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Article

Detecting Intermediate-Mass Black Holes out to 20 Mpc with ELT/HARMONI: The Case of FCC 119

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Abstract

Intermediate-mass black holes (IMBHs; $M_{\rm BH} \approx 10^{3-5} {\rm M}_{\odot}$) play a critical role in understanding the formation of supermassive black holes in the early universe. In this study, we expand on Nguyen et al. simulated measurements of IMBH masses using stellar kinematics, which will be observed with the High Angular Resolution Monolithic Optical and Near-infrared Integral (HARMONI) field spectrograph on the Extremely Large Telescope (ELT) up to the distance of 20 Mpc. Our sample focuses on both the Virgo Cluster in the northern sky and the Fornax Cluster in the southern sky. We begin by identifying dwarf galaxies hosting nuclear star clusters, which are thought to be nurseries for IMBHs in the local universe. As a case study, we conduct simulations for FCC 119, the second faintest dwarf galaxies in the Fornax Cluster at 20 Mpc, which is also fainter than most of Virgo Cluster members. We use the galaxy's surface brightness profile from Hubble Space Telescope (HST) imaging, combined with an assumed synthetic spectrum, to create mock observations with the HSIM simulator and Jeans Anisotropic Models (JAM). These mock HARMONI datacubes are analyzed as if they were real observations, employing JAM within a Bayesian framework to infer IMBH masses and their associated uncertainties. We find that ELT/HARMONI can detect the stellar kinematic signature of an IMBH and accurately measure its mass for $M_{\rm BH} \gtrsim 10^5~{\rm M}_{\odot}$ out to distances of $\sim 20~{\rm Mpc}$.

Keywords: galaxies; individual; FCC 119 – galaxies; supermassive black holes – galaxies; nuclei – galaxies; kinematics; dynamics – galaxies; evolution – galaxies; formation

1. Introduction

Intermediate-mass black holes (IMBHs; $M_{BH} \approx 10^{3-5} \rm M_{\odot}$) in dwarf galaxies ($M_{\star} \lesssim 10^9 \rm M_{\odot}$) represent a crucial but largely missing component of the black hole (BH) mass spectrum [60,90]. Studying the correlations between IMBH mass and the macroscopic properties of their host galaxies—such as central stellar velocity dispersion (σ_{\star}) [45,51] or bulge stellar mass ($M_{\rm bulge}$) [75]—provides valuable insights when compared to their supermassive black hole (SMBH; $M_{\rm BH} \approx 10^{6-10} \rm M_{\odot}$) counterparts in more massive galaxies ($M_{\star} \gtrsim 10^{10} \rm M_{\odot}$). Additionally, determining the occupation fraction—the proportion of dwarf galaxies hosting IMBHs—can constrain potential formation pathways of these black holes in the early universe [61]. Despite extensive searches [61], the presence of IMBHs remains elusive, and their origins continue to be an open question in astrophysics [69,94].

Despite the fact that dynamical methods using stellar or gas motions in the centers of nearby dwarf galaxies are among the most precise techniques for searching for strong evidence of IMBHs [1,34,36,96,97,99–105,131,132,135,148], they remain observationally challenging due to the insufficient spatial resolution of current facilities [99,122] to resolve the black hole's sphere of influence (SOI), defined as $r_{SOI} = GM_{BH}/\sigma_{\star}^2$, where G is the gravitational constant. For example, for an IMBH

with $M_{\rm BH} \approx 10^5~\rm M_{\odot}$ and $\sigma_{\star} \approx 30~\rm km~s^{-1}$ at the distance of the Local Group ($D \approx 3.5~\rm Mpc$), the corresponding $r_{\rm SOI} \approx 0''.028$ [e.g., see Eq. 1 of 98], which is smaller than the diffraction limits of current 8–10 meter class ground-based telescopes equipped with adaptive optics (AO), which achieve point-spread functions (PSFs) with typical full widths at half maximum (FWHM) in the range of FWHM_{PSF} $\approx 0''.05$ –0''.1. This limitation raises the question of whether the current lack of strong observational evidence for IMBHs reflects their true absence, or simply the restrictions imposed by existing instrumentation.

The next generation of giant ground-based telescopes may provide a solution to this problem. In particular, the Extremely Large Telescope (ELT) is poised to overcome current resolution challenges with its 39 meter primary mirror, equipped with the Multi-AO Imaging Camera for Deep Observations (MICADO; FWHM_{PSF} \approx 0'.'01) imager [32,33] and the High Angular Resolution Monolithic Optical and Near-infrared Integral (HARMONI; FWHM_{PSF} \approx 0'.'012) field spectrograph, which also offers a high-spectral resolution ($R \approx 17,400$) [136,137]. However, to locate potential IMBHs, we must first identify suitable targets. According to [106], bright nuclear star clusters (NSCs) within 10 Mpc are promising hosts of IMBHs. These IMBHs can be dynamically detected if their masses are at least \sim 0.5% of their NSC mass ($M_{\rm NSC}$), although the effects of dynamical mass segregation must be carefully considered before drawing strong conclusions.

In this study, we extend the distance limit of the HARMONI IMBH survey [106] to include the Virgo ($D \approx 16.5$ Mpc) and Fornax ($D \approx 20$ Mpc) clusters by simulating $M_{\rm BH}$ measurements for FCC 119, the second faintest nucleated member of the Fornax Cluster and significantly fainter than nearly all members of the Virgo Cluster. These measurements are derived from high-spatial-resolution integral-field unit (IFU) stellar kinematics extracted from mock HARMONI data, generated using the HARMONI Simulator (HSIM¹) [155]. We also investigate the observational limits in terms of the minimum NSC central surface brightness (SB), IMBH mass, and on source exposure time that HARMONI can effectively detect.

We describe the criteria used to finalize the extended HARMONI IMBH sample and its properties, as well as the characteristics of the Virgo and Fornax Clusters, which contain a significant portion of our sample, in Section 2. In Section 3, we present the stellar mass model of FCC 119. Section 4 provides a detailed description of our HARMONI IFS simulations. Finally, we discuss the results in Section 5 and summarize our findings in Section 6.

We adopt a flat Universe with a Hubble constant of $H_0 \approx 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, a matter density of $\Omega_{\text{m},0} \approx 0.3$, and a dark energy density of $\Omega_{\Lambda,0} \approx 0.7$. All photometric measurements are reported in the AB magnitude system [108] and corrected for foreground extinction [119] using a interstellar extinction law [24]. Furthermore, all kinematic maps are aligned such that the galaxy's major and minor axes correspond to the horizontal and vertical directions, respectively.

2. Sample Selection

2.1. Sample Selection

We compiled a list of known nucleated dwarf galaxies from previous NSC surveys, including both photometry and spectroscopy [5,13,22,42,43,53,54,56,58,59,65,67,95,112,115,126], to create a parent sample for selecting the best candidates likely to host IMBHs beyond the current HARMONI IMBH sample within 10 Mpc [106], extending the survey out to 20 Mpc. The selection criteria are listed in the upper part of Table 1.

v3.11: https://github.com/HARMONI-ELT/HSIM

Table 1. Selected criteria for dwarf galaxies

 $\begin{array}{ccc} \text{ELT Observability:} & |\delta + 24^{\circ}| < 45^{\circ} \\ \text{Galaxy zone of avoidance:} & |b| > 8^{\circ} \\ \text{Central stellar-velocity dispersion:} & \sigma_{\star,c} < 75 \text{ km s}^{-1} \end{array}$

2.2. Sample Properties

We present in Figure 1 our selected targets on the distance–magnitude plane and compare them with previous galaxy surveys, including MMBH [98], MASSIVE [84], ATLAS^{3D} [22], and the HARMONI IMBH survey [106]. We calculated the absolute magnitude in the K_s -band using the relation $M_K = K_s - 5\log_{10}D - 25 - 0.11A_V$, calibrated from a sample of 260 elliptical galaxies within 48 Mpc in the ATLAS^{3D} project [22]. Here, K_s magnitudes are sourced from the NED-D² and HyperLeda³ databases, while A_V represents the Galactic extinction in the Landolt V-band [119], adopting the [25] reddening relation with $R_V = A_V / E(B - V) = 3.1$.

We summarized in Table 2 our expanded HARMONI IMBH sample, spanning a distance range of $10 \lesssim D \lesssim 20$ Mpc and an absolute K_s -band magnitude range of $-23.7 \lesssim M_K/\text{mag} \lesssim -17.4$. It overlaps with the ATLAS^{3D} survey by 2.2 mag ($M_K < -21.5$ mag) and focuses more on lower-mass galaxies. Additionally, our sample is distinct from the 10 Mpc HARMONI IMBH sample [106], although both samples cover nearly the same stellar mass range of 10^8 to 5×10^{10} M $_\odot$. Our selected galaxies also exhibit central stellar velocity dispersions ranging from 16 to 75 km s $^{-1}$.

