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Article

Hardware in the Loop Testing Of Spacecraft Relative Dynamics And Tethered Satellite System On A Tip-Tilt Flat-Table Facility

Giuseppe Governale ^{1*} , Armando Pastore ¹ , Matteo Clavolini ¹ , Mattia Li Vigni ¹ , Christian Bellinazzi ¹ , Catello Leonardo Matonti¹ , Riccardo Apa ¹ , Stefano Aliberti ¹  and Marcello Romano ² 

¹ Department of Mechanical and Aerospace Engineering, Politecnico di Torino, Corso Duca degli Abruzzi 24, 10129, Turin, Italy

² Head, Chair of Astrodynamics, Technical University of Munich (TUM), Willy-Messerschmitt-Str. 11, 82024 Taufkirchen, Germany; marcello.romano@tum.de

* Correspondence: giuseppe.governale@polito.it

Abstract

This article presents a compact tip-tilting platform designed for hardware-in-the-loop emulation of spacecraft relative dynamics and a physical setup for testing tethered systems. The architecture consists of a granite slab, supported by a universal joint and two linear actuators to control its orientation. This configuration allows a Floating Spacecraft Simulator to move on the surface in a quasi-frictionless environment under the effect of the gravitational acceleration. The architecture includes a dedicated setup to emulate tethered satellite dynamics, providing continuous feedback on the tension along the tether through a mono-axial load cell. By adopting the Buckingham “ π ” theorem, the dynamic similarity is introduced for the ground-based experiment to reproduce the orbital dynamics. Proof-of-concept results demonstrate the testbed’s capability to accurately reproduce the Hill-Clohessy-Wilshire equations. Moreover, the results of the deployed tethered system dynamics are presented. The paper also details the system architecture of the testbed and the methodologies employed during the experimental campaign.

Keywords: Hardware-In-the-loop Testbed; Floating Spacecraft Simulator; Tethered Satellite Systems; Real-time Control Algorithms

1. Introduction

Ground-based testbeds play a crucial role in validating and testing all spacecraft at the hardware and software level and for operational procedures before in-orbit commissioning. Ground tests help identify potential problems and optimize system performances in a controlled setting, reducing mission risks and costs [1–5]. In this context, simulations in environments that emulate space conditions are essential, especially when leading experiments on the attitude and orbital dynamics of space systems, which are required to ensure mission success [6].

There are currently different methods for emulating microgravity conditions, such as drop towers, parabolic flights, and sounding rockets, which still offer reduced gravity conditions, respectively, of up to 10 seconds, 25 seconds, and more than 1 minute [7–10].

Alongside these facilities, the space-flight environment can offer longer microgravity conditions. The International Space Station often carries special laboratories for conducting scientific research, in similarity to the Space Shuttle until its retirement [10].

Another method to emulate microgravity is represented by the air-bearing platforms, which rely on reducing the friction on the involved table [7,11]. These systems are used for validating Guidance, Navigation, and Control (GN&C) algorithms in proximity operations, formation flying, rendezvous,

and docking manoeuvres [12–16]. Air-bearing testbeds are based on having a reduced-scale Floating Spacecraft Simulator (FSS) moving on an accurately levelled and smooth surface. In this paper, an air-bearing platform for emulating orbital dynamics is presented. The testbed is based on controlling the attitude of the granite slab, which can be tip-tilted using two linear actuators and one universal joint. This setup allows the non-actuated FSS to respond to the granite surface inclination with reactivity, effectively replicating accelerations being similar to those a satellite would experience in orbit. This work builds up on the prototype proposed by Ogundele et al. [17].

The main advantage of the testbed is that continuous accelerations can be reproduced on the non-actuated FSS. Unlike state-of-art setups, in which thrusters are onboard the floating platform and impulsive forces are generated, the proposed configuration allows the emulation of non-impulsive forces, effectively simulating disturbances that are continuously acting on satellite systems in orbit.

The present work also introduces the integration of a system specifically designed for testing control algorithms oriented to tethered satellite systems. In detail, the floating platform is also equipped with a tether which is connected to a load cell to measure its tension. The load cell is mounted on a vertical support on a linear guide, enabling dynamic simulations of tethered satellite systems during different operational phases. This setup is developed with a different approach with respect to previous research on testing tethered satellite systems on ground-based frictionless tables [18–20]. The proposed architecture leverages the capability of a novel tip-tilting facility to reproduce continuous and precise accelerations.

The paper is organised as follows: The architecture of the testbed is introduced in section 2. The mathematical and on-orbit modelling, on which the testbed is based, are described in section 3. To demonstrate the testbed working capabilities, a test campaign is conducted, described in section 4. Numerical results of the experimental campaign are presented and discussed in section 5. Finally, conclusions are resumed and drawn in section 6.

2. System Architecture

This section introduces the testbed, which is subdivided into the following subsystems: the Floating Spacecraft Simulator (FSS), an air-bearing scaled and simplified model that represents the satellite; the Floating Platform (FP), which includes the granite slab, the actuator system, and the laser sensors, standing as central stage on which the FSS operates; the Motion Capture System (MCS), the Tether Testing System (TTS), and the Master Console (MC). The principal units of the testbed are represented in the breakdown tree in Figure 1, while Figure 2 shows the setup of the testbed.

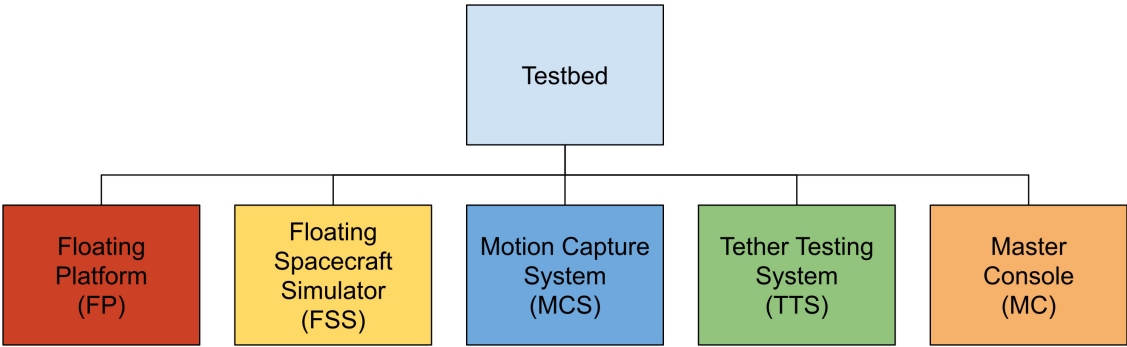


Figure 1. Testbed’s units.

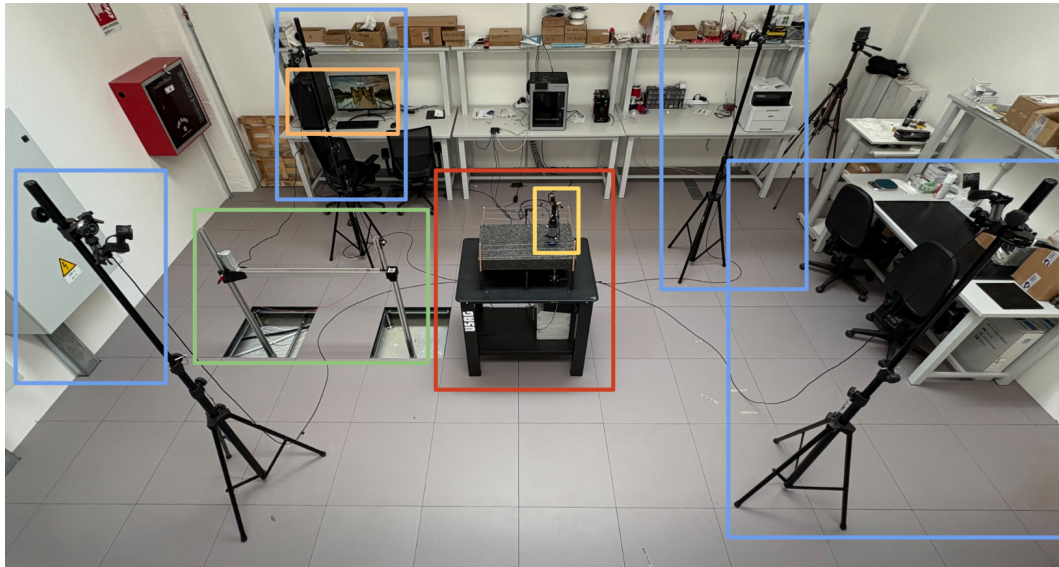


Figure 2. Laboratory setup. The subsystems of the testbed are framed with the corresponding colours, referenced in Figure 1.

2.1. Floating Platform

The FP is shown in Figure 3. It features a polished granite slab (1), tilted by two linear actuators (2). Laser sensors (3) can be used to measure the relative attitude of the granite surface. The proper kinematics of the mechanism is then ensured by one spherical joints' couple, which links the linear actuators to the slab, and a universal joint (4). This setup enables the slab to be tilted around two axes, allowing a rigid body positioned on the granite table to experience a simulated gravitational force.

Granite is identified as the most suitable material for the slab, as an extremely flat and smooth surface is required for air-bearing platforms. The granite slab, manufactured by Zali Precision, is polished with high precision, resulting in very low surface roughness and high surface flatness. Moreover, it has great rigidity and high thermal stability. In this way, it does not warp or expand with temperature changes. Furthermore, unlike metal, granite has good damping properties, helping to absorb vibrations to avoid oscillations.

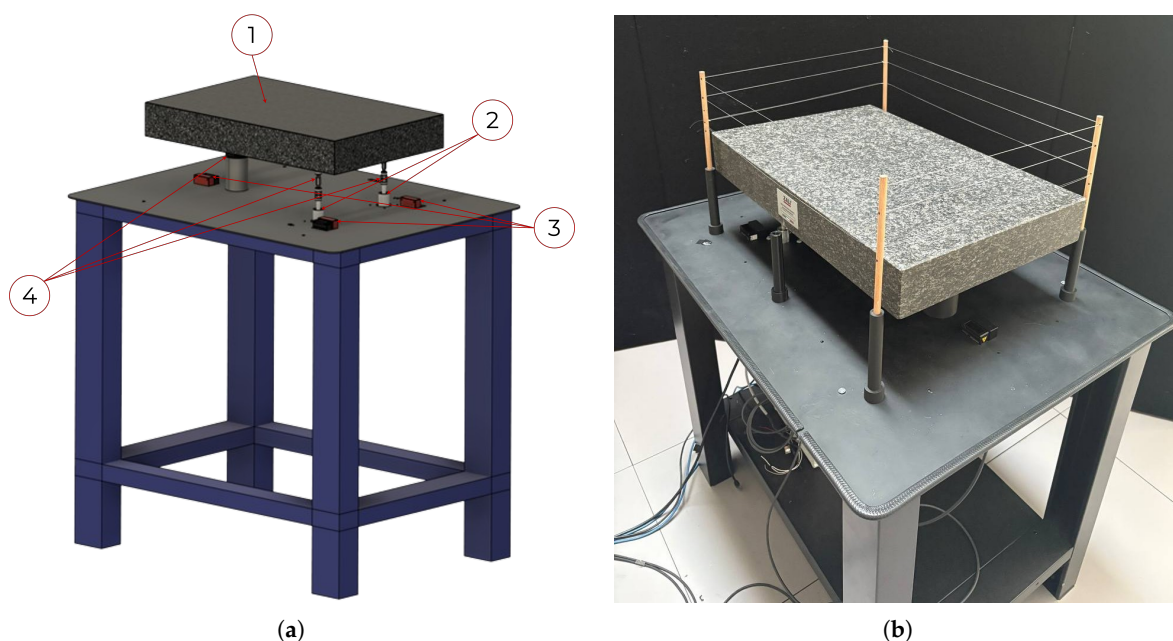


Figure 3. 3(a) 3D model of the FP. 3(b) Physical apparatus of the FP.

