
Are Some Voltage State Orbits, Computed in a Discrete-Time Hopfield Neural Network, Which Correspond to Bifurcation Values, Fractals? Exists These Orbits in Reality or They Exists Just in Theory i.e.Has These Orbits a Real Couterpart?

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Article

Are Some Voltage State Orbits, Computed in a Discrete-Time Hopfield Neural Network, Which Correspond to Bifurcation Values, Fractals? Exists These Orbits in Reality or They Exists Just in Theory i.e.Has These Orbits a Real Couterpart?

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Abstract: In case of a discrete-time Hopfield neural network of two neurons with two delays and no self-connections at 23 bifurcation values 23 voltage trajectories appear. Among the 23 voltage trajectories 14 voltage trajectory are not what we call orbits in classic sense. The geometrical aspect of these trajectories suggest that they are fractals. Our conjecture is that the 14 trajectories in discussion are fractals having real couterpart. In case of a discrete-time Hopfield neural network of five neuron with delay and ring architecture at 9 bifurcation values 9 voltage trajectories appear. Among the 9 voltage trajectories we find 5 voltage trajectory which are not what we call orbits in classic sense. The geometrical aspect of these trajectories suggest that they are fractals. Our conjecture is that the 5 trajectories in discussion are fractals having real counterparts.. In case of a discrete-time Hopfield neural network of two neurons with a single delay and self- connections the computed trajectories are what we call orbits in classic sense. In case of a discrete-time Hopfield neural network of two neurons with two delays and self- connections the computed trajectories are what we call orbits in classic sense.

Keywords: fractals; discrete- time Hopfield neural networks; two neurons with delay and no self-connection; five neurons with delay and ring architecture

1. Introduction

There is some disagreement among mathematician about how the concept of fractal should be formally defined. Mandelbrot himself summarized it is as "beautiful damn hard increasingly useful. That's fractals'. [1] More formally, in 1982Mandelbrot defined fractal as follows: "A fractal is by definition a set for which the Hausdorff–Besicovitch dimension strictly exceeds the topological dimension." [2]

Later, seeing this as too restrictive, he simplified and expanded the definition to this: "A fractal is a rough or fragmented geometric shape that can be split into parts, each of which is (at least approximately) a reduced-size copy of the whole." [1] Still later, Mandelbrot proposed, "to use fractal without a pedantic definition, to use fractal dimension as a generic term applicable to all the variants". [3] The consensus among mathematicians is that theoretical fractals are infinitely self-similar iterated and detailed mathematical constructs, of which many examples have been formulated and studied [1–3]. Fractals are not limited to geometric patterns, but can also describe processes in time [4–9]. Fractal patterns with various degrees of self-similarity have been rendered or studied in visual physical, and aural media and found in nature, [10–13] technology, [14–17] art, [18,19] and architecture. [20] In this context fractal geometry lies within the mathematical branch of measure theory [3] presenting different measures and algorithms for the computation of different types of fractal dimensions of geometrical patterns; box counting dimension, correlation dimension

,generalized or Renyi dimensions, Higuchi dimension, Lyapunov dimension, Multifractal dimensions, Hausdorff dimension, Packing dimension, Assouad dimension, Local connected dimension, Degree dimension, Parabolic Hausdorff dimension. These fractal dimensions strictly exceed the topological dimension and although for compact sets with exact affine self-similarity all these dimensions coincide; in general, they are not equivalent. When fractal describe processes in time, then for the analysis of statistical data, fractal analysis use also tools for: investigate statistical self-similarity, find probability density function (PDF), fractal scaling in time, time series analysis, fractal kinetic description ,power law and fractal statistics. [4–9,16,21–33]

Beside the mathematical description of the voltage dynamics of neural cells and systems using fractals, there exist a large amounts of papers, which use classical theory of electrical circuits or fractional order derivative circuits for describe the neural cells dynamics.

Papers [34–41] use classical theory of electricity (integer order differential equations) for the description of the ion transport through the nervous cell membrane. Papers [42,43] use time fractional order circuits [44] for the description of the ion transport through the nervous cell membrane. In the papers [45,46] it was shown that: Mathematical descriptions of the ion transport, across passive or active biological neuron membrane, voltage propagation along neuron axons and dendrites having passive or active membrane and ion transport in biological neuron networks, using classic Caputo or Riemann-Liouville fractional order derivatives, is nonobjective. The no objectivity is originated in the incompatibility of the classic Caputo and Riemann-Liouville fractional order derivatives with the understanding of the time evolution, used in case of the mathematical description of real word phenomena.