Table 2. Properties of the selected sample

Distant range: $10 < D \lesssim 20 \,\mathrm{Mpc}$ Galaxies K_s -band absolute magnitude: $-23.7 \lesssim M_K/\mathrm{mag} \lesssim -17.4$ Galaxy stellar mass range: $10^8 \lesssim M_{\star,\mathrm{gal}}/M_{\odot} \lesssim 5 \times 10^{10}$ Effective: $R_{\mathrm{e,gal}} \lesssim 7 \,\mathrm{kpc}$ NSC mass range: $1.6 \times 10^5 < M_{\mathrm{NSC}}/M_{\odot} < 2.3 \times 10^8$ NSC effective radius: $0.5 < R_{\mathrm{e}} < 62 \,\mathrm{pc}$

This expanded HARMONI IMBH sample comprises 85 galaxies, including 29 and 20 members of the Virgo (prefix VCC) and Fornax (prefix FCC) Clusters, respectively, along with 36 isolated dwarf galaxies or members of other groups, as detailed in Table 4. It spans a diverse range of galaxy morphologies, consisting of 27% ellipticals, 29% lenticulars, 39% spirals, and 5% irregulars. These galaxies host luminous NSCs with masses in the range $1.6 \times 10^5 < M_{\rm NSC} < 2.3 \times 10^8 \, {\rm M}_{\odot}$ and effective radii spanning $0.5 < R_{\rm e} < 62$ pc, as shown in Figure 2 for the M_{\star} – $M_{\rm NSC}$ and M_{\star} – $R_{\rm e}$ relationships.

NASA/IPAC Extragalactic Database: https://ned.ipac.caltech.edu

³ HyperLeda: https://leda.univ-lyon1.fr

Table 3. Full list and properties of our extended HARMONI IMBH sample hosting NSCs within the distant range of 10–20 Mpc

No.	Object	RA	Decl.	D	Hubble	$M_{\star,\mathrm{gal}}$	R _{e,gal}	M _{NSC}	R _{e,NSC}	$\sigma_{\star, \mathrm{NSC}}$	Ref.
140.	name	(h:m:s)	(d:m:s)	(Mpc)	type	(M_{\odot})	(kpc)	(M_{\odot})	(pc)	(km s^{-1})	TCI.
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)
1	ESO301-IG11	03:23:54.21	-37:30:33.0	19.94	9.9	9.61	2.11	5.79	3.7	28.9	[54]
2	IC1933	03:25:39.97	-52:47:07.5	18.98	6.1	8.35	1.85	5.21	6.2	42.7	[54]
3	MCG-1-03-85	01:05:04.88	-06:12:44.6	11.38	7.0	9.66	4.04	6.40	2.9	42.0	[5,54]
4	NGC0428	01:12:55.75	+00:58:53.5	16.54	8.6	9.54	2.63	6.41	1.2	56.9	[54]
5	NGC0864	02:15:27.64	+06:00:09.4	19.41	5.1	10.10	0.21	7.86	2.7	26.9	[54,65]
6	NGC1042	02:40:23.97	-08:26:00.7	19.59	6.0	9.32	4.59	6.33	1.3	65.3	[54]
7	NGC1073	02:43:40.52	+01:22:34.0	13.80	5.3	9.69	0.16	6.53	0.5	24.8	[54,65]
8	NGC1493	03:57:27.46	-46:12:38.6	15.54	6.0	9.77	4.85	6.55	3.6	51.2	[54]
9	NGC1518	04:06:49.82	-21:10:23.5	13.50	8.2	9.10	0.83	5.35	3.4	36.6	[54]
10	NGC1559	04:17:35.77	-62:47:01.2	15.21	5.9	9.77	3.50	6.66	1.5	72.6	[5,54,130]
11	NGC2835	09:17:52.91	-22:21:16.8	10.86	5.0	9.70	2.80	6.61	3.4	70.7	[5,54,130]
12	NGC3346	10:43:38.91	+14:52:18.9	17.84	5.9	9.51	3.34	6.37	1.1	49.6	[54]
13	NGC3423	10:51:14.33	+05:50:24.1	11.27	6.0	9.45	0.12	6.29	2.2	54.6	[54]
14	NGC3455	10:54:31.08	+17:17:04.6	15.79	3.7	10.57	1.44	6.81	2.2	46.2	[54]
15	NGC3666	11:24:26.07	+11:20:32.0	19.32	5.2	9.94	0.10	7.81	9.6	60.6	[54,65]
16	NGC4212	12:15:39.36	+13:54:05.4	17.30	4.9	10.40	4.56	6.70	12.3	61.0	[54,59]
17	NGC4487	12:31:04.46	-08:03:14.1	19.59	6.0	10.38	3.50	6.93	1.1	51.0	[54]
18	NGC4496A	12:31:39.21	+03:56:22.1	15.49	7.6	9.65	3.10	6.23	0.7	74.9	[54,65,130]
19	NGC4504	12:32:17.41	-07:33:48.9	14.51	6.0	10.54	2.10	6.85	3.9	53.0	[54]
20	NGC4517	12:32:45.59	+00:06:54.1	10.67	6.0	10.08	6.50	5.85	2.3	43.8	[54,65]
21	NGC4592	12:39:18.74	-00:31:55.0	15.21	8.0	10.28	1.84	5.80	1.1	42.6	[54]
22	NGC4635	12:42:39.23	+19:56:43.7	13.54	6.6	9.01	1.40	6.79	4.3	34.7	[54]
23	NGC4771	12:53:21.26	+01:16:09.5	16.07	6.2	10.32	2.73	6.63	1.9	64.4	[54]
24	NGC4900	13:00:39.26	+02:30:02.7	13.73	5.2	10.12	2.59	7.24	3.1	67.1	[54]
25	NGC4904	13:00:58.64	-00:01:39.4	16.94	5.8	9.92	2.06	6.23	2.5	62.5	[54]
26	NGC5054	13:16:58.49	-16:38:05.5	18.62	4.2	10.43	0.11	7.02	8.3	57.4	[5,54]
27	NGC5300	13:48:16.04	+03:57:03.1	16.74	5.2	9.97	2.68	7.21	20.6	54.1	[54]
28	NGC5334	13:52:54.48	-01:06:52.0	19.63	5.2	9.79	2.96	6.85	11.9	49.5	[54]
29	NGC7090	21:36:28.81	-54:33:25.3	11.86	5.0	9.64	4.56	6.48	1.1	54.3	[54]
30	NGC7424	22:57:18.37	-41:04:14.1	11.48	6.0	9.75	-	6.00	6.8	15.6	[54,150]
31	NGC7713	23:36:14.99	-37:56:17.1	10.29	6.7	9.90	2.17	6.06	1.4	42.4	[54]
32	UGC08041	12:55:12.65	+00:06:59.9	17.14	6.9	8.83	_	6.74	10.3	34.4	[5,54]
33	UGC08516	13:31:52.59	+20:00:04.6	14.50	5.9	8.95	1.29	6.37	11.2	34.1	[54]
34	UGC09215	14:23:27.13	+01:43:34.4	19.96	6.3	9.77	2.49	7.24	0.8	41.3	[54]
35	VCC0033	12:11:07.76	+14:16:29.3	15.00	-4.0	9.62	0.72	6.86	7.2	38.9	[28,80]
36	VCC0140	12:15:12.56	+14:25:58.4	17.15	-2.2	9.79	2.10	7.09	7.8	48.7	[28,80]
37	VCC0437	12:20:48.82	+17:29:13.4	16.90	-3.2	9.90	2.96	6.94	61.8	70.7	[28,80]
38	VCC0543	12:22:19.53	+14:45:38.8	16.04	-2.7	9.74	2.19	6.95	10.7	19.2	[28,46,58,80,112,113]
39	VCC0698	12:24:05.02	+11:13:05.0	18.71	-2.0	9.98	1.45	6.41	-	62.0	[13,46,58,113]
40	VCC0751	12:24:48.36	+18:11:42.4	15.59	-2.7	9.63	1.79	6.98	22.0	65.0	[28,80]
41	VCC0856	12:25:57.93	+10:03:13.6	16.83	-3.3	9.35	1.22	5.93	_	33.0	[46,48,58,113]
42	VCC0939	12:26:47.23	+08:53:04.6	16.83	6.3	9.48	4.30	6.55	4.8	45.8	[5,54]
43	VCC1087	12:28:14.88	+11:47:23.4	16.67	-4.1	9.51	1.53	6.03	_	45.0	[13,46,58,113,124]

Table 4. Full list and properties of our extended HARMONI IMBH sample hosting NSCs within the distant range of 10–20 Mpc