The linear actuators, LGA56 – Captive linear actuator – NEMA 23 by Nanotec, are powered by a power supply unit and actuated via motor controllers, C5-E series by Nanotec, which communicate with the MC via Modbus TCP protocol. The testbed is also equipped with three laser sensors, optoNCDT-1220 of Micro-Epsilon series, working in VIS-RF (VISible RadioFrequency, red semiconductor laser). Lasers are used to measure the relative displacement of three points on the tilting surface with respect to the reference plane. Those three points can be used to evaluate the attitude of the table. The motors and lasers are placed on a steel base attached to the bottom of the support frame, as shown in [Figure 3](#). The specifics of each equipment are reported below in [Table 1](#).

Table 1. Floating Platform components’ properties.

GRANITE SLAB	Value	Unit of Measure
Manufacturer	Zali Precision	
Size	630 × 400 × 80	[mm × mm × mm]
Mass	60	[kg]
Planarity	2.5	[μm]
ACTUATOR		
Manufacturer	Nanotech	
Quantity	2	
Maximum speed	30	[mm/s]
Step resolution	0.01	[mm/step]
Motor holding torque	1.12	[Nm]
Stroke	38.1	[cm]
Step angle	1.8	[deg]
LASER SENSOR		
Manufacturer	Micro-Epsilon	
Quantity	3	
Power voltage	24	[V]
Power consumption	<2	[W]
Start of Measuring Range SMR	5	[cm]
End of measuring Range EMR	15	[cm]
Measuring range MR	10	[cm]
Weight	30	[g]
Wavelength	670	[nm]

2.2. Floating Spacecraft Simulator

The Spacecraft Simulator developed at Politecnico di Torino [\[21\]](#) is a revised version of that designed and tested at the Naval Postgraduate School of Monterey in 2022 [\[22\]](#). The FSS is always characterized by a pneumatic subsystem and an air-bearing, shown in [Figure 4](#). It can move on the monolithic granite slab depending on the adopted control strategy. The structure is composed by a 3D printed tank container, an air bearing, and a 3D printed component to fix the tether anchor point, to let the markers being placed to track the FSS through the optical capture system. The pneumatic subsystem, schematized in [Figure 5](#), includes valves and gauges to regulate and measure air pressure, an air-tank, and a manual valve that allows compressed airflow to reach the bearing. The prenumatic system is needed to create an air film between the FSS and the matching plane to establish a quasi-frictionless environment. The air-tank operates at 200 bar (3000 psi) and has a volume of 0.21 L.



Figure 4. Air tank container with the air-bearing laid on a microfiber cloth, reflective markers (left) and pneumatic system (right).

It has a built-in regulator that delivers air at a pressure of 60 bar (800 psi), with maximum utilisation pressure of about 310 bar (4500 psi). Since the air bearing works at an operating pressure of 4.1 bar, an additional 90° pressure regulator is fed by the air directly coming from the high-pressure one. The output pressure can be monitored by a low-pressure gauge. The pressure relief knob is attached to the low-pressure regulator to regulate it to a desired value, and a manual valve that allows the airflow to be simply shut off or enabled. The air-bearing is flat and round, with a diameter of 65 mm. The entire apparatus weighs approximately 1 kg, taking into account both the tank container and the air bearing, whereas the air-tank has an approximate operative time equal to 1 hour, between recharges, for supporting experiments’ phases.

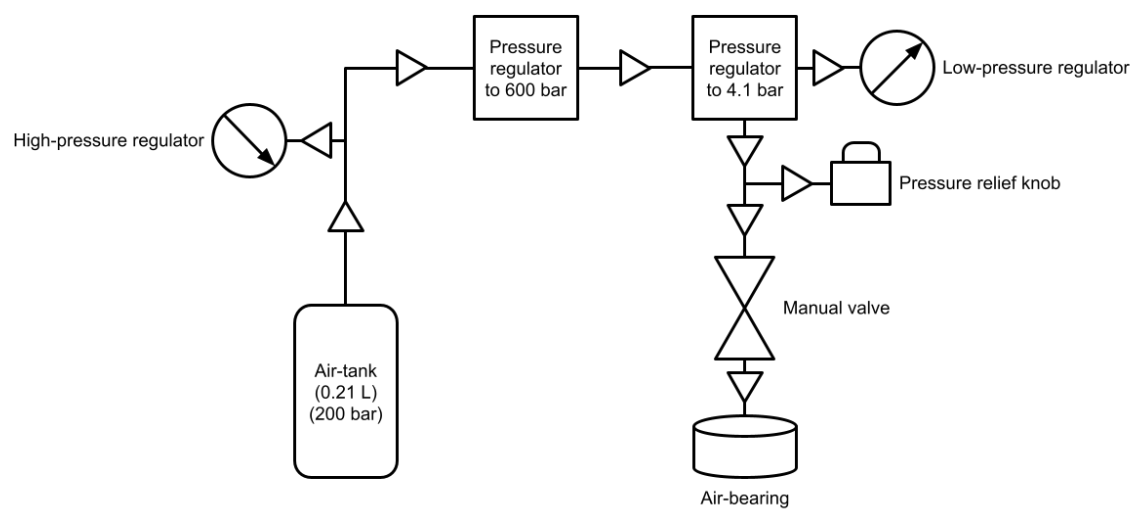


Figure 5. Detailed schematics of the FSS pneumatic system.

As already described, around the 3D printed container are fixed four high-reflective markers to let the motion capture system track the FSS on the granite monolith. Eventually, the 90° boulder CO₂ low-pressure regulator allows the insertion of an external gauge and guarantees the connection between the opening valve and the below-standing flat air-bearing through the air-hose. The air-tank is refilled by an air compressor that harnesses an Oil-Water separator, a carbon filter rod and one sponge to reduce moisture in the tank as more during recharge.

2.3. Motion Capture System

The MCS is essential for accurately tracking the floating system, which would let precise position control. The system comprises four InfraRed (IR) cameras, OptiTrack Flex-13 series, mounted on tripods. These cameras would enable to observation of the tracked volume, where the table system is positioned, on which the FSS moves. Additional cameras can be easily integrated if needed to expand the tracking volume. Furthermore, tripods can be easily moved to positions that frame different desired areas. To track the position of objects within the volume, special markers are adopted. These markers (bolded points in Figure 6) are placed on the FSS system, allowing its position to be continuously monitored. The information on the FSS displacements and velocities are extrapolated by numerically plotting them over a short period.

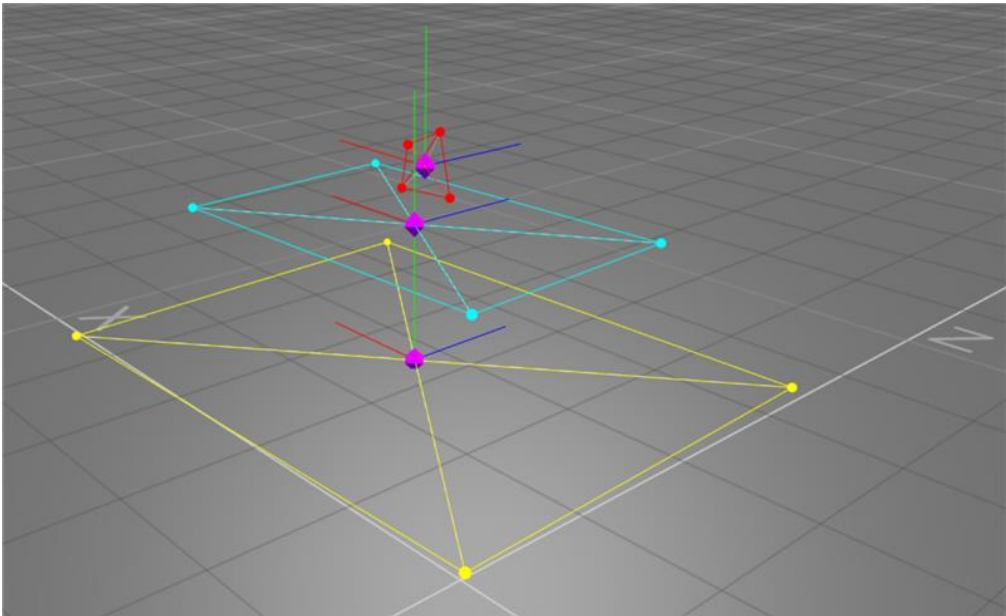


Figure 6. MCS visualization output, evidencing the FSS (red), the granite table (cyan) and the stainless-steel table (yellow).

Table 2. Characteristics of the MCS.

MCS	Value	Unit of Measure
Manufacturer	OptiTrack	
Cameras quantity	4	
Resolution	1.3	[Mega Pixel]
Latency	8.3	[ms]
Position accuracy	± 0.20	[mm]
Attitude accuracy	± 0.50	[mm]
Frame rate	120	[FPS]
Cover area	9 × 9	[m × m]
Tripod height	215	[cm]

2.4. Tether Testing System

The TTS is studied to carry out tests in both deployed (see [Figure 7](#) and [Figure 8](#)) and under deployment scenarios (see [Figure 9](#)). Two different configurations are proposed, accounting for the two scenarios.

The deployed configuration is studied by a setup comprising an aluminium support structure, a linear guide with an actuated prismatic joint and a load cell, of the LSB-200 series by Futek, for real-time tension measurement. Specifically, two aluminium rods and the linear guide form a U-shaped support structure where an extra aluminium profile is fixed with the load cell attached to it. The tether is then connected to the FSS and the load cell with two knots. In this case, the guide does not allow relative motion, though it ensures alignment and structural support.

