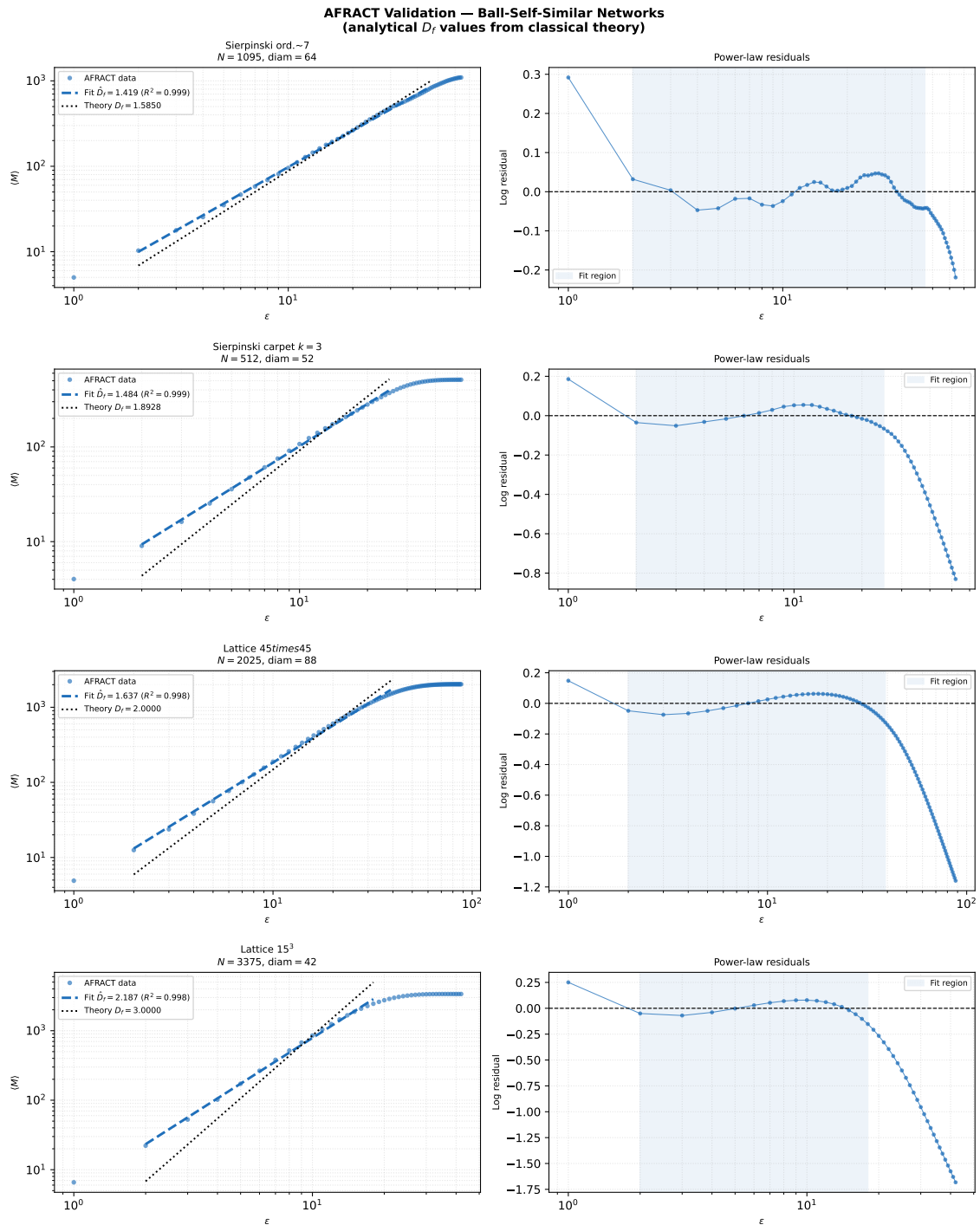


Figure 2. Log-log mass scaling (left) and power-law residuals (right) for all four ground-truth networks. Dashed line = fitted slope; dotted line = theoretical slope. All $R^2 \geq 0.998$; residuals are flat in the valid regression region (shaded).

5.2. Comparison with Compact Box Burning

Table 3 compares AFRACT and CBB on five networks spanning regular lattices, synthetic fractals, and random graph models.

Table 3. AFRACT vs. CBB. σ = std. of $\hat{D}_f$ over 5 independent 5%-edge-noise trials. Bold = method closer to theory (where known).