44	VCC1125	12:28:43.31	+11:45:18.1	16.04	-1.9	9.93	7.15	7.32	38.1	65.3	[28,80]
45	VCC1192	12:29:30.25	+07:59:34.3	16.50	-4.8	9.27	0.60	5.79	_	58.2	[46,58,113,151]
46	VCC1250	12:29:59.08	+12:20:55.2	17.62	-2.9	9.57	1.35	6.31	_	69.0	[13,46,58,113]
47	VCC1261	12:30:10.33	+10:46:46.1	18.91	-4.8	9.69	1.71	6.23	_	50.0	[58,112,113]
48	VCC1355	12:31:20.19	+14:06:54.7	16.90	-5.0	9.75	1.77	6.80	23.8	64.8	[28,80]
49	VCC1422	12:32:14.21	+10:15:05.3	15.35	-4.8	9.58	1.82	6.13	_	36.0	[13,58,113]
50	VCC1431	12:32:23.39	+11:15:46.7	17.50	-4.8	9.97	2.43	7.11	_	55.0	[28,80]
51	VCC1488	12:33:13.44	+09:23:50.5	18.30	-3.5	9.62	0.70	6.83	1.3	18.6	[28,80]
52	VCC1489	12:33:13.90	+10:55:42.7	18.82	-5.0	9.68	0.34	6.66	5.6	25.5	[28,80]
53	VCC1528	12:33:51.62	+13:19:20.8	15.41	-4.4	9.88	1.84	7.05	7.2	51.9	[28,80]
54	VCC1539	12:34:06.74	+12:44:29.8	16.35	-5.0	9.85	0.69	6.66	32.1	28.5	[28,80]
55	VCC1545	12:34:11.53	+12:02:56.3	17.60	-5.0	9.94	1.33	6.96	11.5	45.6	[28,80]
56	VCC1661	12:36:24.78	+10:23:04.8	17.76	-5.0	9.67	1.08	6.63	51.4	64.9	[28,80]
57	VCC1695	12:36:54.85	+12:31:12.3	16.50	-3.0	9.76	2.05	6.85	10.6	54.5	[28,80]
58	VCC1861	12:40:58.56	+11:11:04.2	15.23	-4.8	9.87	2.54	6.99	60.8	58.6	[28,80]
59	VCC1871	12:41:15.73	+11:23:14.1	15.49	-5.0	9.35	0.52	5.76	_	51.0	[13,46,58,113]
60	VCC1883	12:41:32.75	+07:18:53.6	16.60	-2.0	10.22	2.01	6.59	_	60.8	[46,58,113,151]
61	VCC1895	12:41:51.98	+09:24:10.4	15.80	-3.0	9.63	0.87	6.87	2.0	25.6	[28,80]
62	VCC2019	12:45:20.42	+13:41:33.6	17.06	-3.9	9.00	1.27	5.86	_	37.0	[46,58,112,113]
63	VCC2048	12:47:15.30	+10:12:12.9	16.60	-4.5	9.73	2.77	7.08	21.1	67.5	[28,80]
64	VCC2050	12:47:20.64	+12:09:59.1	17.01	-4.0	9.63	0.78	6.83	6.6	26.0	[28,80]
65	FCC100	03:31:47.58	-35:03:06.7	19.48	-2.2	8.70	1.10	6.40	_	35.0	[35,39,129]
66	FCC106	03:32:47.65	-34:14:19.3	20.00	-2.8	8.90	1.00	6.70	_	40.0	[114,129,141]
67	FCC119	03:33:33.84	-33:34:23.9	20.00	-2.9	9.00	1.70	6.81	9.7	20.0	[42]
68	FCC135	03:34:30.86	-34:17:51.0	18.73	-3.0	8.70	1.41	6.30	_	24.0	[39,110,129]
69	FCC136	03:34:29.48	-35:32:47.0	18.80	-2.4	9.10	1.29	6.60	_	64.3	[35,129]
70	FCC148	03:35:16.82	-35:15:56.4	20.00	-2.2	9.76	2.70	8.37	22.0	56.0	[42,83,146]

A key characteristic of our sample is its predominance of more massive NSCs, which are larger in both mass and size compared to those in the previous HARMONI IMBH sample [106]. This distinction arises because our survey extends to greater distances, leading to the omission of some smaller and fainter NSCs present in previous surveys. Despite our thorough efforts to select the most suitable targets from all available databases and literature, our sample represents only a small subset of the total potential candidates. A comprehensive, high-resolution, and deep imaging survey remains necessary to construct a complete census of NSCs and uncover the elusive IMBH population in nearby nucleated dwarf galaxies.

2.3. Searching Tip-Tilt Stars and Observing Strategies

The AO system is crucial for achieving high spatial resolution observations with the ELT. In LTAO mode, the system is supported by laser guide stars (LGS) and operates with a natural guide star (NGS) that can be up to 10,000 times fainter than those required by the Gemini and VLT AO systems. The NGS serves as a reference star to correct tip/tilt distortions and must have an H-band magnitude of $m_H < 20.4$ ABmag [136].

Since the ELT will reach the lowest sensitivity limits for NGSs, current NGS search tools are inadequate. To identify suitable NGSs, we utilized the Gaia Data Release 3^4 , which provides an all-sky catalog of $\approx 10^9$ point sources with a magnitude limit of $m_V \lesssim 20$ Vegamag [111]. We employed the CONE_SEARCH_ASYNC function from the astroquery package⁵ [55] to retrieve stellar coordinates and photometric information, which will then converted to the H-band ABmag using the following relation [16]:

$$G - H = a + b(G_{BP} - G_{RP}) + c(G_{BP} - G_{RP})^{2},$$
(1)

Here, the *G*-band magnitude and the integrated blue-to-red passband color ($G_{BP} - G_{RP}$) were retrieved from Gaia, while the coefficients *a*, *b*, and *c* were taken from Table 5.9 of [16].

For each galaxy, we selected one star based on the following criteria: (1) it lies within an annulus of 12'' < r < 60'' from the galaxy center, and (2) it has an H-band magnitude of $m_H < 20.4$ ABmag. The selected stars are marked with white crosses, with their H-band magnitudes labeled next to them in Figure 3.

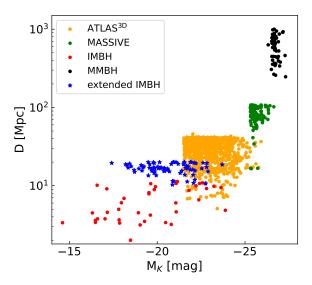


Figure 1. The distribution of our extended-HARMONI IMBH sample on distance vs. *Ks*-band absolute magnitude comparing with the previous surveys: HARMONI MMBH [98], MASSIVE [84], ATLAS^{3D} [22] and HARMONI IMBH [106].

⁴ https://www.cosmos.esa.int/web/gaia/dr3

⁵ https://github.com/astropy/astroquery

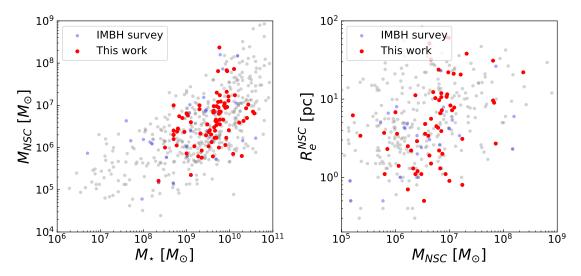


Figure 2. Left: The M_{\star} – $M_{\rm NSC}$ relation shows the correlation between galaxy mass and NSC mass across various Hubble types [41,54,109,116,126]. **Right:** The M_{\star} – R_{e} relation illustrates the connection between galaxy mass and the NSC's effective radius [5,28,53,107]. For comparison, we included the 10 Mpc HARMONI IMBH sample [106].

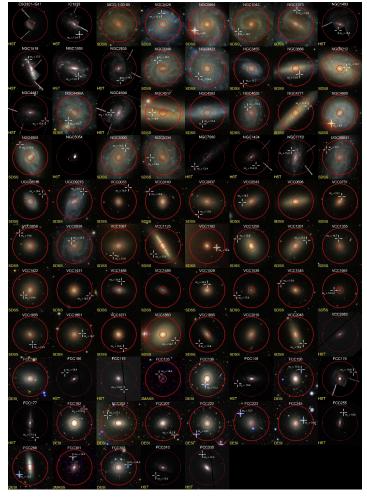


Figure 3. Images from the HST, the Dark Energy Spectroscopic Instrument (DESI) Legacy Surveys DR10, The Two Micron All-Sky Survey (2MASS), and Sloan Digital Sky Survey (SDSS) of our extended HARMONI IMBH sample. Two red circles indicate radii of 12'' (inner) and 60'' (outer) from the galaxy center, marking the optimal region for selecting NGSs for AO correction. These stars are indicated by white crosses, with their respective *H*-band magnitudes estimated from Gaia.

2.4. The Simulated Target: FCC 119

We aim to determine the required exposure time (sensitivity) for observing faint dwarf galaxies with ELT/HARMONI, ensuring to resolve the stellar kinematic signature of central IMBHs and accurately measure their masses dynamically. At \approx 20 Mpc, galaxies in the Fornax Cluster exhibit lower SB and smaller angular sizes compared to those in the Virgo Cluster ($D \approx 16.5$ Mpc), making difficulty in measuring their IMBH. To assess this, we selected FCC 119 ($m_B = 15.44$ mag), the second faintest galaxy in the Fornax Cluster [141], which is fainter than nearly all Virgo Cluster galaxies in our sample [28].

FCC 119 is an S0 dwarf galaxy located at $03^h33^m33^s7-33^\circ34'18''$. Its stellar mass is estimated to be $M_{\star}\approx (1.06-1.4)\times 10^9~{\rm M}_{\odot}$, with an effective radius of $R_{\rm e}=1.33~{\rm kpc}$ [81,125,141], based on r-band observations from OmegaCAM at the ESO/VLT Survey Telescope [VST; 77,118]. The galaxy has a SB of 20.1 mag arcsec $^{-2}$ at a radius of 1" in the z-band [141]. FCC 119 also hosts a bright NSC with a mass of $M_{\rm NSC}\approx 6.5\times 10^6~{\rm M}_{\odot}$ [42], as observed with the HST/Advanced Camera for Surveys (ACS; Figure 4) [73,141].

