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Article

On the Foundations of Quantum Mechanics

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Abstract

Matrix mechanics describes the spectral emissions of hydrogen gas, but quantum mechanics concerns the emission and absorption of radiation by a single atom. To obtain the spectral emissions of a single hydrogen atom we introduce thermal energy in infinitesimal increments registering the spectral lines one by one until a complete spectrum is obtained. We describe the gradual increase of the internal *quantum mechanical* energy in a non-conservative system by integrating between two specific endpoints, the atomic orbitals, and two specific times, the period τ , which are the conditions necessary for Hamilton's principle. Wave mechanics is shown to be invalid for non-conservative systems. A complete quantum mechanics is formulated using energy-time conjugate variables as the time integral of a Lagrangian independently of Hamiltonian formulations that use position-momentum conjugate variables.

Keywords: matrix mechanics; continuous time; energy conservation; Hamilton's principle; Lagrangian; action function

1. Introduction

Although matrix mechanics and wave mechanics are believed to be equivalent formulations of quantum mechanics, they are not equal participants. They are based on widely separated initial assumptions, distinct experiments, and different mathematical models. Contributions to the foundations by wave mechanical models predominate, while matrix mechanics has not advanced significantly in its contributions to the foundations beyond initial assessments 100 years ago. The reasons given for the imbalance are its lack of visualizability and ease of calculation, properties not of matter but of the formulation. The relative importance of a theory should not depend upon how it was formulated. Rather than ignoring matrix mechanics during discussions of the foundations we will reformulate it to obtain insights that complement the contributions of wave mechanics.

2. Matrix Mechanics Reformulated

Classical mechanics is the study of the motion of objects and their causes; in other words, the equations of motion. Ideally the objective of quantum mechanics should be to find out the subatomic equations of motion. This is clearly what wave mechanics does when it plots the probability densities for the electron orbitals of a hydrogen atom. However, matrix mechanics makes no attempt at all for it describes the properties not of an atom but of the simultaneous emission of radiation by many hydrogen atoms. Measurements are performed in order to obtain matrices that describe black body radiation in terms of the frequency and transition probability of the spectral lines, which are the observables. The set of all possible observables are understood to be the set of eigenvalues of the matrix. Thus the purpose of matrices is to give a complete set of the observables of hydrogen gas for any arbitrary electron transition. Instead we will use a single hydrogen atom and quantum system to obtain the observables of the spectrum.

There are difficulties involved in deriving the laws of mechanics for the single hydrogen atom as noted by Nobel laureate Richard Feynman [1].

But suppose we look at the whole hydrogen atom as a "particle." If we didn't know that the hydrogen atom was made out of a proton and an electron, we might have started out

and said: “Oh, I know what the base states are—they correspond to a particular momentum of the hydrogen atom.” No, because the hydrogen atom has internal parts. It may, therefore, have various states of different internal energy, and describing the real nature requires more detail.

The internal energy of a hydrogen atom may be described quantum mechanically by the Hamiltonian matrix H_{ij} because it has an infinite number of matrix elements and there are an infinite number of possible energy levels that an electron can transition between. However, it is impossible to determine the transition energies of all states. As energy states go higher, the energy differences between successive states become smaller, and the spectral lines converge and become continuous. Heisenberg chose to disregard spectral lines that cannot be observed stating that quantum mechanics should be “founded exclusively upon relationships between quantities which in principle are observable” [2]. We will include all electron transitions in a theory of the hydrogen atom, irrespective of whether the corresponding spectral lines are observable.