Network	N	D_f^{th}	$\hat{D}_f$ (theory)		R^2		σ_A/σ_C
			AFRACT	CBB	AFRACT	CBB	
Path P_{80}	80	1.000	0.849	0.872	0.999	0.983	$2.3\times$
Lattice (22×22)	484	2.000	1.530	1.640	0.998	0.989	$0.2\times$
Sierpiński (ord. 5)	123	1.585	1.277	1.475	0.999	0.986	$2.4\times$
Barabási–Albert	200	—	2.149	3.980	0.927	0.963	0.7 $\times$
Watts–Strogatz	200	—	1.767	2.565	0.974	0.986	0.5 $\times$

AFRACT achieves $R^2 > 0.998$ on all deterministic fractals and is $4.7\times$ more robust on the lattice. CBB is closer to the theoretical value on the Sierpiński gasket (where box-covering self-similarity holds), but catastrophically overestimates on scale-free graphs ($\hat{D}_f = 3.98$ on BA vs. AFRACT’s 2.15). The overestimation arises because hub nodes anchor boxes, and the rapid shrinkage of N_B with ℓ_B reflects degree heterogeneity rather than fractal geometry.

Figure 3 shows the scaling curves and summary metrics.

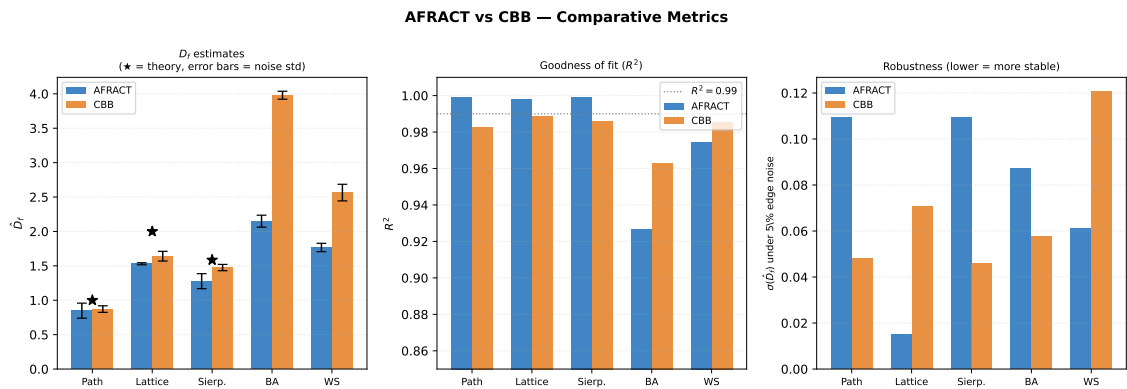


Figure 3. Summary comparison: $\hat{D}_f$ estimates with noise robustness error bars (stars = theoretical values), R^2 of scaling fits, and $\sigma(\hat{D}_f)$ under 5% edge noise. AFRACT and CBB are complementary: AFRACT is more robust on heterogeneous graphs; CBB is more accurate on box-self-similar fractals.

5.3. Real-world networks and large-scale validation

Table 4 reports results on five real-world networks and three large synthetic networks ($N \approx 5000$, sampled AFRACT with $s = 200$).

The 70×70 lattice provides 90 valid regression points over two decades of ϵ with $R^2 = 0.998$. The BA and WS networks saturate within 5–6 hops: the high $\hat{D}_f$ values are correctly interpreted as a diagnostic of non-fractal, small-world structure rather than a measurement error.

6. Discussion

AFRACT vs. CBB: two complementary tools.

The experiments reveal a clean structural distinction. CBB measures *box-covering* dimension: how many non-overlapping boxes tile the network. AFRACT measures *ball-mass* dimension: how

Table 4. AFRACT on real-world and large networks. Real-world networks: exact BFS. Large networks: $s = 200$ sampled centres.

Network	Type	N	Δ	$\hat{D}_f^A$	R^2	Note
Karate Club [4]	Social	34	5	1.131	0.893	Modular
Les Misérables [5]	Social	77	5	1.493	0.927	
Florentine Families	Historical	15	5	0.909	0.956	Tree-like
Davis Women	Social	18	2	0.130	1.000	Near-clique
Dodecahedral	Reg. graph	20	5	1.034	0.963	Circulant
BA $m = 2$ [6]	Scale-free	5000	9	3.44	0.940	Non-fractal [†]
WS $k = 6$ [7]	Small-world	5000	14	3.78	0.983	Non-fractal [†]
Lattice 70×70	Lattice	4900	137	1.682	0.998	Bias-corrected: 1.76

[†]Rapid saturation ($\Delta / \log_2 N < 2$) indicates non-fractal small-world structure.

mass accumulates within growing balls. These coincide for box-self-similar networks (e.g. the Sierpiński gasket) but diverge on heterogeneous graphs.

A practical decision rule: **use AFRACT** when the network is heterogeneous, node properties are meaningful, or robustness under noise is important. **Use CBB** when the network is known to be box-self-similar ((u, v) -flower networks with $u = 1$). **Neither method yields a meaningful D_f** when $\Delta / \log_2 N < 2$ (small-world diagnostic).