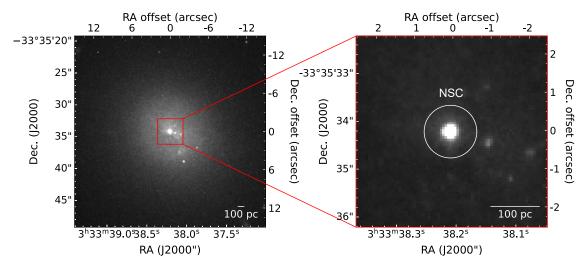


Figure 4. Left: The HST/ACS WFC F850LP image of FCC 119 ($15'' \times 15''$ or $1.5 \times 1.5 \text{ kpc}^2$) shows its large-scale morphology. **Right:** A zoom-in to the central region ($5'' \times 5''$) highlights its NSC.

The galaxy's kinematic inclination was determined from optical data obtained with the VLT/Multi Unit Spectroscopic Explorer (MUSE), yielding $i=64^{\circ}.9\pm5^{\circ}.4$. Stellar kinematic analysis of these data revealed rotation with $V\approx 10$ km s⁻¹ and a velocity dispersion of $\sigma_{\star}\approx 20$ km s⁻¹ [37,71]. Additionally, FCC 119 has a relatively young NSC, with an estimated age of 1.7 ± 0.5 Gyr and a metallicity of -1.02 ± 0.02 dex [42], while its outer regions are older, with ages increasing from 5 to 10 Gyr [42,71]. The presence of emission lines exclusively within the NSC suggests ongoing star formation concentrated in the nuclear region.

2.5. Fornax & Virgo Clusters

2.5.1. Virgo Cluster

Although the Virgo Cluster is located in the Northern Hemisphere, some targets can still be observed with the ELT (Table 1). As the most extensively studied cluster to date, Virgo contains \approx 2000 members [9] at a distance of \approx 16.5 Mpc [89]. It is the most massive galaxy cluster in the Local Supercluster, with an estimated virial mass of $M_{\rm vir}\approx(1.4-6.3)\times10^{14}~\rm M_{\odot}$ [74,88,142]. This cluster exhibits a strong segregation of Hubble types: its center is dominated by early-type galaxies (ETGs), while dwarf elliptical/dwarf S0 (dE/dS0) galaxies dominate the overall census, comprising \approx 75% of the members [10,44]. Several large surveys have attempted to study the Virgo Cluster in multiple wavelengths, including X-ray [12,49,139], optical [28,47], near-infrared [79,87,123], far-infrared [14,30], and radio [6,152,154]. These surveys have provided the groundwork for investigating dwarf galaxies and predicting the presence of IMBHs [50,58,59,149].

2.5.2. Fornax Cluster

Fornax is an excellent galaxy cluster for studying galaxy evolution in the Southern Hemisphere. With a virial mass of $M_{\rm vir}=7\times10^{13}~{\rm M}_{\odot}$ [38], it is the second largest cluster on the sky after Virgo, which contains numerous ETGs at a distance of 20 Mpc [11]. This cluster has been investigated in multi-wavelength surveys: X-ray [128], optical [70,73,78,92,117], NIR [31], and radio [82]. The distance to the cluster currently prevents dynamical measurements of its dwarf members' $M_{\rm BH}$ because their $r_{\rm SOI}$ is far below the resolving power of current telescopes, except for FCC 213 [52,66] and FCC 47 [133]. However, given the significant number of its dwarf ETGs and GCs, Fornax represents a promising location for searching for IMBHs in the local volume with the ELT.

3. Stellar Mass Model of FCC 119

3.1. HST Observations

We used Hubble Space Telescope (HST) Advanced Camera for Surveys (ACS) Wide Field Channel (WFC) F850LP observations of the Fornax Cluster (PI: Andres Jordan) from the Hubble Legacy Archive (HLA) to constrain the radial SB profile and stellar mass model for FCC 119. The image has a pixel scale of 0.05 and a total exposure time of 1220 seconds, composed of three frames.

3.2. HST Point Spread Function (PSF) Simulations

We generated the F850LP PSF using the Tiny Tim package [76] to accurately derive the intrinsic SB of FCC 119. This software constructs model PSFs tailored to the specific instrument, detector chip, chip position, and filter used in the observations. To ensure consistency with the processing of the HST images, we created individual model PSF for each exposure frame (i.e., three PSFs for three exposures), incorporating sub-pixel offsets following a four-point box dither pattern. Additionally, each model PSF was convolved with an appropriate charge diffusion kernel to account for electron leakage into adjacent CCD pixels. The three individual PSFs were then combined and resampled onto a final grid with a pixel size of 0½05 using Drizzlepac/AstroDrizzle [3], producing the final PSF, which has a FWHM of 0½11.

3.3. F850LP Surface Brightness Profile

To mitigate contamination from foreground stars, we first masked them in the HST image using SExtractor⁶ [7], a tool designed to extract light sources from astronomical images and classify them as galaxies or point sources (stars) based on the "stellarity index" in the CLASS_STAR parameter. We set a seeing FWHM of 0.2 and a detection threshold (Detect_thresh) of 20 mag. Objects with a stellarity index greater than 0.5 were classified as stars and subsequently masked from the image.

Second, we extracted the radial SB profile along the galaxy's semi-major axis from the masked image using the ellipse task with the Image Reduction and Analysis Facility (IRAF) [72]. This task measures the integrated flux (in counts $\rm s^{-1}$) within concentric annular rings while accounting for variations in position angle and ellipticity. A photometric zero-point of 24.87 mag was applied, calculated using the ACS Zeropoints Calculator from the acstools package⁷. Additionally, a solar absolute magnitude of 4.5 mag [153] and a Galactic extinction value of 0.017 mag [119] were adopted. During this process, we deconvolved the HST image using the final Tiny Tim PSF (Section 3.2). The integrated flux within each annulus was then converted to mag arcsec⁻².

Next, we fitted the radial SB profile using a combination of a core-Sérsic [57,140] and a Sérsic [120] function, which describe the light profiles of the NSC and the disk component, respectively. The core-Sérsic function has the following analytical form:

$$I(r) = I' \left[1 + \left(\frac{R_b}{r} \right)^{\alpha} \right]^{\gamma/\alpha} \exp \left[-b_{n_{\text{NSC}}} \left(\frac{r^{\alpha} + R_b^{\alpha}}{R_e^{\alpha}} \right)^{1/(\alpha \, n_{\text{NSC}})} \right]$$
(2)

⁷ acstools v3.7.0: https://pypi.org/project/acstools/



⁶ SExtractor: https://github.com/astromatic/sextractor/

where
$$I' = I_b 2^{-\gamma/\alpha} \exp\left[b_{n_{\rm NSC}} 2^{1/(\alpha\,n_{\rm NSC})} \left(\frac{R_b}{R_e}\right)^{1/n_{\rm NSC}}\right]$$
. Here, R_b is the break radius, marking the point at which the radial SB profile transition from the outer Sérsic to the inner power law regimes, while the sharpness of this transition is controlled by α . I_b is the intensity at R_b , which will be converted into SB values (μ_b) . The shape of the outer Sérsic part $(r > R_b)$ is determined by its Sérsic index $(n_{\rm NSC})$ and effective radius (R_e) , while the inner-power law part $(r < R_b)$ is sharped by the power-law index γ . And the constant $b_{n_{\rm NSC}}$ is approximated as $b_{n_{\rm NSC}} = 2n_{\rm NSC} - \frac{1}{3} + \frac{4}{405n_{\rm NSC}} + \frac{46}{25515n_{\rm NSC}^2} + O\left(\frac{1}{n_{\rm NSC}^3}\right)$ [27].

Given that the observable scale of the HST image is 5 times coarser than that of HARMONI, detailed stellar information near the IMBH's SOI is unavailable. To address this, [106] assumed a core-Sérsic function with a power-law index of $\gamma=0.1$ for the NSC. We also adopt this assumption in this work.

The Sérsic function [120] describes the SB profile of the galaxy's disk component, characterized by the half-light radius ($R_{e,\text{disk}}$), Sérsic index (n_{disk}), and intensity ($I_{e,\text{disk}}$) at $R_{e,\text{disk}}$, expressed as:

$$I(r) = (I_{e,\text{disk}}) \exp\left\{-b_{n_{\text{disk}}} \left[\left(\frac{r}{R_{e,\text{disk}}}\right)^{1/n_{\text{disk}}} - 1 \right] \right\}$$
 (3)

where, $b_{n_{\rm disk}}$ is determined so that $\gamma(2n_{\rm disk};b_{n_{\rm disk}})=\frac{1}{2}\Gamma(2n_{\rm disk})$, where Γ and γ respectively refer to the Gamma and lower incomplete Gamma function [26]. And, we adopt $b_{n_{\rm disk}}\approx 2n_{\rm disk}-1/3$.