A matrix is a snapshot of an atom’s energy state at a particular point in time. The diagonalized matrix describing the hydrogen spectrum due to black body radiation is not the result of a single hydrogen atom, rather it is due to many atoms acting in unison. It is the result of many events, where each one occurs at an independent point in space and time. It is nevertheless possible for a solitary, electrically neutral hydrogen atom, isolated in space and time, to emit the entire hydrogen spectrum if we apply heat incrementally and describe the energy states sequentially in a series. Initially the states are of infinitesimal energy and undetectable, but they eventually register lines in the infrared, the visible spectrum, and beyond to the ultraviolet spectrum. The complete spectrum as viewed by a single atom appears in the manner of a motion picture with each frame consisting of a Hamiltonian matrix. We observe the spectrum as a continuously increasing background thermal energy that surrounds the hydrogen atom and is periodically interrupted by spectral lines of increasing energy in the Lyman, Balmer, Paschen, and Brackett series. Thus the discrete energy eigenvalues of the atom appear sequentially as diagonalized elements of the Hamiltonian matrix. In other words, the complete spectral series of hydrogen is generated by integrating the emissions of a single atom over a sufficiently long time period.

3. Wave Mechanics

We proceed next to examine the methods used to derive a second formulation of the equations of motion based on the Schrödinger wave equation, $i\hbar d/dt|\Psi(t)\rangle = \hat{H}|\Psi(t)\rangle$. He arrived at his equation by integrating over the entire coordinate space to obtain a wave function [3].

“We now seek a function ψ such that for any arbitrary variation of it the integral . . . taken over the whole coordinate space is stationary, ψ being everywhere real, single-valued, finite, and continuously differentiable up to the second order. *The quantum conditions are replaced by this variation problem.*” p. 2

He seeks to derive a theory based on an action functional and variational principle, a method which is central to all of physics. A functional $\psi(r,t)$ is introduced (the energy) whose variation over the coordinate space is stationary.

“There arises, however, an urgent need for the extension of the theory for non-conservative systems, because *it is only in that way that we can study the behavior of the system under the influence of prescribed external forces, e.g., a light wave, or a strange atom flying past.*” p. 103

Wave mechanics is derived with respect to a conservative electromagnetic system; that is, a proton- electron pair with a fixed potential. The probability density function of the orbiting electron for a fixed potential is defined by the square of the absolute value of the wave function at each point in coordinate space. Extending the theory to non-conservative systems would allow external forces to be introduced thereby enabling greater flexibility in the number and variety of applications. He concludes the first draft of his paper by repeating this same concern.

“In the pair of equations [given in the text], we have before us only the substitute . . . for a real wave equation . . . which, however, I have not succeeded in forming for the non-conservative case.” p. 123

Schrödinger reaffirms here the need for a quantum theory of non-conservative systems which would include the “external forces” on a hydrogen atom and enable it to be merged with classical theory. His theory as it stands is incomplete. Years later Einstein realized as much [4].

“The light quantum has a definite localization and a definite color. Naturally one cannot do justice to this by means of a wave function. Thus I incline to the opinion that the wave function does not (completely) describe what is real, but only a to us empirically accessible maximal knowledge regarding that which really exists. . . . This is what I mean when I advance the view that quantum mechanics gives an incomplete description of the real state of affairs.”

We postulate here that matrix mechanics provides the required “missing piece” of the puzzle that is needed to define a complete theory that is also real. It is mathematically equivalent to wave mechanics because it gives the same eigenvalues for the energy, but it also includes physical variables such as frequency that are lacking in the wave mechanical model.

4. Energy Conservation

It is well known that energy is not always conserved in quantum mechanics [5]. “Wave functions collapse when they are observed. This collapse process is unpredictable, and doesn’t obey the Schrödinger equation. As a result, the average energy is not conserved in the process of quantum measurement.” In fact the original proof of energy conservation given by Heisenberg in his study of the anharmonic oscillator is not rigorous. He demonstrated energy conservation for a conservative quantum system; however, the atom absorbs thermal energy from the environment as a continuous series of infinitesimal transitions during black body radiation in an open, non-conservative system. Neither he nor Schrödinger proved that energy is conserved for open, non-conservative quantum systems.