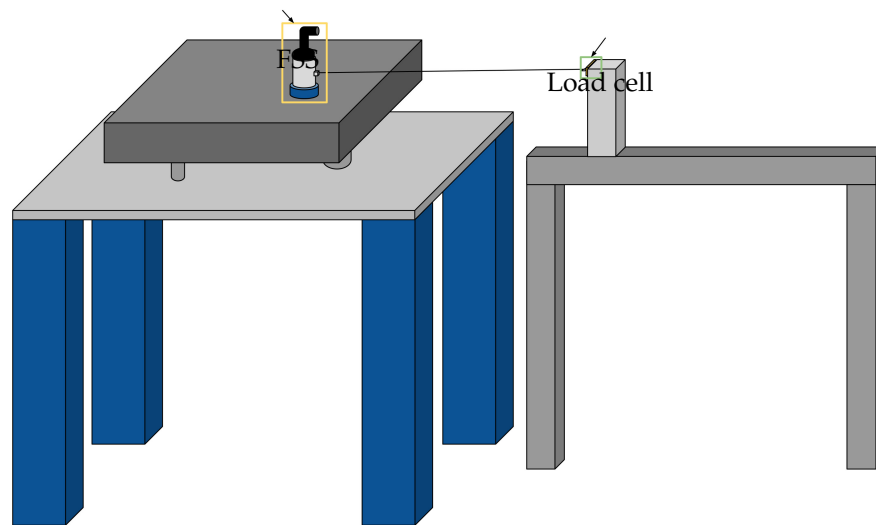


Figure 7. Illustration of the TTS-deployed (static) architecture.



Figure 8. TTS-deployed (static) architecture.

The setup for the deploying case needs an actuated spool to provide tension control during the deployment. The setup for this configuration (Figure 9) derives from the desire of keeping the FSS as lightweight as possible, with the aim of maximizing its usage time between two recharges. As a result, both the actuated spool and the load cell, features shown in Table 3, have been mounted on the vertical support, with the tether being wrapped on the actuated spool. Its free-end extends toward the FSS, which houses a passive spool, before being routed back to the vertical support, where it is secured to the load cell. With this setup, it is possible to perform a controlled deployment, where a constant value of tension is achieved by the combined action of the linear guide and the actuated spool. From a dynamic standpoint, the presence of the double cable attached to the FSS should not affect its dynamic response as long as one keeps in mind that the true force acting on the FSS is twice the value measured by the load cell. The deployment scenario is provided as key-finding for future integration to the proposed testbed.

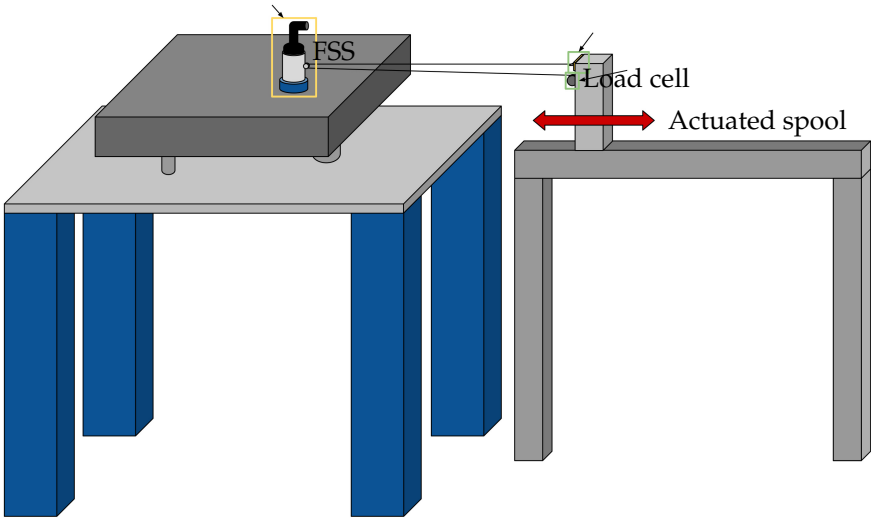


Figure 9. Illustration of the TTS-deployment (dynamic) architecture.

Table 3. Tether subsystem’s properties.

LOAD CELL	Value	Unit of Measure
Manufacturer	Futek	
Quantity	1	
Max load	2	[mN]
Natural frequency	140	[Hz]
Mass	0.0193	[kg]
Tether length	1	[m]

3. Mathematical Models for On-orbit Dynamics

The testbed is developed to study the relative dynamics of two satellites in a closed-proximity configuration, generally defined as target and chaser.

3.1. In-Orbit Dynamical Model

Under the hypothesis that the target moves on a circular orbit and its relative distance is smaller than the target’s orbital radius, the Hill-Clohesy-Wilshire (HCW) linearized equations of motion hold true. The HCW equations are written in the Local-Vertical-Local-Horizontal (LVLH) reference frame

shown in Figure 10. The framework at hand is a Cartesian Coordinate System (CCS) centred in the target's (T) centre of mass, where the $\hat{\xi}$ axis lays on the radial direction from the primary body to the target; the $\hat{\eta}$ axis lays on the along-track direction; the $\hat{\zeta}$ axis completes the right-hand triad being perpendicular to the orbital plane, in the same direction of the angular momentum, whereas R the primary body radius. The subsequent Equations of Motion (EoM) are written in Equation 1,

$$\begin{aligned}\ddot{\xi} - 2\Omega\dot{\eta} - 3\Omega^2\xi &= \frac{F_{\xi}}{M} \\ \ddot{\eta} + 2\Omega\dot{\xi} &= \frac{F_{\eta}}{M} \\ \ddot{\zeta} + \Omega^2\zeta &= \frac{F_{\zeta}}{M}\end{aligned}\quad (1)$$

where $\chi = [\xi, \eta, \zeta]^T$ represents the components of the position vector of the chaser Center of Mass (CoM) with respect to the target's CoM, expressed in LVLH ccs. Dot notation stands for the derivative with respect to the time τ .

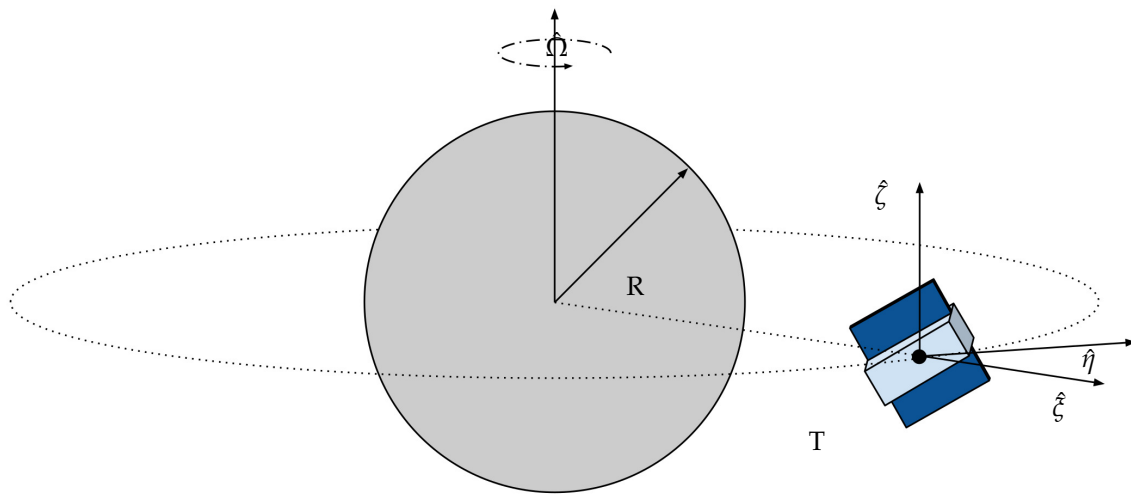


Figure 10. Local-Vertical Local-Horizontal Cartesian Coordinate System with origin at the Center of Mass of the Target (T).

$\Omega = \sqrt{\frac{\mu}{R^3}}$ is the orbital angular velocity of the target spacecraft, $F_{\chi} = [F_{\xi}, F_{\eta}, F_{\zeta}]^T$ is the resultant of forces acting on the chaser and M represents its mass.

3.2. Buckingham theorem and scaled orbital dynamics

To replicate the relative orbital dynamics, on a smaller and compact testbed, it is necessary to scale the HCW equations [23]. Based on [17,24], the dynamic similarity is ensured between the two systems by applying the Buckingham π theorem and the *principle of similarity*. The Buckingham π theorem states that if a phenomenon is described by p variables, with q fundamental units involved, the entire system can be expressed using $r = p - q$ dimensionless parameters, called π variables. By imposing equality between the π variables of the real case and those of the scaled case, the *principle of similarity* ensures dynamic similarity between the two [25,26]. The system involves a total of 15 variables:

$$\left[\xi \quad \eta \quad \zeta \quad \dot{\xi} \quad \dot{\eta} \quad \dot{\zeta} \quad \ddot{\xi} \quad \ddot{\eta} \quad \ddot{\zeta} \quad F_{\xi} \quad F_{\eta} \quad F_{\zeta} \quad M \quad \tau \quad \Omega \right]$$

A new set of variables is introduced for the scaled system:

$$\left[x \ y \ z \ \dot{x} \ \dot{y} \ \dot{z} \ \ddot{x} \ \ddot{y} \ \ddot{z} \ F_x \ F_y \ F_z \ m \ t \ \omega \right]$$

It is possible to recognize that there are three fundamental units: length, mass, and time. As a consequence, a total of 12 π variables can be identified and used to establish the similarity conditions between the two systems. The final result is shown in the system of equations [Equation 2](#). In detail, the scaling factors for the three fundamental quantities, namely λ_L , λ_m , and λ_t , are defined as the ratio between the real case and the scaled variables.

$$\begin{aligned} \lambda_t &= \frac{\tau}{t} \\ \lambda_m &= \frac{M}{m} \\ \lambda_L &= \frac{\xi}{x} = \frac{\eta}{y} = \frac{\zeta}{z} \end{aligned} \quad (2)$$

The scaling parameters for the derived quantities involved, such as velocity, acceleration, force, and mean motion, are then evaluated as a function of the fundamental ones. According to the Buckingham theorem, all scaling factors can be written as a function of the three factors referring to the fundamental units. In other words, when a test shall be performed, it is sufficient to set only λ_m , λ_L and λ_t to obtain all the scaled quantities involved with the phenomenon. By substituting relations within [Equation 2](#) in [Equation 1](#), the system of [Equation 3](#) is obtained in a testbed-fixed reference frame coordinates that will be defined in the next paragraph.

$$\begin{aligned} \ddot{x} - 2\omega\dot{y} - 3\omega^2x &= \frac{F_x}{m} \\ \ddot{y} + 2\omega\dot{x} &= \frac{F_y}{m} \\ \ddot{z} + \omega^2z &= \frac{F_z}{m} \end{aligned} \quad (3)$$

where $\underline{F} = \underline{F}_{ctrl} + \underline{F}_{disp}$ is the sum of the control action and the disturbances.