We employed a non-linear least-squares fitting algorithm using the Python MPFIT⁸ function [86] to iteratively fit a combined core-Sérsic and Sérsic function to the intrinsically deconvolved radial SB profile of FCC 119. The best-fit parameters include those for the core-Sérsic function: $n_{\rm NSC}=1.86\pm0.41$, $\alpha=4.25\pm0.35$, $R_{\rm e,NSC}=0.100\pm0.006$, $R_b=0.000$, $R_b=0.000$, and $\mu_b=16.5\pm0.4$ mag arcsec⁻²; and for the Sérsic function: $n_{\rm disk}=0.8\pm0.2$, $R_{\rm e,disk}=11.000$, and $\mu_{\rm e,disk}=21.2\pm0.8$ mag arcsec⁻². Based on these best-fit parameters, we extrapolated the SB profile toward the galaxy center down to a scale of 0.0000, as required for this extended HARMONI IMBH sample, and present the results in Figure 5.

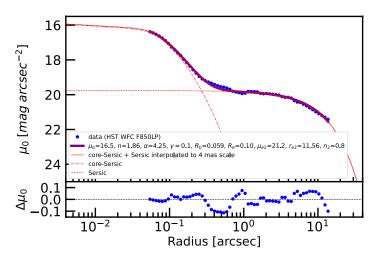


Figure 5. Top: The HST ACS/WFC F850LP radial SB profile of FCC 119 (blue dots) is shown simultaneously with the best-fit model (purple-thick line), which is a combination of a core-Sérsic (red-dotted line) and a Sérsic (red-dashed line) function, extrapolated into both the unresolved region of 0.04 and the extended region beyond 20. The best-fit parameters of the best-fit model are shown in the legend. **Bottom:** The residual (data-model) between the SB and the best-fit combined function.

⁸ http://purl.com/net/mpfit.

3.4. MGE Mass Model

We used the best-fitting parameters from the combined core-Sérsic and Sérsic functions to construct an MGE model using the mge.fit_1d procedure from the MGEFIT package⁹ [17,40]. The input logarithmically sampled Sérsic profile was sampled over the range 0.04– $400R_e$ and fitted with nine Gaussians to ensure an accurate representation of the galaxy SB distribution. To account for ellipticity, defined as $\epsilon = 1 - b/a = 1 - q'$, which quantifies the elongation of the galaxy along its major axis, we used the FIND_GALAXY procedure from the MGEFIT package to measure ϵ . FCC 119 shows a marked change in ellipticity at a radius of r = 6''. Beyond this radius, we find $\epsilon(r > 6'') \approx 0.15$, while inside this radius the morphology is nearly spherical, with $\epsilon(r < 6'') \approx 0$. This ellipticity information was incorporated into the HST/WFC F850LP MGE model presented in Table 5 (Column 4) in terms of axis ratio q' for each Gaussian.

Finally, we created the stellar mass model from this MGE model by scaling its surface luminosity (Column 2 of Table 5) with a constant M/L. Here, we estimated the M/L_{F850LP} of FCC 119 utilizing the empirical correlation between color and M/L from [144]:

$$\log_{10}(M/L_z) = a_z(g-z) + b_z \tag{4}$$

where $a_z=1.25$ and $b_z=-1.11$ are given in Table 3 of [144]. We adopted the color value of $g-z=0.63\pm0.12$ mag [141] to infer $M/L_{\rm F850LP}\approx M/L_z\approx0.6$ (M_{\odot}/L_{\odot}) for FCC 119. We applied this estimated $M/L_{\rm F850LP}$ for our mass model, which yields an NSC mass of $M_{\rm NSC}=6.5\times10^6$ M $_{\odot}$ and a total galaxy mass of $M_{\star}=1.15\times10^9$ M $_{\odot}$. These values are consistent with the estimates previously discussed in Section 2.4 within 8%.

In this work, we ignored the variation in M/L from stellar populations [91,99,134] and the distribution of dark matter halos [93].

4. HARMONI IFS Simulations

4.1. Jeans Anisotropic Model (JAM)

Our extended HARMONI IMBH sample consists of galaxies with NSCs that may exhibit rotational signatures [121], making them well-suited for a cylindrical projection. To model their kinematics, we employed JAM, assuming a velocity ellipsoid aligned with cylindrical coordinates ($\sigma_z \neq \sigma_r = \sigma_\theta$) in the meridional plane [18]. This was implemented by setting the model keyword align = 'cyl' in the JAM_AXI_PROJ procedure within the JAMPY package¹⁰ [19].

4.2. MARCS Synthetic Library of Stellar Spectra

To simulate the HARMONI IFS, we used the [85] stellar population synthesis (SPS) spectra, which are based on the Model Atmospheres with a Radiative and Convective Scheme (MARCS) synthetic spectra [62]. The MARCS library provides \approx 52,000 atmospheric models at a high spectral resolution of $R\approx$ 20,000, corresponding to an instrumental broadening of $\sigma_{\rm instr}=6.4~{\rm km~s^{-1}}$. It offers fine spectral sampling with 100,724 flux points ($\Delta\lambda\approx$ 0.065 Å) and covers a broad wavelength range of 0.13–20 μ m, making it well-suited for generating mock IFS data to detect the stellar kinematic signatures of IMBHs.

4.3. HARMONI Instrument and HSIM Simulator

HARMONI is the first IFS for the ELT, covering optical and NIR wavelengths (0.46–2.46 μ m) and offering 13 spectral gratings with three different spectral resolving powers: $R \approx 3,300$ (low), $\approx 7,100$ (medium), and $\approx 17,400$ (high). The instrument provides four observational scales: 4×4 , $10 \times 10, 20 \times 20$, and 30×60 mas², corresponding to fields of view (FoVs) of 0.86 \times 0.61, 2.14 \times 1.52, 4.28 \times 3.04, and 9.12 \times 6.42, respectively [155]. In this study, we conducted our HARMONI IFS

⁹ v5.0.15: https://pypi.org/project/mgefit/

jampy v7.2.4: https://pypi.org/project/jampy/

Table 5. The HST/WFC F850LP MGE model of FCC 119

j	$\Sigma_{\star,j}/({\rm L}_{\odot}~{ m pc}^{-2})$	$\sigma_j(arcsec)$	$q_j' = b_j/a_j$
(1)	(2)	(3)	(4)
1	9428.31	0.060	1.00
2	2535.54	0.131	1.00
3	175.71	0.346	1.00
4	91.52	2.679	1.00
5	243.69	6.034	1.00
6	187.84	9.868	0.85
7	68.45	13.762	0.85
8	10.19	17.689	0.85
9	0.37	21.999	0.85

Notes: Columns in order: Gaussian component number, luminosity surface density, dispersion along the major axis, and axial ratios.

Table 6. HSIM simulation via HSIM.

HARMONI grating	Exposure time (NDIT)	Sensitivity (NDIT)		
(1)	(2)	(3)		
H-high	16	12		
K-short	30	24		
K-long	20	16		

Notes: The total exposure time (and sensitivity) on source is determined as DIT×NDIT with DIT = 900 seconds. These estimated time are the science time on sources without accounting for the target acquisition, overhead, and LTAO setup.

simulations using the high-spectral-resolution mode with an intermediate spatial resolution of 10×10 mas² spaxel [98,106].

The HARMONI IFS cubes were generated using HSIM [155], which applies observational effects—including instrumental noise, atmospheric conditions, celestial sources with physical properties, realistic detector characteristics, and readout noise—to initially noiseless input data cubes (Section 4.4.1). This produces realistic simulated observations [98,106] and ensures that the resulting data closely resemble actual astronomical observations.

It should be noted that the final performance of HARMONI remains subject to possible descoping in the coming years, which could reduce its spatial resolution and spectral resolution. Such changes in the telescope's specifications would directly affect the IMBH detection limits. Therefore, the results presented here should be regarded as best-case estimates, based on the ELT's current planned design.

4.4. HARMONI Simulations

We employed the high-spectral-resolution gratings of HARMONI, including the H-high (1.538–1.678 μ m), K-short (2.017–2.201 μ m), and K-long (2.199–2.400 μ m) bands, which include several strong spectral lines that have been widely used in previous stellar kinematic studies [29,106].

The *H*-high band contains numerous metallicity indicators, such as Mg I (λ 1.487, 1.503, 1.575, 1.711 μ m), Fe I (λ 1.583 μ m), and Si I λ 1.589 μ m, along with CO absorption lines, including ¹²CO(3-0) λ 1.561 μ m, ¹²CO(4-1) λ 1.578 μ m, ¹²CO(5-2) λ 1.598 μ m, ¹²CO(6-3) λ 1.619 μ m, and ¹²CO(7-4) λ 1.640 μ m. In addition, the *K*-long band features prominent CO molecular absorption lines, including ¹²CO(2-0) λ 2.293 μ m, ¹²CO(3-1) λ 2.322 μ m, ¹²CO(4-2) λ 2.351 μ m, and ¹²CO(5-3) λ 2.386 μ m, as well as atomic absorption lines, such as Na I λ 2.207 μ m, Ca I λ 2.263 μ m, and Mg I λ 2.282 μ m. Furthermore, the *K*-short grating contains additional atomic lines, including Si I (λ 2.062, 2.083, 2.092, 2.137, 2.179, 2.183, 2.188 μ m) and Fe I (λ 2.070, 2.088 μ m). All these spectral features are illustrated in Figure 6.