Hopfield paper [47] claims to be mathematical descriptions of electrical phenomena appearing in nervous system. Later in [48–50] this model was extended to optimization and cryptography. In papers [51–71], theoretical studies of the Hopfield neural system are carried out in the continuous time and discrete time versions. Papers reveal configuration of steady states , local exponential stability of steady states in Lyapunov sense, regions of attraction of exponentia stable steady states, bifurcation properties etc. Several example of low dimensional (two, fife neurons) neural network are presented which exhibit extremely interesting orbits and bifurcation diagrams.

Although for some values of the bifurcation parameter the orbits completely disintegrate and no longer resemble what we call orbit in classical sense, their fractal quality is not analyzed and highlighted. Also, no reference is made to the question: whether or not these completely dismembered orbits have a real counterpart? In the case of strange bifurcation diagrams, their approach with fractals, as far as we know, is not done.

In this paper, we will present in case of four neural networks, examples of strange orbits and bifurcation diagrams, and try to bring arguments to support that they are fractals. We also hypothesize that they do not have a real counterpart. They exists just in theory.

2. Case of a Discrete-Time Hopfield Neural Network of Two Neurons with Two Delays and No Self- Connections

In [67] an extended analysis of the chaotic dynamics of the neural network defined by

$$x_{p+1} = a_1 \cdot x_p + T_{12} \cdot g_2(y_{p-k_2}) , y_{p+1} = a_2 \cdot y_p + T_{21} \cdot g_1(x_{p-k_1}) \quad (2.1)$$

was undertaken. The external inputs are equal to zero and the system parameters are regarded as numerical expresion of the state of network..The network state can change due to a disease for example. Therefore the network voltage dynamics change too. In [67] this kind of change is analyzed with respect to the product of the interconnection coefficients as aggregated characteristic parameter for the system. The main result obtained in [67], is that, if the magnitudes of the interconnection coefficients are large enough, then the neural network exhibits Marotto's chaotic behavior.

For illustrate what this means the following numerical example was considered:

$$a_1 = 1/4, a_2 = 3/4, T_{12} = 1, g_1(t) = \tanh(t) \text{ and } g_2(t) = \sin(t), b = T_{21}k_1 = 1 \text{ and } k_2 = 2.$$

For different values of b , trajectory of the neurons voltages were computed and represented. We will present some of these trajectories.

For $b = -0.35$, the null solution is asymptotically stable, and the trajectories converges to the origin.

For $b = -0.36$, an asymptotically stable cycle (1-torus, drift ring) is present, and the trajectories converges to this cycle. At $b = b_1 = -0.35635$ a supercritical Neimark–Sacker bifurcation takes place.

For $b = -1.26$, there is only one asymptotically stable limit cycle (1-torus), symmetrical to the origin.

For $b = -1.27$, there are two stable limit cycles, not symmetrical to the origin. For each plot, the first 10^6 iterations have been dropped, and the next 10^4 iterations have been plotted. At around $b = -1.265$, a bifurcation phenomenon takes place, which determines the appearance of two stable limit cycles (1-tori) close to the origin.

For the localization of bifurcations values the Largest Lyapunov Characteristic Exponent (LCE) and the Lyapunov spectrum (LCEs) for system (2.1) was used. For the computation of the Lyapunov spectrum, for each b value (step size 0.01 for b), the initial conditions were reset and 10^5 time-steps were iterated before calculating the LCEs (which were computed over the next 10^5 time steps). The Lyapunov spectrum was computed using the Householder QR based (HQRB) method. The obtained LCE and LCEs are represented in Figure 2.5 and Figure 2.6. respectively.

The trajectories presented in Figures 2.1–2.4 are what we call orbits in classical sense. However, they become more and more complex for various values of $b \in (-1.5, -1.4)$. This can be seen on the following figures.