3.3. Testbed system and models

The novel idea behind the proposed testbed is that the plane on which the FSS can freely move is actuated, generating accelerations on the FSS by tip-tilting the matching plane. This section explores the connection between the actuators' displacements and the acceleration induced on the FSS due to the resulting plane inclination. For the purpose of this calculation, a new reference frame is introduced. As shown in [Figure 11](#), which illustrates the plant of the granite, P_1 and P_2 are the points where the two linear actuators are mounted, P_0 is the point where the pillar is mounted, d is the distance between the two actuators, l the distance between any of the actuators and P_0 , and p and q the sides of the slabs. The new CCS is fixed to the granite slab, centred on the rotation pole of the system P_0 , and orientated with \hat{x} axis parallel to p , \hat{y} parallel to q , and \hat{z} axis completes the right-hand triad.

It is now essential to determine the relationship between the actuators' displacements, $[z_1, z_2]$ and the acceleration induced on the FSS. This relationship has already been established by [\[17,24\]](#), and only the most relevant results are discussed here.

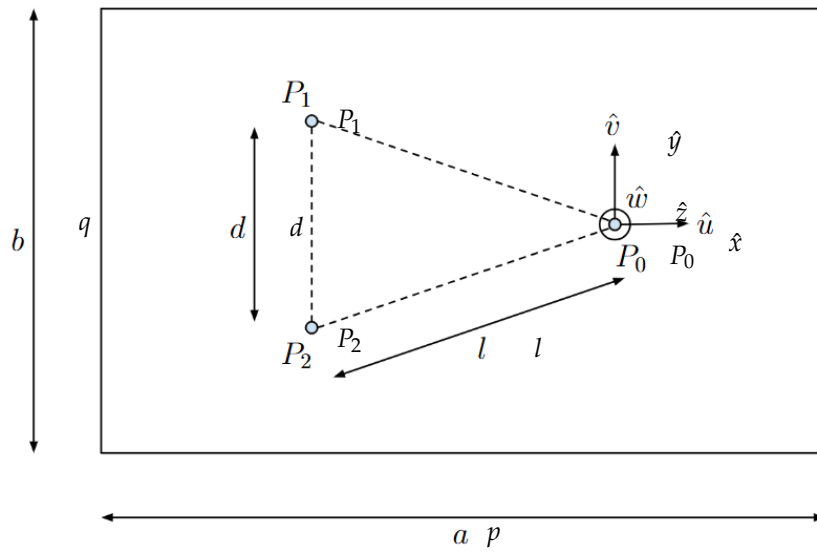


Figure 11. Granite monolith geometrical characteristics.

The position of points P_1 and P_2 are respectively

$$\begin{aligned} \mathbf{r}_{P_1} &= [x_1, y_1, z_1]^T \\ \mathbf{r}_{P_2} &= [x_2, y_2, z_2]^T \end{aligned} \quad (4)$$

Using a Euler angle and axis parametrization, when a small displacement z_1 is applied, a rotation of $\alpha \simeq \frac{z_1}{d}$ is made around the vector $\mathbf{e}_1 = \frac{1}{l}[-x_2, -y_2, 0]^T$, see [Figure 12](#).

The new attitude of the matching plane is expressed by the Direction Cosine Matrix (DCM) in [Equation 5](#).

$$\mathbf{C}_1 = \frac{1}{ld} \begin{bmatrix} 1 & 0 & y_2 z_1 \\ 0 & 1 & -x_2 z_1 \\ -y_2 z_1 & x_2 z_1 & 1 \end{bmatrix} \quad (5)$$

When a displacement z_2 is applied, a rotation of $\beta \simeq \frac{z_2}{d}$ around $\mathbf{e}_2 = \frac{1}{l}[x_1, y_1, 0]$ is triggered and the subsequent DCM is obtained in [Equation 6](#), see [Figure 13](#).

$$\mathbf{C}_2 = \frac{1}{ld} \begin{bmatrix} 1 & 0 & -y_1 z_2 \\ 0 & 1 & x_1 z_2 \\ y_1 z_2 & -x_1 z_2 & 1 \end{bmatrix} \quad (6)$$

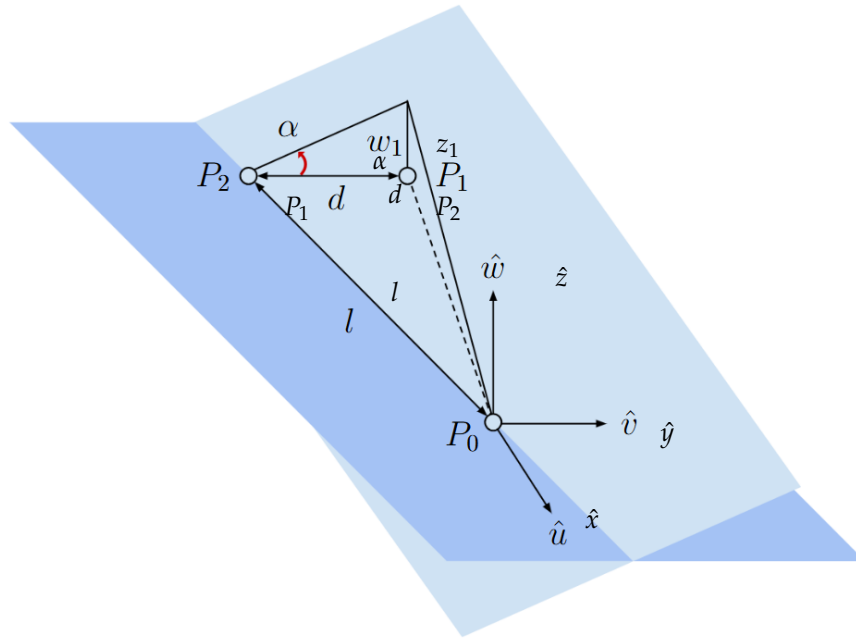


Figure 12. Schematic view of first linear actuator displacement.

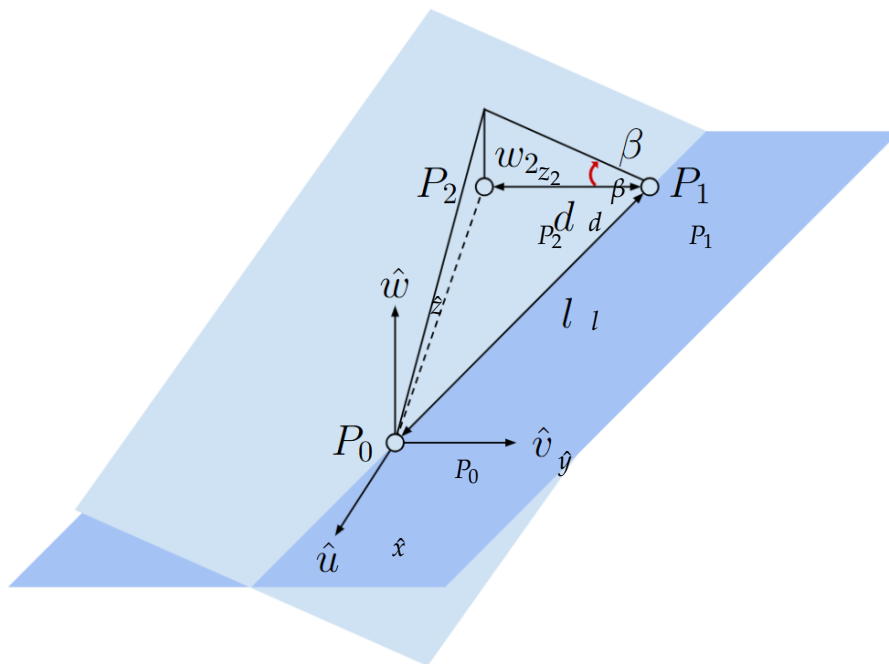


Figure 13. Schematic view of second linear actuator displacement.

Finally, the attitude of the table after two generic rotations can be expressed as the matrix multiplication of the two DCMs. Assuming the hypothesis of small displacement holds, the second or higher order terms are neglected. The resulting matrix C is obtained in [Equation 7](#).

$$C = C_1 C_2 = \frac{1}{ld} \begin{bmatrix} 1 & 0 & -y_1 z_2 + y_2 z_1 \\ 0 & 1 & x_1 z_2 - x_2 z_1 \\ y_1 z_2 - y_2 z_1 & -x_1 z_2 + x_2 z_1 & 1 \end{bmatrix} \quad (7)$$

As a consequence, the gravitational acceleration experienced by the FSS in the reference frame of the table is determined by Equation 8.

$$C[0, 0, -g]^T = -\frac{g}{ld} \begin{bmatrix} -y_1 z_2 + y_2 z_1 \\ x_1 z_2 - x_2 z_1 \\ 1 \end{bmatrix} \quad (8)$$

As shown in Equation 8, the acceleration along \hat{z} is dependent only on the gravitational acceleration, which is balanced by the plane's normal force, while the other two components are related to the movement of actuators. For this reason, Equation 9 can be written.

$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \end{bmatrix} = \frac{g}{ld} \begin{bmatrix} -y_2 & y_1 \\ x_2 & -x_1 \end{bmatrix} \begin{bmatrix} z_1 \\ z_2 \end{bmatrix} \quad (9)$$

By inverting Equation 9, it is possible to write the table's steering law in Equation 10, which represents the required displacements z_1, z_2 that are necessary to obtain the desired acceleration components $\ddot{x}_{des}, \ddot{y}_{des}$.

$$\begin{bmatrix} z_1 \\ z_2 \end{bmatrix} = \frac{ld}{g} \begin{bmatrix} -y_2 & y_1 \\ x_2 & -x_1 \end{bmatrix}^{-1} \begin{bmatrix} \ddot{x}_{des} \\ \ddot{y}_{des} \end{bmatrix} \quad (10)$$

Knowing the relations between the displacements of both the actuators and the accelerations experienced by the FSS due to tilting the granite slab, it is possible to obtain a surface-map of the reproducible acceleration. As shown in Figure 14, the acceleration a , which is the norm of the vector $[\ddot{x}, \ddot{y}]^T$, can be plotted as a function of z_1 and z_2 , individuating the accelerations that can be reproduced on the table. In detail, knowing the characteristics of the actuators, resumed in Table 1, and imposing a security coefficient that limits their strokes to 30 mm (± 15 mm), the maximum and minimum reproducible accelerations are determined, reported in Table 4.