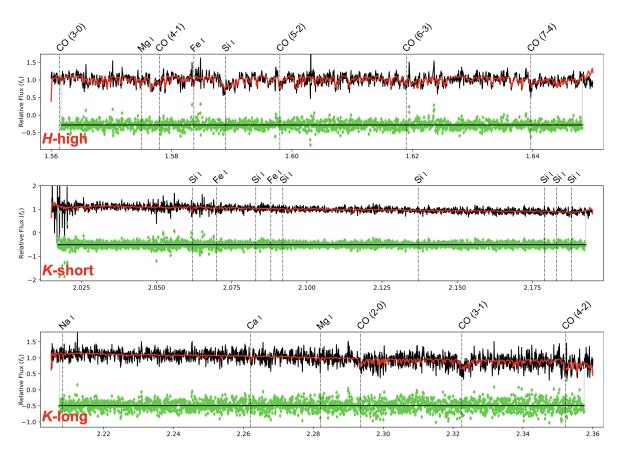


Figure 6. The central Voronoi-binned spectra of FCC 119 are presented for the three mock IFS data cubes generated by HSIM (black), along with their best-fit pPXF models (red). The residuals, calculated as (data-model), are shown as green points.

4.4.1. Creation of the Input-Noiseless Cubes for HSIM

We assumed that the line-of-sight velocity distribution (LOSVD) in the nucleus of FCC 119 is described well by a Gaussian profile. Accordingly, we computed the 2D first (V) and second ($V_{\rm rms}$) moments of the stellar component using the JAM_AXI_PROJ routine from the JAMPY package (Section 4.1) [19]. The velocity dispersion was then derived as $\sigma_{\star} = \sqrt{V_{\rm rms}^2 - V^2}$. In our JAM modeling, we integrated the MGE model (Table 5) with a constant $M/L_{\rm F850LP} \approx 0.6~{\rm M}_{\odot}/{\rm L}_{\odot}$ (Section 3), isotropic stellar orbits ($\beta_z = 0$), and an inclination angle of $i = 65^{\circ}$. To assess the impact of a central IMBH on the simulated IFS and stellar kinematics, we considered two scenarios: (i) a model without a BH, and (ii) a model including an IMBH with a mass equivalent to 5% of the NSC mass in FCC 119, corresponding to $M_{\rm BH} = 3.25 \times 10^5~{\rm M}_{\odot}$.

We generated the input-noiseless cubes for HSIM within a FoV of $0''.4 \times 0''.4$ and a pixel size of 2×2 mas², following the below steps:

- (i) We rebinned the MARCS SPS spectra with an assumed stellar population (Section 4.2) onto a logarithmic scale using velscale = 0.5 km s⁻¹, ensuring a constant $\Delta \log \lambda$ interval.
- (ii) We constructed a Gaussian LOSVD kernel for each spatial position (x,y) with velocity V(x,y) and velocity dispersion $\sigma_{\star}(x,y)$ using JAM, given specific dynamical parameters. The velocity was sampled at $\Delta V=0.5\,\mathrm{km\,s^{-1}}$.
- (iii) We convolved the rebinned spectrum from step (i) with the Gaussian LOSVD from step (ii) in log-wavelength space, yielding the broadened stellar spectrum at position (x, y).
- (iv) The spectrum was rebinned onto a linear wavelength grid with a sampling step at least twice as fine as the HARMONI spectral resolution (e.g., K-long has $\Delta\lambda\approx0.5\text{Å}$). This process involved thorough integration across pixels to prevent any loss of information during HSIM simulating process.
- (v) Finally, the spectrum at (x, y) was then scaled to match the SB predicted by the MGE photometric model (Table 5) for that position. This scaling factor was obtained by comparing the integrated

flux of the original spectrum within the corresponding photometric band to the MGE model's SB, using the ppxf_util.mag_spectrum function from the Penalized Pixel-Fitting (pPXF) method¹¹ [20].

4.4.2. HSIM Output Datacubes

We used the input-noiseless cubes created in Section 4.4.1 as inputs for HSIM to simulate HAR-MONI IFS observations in the H-high, K-short, and K-long gratings. The simulations were conducted under median observational conditions at the Armazones site, utilizing LTAO mode with a NGS of 17.5 mag in the H-band, located within a 30'' radial distance. We assumed an optical zenith seeing (0.5 μ m) with a FWHM of 0'.64 and an airmass of 1.3.

The exposure time for each simulation was carefully calibrated to achieve a minimum S/N of 3 per spaxel across the simulated FoV at the observed stellar features (before voronoi binning). To replicate real-time observations, we incorporated multiple exposure frames and dithering, with a Detector Integration Time (DIT) of 900 seconds per exposure. The total exposure time was determined by the number of exposures (NDIT), following the relation DIT×NDIT.

We obtained unexpected results regarding the exposure time required for the K-short grating to achieve stellar kinematics of comparable quality to the K-long. Given its higher flux level, the exposure time for K-short should theoretically be shorter than that for K-long. However, we found that K-short requires a longer exposure time (7.5 hours) compared to K-long (5 hours), arising because the K-short grating lacks strong stellar absorption features, such as the CO bandheads, which are present in the K-long range and are critical for detecting IMBH-induced stellar kinematics [106]. Furthermore, the atomic stellar features present in K-short are relatively weak and can be easily blended into spectral noise if the BH mass is large ($\gtrsim 10^6~{\rm M}_\odot$) or if the exposure time is insufficient. This result highlights the limitations of the K-short IFS for detecting and accurately measuring the stellar kinematic signatures of central BHs.

5. Results

5.1. Stellar Kinematics Extraction

We extracted the stellar LOSVD from the mock IFS data cubes generated in Section 4.4.2 using the adaptive Voronoi binning technique via the vorbin procedure [21] and the pPXF method [20]. The Voronoi binning optimally enhances the spectral S/N to a target threshold by setting targetSN = 25, while pPXF fits the binned spectra using the MARCS SPS templates (see Section 4.2). The pPXF fit was constrained to recover the velocity (V) and velocity dispersion (σ_{\star}) by setting moments = 2. Since we did not account for the continuum and sky background, we excluded the additive and multiplicative polynomials by setting mdegree=0 and degree=-1. The root-mean-square velocity was then computed as $V_{\rm rms} = \sqrt{V^2 + \sigma_{\star}^2}$. During the fitting process, we accounted for the instrumental broadening of the HARMONI IFS by convolving the stellar templates with the constant instrumental dispersion difference between the data and the templates.

We enhanced the fidelity of our fits by incorporating 13 MARCS SPS templates with ages ranging from 3 to 15 Gyr and z004 metallicity. [106] demonstrated that using either the full grating spectrum or an optimally selected spectral range resulted in only a 5% difference in the recovered kinematics. In this study, we selected the spectral ranges containing the strongest absorption features (as described in Section 4.4 and shown in Figure 6) to extract the stellar kinematics. The selected ranges are $\lambda 1.56-1.65$ μ m, $\lambda 2.02-2.19$ μ m, and $\lambda 2.21-2.36$ μ m for the H-high, K-short, and K-long gratings, respectively.

The best-fit pPXF models for the central bin, corresponding to the NSC center (x,y) = (0,0), and the optimal spectral ranges of the H-high, K-short, and K-long data cubes are shown in Figure 6. Figure 7 further presents the resulting stellar LOSVD maps, which reveal significant differences in σ_{\star} (and hence $V_{\rm rms}$) between the two scenarios: one without a BH and the other with an IMBH of $M_{\rm BH}$

v3.1.5: https://pypi.org/project/vorbin/



¹¹ v9.2.1: https://pypi.org/project/ppxf/

 $= 3.25 \times 10^5 \ {\rm M_{\odot}}$. In the absence of a BH, both σ_{\star} and $V_{\rm rms}$ exhibit a central drop, whereas the IMBH models show a pronounced central peak in these kinematic quantities. This contrast represents a clear kinematic signature of a central massive object.

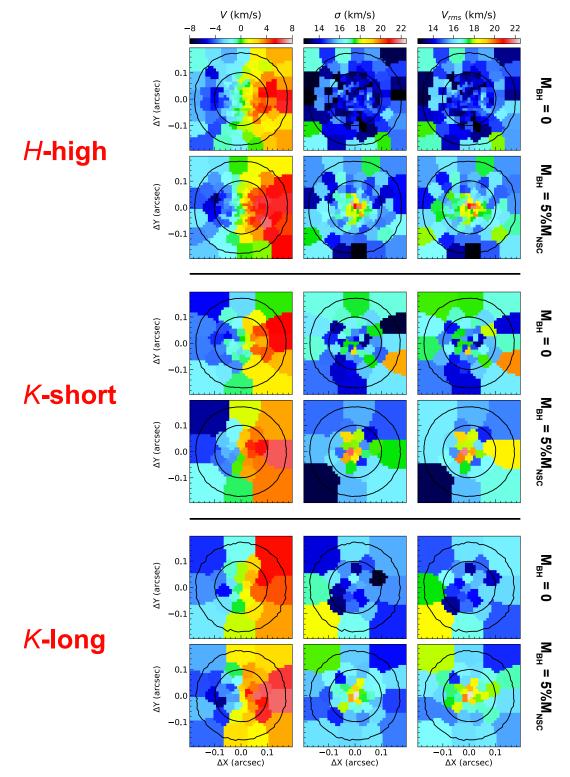


Figure 7. The 2D stellar kinematic maps of FCC 119, derived from the three mock IFS gratings, show the velocity (V, left), velocity dispersion (σ_{\star} , middle), and root-mean-square velocity ($V_{\rm rms}$, right). The black contours in each map represent isophotes extracted from the collapsed mock IFS cubes, decreasing in brightness by 1 mag arcsec⁻² from the center outward. For each mock IFS grating, the kinematic maps were simulated assuming no BH (top) and including an IMBH (bottom).