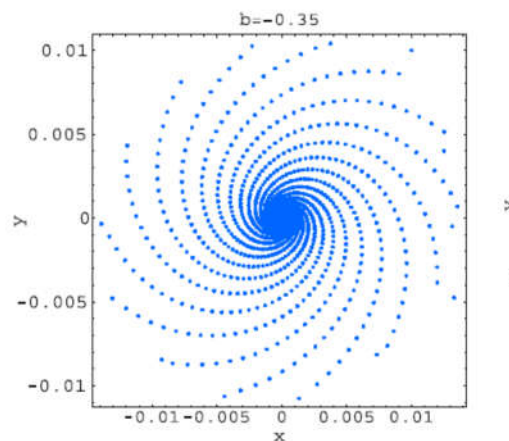


Figure 2.1. Trajectory for $b = -0.35$.

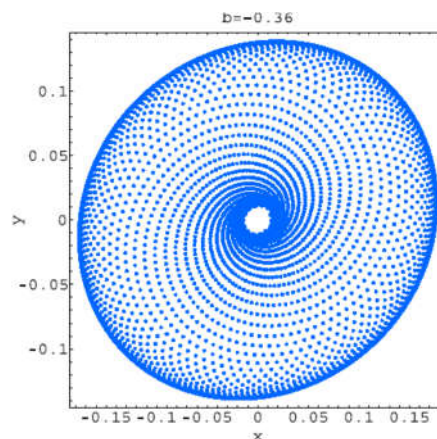


Figure 2.2. Trajectory for $b = -0.36$.

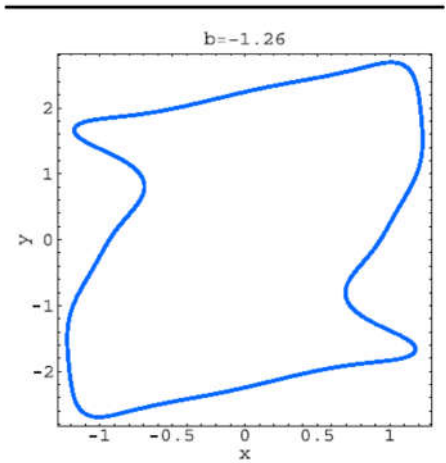


Figure 2.3. Trajectory for $b = -1.26$.

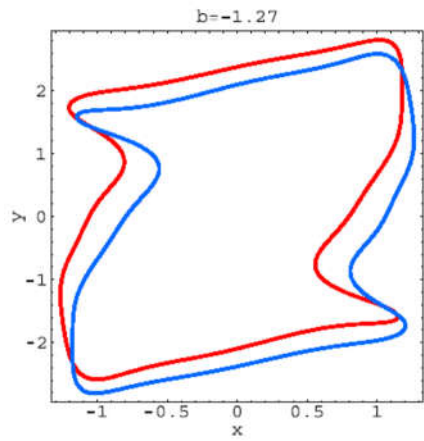


Figure 2.4. Trajectory for $b = -1.27$.

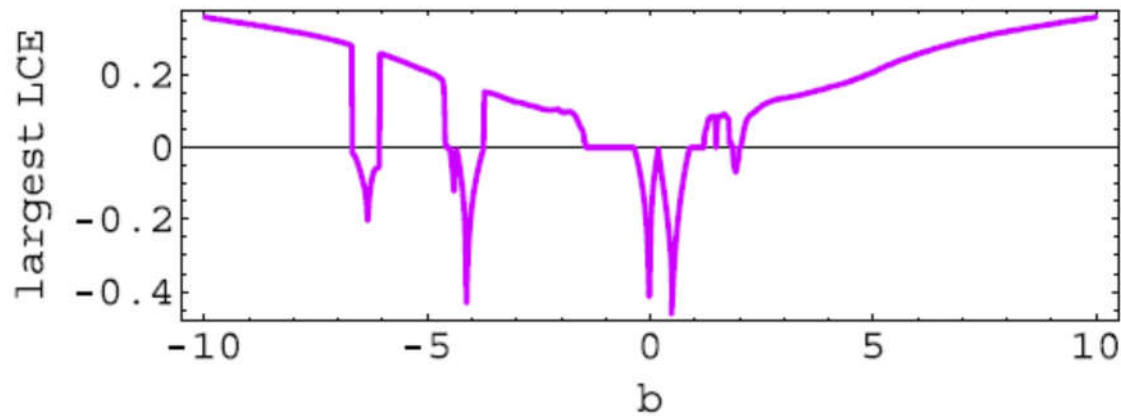


Figure 2.5. Largest Lyapunov Characteristic Exponent.