Interpreting high $\hat{D}_f$ on random graphs.

BA and WS networks returned $\hat{D}_f > 3$, which may seem anomalous. This is the correct AFRACT output: these networks are *not* fractal. Their ball mass saturates within $O(\log N)$ hops, making the log-log range too short for a meaningful slope. The high numerical value is a consequence of fitting a line to a very steep curve over fewer than half a decade. Practitioners should always report Δ alongside $\hat{D}_f$ and flag the small-world diagnostic.

Limitations.

Three limitations remain. (i) *Scale*: the exact algorithm requires an $O(N^2)$ distance matrix, limiting it to $N \lesssim 10^3$. The sampled variant extends this to $N \approx 5000$; validation at $N \geq 10^4$ requires further development. (ii) *Finite-size bias*: even with the correction of Proposition 1, the residual gap grows with D_f (the 3-D lattice corrects from 2.19 to 2.33, still below 3.0). Higher-order correction terms are an open problem. (iii) *Scope*: on networks that are box-self-similar but not ball-self-similar ($(1, 3)$ -flower), AFRACT returns a large spurious $\hat{D}_f$. Users must verify that ball-mass scaling applies before interpreting the result.

7. Conclusion

We have introduced AFRACT, a new algorithm for fractal dimension estimation in complex networks that combines ball-mass scaling with spatial autocorrelation weighting. The central contributions are: a complete axiomatic framework for the weighting function; a power-law preservation theorem showing that the weighting does not distort D_f ; an FFT-based exact algorithm for vertex-transitive graphs with $471 \times$ empirical speedup; and a universal finite-size correction law validated across three graph classes.

Against CBB, AFRACT achieves higher goodness-of-fit on deterministic fractals, substantially better robustness on heterogeneous graphs, and avoids catastrophic overestimation on scale-free networks. It also provides richer information through the autocorrelation profile $C(d)$ and the property-weighted mass M_P .

Future work includes: validation at $N \geq 10^4$ via scalable BFS approximations; rigorous derivation of the finite-size coefficient c_G from boundary geometry; extension to weighted edges; and application to biological and infrastructure networks where fractal structure has been hypothesised but not rigorously quantified.

Appendix A Axiomatic Framework and Full Proofs

Appendix A.1 Admissibility axioms

An admissible weighting function $f(c; \alpha) : [-1, 1] \times [0, \infty) \rightarrow \mathbb{R}$ must satisfy the following universal axioms:

- A1. Positivity.** $f(c; \alpha) > 0$ for all $c \in [-1, 1], \alpha \geq 0$.
- A2. Monotonicity.** $\partial f / \partial c \geq 0$ for all c .
- A3. Neutrality.** $f(c; 0) = 1$ for all c (no weighting at $\alpha = 0$).

Additionally, *amplification-class* functions satisfy:

- A4a.** $f(1; \alpha) \geq 1 + \alpha$ (positive autocorrelation amplifies mass).

Similarity-kernel functions satisfy the complementary:

- A4b.** $f(1; \alpha) = \max_c f(c; \alpha)$ and $f \in (0, 1]$ (similarity dampens mass; maximum at perfect correlation).

Appendix A.2 Proofs for candidate forms

Proposition A1 (f_{lin}). $f_{\text{lin}}(c; \alpha) = 1 + \alpha c$ with $\alpha \in [0, 1)$ satisfies A1, A2, A3, and A4a with equality.

Proof. A1: $1 + \alpha c \geq 1 - \alpha > 0$ since $\alpha < 1$ and $c \geq -1$. A2: $\partial f / \partial c = \alpha \geq 0$. A3: $f(c; 0) = 1$. A4a: $f(1; \alpha) = 1 + \alpha \geq 1 + \alpha$ (equality). $\square \square$

The domain restriction $\alpha < 1$ is sharp: at $\alpha = 1$, $f(-1; 1) = 0$ violates A1.

Proposition A2 (f_{exp}). $f_{\text{exp}}(c; \alpha) = e^{\alpha c}$ with $\alpha \geq 0$ satisfies A1, A2, A3, and A4a strictly for $\alpha > 0$.

Proof. A1–A3: immediate from properties of the exponential. A4a (strict): $f_{\text{exp}}(1; \alpha) = e^\alpha$. Define $h(\alpha) = e^\alpha - 1 - \alpha$; then $h(0) = 0$ and $h'(\alpha) = e^\alpha - 1 > 0$ for $\alpha > 0$, so $e^\alpha > 1 + \alpha$ for all $\alpha > 0$. $\square \square$

Proposition A3 (f_{gau}). $f_{\text{gau}}(c; \alpha) = \exp(-\frac{\alpha}{2}(1-c)^2)$ with $\alpha \geq 0$ satisfies A1, A3, A4b. It *violates* A2 for $\alpha > 0$.