A central decrease in σ_{\star} within resolved nuclei is a common feature in systems that either lack a central BH or host a low-mass BH [106]. This behavior occurs across a broad range of Sérsic indices (n) and power-law indices (γ) in the core-Sérsic profile of an NSC, as theoretically predicted by [138] and numerically demonstrated by [106] using HARMONI IFS stellar kinematics simulations. In contrast, the presence of an IMBH induces a central increase in σ_{\star} due to the IMBH's gravitational dominance within its SOI. The difference in $V_{\rm rms}$ between scenarios with and without an IMBH is \approx 6 km s⁻¹, corresponding to 30% of the nuclear σ_{\star} of FCC 119 ($\sigma_{\star} \approx$ 20 km s⁻¹) [37,71]. This difference is substantial enough to reliably distinguish between the two cases.

5.2. The Black Hole Mass Recovering

We combined the MGE model (Section 3.4) and the mock stellar kinematic measurements (Section 5.1) to reconstruct the LOSVD of FCC 119, including V and σ_{\star} (or $V_{\rm rms}$), using JAM (Section 4.1). The fitting was performed within a Bayesian inference framework [23], employing a Markov Chain Monte Carlo (MCMC) simulation via the adaptive Metropolis algorithm implemented in AdaMet ¹³ [63]. We assumed a uniform Gaussian error distribution, making the posterior probability proportional to the logarithm of the likelihood, $\ln P({\rm data}|{\rm model}) \propto -0.5\chi^2$, where the chi-squared χ^2 is given by:

$$\chi^2 \equiv \sum_i \left(\frac{V_{\text{rms},i} - \overline{V}_{\text{rms},i}}{\sigma_i} \right)^2, \tag{5}$$

Here, $V_{\text{rms},i}$ and σ_i represent the root-mean-square velocity and its associated error for bin i extracted from the mock data. The model root-mean-square velocity, $\overline{V}_{\text{rms},i}$, was computed using JAM and convolved with the LTAO PSF, which has a FWHM_{PSF} = 0.012 [98,106,137]. The JAM model explores a four-parameter space, including the black hole mass (M_{BH}) in logarithmic scale, as well as the inclination (i), the mass-to-light ratio in the F850LP filter (M/L_{F850LP}), and the anisotropy parameter (β_z) in linear scale.

We performed a total of 3×10^5 calculations and excluded the first 20% of the MCMC steps as the burn-in phase. The best-fit parameters and their uncertainties were then derived from the probability distribution function (PDF) of the remaining 80% of the calculations. We used the same initial guesses for the free parameters in all fittings, which are: $i = 65^{\circ}$, $M/L_{\rm F850LP} = 0.6$ (M_{\odot}/L_{\odot}), $\beta_z = 0$, and $M_{\rm BH}$ either $0~M_{\odot}$ or $3.25 \times 10^5 M_{\odot}$, identical to the input values for HSIM. We also defined the search ranges as follows: $i~(33^{\circ}-90^{\circ})$, $\log_{10}M_{\rm BH}~(0-7M_{\odot})$, $M/L_{\rm F850LP}~(0-2~M_{\odot}/L_{\odot})$, and $\beta_z~(-1~{\rm to}~0.99)$

We present the AdaMet MCMC results for the H-high, K-short, and K-long bands, corresponding to both input $M_{\rm BH}$ values, in Figures 8 and 9. The 2D distribution of each parameter pair is shown in the corner plots after marginalizing over the other parameters, where each point represents a JAM model. The color of these scatter points, ranging from white to blue, indicates their likelihood at different confidence levels (CLs), from 1σ to 3σ . Black points represent CLs greater than 3σ , while white points denote the highest likelihood (best fit) within the 1σ CL. The top histogram illustrates the 1D distribution of each parameter and is used to estimate the best-fit values and associated uncertainties. An anti-correlation between $M/L_{\rm F850LP}$ and $M_{\rm BH}$ is evident in their 2D PDFs, arising from the interplay between the gravitational potentials of the IMBHs and their host galaxies, where larger BHs correspond to lower $M/L_{\rm F850LP}$ values and vice versa. This behavior highlights the high spatial resolution of our simulations, enabled by the 10 mas observational scale, which is sufficient to resolve the SOIs of these IMBHs at the distance of the Fornax Cluster.

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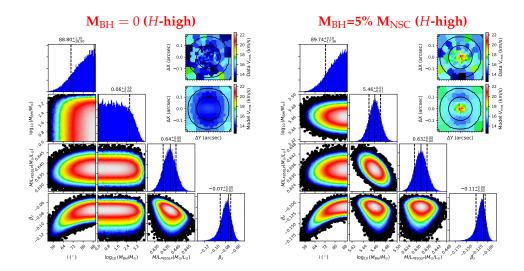


Figure 8. The PDF of four parameters (i, $M_{\rm BH}$, $M/L_{\rm F850LP}$, β_z) constrained using JAM within the AdaMet MCMC optimization framework, derived from mock H-high IFS data cubes. The figure presents results for two different $M_{\rm BH}$ cases: no BH (left) and $5\%M_{\rm NSC}$ (right). Each data point in the 2D scatter plots represents a single model, with colors indicating the log-likelihood. The 1D histograms above display the corresponding projected 1D distributions. In the upper-right corner, a comparison is shown between the mock kinematic maps and their best-fitting model.

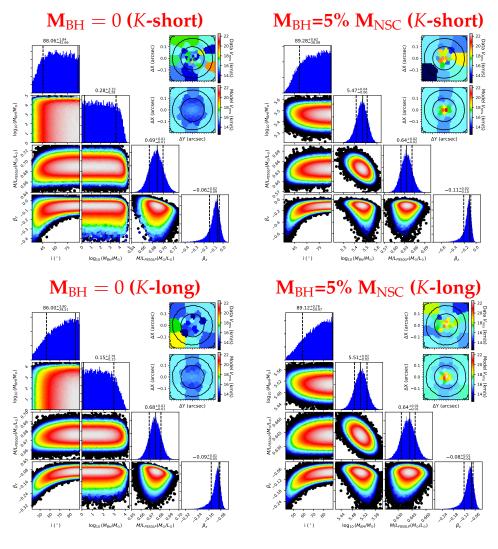


Figure 9. Same as Figure 8, but for the *K*-short and *K*-long gratings.

The top-right corner of these figures compares the $V_{\rm rms}$ map extracted from the HSIM datacubes (top) with that derived from the best-fit JAM model (bottom), both using the same color scale. The best-fit parameter values, along with their associated 1σ and 3σ uncertainties, are listed in Table 7. Our recovered values for both $M/L_{\rm F850LP}$ and $M_{\rm BH}$ closely match the input values, with differences of <15%. Notably, the uncertainties from the MCMC fits remain relatively small (<10% at the 3σ CL). This precision is attributed to the exceptional spatial and spectral resolutions of our IFS simulations, which effectively resolve the BH's SOI and capture subtle increasing in $V_{\rm rms}$ induced by the presence of a small IMBH.

Initial Input			$M_{\mathrm{BH}} =$	0	$M_{\rm BH} = 5\% M_{ m NSC}$			
parameters	JAM guess	best-fit	(16-84)%	(0.14-99.86)%	best-fit	(16-84)%	(0.14-99.86)%	
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	
		H-high						
<i>i</i> (°)	65	89	>19	>36	90	>17	>34	
$\log_{10}(M_{\rm BH}/{ m M}_{\odot})$	(*)	0.06	< 2.58	< 3.68	5.460	± 0.010	± 0.031	
$M/L_{\rm F850LP}~({ m M}_{\odot}/{ m L}_{\odot})$	0.6	0.637	± 0.003	± 0.009	0.633	± 0.004	± 0.011	
eta_z	0	-0.074	± 0.010	± 0.039	-0.110	± 0.011	± 0.044	
	K-short							
<i>i</i> (°)	65	88	>43	>56	89	>38	>56	
$\log_{10}(M_{\rm BH}/{ m M}_{\odot})$	(*)	0.28	< 3.33	< 4.37	5.472	± 0.050	± 0.150	
$M/L_{\rm F850LP}~({ m M}_{\odot}/{ m L}_{\odot})$	0.6	0.689	± 0.013	± 0.039	0.636	± 0.017	± 0.050	
eta_z	0	-0.057	± 0.052	± 0.240	-0.114	± 0.062	± 0.340	
	K-long							
i (°)	65	86	>30	>48	89	>30	>50	
$\log_{10}(M_{\rm BH}/{ m M}_{\odot})$	(*)	0.15	< 2.79	< 3.92	5.507	± 0.022	± 0.065	
$M/L_{\rm F850LP}~({ m M}_{\odot}/{ m L}_{\odot})$	0.6	0.675	± 0.006	± 0.019	0.643	± 0.008	± 0.022	
eta_z	0	-0.093	± 0.029	± 0.130	-0.080	± 0.025	±0.230	

Table 7. Best-fit JAM models and uncertainties.