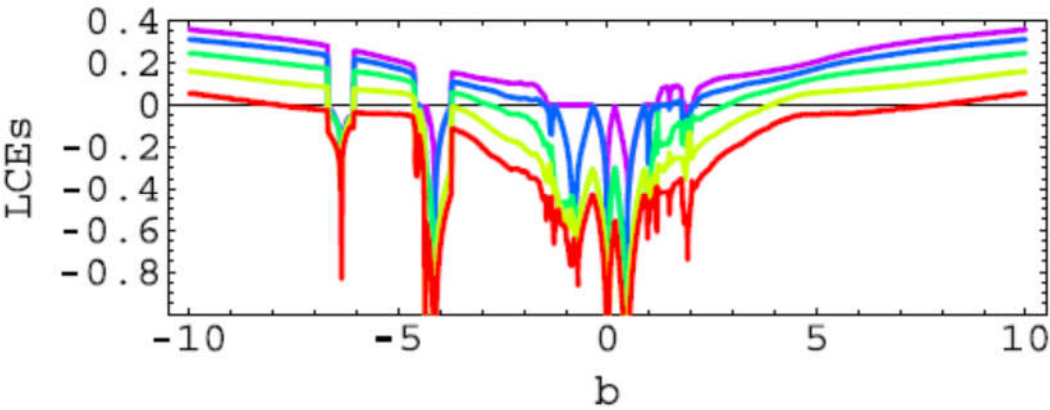


Figure 2.6. Largest Lyapunov Characteristic Spectrum.

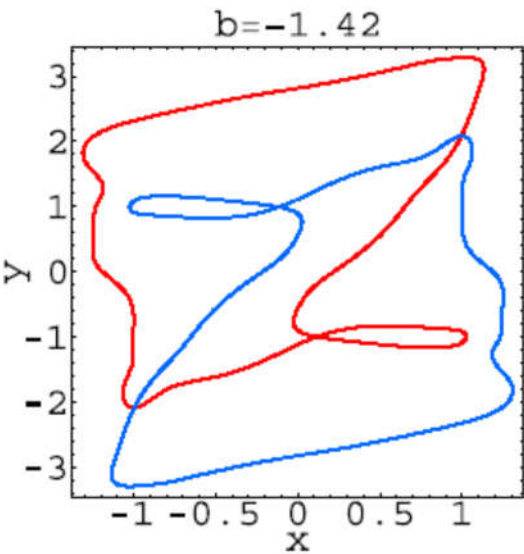


Figure 2.7. Trajectory for $b=-1.42$.

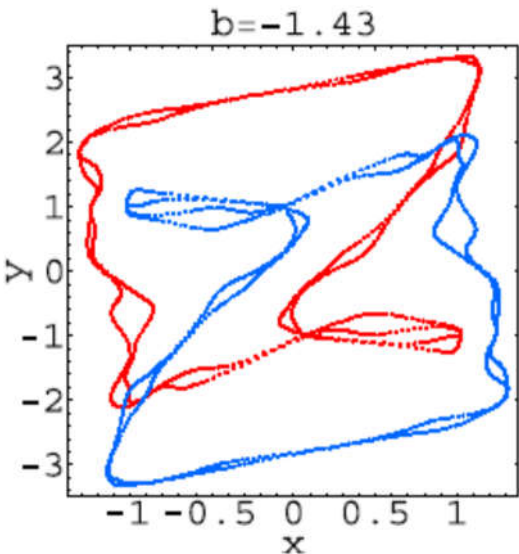


Figure 2.8. Trajectory for $b=-1.43$.

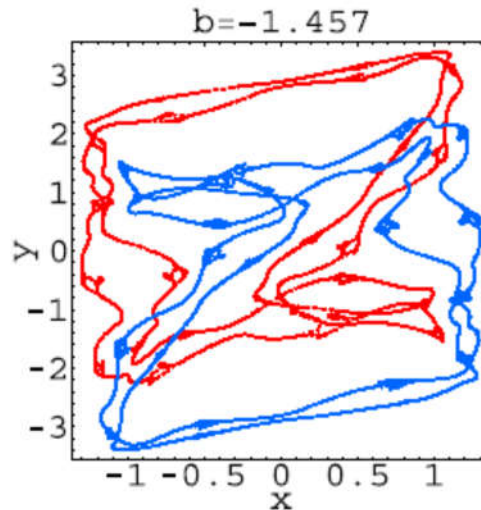


Figure 2.9. Trajectory for $b=-1.457$.