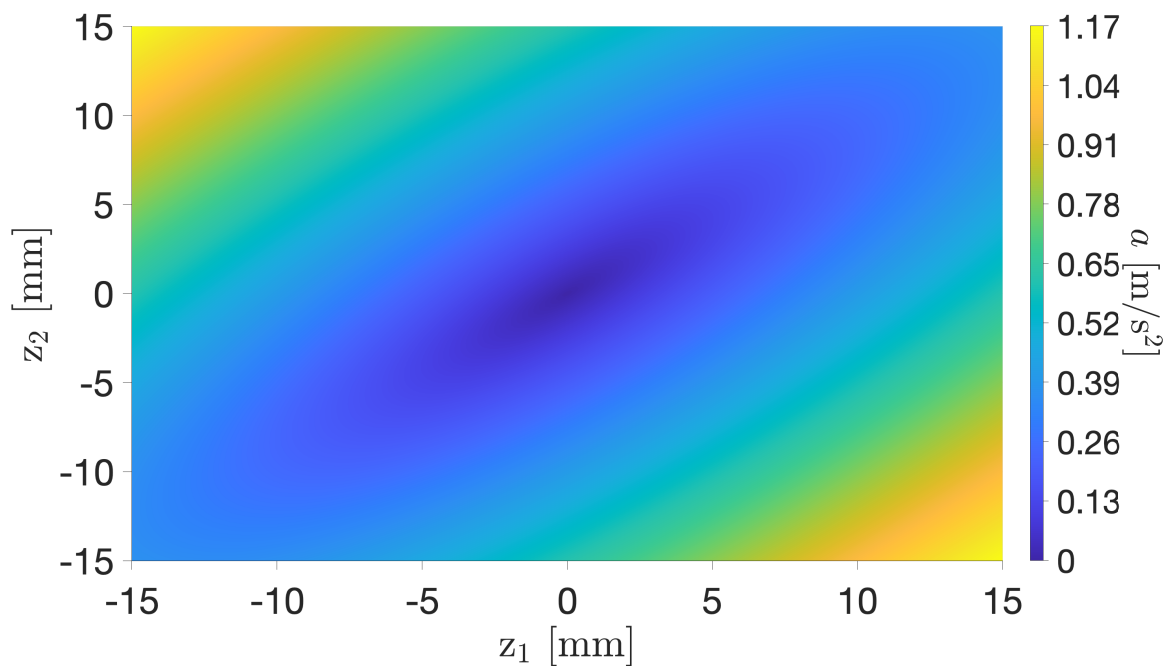


Figure 14. Map of reproducible acceleration on the monolithic granite table.

Table 4. Maximum and minimum acceleration reproducible on the matching plane.

Maximum acceleration	1.1688	m/s ²
Minimum acceleration	0.0025	m/s ²

The map in figure [Figure 14](#) refers to an ideal case where the supporting structure where the actuators are mounted is perfectly flat. However, taking into consideration [18], the testbed may present a slight residual inclination. That may be caused by all the necessary manufacturing processes that might have plastically deformed the structure at the micrometric level. For this reason, a worst-case scenario with an expected residual acceleration of the order of magnitude of the minimum reproducible value of 0.0025 m/s² is considered.

4. Test Campaign

The test campaign is crucial for the testbed. In detail, the campaign is planned not only to conduct a test related to the dynamic control of a tethered system, but also to validate the apparatus. In detail, to assess the capabilities and the performance of the tip-tilting testbed, a first test is conducted to verify that the system is able to work in closed-loop, autonomously responding to inputs on the system, while a second test is planned to assess the performance of reproducing the orbital dynamics of Hill-Clohessy-Wiltshire equations.

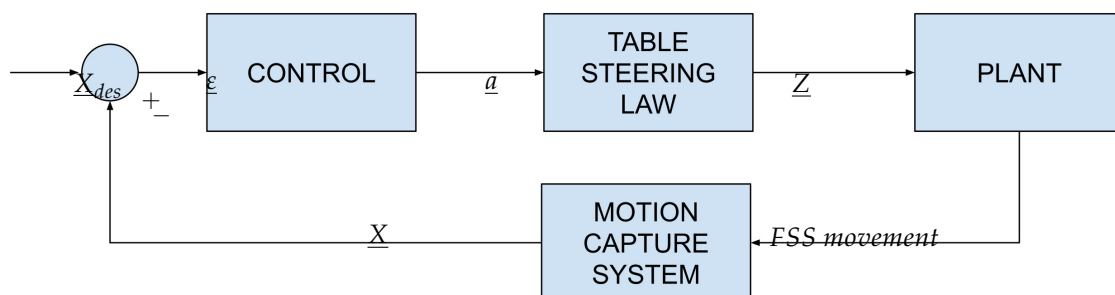
4.1. Closed-Loop Calibration

The first test aims to control the granite table's attitude, achieving a horizontal configuration. In order to do that, a closed-loop approach is implemented via a Proportional Integrative Derivative (PID) controller [27]. By tracking the FSS's position and velocity with the MCS, the goal is to guide it towards the middle point of the granite surface with zero velocity. As a result, the control action imposed by the table is evaluated by [Equation 11](#):

$$\underline{a} = \underline{K}_P(\underline{X}(t) - \underline{X}_{des}) + \underline{K}_D(\dot{\underline{X}}(t) - \dot{\underline{X}}_{des}) + \underline{K}_I \int_0^t (\underline{X} - \underline{X}_{des}) dt \quad (11)$$

where $\underline{X}_{des} = [x_{des}, z_{des}]^T$ and $\dot{\underline{X}}_{des} = [\dot{x}_{des}, \dot{z}_{des}]^T$ represent the desired final position and velocity. \underline{K}_P , \underline{K}_D , \underline{K}_I are the gain matrices needed to compute the PID controller.

The acceleration calculated in real-time with [Equation 11](#) is then transformed into information on the displacement and fed into the plant. The evolution of the dynamics in terms of position and velocity is then extracted by the motion capture system and used as feedback for the control law. The schematic in [Figure 15](#) summarizes the process.

**Figure 15.** Flow chart of the closed-loop calibration algorithm.

4.2. Emulation of Linearized Relative Dynamics

The objective of this test is to evaluate how accurately the testbed can replicate the orbital relative dynamics governed by the HCW equations of motion. For this case study, orbital perturbations are

neglected, and the motion is analysed solely within the orbital plane. As a consequence, the system in Equation 1 is reduced to 2 degrees of freedom and the external forces are set to zero, obtaining the set of coupled differential equations in Equation 12 with the initial condition of $\chi_0 = [\xi_0, \eta_0, \dot{\xi}_0, \dot{\eta}_0]^T$.

$$\begin{aligned}\ddot{\xi} - 2\Omega\dot{\eta} - 3\Omega^2\xi &= 0 \\ \ddot{\eta} + 2\Omega\dot{\xi} &= 0\end{aligned}\tag{12}$$

To set up the experiment, it is necessary to relate the LVLH CCS with the table-fixed CCS. Without loss of generality, the $\hat{\xi}$ and $\hat{\eta}$ are set to be coincident with \hat{x} and \hat{y} . Secondly, scaling factors are to be selected. In particular, after propagating the real system trajectory and defined $\xi_{max} = \max(\xi(t))$ and $\eta_{max} = \max(\eta(t))$, they can be used to evaluate the λ_L . Specifically, to guarantee that the FSS does not exceed the table's boundaries, the constraint in Equation 13 shall be respected.

$$\lambda_L \geq \max\left(\left|\frac{\xi_{max}}{p}\right|, \left|\frac{\eta_{max}}{q}\right|\right)\tag{13}$$

Additionally, a constraint must also be set on acceleration. Indeed, as seen in subsection 3.3, a residual acceleration of the order of magnitude of 0.0025 m/s², see Table 4, is considered. To replicate the relative dynamics governed by the HCW equations of motion (EoM) in the absence of control, it is crucial to ensure that residual accelerations do not significantly alter the FSS dynamics. Moreover, the scaled accelerations should also not be greater than the maximum reproducible value in Table 4. To mitigate these risks, the constraints in Equation 15 are enforced, ensuring that the scaled HCW acceleration is dominant over unwanted perturbations.

$$a_{min} = \frac{\|\ddot{\chi}_{min}\|\lambda_t^2}{\lambda_L} \geq 10 \cdot 0.0025 \text{ m/s}^2\tag{14}$$

$$a_{max} = \frac{\|\ddot{\chi}_{max}\|\lambda_t^2}{\lambda_L} \leq 1.169 \text{ m/s}^2\tag{15}$$

As the scaling factors are chosen, the equations of HCW dynamics are scaled, allowing to write Equation 16.

$$\begin{aligned}\ddot{x} - 2\omega\dot{y} - 3\omega^2x &= 0 \\ \ddot{y} + 2\omega\dot{x} &= 0\end{aligned}\tag{16}$$

A new Matlab/Simulink script is developed, and the workflow is shown in Figure 16.

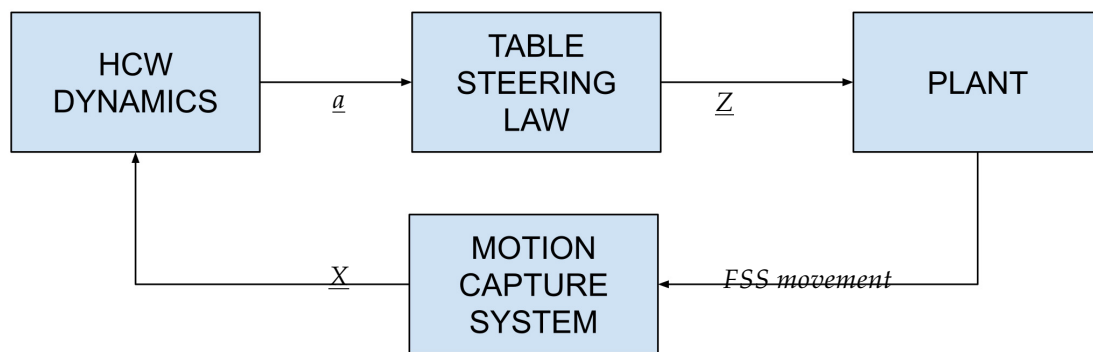


Figure 16. Flow chart for the HCW emulation test.

4.3. Fully-Deployed Tether Dynamics

The objective of the third experiment is to test the dynamics and control algorithm of a novel cross-track tethered satellite system. Unlike a conventional tethered spacecraft that primarily exploits the gravity gradient torque for radial alignment, the system of interest consists of two satellites connected by a tether that is extended along the cross-track direction (\hat{c} axis in the LVLH reference frame, Figure 17). In this configuration, the tethered system maintains a perpendicular orientation to the orbital plane. This orientation is particularly suited for remote sensing applications, especially for radar imaging, since multiple radar instruments can be positioned along the tether, obtaining a distributed radar configuration [28]. A key advantage of this architecture is its potential to minimize propellant consumption for formation flying [28]. While similar tethered configurations have been proposed for applications such as position keeping and rendezvous maneuvers [29], the present configuration introduces a novel tethered system concept, where the tether is maintained in cross-track orientation by generating control forces, using thrusters, or aerodynamic drag plates if the spacecraft is flying at Very-Low Earth Orbits (VLEOs) altitude [30]. The TSS is composed of two CubeSat platforms which are connected by a tether with a length that can reach hundreds of meters. If the spacecraft operates in VLEO, each satellite is equipped with two triangular-shaped drag plates designed to interact with the residual atmospheric environment [31], generating control forces that mitigate external disturbances and maintain the tether in a tensioned state, thereby stabilizing the formation. If no residual drag can be exploited, thrusters can be used to stabilize the spacecraft system. The system dynamics is analysed within the LVLH reference frame. The main goal is to assess the testbed's ability to accurately replicate the orbital dynamics of a TSS in a cross-track configuration.