In order to prove that energy is conserved for an open, non-conservative quantum system the energy absorbed must be identical to the energy emitted. In fact due to energy conservation energy input equals energy output for all isolated material systems, whether the system is a hydrogen atom, a machine, a hydroelectric power plant, or a galaxy. The quantum system conserves energy for changes in electromagnetic potential in the same way that a gravitational system conserves energy for changes in gravitational potential. Energy conservation causes matrix mechanics to derive physical meaning from energy absorption and it necessarily equates with energy emission and the physical meaning attributed to wave mechanical formulations.

We shall illustrate energy conservation in a quantum system formally by using the conjugate variables energy and time to show that matrix and wave mechanics work together to describe the emission and absorption of radiation. We have described energy absorption in real space and real time with matrices as incremental increases. It can also be described in real time by one rotation of the wave function. Emission follows by means of a second rotation of the wave function demonstrating that absorption and emission are part of the same quantum process. We interpret the first rotation of the wave function not as a rotation in *abstract space*, but as the change in phase of electromagnetic fields from 0 to 2π in *real space* with phase information given by the off-diagonal matrix elements. Thus one “rotation” of the wave function is interpreted as one full cycle of an electromagnetic wave and an increase in the electron’s energy from the ground state to an excited state. The second rotation occurs as the electron returns to the ground state and is interpreted as a localization of electromagnetic field energy and emission of a photon. The dual wave-particle nature of the photon is thereby realized in a continuous transformation of fields. When a continuous function is integrated between two fixed endpoints, the orbitals, and two specified times, the period, the result is a relativistic invariant, the time integral of a Lagrangian $\int L dt$, or action function.

5. The Lagrangian in Quantum Mechanics

The photons of a spectral line are emitted when electrons “jump” from the ground state to an excited state and return to the ground state. The paths of excitation and decay occur between two specific endpoints, the atomic orbitals, and two specific times, the period τ , which are the physical conditions that define Hamilton’s principle of least action [6]. *Thus the physical characteristics of a quantum system naturally define the correct equations of motion needed to describe it.* The principle of least action simply means that of all the paths that an electron can take between two atomic orbitals, the ones actually taken are found by computing the action for each of the trajectories, and selecting the ones that have the least action. The paths selected will be the ones that have an action equal to the reduced Planck’s constant \hbar .

The action is the time integral of the Lagrangian $\int L dt$, where $L=T-V$, T is the kinetic energy, and V is the potential energy. We compute the action of an electron transitioning from the ground state to an excited state by using generalized coordinates, three to describe its position on the electron shells $R_1 = (x_1, y_1, z_1)$ and $R_2 = (x_2, y_2, z_2)$, and three to describe its trajectory.

$$S[r(t)] = \int_{R_1}^{R_2} \int_{t_1}^{t_2} (T-V) dt \quad (1)$$

The action, $S[r(t)]$, is a functional that describes the absorption process in four dimensions. It has as its argument an infinite number of functions, the possible electron paths $r(t)$. We know that it is relativistically correct because gravitational fields influence the action identically to the way that they influence clocks; that is, by causing time dilation [7].

Emission initiates from the excited state $R_2 = (x_2, y_2, z_2)$ at time t_2 and it finalizes at the ground state $R_1 = (x_1, y_1, z_1)$ at time t_1 . Each of the electron shells R_2 and R_1 determines a locus of points where the fields vanish. The Lagrangian density over the region of space-time between the excited and ground states is stationary for all small variations of the coordinates inside the region. Changes in action during emission are evaluated by integrating the Lagrangian density four-dimensionally thereby yielding a relativistic formulation.

$$S[\phi_i(t)] = \int_{R_2}^{R_1} \int_{t_2}^{t_1} L(\phi_i, \phi_{i,\mu}) d^3x dt = E\tau \quad (2)$$