Here for each plot, the first 10^6 iterations of system (2.1) have been dropped, and the next 10^4 iterations have been plotted. These trajectories correspond to the bifurcation values: $b=-1.42$: 1-tori; $b=-1.43$: 2-tori; $b=-1.457$: 2-tori declining into strange attractors; $b=-1.46$: strange attractors; $b=-1.47$: two stable period-52 orbits; $b=-1.473$: strange attractors; $b=-1.48$: strange attractors; $b=-1.485$: two stable period-39 orbits; $b=-1.49$: strange attractors.

It seems that from bifurcation point of view everything is clear but from the point of view of the trajectory geometry is not at all clear. That is because it can be seen that “Trajectories” presented in Figures 2.10–2.15. are not at all what we call orbits in classic sense. Their geometry moreover is fractal. Are “Trajectories” presented in Figures 2.10–2.15 objects of Fractal Geometry? Have they real counterpart or they exist just in theory? There are similar behavior mentioned in [10–12,16,21–25,27–33] or elsewhere in specialized literature? These are our main questions.

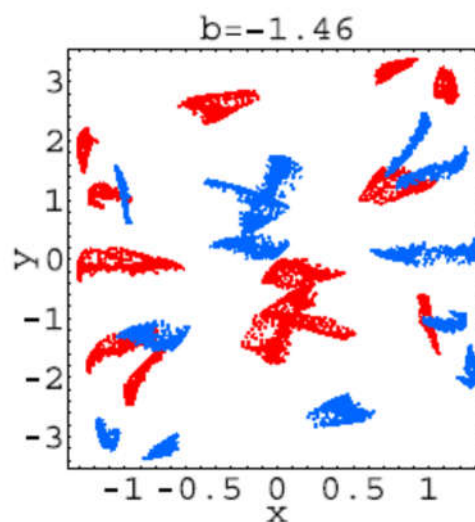


Figure 2.10. Trajectory for $b=-1.46$.

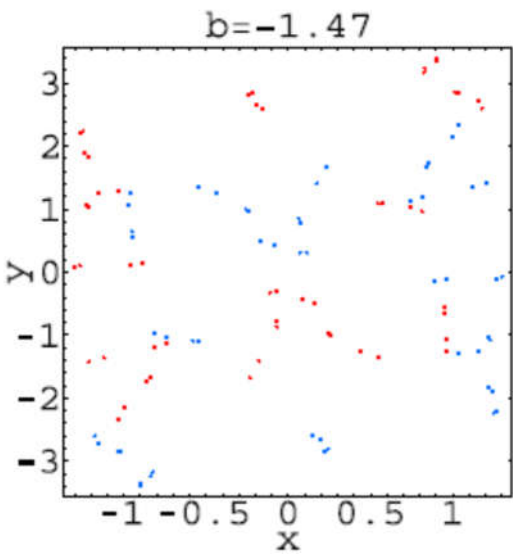


Figure 2.11. Trajectory for $b=-1.47$.

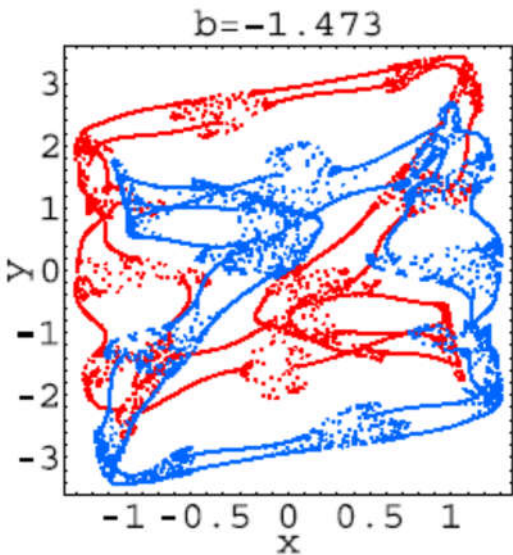


Figure 2.12. Trajectory for $b=-1.473$.