Notes: Columns 1–2: The free parameters of JAM, their initial guessed values. The asterisk (*) in $M_{\rm BH}$ in Column 2 indicates that the initial guesses for the $M_{\rm BH}$ either zero or $M_{\rm BH}=3.25\times10^5~\rm M_{\odot}$. Columns 3–5: the best-fitting parameters for the case without a central BH ($M_{\rm BH}=0~\rm M_{\odot}$), their 1σ and 3σ uncertainties. Columns 6–8: Same as Columns 4–5 but for the input $M_{\rm BH}=3.25\times10^5~\rm M_{\odot}$.

In addition, for a scenario without a central IMBH ($M_{\rm BH}=0$), the MCMC yields an upper limit of $M_{\rm BH}\approx 10^4~{\rm M}_{\odot}$. This indicates HARMONI's ability to detect IMBHs above this mass threshold at a distance of 20 Mpc. Conversely, simulations with an IMBH of $M_{\rm BH}\approx 3\times 10^5~{\rm M}_{\odot}$ and $\sigma_{\star}=20~{\rm km~s}^{-1}$ at a distance of 20 Mpc give an $r_{\rm SOI}\approx 0.0000$. This radius is three times larger than the simulated spaxel size of 0.00000. This establishes a fundamental lower limit for $M_{\rm BH}$ measurements using dynamical techniques with ELT/HARMONI in the future.

The recovered β_z values exhibit a broad range from -0.5 to -0.03, significantly deviating from the input value of $\beta_z = 0$. This suggests that stars within the nucleus of FCC 119 predominantly follow tangential orbits. These discrepancies may arise from the scattered nature of the $V_{\rm rms}$ measurements at the FoV edges. The faintness of FCC 119, combined with a steep decline in S/N toward the FoV boundaries, affects the accuracy of spectral fitting in these regions.

Additionally, the inclination i is not well constrained, varying between 33° and 90°. This is likely due to the kinematic properties of FCC 119, where velocity dispersion dominates and rotational motion is minimal ($V \approx \pm 8 \text{ km s}^{-1}$). As discussed in [106], NSCs with little to no rotation exhibit kinematic characteristics similar to nearly spherical systems. In such cases, galaxies appear similar over a wide range of inclinations, making it difficult to precisely constrain i in our models.

We compared the input $M_{\rm BH}$ values used to generate the mock IFS data cubes with the corresponding recovered values in Figure 10. The recovered data points closely align with the line of equality, demonstrating the robustness of our method for constraining IMBH masses using HARMONI observations.

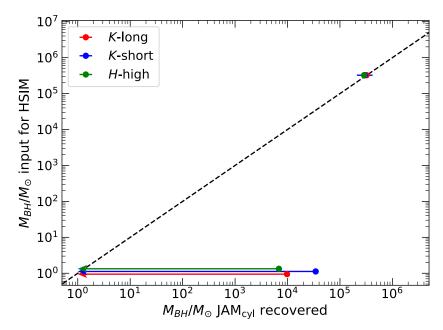


Figure 10. A comparison of the input M_{BH} used in HSIM and the recovered values within 3σ CL. The black dashed line represents the line of equality between the input M_{BH} in HSIM and our recoveries.

5.3. Dynamical Mass-Segregation in FCC 119 NSC?

In a stellar system, stars will reach an energy equipartition state through two-body interactions, causing higher-mass stars decrease velocities and move toward the center of the system, while lower-mass stars gain higher velocities and drift outward (i.e., mass-segregation process) [127]. Such massive star's deaths leads to a gradual accumulation of dark remnants (e.g., neutron stars, stellar-mass black holes) and cause an increase in the M/L toward the central region [8]. The central peak-up of M/L in a star cluster system can produce a stellar kinematics that mimic the presence of a central IMBH and potentially leading to an ambiguous IMBH detection [106].

To assess the impact of this mass-segregation phenomenon on the FCC 119's NSC, we evaluated the relaxation time (t_{relax}) as a function of radius [4,143]:

$$t_{\rm relax}(r) \approx (1.4 \times 10^8 {\rm yr}) \,\sigma_{20}^3(r) \,\rho_5^{-1}(r) (\ln \Lambda_{10})^{-1},$$
 (6)

5.4. Constraint on Nearby NSC Brightness and Sensitivity

We examined the sensitivity of the ELT to low-SB galaxies, focusing on FCC 119, located at a distance of 20 Mpc, which has a SB of 20.1 mag $arcsec^{-2}$ at a radius of 1" in the z-band [141]. To assess this, we generated mock high-spectral-resolution HARMONI IFS cubes with varying exposure times to determine the minimum integration time required to achieve a sufficient S/N ratio for robust kinematic LOSVD measurements with each grating. Our simulations indicate that the required exposure times are at least 3, 6, and 4 hours for the *H*-high, *K*-short, and *K*-long bands, respectively, as summarized in Table 6. These sensitivity estimates account only for the on-source science time and do not include

additional time for target acquisition, instrument overheads, and LTAO setup. Consequently, the total required integration time in practice is expected to be at least twice as long. Furthermore, the SB profile was derived from HST imaging and extrapolated down to a spatial scale of 0.01 toward the galaxy's central region, potentially necessitating longer exposure times than initially estimated.

Thus, we established a limiting distance of 20 Mpc for conducting HARMONI observations within a reasonably short exposure time. At this distance, the NSC must have a central SB as low as $\mu_c \approx 16$ mag arcsec⁻², decreasing to approximately 18 mag arcsec⁻² at a radius of 0'.'2 (see Figure 5). [106] reported a similar SB decline in the NSCs of two much closer galaxies, NGC 300 and NGC 3115 dw01, where the central SB is $\mu_c \approx 15$ mag arcsec⁻², dropping to about 17 mag arcsec⁻² at a radius of 0'.'2. This consistency reinforces our observational lower limit. Notably, our findings indicate a sensitivity approximately two orders of magnitude lower than previous limits for 8–10 meter class telescopes, which are capable of observing nuclei with a central SB above $\mu_c \approx 14$ mag arcsec⁻², decreasing to about 16 mag arcsec⁻² at a radius of 1'.'5.

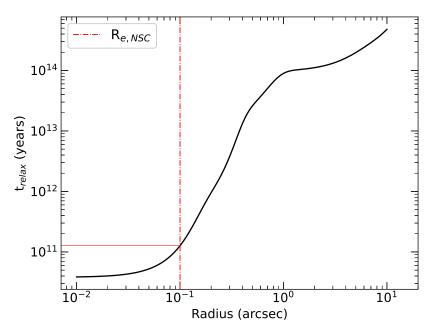


Figure 11. Relaxation time (t_{relax}) as a function of radius plots along the semi-major axis of FCC 119 (black-solid line). The red dash-dotted line indicates the effective radius of its NSC ($R_{e,NSC}$)

6. Conclusions

We expanded the current HARMONI IMBH sample, which was limited within 10 Mpc [106] to 20 Mpc by exploring the capabilities of HARMONI for dynamically measuring the masses of IMBHs in the second faintest dwarf and nucleated member of the Fornax Cluster, FCC 119, which is also much fainter than almost all members of the Virgo Cluster. Based on our findings, our conclusions are summarized as follows:

- (i) We compiled an expanded sample of 85 dwarf galaxies hosting bright NSCs, with masses ranging from 1.6×10^5 to 2.3×10^8 M $_\odot$ and effective radii between 0.5 and 62 pc, located at distances of 10–20 Mpc, representing promising sites for IMBHs.
- (ii) The sample spans a variety of galaxy types, including ellipticals (27%), lenticulars (29%), spirals (39%), and irregulars (5%). It comprises 29 members of the Virgo Cluster, 20 members of the Fornax Cluster, and 36 isolated galaxies.
- (iii) We performed HARMONI simulations for FCC 119 using the H-high, K-short, and K-long gratings with a spaxel scale of 0.01 \times 0.01, requiring on-source exposure times of 4, 7.5, and 5 hours, respectively. The consistent stellar kinematic measurements obtained across the mock IFS data cubes validate our methodology and observational strategy, demonstrating the reliability of extending the IMBH sample and highlighting HARMONI's capabilities.

- (iv) The recovered $M_{\rm BH}$ and stellar $M/L_{\rm F850LP}$ closely match the input simulation values, with uncertainties \lesssim 5%, demonstrating the robustness and reliability of our measurement technique using HARMONI data.
- (v) We determined that HARMONI can effectively observe NSCs at the distance of the Fornax Cluster with a central surface brightness as low as $\mu_c \approx 16$ mag arcsec⁻² (decreasing to ~ 18 mag arcsec⁻² at 0'.'2) and can dynamically resolve IMBHs with masses down to $\sim 10^5$ M $_{\odot}$.

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Software: Python 3.12: [145], Matplotlib 3.6: [68], NumPy 1.22: [64], SciPy 1.3: [147], photutils 0.7: [15], AstroPy 5.1 [2], AdaMet 2.0 [23], JamPy 7.2 [19], pPXF 8.2 [20], vorbin 3.1 [21] MgeFit 5.0 [17], and HSIM 3.11 [155].

Data Availability

All data and software used in this paper are public. We provided their links in the text when discussed.

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