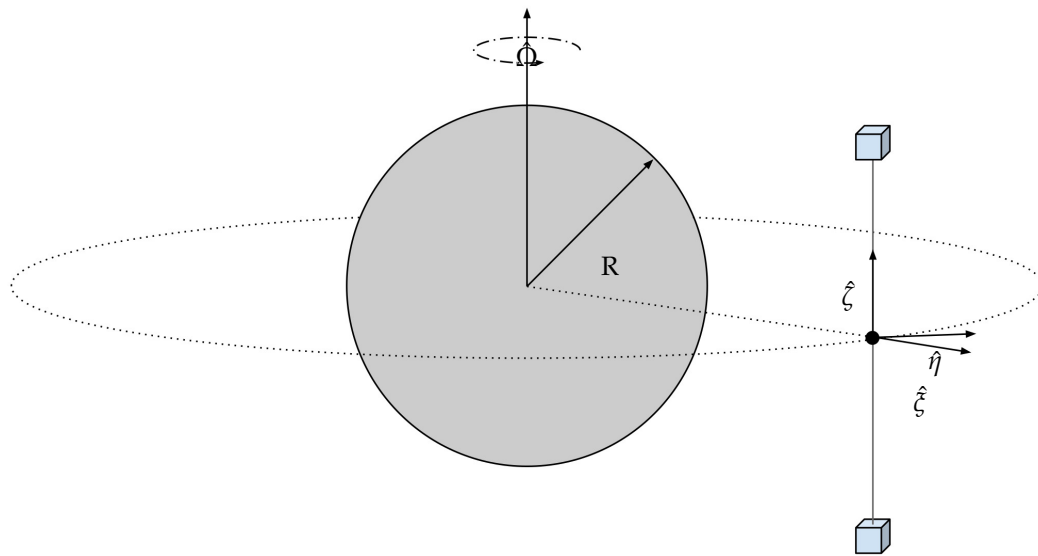


Figure 17. Illustration of a stabilized TSS in a cross-track configuration.

Specifically, considering the system in circular orbit at an altitude of 500 km, where a tether of 100 m is connected to two CubeSat platforms of 20 kg and is kept in tension by thruster propulsion, and with an initial off-nominal condition where the two satellites are misaligned with respect to the across-track direction, the controller is tasked with correcting the positional error and restoring proper cross-track alignment. Additionally, once the satellites are correctly positioned, the controller is designed to maintain constant tension along the tether.

For the experimental campaign, the TTS-deployed configuration is used (see Figure 7). The setup allows for the simulation of half of the cross-track tethered satellite, where the dynamics of just one CubeSat is analysed. The CCS is centred on the load cell, at one end of the testbed tether, which corresponds to half of the TSS. The CCS, with respect to the reference frame in Figure 11, is oriented with $\hat{\xi}$, $\hat{\zeta}$, and $\hat{\eta}$ axes concordant, respectively, with \hat{x} , \hat{y} , and \hat{z} vectors. As the radial and cross-track directions lie down on the matching plane, the dynamics along $\hat{\eta}$ is not considered, and the HCW equations of motion are reduced to two degrees of freedom (Equation 17):

$$\begin{aligned}\ddot{\xi} - 3\Omega^2\xi &= \ddot{\xi}_{ctrl} + \frac{T_{\xi}}{m} \\ \ddot{\zeta} + \Omega^2\zeta &= \ddot{\zeta}_{ctrl} + \frac{T_{\zeta}}{m}\end{aligned}\tag{17}$$

where $\ddot{\xi}_{ctrl}$ and $\ddot{\zeta}_{ctrl}$ are the control actions, while T_{ξ} and T_{ζ} are the components of the tension along the tether.

Given the order of magnitude of the acceleration involved and the brief time interval in which this phenomenon unfolds, a simplified model, where disturbances are neglected, is assumed.

The along-track motion on η is neglected, and the accelerations due to the control ($\ddot{\xi}_{ctrl}$ and $\ddot{\zeta}_{ctrl}$) and the presence of the tether tension (T_{ξ} and T_{ζ}) are considered. During the experiment, a physical tether connected to the FSS introduces the tension, with the architecture described in subsection 2.4, while the granite matching plane actuation replicates the joint effect of the HCW and the control. Also in this case, the system requires scaling, and the same approach of section subsection 4.2 is implemented.

4.3.1. LQR Controller

A Linear Quadratic Regulator (LQR) controller [27] is designed based on the linearized HCW EoM and a linearized model for the effect of tether on the dynamics. Specifically, the system is represented by the following system:

$$\dot{\underline{X}} = \underline{A} \underline{X} + \underline{B} \underline{F}_{ctrl} + \underline{L} \underline{T} \quad (18)$$

where \underline{A} and \underline{B} represents the state and input matrices obtained by writing Equation 17 in matrix form, \underline{L} is the disturbance matrix, related to the effect of the tension $\underline{T} = [T_x, T_z]^T$ on the dynamics and $\underline{F}_{ctrl} = [F_x, F_z]^T$ is the control input. In particular, to design the controller, a linear, elastic model is assumed for the tether, hence giving:

$$\underline{T} = \frac{ES}{L_t} \underline{\Delta L} \quad (19)$$

where E is the Young's modulus, S and L_t represent the cross-sectional area and the nominal length of the tether respectively, and $\underline{\Delta L} = [\Delta x, \Delta z]^T = [x - x_{eq}, z - z_{eq}]^T$ is the linear deformation. The optimal LQR controller is in the form $\underline{F}_{ctrl} = -\underline{K}(\underline{X} - \underline{X}_{eq})$, where \underline{K} is built as to minimize the following cost function:

$$J_{LQR} = \int_0^\infty (\underline{X}^T \underline{Q} \underline{X} + \underline{F}^T \underline{R} \underline{F}) dt \quad (20)$$

\underline{Q} and \underline{R} are two user defined weight functions. For the sake of completeness, two different \underline{K} matrices are defined, as there are two possible states the system can be in. When the cable is slack, the spacecraft's dynamics can be described purely by the HCW EoM, while when there is tension in the tether, the elastic force must also be added. The matrices \underline{A} , \underline{B} and \underline{L} for the two cases are obtained as shown in Equation 21 and Equation 22 respectively.

$$\underline{A} = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 3\omega^2 & 0 & 0 & 0 \\ 0 & -\omega^2 & 0 & 0 \end{bmatrix}, \quad \underline{B} = \frac{1}{m} \begin{bmatrix} 0 & 0 \\ 0 & 0 \\ 1 & 0 \\ 0 & 1 \end{bmatrix}, \quad \underline{L} = 0_{4 \times 2} \quad (21)$$

$$\underline{A} = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 3\omega^2 & 0 & 0 & 0 \\ 0 & -\omega^2 & 0 & 0 \end{bmatrix}, \quad \underline{B} = \frac{1}{m} \begin{bmatrix} 0 & 0 \\ 0 & 0 \\ 1 & 0 \\ 0 & 1 \end{bmatrix}, \quad \underline{L} = -\frac{ES}{mL_t} \begin{bmatrix} 0 & 0 \\ 0 & 0 \\ 1 & 0 \\ 0 & 1 \end{bmatrix} \quad (22)$$

5. Numerical and Experimental Results

In this section, the numerical results of the test campaign are reported and discussed.

5.1. Closed-Loop Calibration

To test the calibration loop described in subsection 4.1, the gains are defined after a few iterations of trial and error, as in Equation 23:

$$\underline{K}_P = 0.3 * \underline{I}_2 \quad \underline{K}_D = 0.6 * \underline{I}_2 \quad \underline{K}_I = 0.03 * \underline{I}_2 \quad (23)$$

where \underline{I}_2 is the identity matrix of second order. The desired final state is set to $\underline{x}_{des} = [0, 0]^T$ and $\dot{\underline{x}}_{des} = [0, 0]^T$ where the notation from subsection 4.1 is kept. The results are shown in Figure 18.

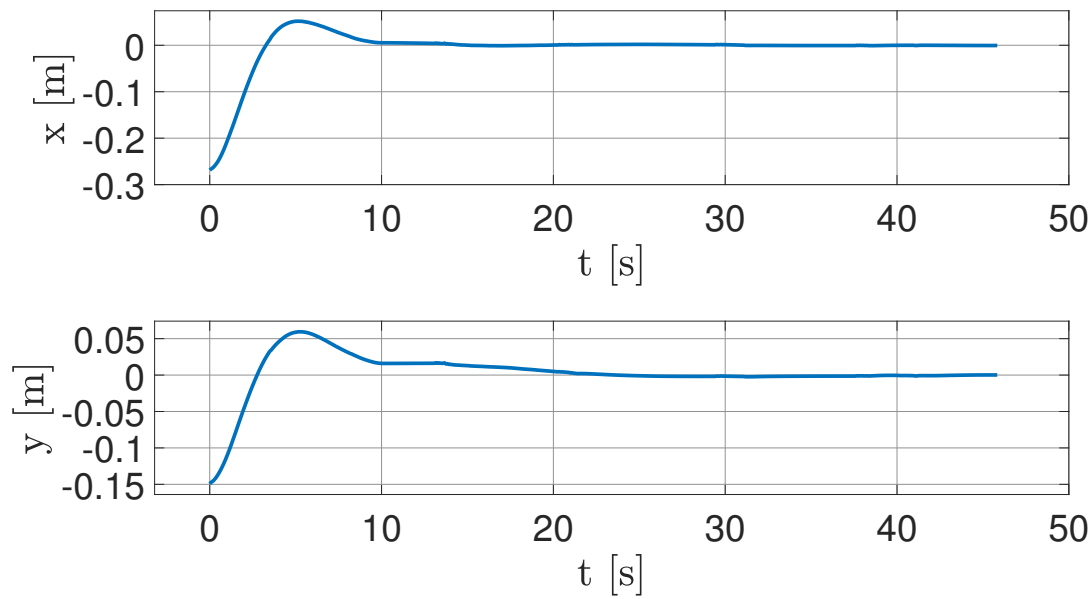


Figure 18. Time evolution of FSS coordinates over time.

The system reaches the desired state within 0.5 millimetres after approximately 20 seconds, with final actuator chattering in the order of their resolution. This behaviour suggests that it is impossible for the setup to achieve a perfect horizontal attitude, due to the limited resolution of the actuators. As a result, the FSS is expected to experience a residual acceleration of approximately $a_{res} \simeq 2.5 \cdot 10^{-3} \text{ m/s}^2$.