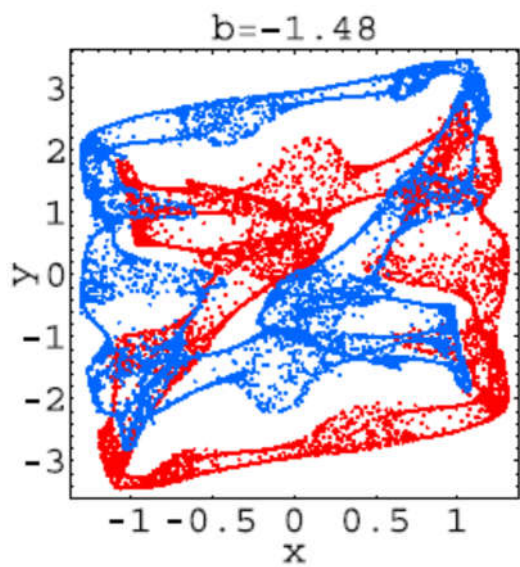


Figure 2.13. Trajectory for $b = -1.48$.

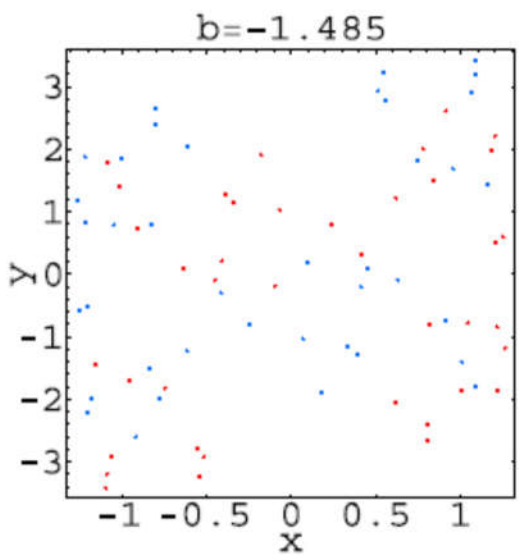


Figure 2.14. Trajectory for $b = -1.485$.

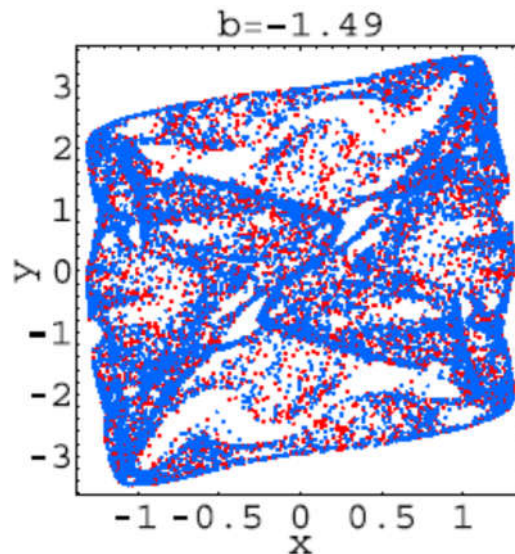


Figure 2.15. Trajectory for $b=-1.49$.

In our opinion Figures 2.10–2.15. are fractals. If not, what are these geometrical figures?

For various values of $b \in (-10, -2)$ the trajectories were also computed using the same procedure i.e. the first 10^6 Iterations of system (2.1) have been dropped, and the next 10^4 iterations have been plotted.

These trajectories corresponds to the bifurcations values: $b = -2$: hyperchaos ($LE2 > 0$, $LE3 < 0$); $b = -3$: hyperchaos ($LE3 > 0$, $LE4 < 0$); $b = -4$: one stable period-2 orbit ($LE1 < 0$); $b = -4.6$: strange attractor developing from the period-2 solution; $b = -5$: hyperchaos ($LE4 > 0$, $LE5 < 0$); $b = -6.5$: two stable period-3 orbits ($LE1 < 0$); $b = -7$: hyperchaos ($LE4 > 0$, $LE5 < 0$); $b = -9$: hyperchaos ($LE5 > 0$).

Our main questions are the same: are “Trajectories” presented in Figures 2.16–2.23 objects of Fractal Geometry? Have they real counterpart or they exist just in theory? If not, what are these geometrical figures? There are similar behavior mentioned in [10–12, 16, 21–25, 27–33] or elsewhere in specialized literature? Our conjecture is that “Trajectories” presented in Figures 2.10–2.23. are fractals and has real counterpart.