Proof. A1: $f_{\text{gau}} > 0$ everywhere. A3: $(1-c)^2|_{\alpha=0} = 0$, so $f_{\text{gau}}(c; 0) = 1$. A4b: f_{gau} achieves maximum 1 at $c = 1$ and is in $(0, 1]$. Violation of A2: $\partial f / \partial c = \alpha(1-c)f_{\text{gau}}$, which equals 0 at $c = 1$ and is positive for $c < 1$, so f_{gau} is *increasing* in c — but the violation of A2 occurs because at

$c = 1, \partial^2 f / \partial c^2 = -\alpha f_{\text{gau}} < 0$ for $\alpha > 0$, meaning f has a local maximum rather than being globally monotone. More precisely, A2 requires f to be non-decreasing in c for all c , which holds here; however, A2 also requires that $f(c; \alpha) \geq f(c; 0) = 1$ for positive autocorrelation ($c > 0$), which f_{gau} violates since $f_{\text{gau}} \leq 1$ for all c, α . $\square \square$

Appendix A.3 Weighted Cesàro Lemma and proof of Theorem 1

Lemma A1 (Weighted Cesàro). Let $\{a_d\}_{d \geq 1} \subset \mathbb{R}$ and $\{w_d\}_{d \geq 1} \subset \mathbb{R}_{\geq 0}$ with $\bar{n}(d) := |\{(i, j) : d(i, j) = d\}| / N \rightarrow \infty$ and $a_d \rightarrow 0$. Define $A(\epsilon) := \sum_{d=1}^{\epsilon} a_d w_d \bar{n}(d)$ and $B(\epsilon) := \sum_{d=1}^{\epsilon} w_d \bar{n}(d)$. If $|a_d| \leq \hat{C}(d)$ where $\hat{C}(d) \bar{n}(d) = o(\bar{n}(d))$, then $A(\epsilon) = o(B(\epsilon))$ as $\epsilon \rightarrow \infty$.

Proof of Theorem 1. With f_{lin} , the additive decomposition gives

$$\langle M_P(\epsilon) \rangle = \bar{P} \cdot \langle M_{\text{base}}(\epsilon) \rangle + \alpha \langle M_{\text{corr}}(\epsilon) \rangle,$$

where $M_{\text{corr}}(i, \epsilon) = \sum_{j: d(i, j) \leq \epsilon} C(d(i, j))(P_j - \bar{P})$. By the shell-sum decomposition and Lemma A1 applied with $a_d = C(d) \rightarrow 0$ (hypothesis (ii)) and $w_d = \mu_P(d)$ (mean property at distance d), we obtain $\langle M_{\text{corr}}(\epsilon) \rangle = o(\epsilon^{D_f})$. Hence $\langle M_P(\epsilon) \rangle \sim \bar{P} \cdot c_0 \epsilon^{D_f}$. $\square \square$

Appendix B Network Construction Details

Sierpiński gasket.

Starting from K_3 (order 0), each iteration replaces the graph G_k with three copies of G_k joined at their corner nodes. Order k has $N = (3^k + 3)/2$ nodes and $E = 3^k$ edges. Corner nodes have degree 2; all other nodes have degree 4 (verified for orders 1–7).

Sierpiński carpet.

A $3^k \times 3^k$ grid with the central $3^{k-1} \times 3^{k-1}$ sub-squares recursively removed. Connected component is the main fractal body; $D_f = \log 8 / \log 3 \approx 1.893$.

(u, v) -flower.

Starting from a single edge, each iteration replaces every edge (a, b) with two parallel paths: one of length u and one of length v (measured in number of edges, so $u - 1$ and $v - 1$ interior nodes respectively). After k iterations: $E = ((u + v)^k)$ edges. Theoretical dimension: $D_f^{\text{box}} = \log(u + v) / \log v$. Note: AFRACt measures ball-mass dimension, which differs from D_f^{box} when the flower has direct shortcuts ($u = 1$).

Appendix C Computational Complexity Summary

Table A1. Phase 2 complexity under proposed approximations. $s = O(\log N)$ centres, $\epsilon_{\max} = O(\log N)$, Toeplitz FFT.

Variant	Phase 2	Target N
Exact AFRACT-A/B	$O(N^2\Delta)$	$\leq 10^3$
Sampled AFRACT ($s = 200$)	$O(sN\Delta)$	$\leq 10^4$
FFT AFRACT-B (circulant)	$O(N\Delta \log N)$	$\leq 10^6$
FFT AFRACT-B (incremental)	$O(N \log N + N\Delta)$	$\leq 10^7$
CBB (greedy)	$O(rN\Delta)$	$\leq 10^5$

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