5.2. Emulation of Linearized Relative Dynamics

The evolution of the dynamics of two satellites starting from different radial positions, with zero velocity, is studied using the HCW equations, as explained in [subsection 4.2](#). In particular, a starting state of $\chi_0 = [\xi_0, \eta_0, \dot{\xi}_0, \dot{\eta}_0]^T = [50, 0, 0, 0]^T$ is chosen. The resulting trajectory and acceleration after propagating for 20% of the orbital period are shown in [Figure 19](#) and [Figure 20](#).

The maximum displacement is $[x_{max}, y_{max}]^T = [153.648, -91.674]^T \text{ m}$, while the acceleration $a = [1.842, 3.551] \cdot 10^{-4} \text{ m/s}^2$. Consequentially, according to the criteria expressed in [subsection 4.2](#), the scaling factors are selected as:

$$\begin{cases} \lambda_L = 700 \\ \lambda_t = 500 \end{cases} \quad (24)$$

The results of emulated trajectories are shown in [Figure 21](#) and [Figure 22](#).

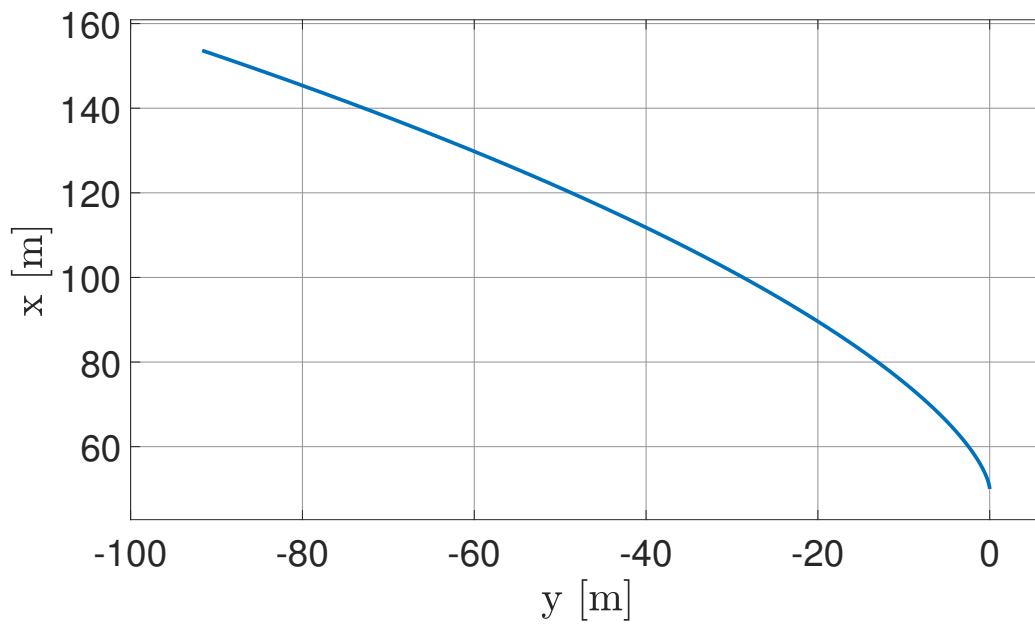


Figure 19. Real-case trajectory propagated in LVLH CCS.

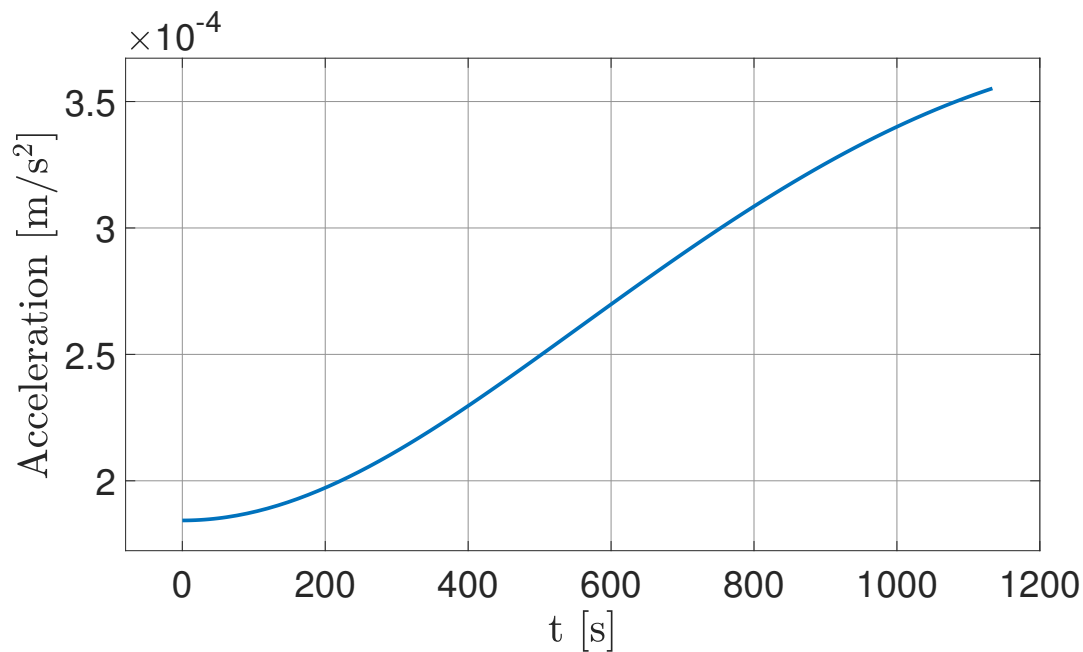


Figure 20. Real-case acceleration.

Figure 21 shows that the experimental results match the numerically propagated trajectory with a reasonable error. By increasing the number of the test and considering the average, spurious errors are mitigated. They are caused by various factors, such as the residual friction between the FSS and the matching plane and minor surface defects.

Figure 22 shows the relative error between the numerically propagated trajectory and the experimental results. This error is computed as the norm of the difference between the vectors' positions at the same instant of time. These results highlight that the testbed can emulate relative dynamics with a maximum error of $\varepsilon_{max} = 7.85$ mm; it corresponds to an error of 5.5 m when scaled back up to the real case scenario, which is in line with the hypothesis of a residual acceleration of approximately 2.25 mm/s².

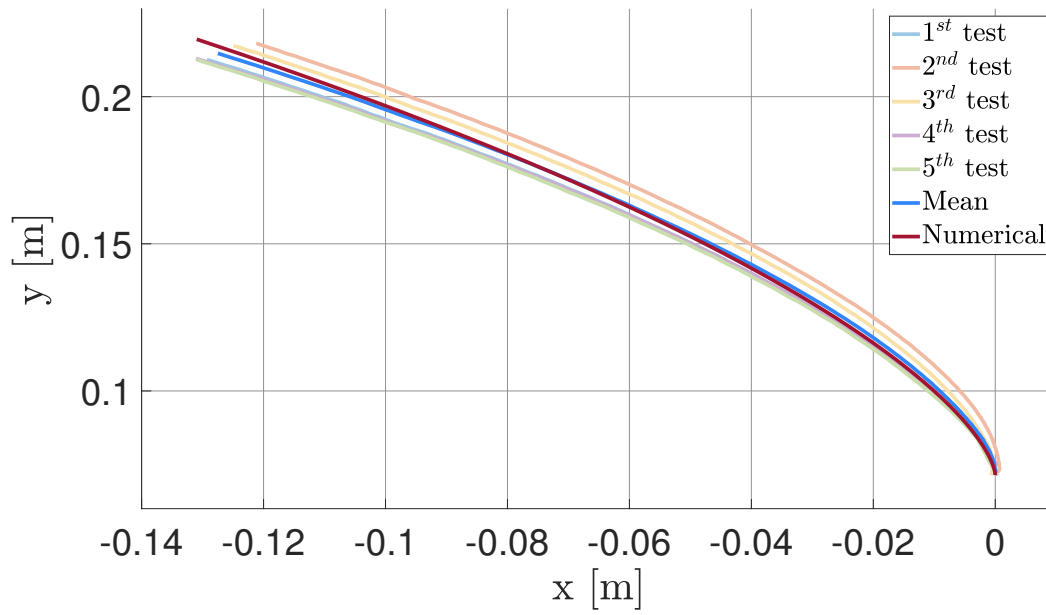


Figure 21. Comparison between HCW analytical solution and experimental trajectories.

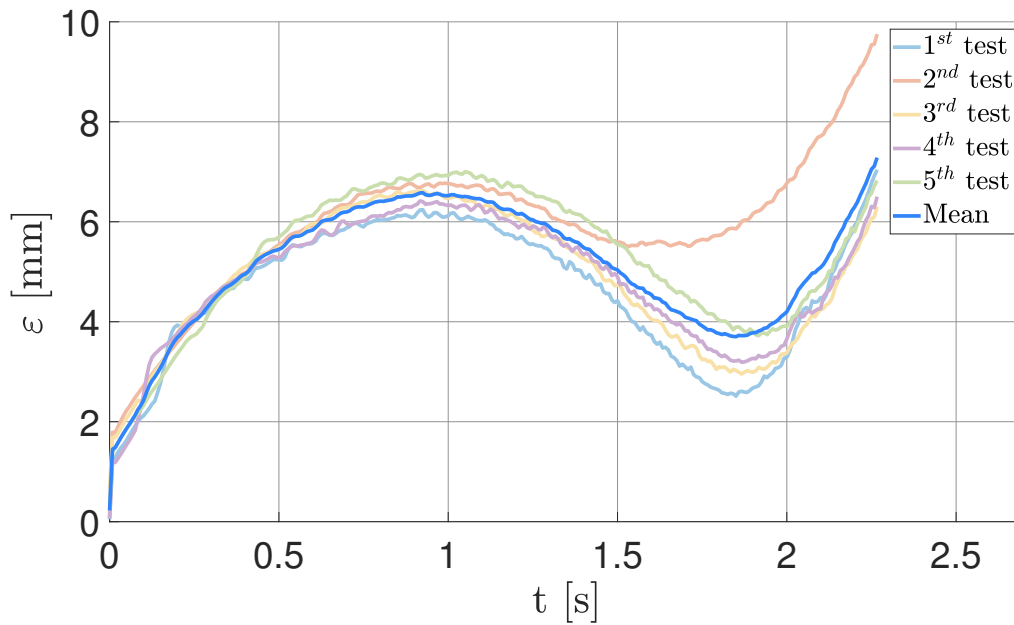


Figure 22. Absolute error between the numerical and experimental trajectories.

5.3. Fully-Deployed Tether Dynamics

The case study for the third test regards a 100 m long tether, starting from an off-nominal slack condition with the goal of testing the ability of the LQR controller (presented in subsection 4.3) to lead the system back to a cross-track configuration, and stabilizing it. Moreover, the load cell included in the setup is used to measure the value of tension during the whole experiment.