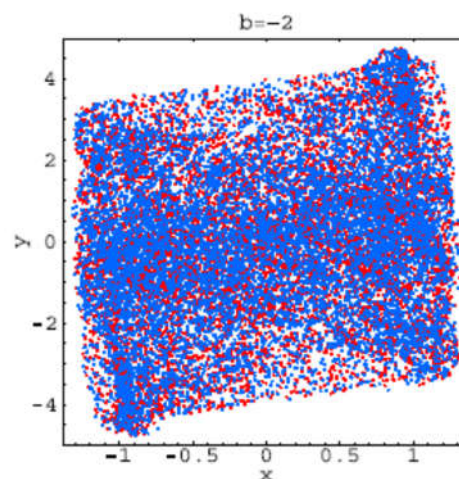


Figure 2.16. Trajectory for $b=-2$.

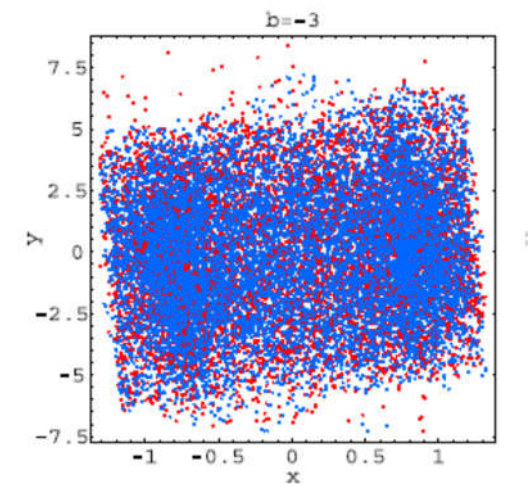


Figure 2.17. Trajectory for b=-3.

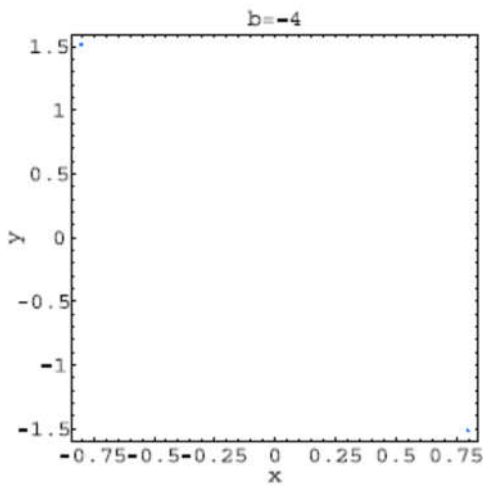


Figure 2.18. Trajectory for b=-4.

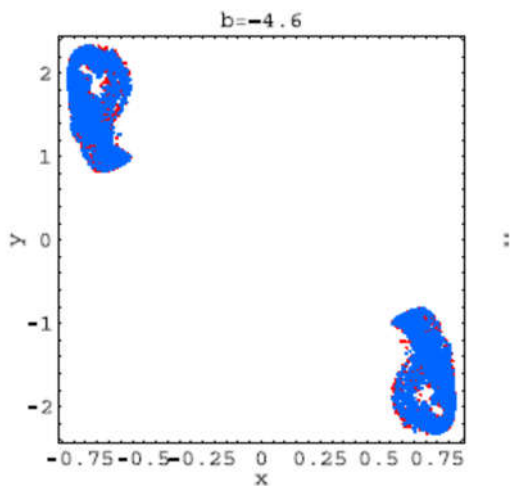


Figure 2.19. Trajectory for b=-4.6.

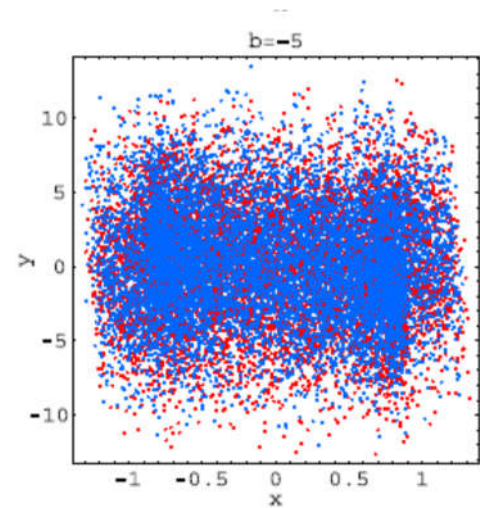


Figure 2.20. Trajectory for b=-5.