Solving (2) we obtain an expression for the localized fields of a photon $E\tau = \hbar$, which is equivalent to the more familiar $E = h\nu$. The emission of a photon is described by the action integral of a Lagrangian density $\int \mathcal{L}(\phi_i, \phi_{i,\mu}) dt$ in the region of space-time between the excited and ground states given by the fields ϕ_i and its first derivatives $\phi_{i,\mu}$. This allows for a complete accounting of the energy interactions that occur during an electron transition within the volume $\int d^3x$ between the electron orbitals where ϕ_i is the current density described *radially* and $\phi_{i,\mu}$ is the electromagnetic field strength described *transversely* that results in the creation of a photon. The photon is described therefore as a four-dimensional localization of fields contained within the volume $\int d^3x$ and the time interval $t_2 - t_1 = \tau$. The discrete and continuous properties of the emission process are described by the Lagrangian density over the time interval of a period which results in a photon being created. A more complete derivation is given elsewhere [8].

6. Conclusions

By his own admission Schrödinger did not complete his derivation of wave mechanics by formulating it in a non-conservative system. As shown in 4.0 that can be accomplished by simply appending the output of matrix mechanics to the input of wave mechanics. The resulting “unified” theory is non-relativistic because the equations of motion are formulated in inertial systems and measurements are performed at single points in time. For a theory to be relativistic it must be formulated in a non-inertial frame so that the equations of motion vary continuously with respect to

velocity and acceleration as shown in 5.0. Dirac recognized already in 1930 the significance and relative importance of a relativistic theory [9]. "Quantum mechanics was built up on a foundation of analogy with the Hamiltonian theory of classical mechanics. Now there is an alternative formulation for classical dynamics provided by the Lagrangian. The two formulations are, of course closely related, but there are reasons for believing that the Lagrangian one is the more fundamental. In the first place the Lagrangian method allows one to collect together all the equations and express them as the stationary property of a certain action function. (This action function is just the time-integral of the Lagrangian.) There is no corresponding action principle in terms of the coordinates and momenta of the Hamiltonian theory. Secondly the Lagrangian method can easily be expressed relativistically, on account of the action function being a relativistic invariant; while the Hamiltonian method is essentially non-relativistic in form, since it marks out a particular time variable as the canonical conjugate of the Hamiltonian function. For these reasons it would seem desirable to take up the question of what corresponds in the quantum theory to the Lagrangian method of the classical theory."

References

1. Feynman, Leighton, Sands *The Feynman Lectures on Physics, Vol III, Chap 8-3 "The Hamiltonian Matrix"* (1963).
2. W. Heisenberg *Z Phys* **33** (1) (1925), in B.L. van der Waerden (ed.), *Sources of Quantum Mechanics* (Amsterdam, 1967). doi: <http://www.mat.unimi.it/users/galgani/arch/heis25ajp.pdf>
3. Schrödinger, E. *Collected Papers on Wave Mechanics* 2nd edn (Blackie, London 1928) https://mwolf.pracownicy.uksw.edu.pl/MK/Schrodinger_Collected_Papers_on_Wave_Mechanics.pdf
4. Einstein, A. Letter to Paul Epstein 11/10/1945 *Sixty-Two Years of Uncertainty Edited by A. I. Miller (Plenum Press, New York, 1990) p. 103.*
5. Carroll, S.; Lodman, J. (2021) "Energy non-conservation in quantum mechanics" <https://doi.org/10.1007/s10701-021-00490-5>
6. Lanczos, C. *The Variational Principles of Mechanics* 4th edn. (Dover, 1970)
7. T. Bothwell *et al.* "Resolving the gravitational red shift within a millimeter atomic sample" *Nature* **602**, Issue 7897, 420 (2022) arXiv:2109.12238
8. Oldani, R. (2024) "Galactic symmetry" doi: preprints.org/manuscript/202009.0215/v8, p. 11
9. Dirac, P.A.M., "The Lagrangian in quantum mechanics" *Phys Zeit Sow* **3**, 1933, p.1. https://www.informationphilosopher.com/solutions/scientists/dirac/Lagrangian_1933.pdf

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