A scaling factor $\lambda_L \geq 50$ is selected as the ratio of half of the nominal length of the real case tether and the 1 m cable used in the experiment. As for the λ_t , again, a constraint regarding the acceleration range is enforced. Particularly, in this case, only a limit on maximum acceleration is enforced, since the system is now controlled in closed loop and, therefore, it is able to keep the FSS in the desired position, with a final continuous chattering of the actuators, in a similar way to subsection 5.1. Specifically, an

acceleration profile as that shown in Figure 23 is found, with a maximum value of $a_{max} = 0.1416 \text{ m/s}^2$, which leads to $\lambda_t \leq 20.3$.

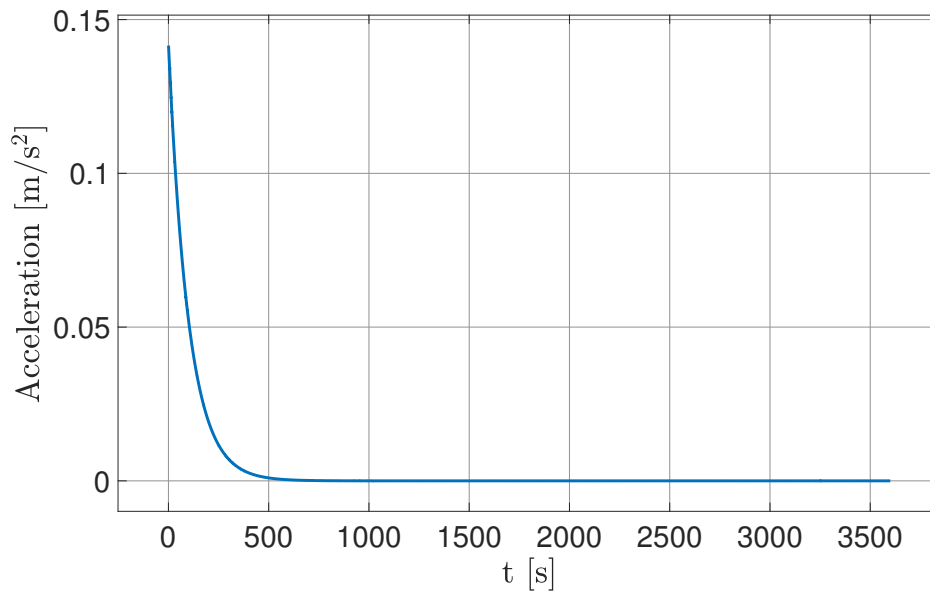


Figure 23. Real-case acceleration for the cross-track TSS

According to these constraints the set of scaling factors is chosen as in Equation 25.

$$\begin{cases} \lambda_L = 50 \\ \lambda_t = 20 \end{cases} \quad (25)$$

The FSS displacements, shown in Figure 24, stabilize at $x = 0 \text{ mm}$ and $y = 1270 \text{ mm}$ with respect to the origin of the CCS placed on the load cell, observing that the steady-state in position is achieved after approximately 20 seconds of simulation.

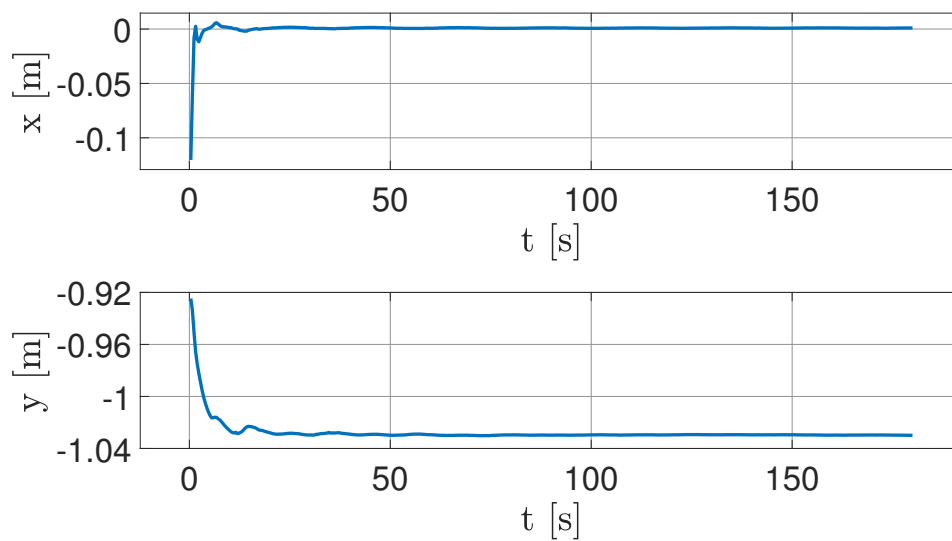


Figure 24. Time evolution of FSS coordinates under tether tension.

In Figure 25, the measured value of tension stabilizes at approximately 1 mN after approximately 80 seconds of the experiment. During an initial phase of uncontrolled rotation of the FSS along its vertical axis, the tether stretches more, increasing the tension. The presence of the tether help stabilize

the FSS, gradually damping the amplitude of oscillation closer to 0 degrees. However, since no attitude control is implemented, the rotation of the FSS around its vertical axis does not stabilize to a constant value. Therefore, the tension of the tether is close to its steady state around 1 mN but experiences minor oscillations caused by the rotation.

As a result, the tether contributes to the stability of the system; however, to achieve a more refined attitude, an active control strategy is required.

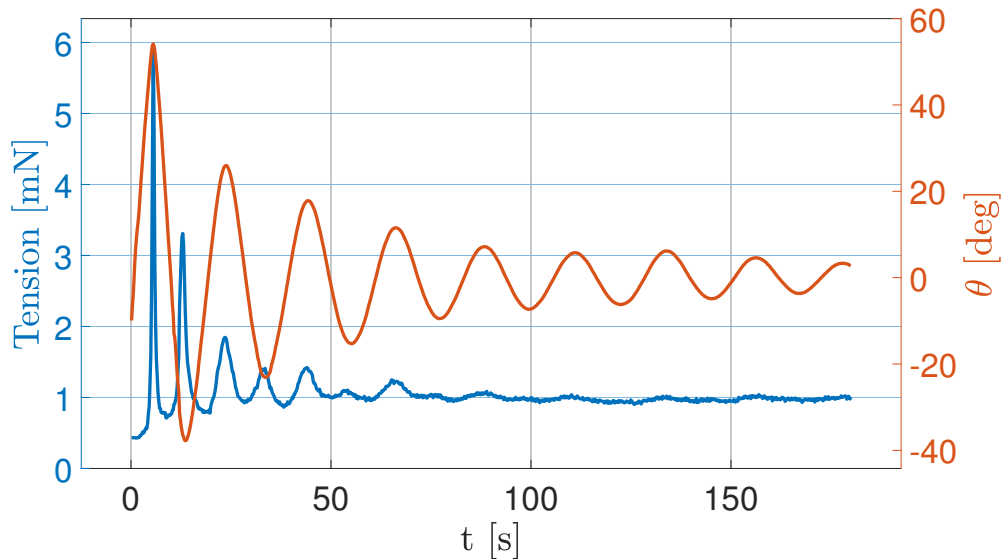


Figure 25. Tension acting on the FSS (blue) and attitude angle θ of the FSS (orange)

6. Conclusions

This paper presents a novel hardware-in-the-loop testbed capable of emulating the relative dynamics between two orbiting spacecrafts, along with an innovative experimental setup for tethered satellite systems. The design is detailed, including the derivation of scaling laws based on the Buckingham theorem and the principle of similarity. Experimental results demonstrate the potential of this approach to replicate orbital dynamics within a laboratory environment through a Hardware-In-the-Loop setup. Experiments are carried out to evaluate the capabilities of the testbed to replicate the dynamics described by the HCW equations of motion, as well as the behaviour of a tethered satellite system in a cross-track configuration. Specifically, results show that the FSS follows a trajectory closely matching the one obtained from numerical propagation of the Hill-Clohessy-Wilshire equation. Furthermore, simulations of tether-assisted stabilization indicate that the combined action of the tether and an LQR controller is effective in stabilizing the tethered spacecraft.

7. Patents

The main principle underlying the tip-tilt flat-table facility described in this paper is based on two existing patent inventions:

- 1) "Method and Apparatus for the Stabilization of a Satellite Formation", [31].

Description: "The invention relates to a method for stabilization of the orbital motion set-up of a group of satellites bound to a tether. To this end, the tension of the tether that connects two or more satellites is controlled in order to keep them in the desired flight formation. This effect is obtained by perturbing a satellite to generate a tension component in the tether, which results in a recoil force that brings the system back into a condition of equilibrium. According to a preferred embodiment, the tension T of the tether is obtained by adjusting the disposition of one or more aerodynamic surfaces provided on the satellites at the ends of the tether".

- 2) *"Dynamically tilting flat table to impart a time-varying gravity-induced acceleration on a floating spacecraft simulator"*, [32].

Description: "Disclosed is a planar test bed comprising a planar surface and further comprising mechanical couplings in mechanical communication with the planar table and the supporting legs. The mechanical couplings are translatable to provide three degrees of freedom for orientation of the planar surface. A processor receives position and velocity information describing an object on the planar surface, and calculates a relative acceleration typically using a function $a_R=f(t,x_R,v_R,\mu t)$. The processor communicates with the mechanical couplings to establish an orientation where a local gravity vector projects onto the planar surface and generates acceleration with magnitude and direction substantially equal to the desired acceleration a_R . The operations occur in cyclic fashion so the desired accelerations and planar orientations are updated as an object transits over the planar surface".

Author Contributions: Conceptualization, A.P., M.C., M.L.V., C.B., G.G., C.L.M., R.A., S.A., M.R.; methodology, A.P., M.C., M.L.V., C.B.; software, A.P., M.C., M.L.V., C.B.; validation, A.P., M.C., M.L.V., G.G., C.B.; formal analysis, A.P., M.C., M.L.V., C.B.; investigation, A.P., M.C., M.L.V., C.B.; writing - original draft preparation, A.P., M.C., M.L.V., C.B.; writing - review and editing, G.G., C.L.M., R.A., S.A., M.R.; supervision, G.G., M.R.; project administration, G.G., M.R.; funding acquisition, C.L.M., R.A., S.A., M.R.; All authors have read and agreed to the published version of the manuscript.

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Conflicts of Interest: The authors declare no conflicts of interest.

Acronyms

The following acronyms are used in this manuscript:

CCS	Cartesian Coordinate System
DCM	Direction Cosine Matrix
EoM	Equation of Motion
FP	Floating Platform
FSS	Floating Spacecraft Simulator
GN&C	Guidance, Navigation and Control
HCW	Hill-Clohessy-Wilshire
HIL	Hardware-In-the-Loop
MC	Master Console
MCS	Motion Capture System
LQR	Linear Quadratic Regulator
LVLH	Local-Vertical Local-Horizontal
PID	Proportional Integrative Derivative
TSS	Tethered Satellite System
TTS	Tether Testing System
VLEO	Very Low Earth Orbit

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