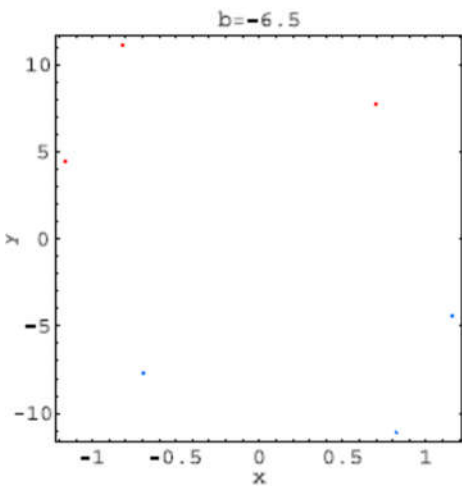


Figure 2.21. Trajectory for b=-6.5.

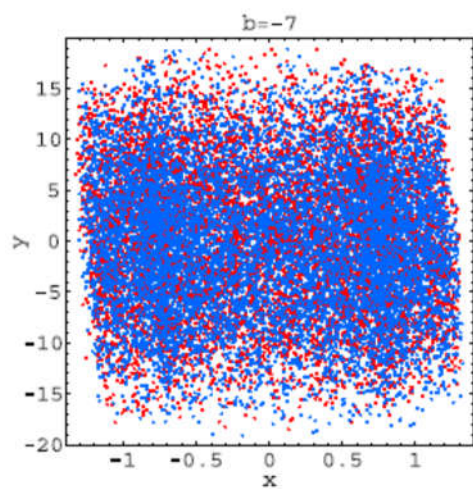


Figure 2.22. Trajectory for b=-7.

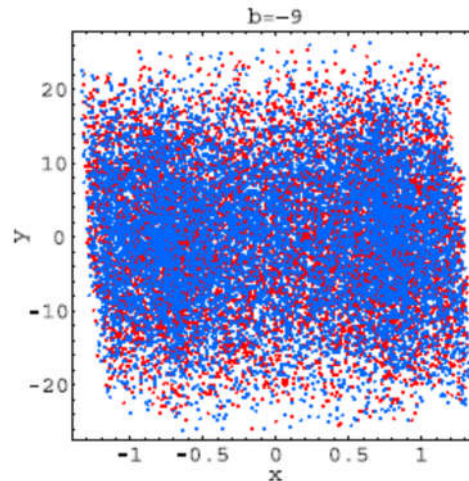


Figure 2.23. Trajectory for $b=-9$.

3. Case of a Discrete-Time Hopfield Neural Network of Two Neurons with a Single Delay and Self-Connections

In [68] an analysis of the dynamics of the neural network defined by

$$\begin{cases} x_{n+1} = ax_n + T_{11}g_1(x_{n-k}) + T_{12}g_2(y_{n-k}) \\ y_{n+1} = ay_n + T_{21}g_1(x_{n-k}) + T_{22}g_2(y_{n-k}) \end{cases} \quad \forall n \geq k \quad (3.1)$$

was undertaken. The external inputs are equal to zero and the system parameters are regarded as numerical expression of the state of network. The network state can change due to a disease for example.

Therefore the network voltage dynamics change too. In [68] this kind of change is analyzed. The conclusion in [68], is: The bifurcation analysis of two-dimensional discrete-time Hopfield neural networks with a single delay reveals the existence of Neimark–Sacker, fold and some codimension 2 bifurcations for certain values of the bifurcation parameters that have been chosen. For illustrate what this means the following numerical example was considered:

$$\begin{cases} x_{n+1} = \frac{1}{2}x_n + \tanh(x_{n-1}) - \sin(y_{n-1}) \\ y_{n+1} = \frac{1}{2}y_n + \alpha \tanh(x_{n-1}) - \sin(y_{n-1}) \end{cases} \quad \forall n \geq 1 \quad (3.2)$$

The origin is asymptotically stable if and only if $\alpha \in (0.75, 1.40693)$. At 1.40693 a supercritical Neimark–Sacker bifurcation occurs.

For $\alpha = 1.4$, the null solution of (3.2) is asymptotically stable. Figure 3.1 illustrate that the trajectories of (3.2) converges to the null solution. These trajectories are what we call orbits in classic sense.

For $\alpha = 1.5$, the null solution is unstable and an asymptotically stable closed invariant curve is present.

The trajectory of (3.2) converges to the asymptotically stable invariant curve as is shown in Figure 3.2. Figure 3.3.

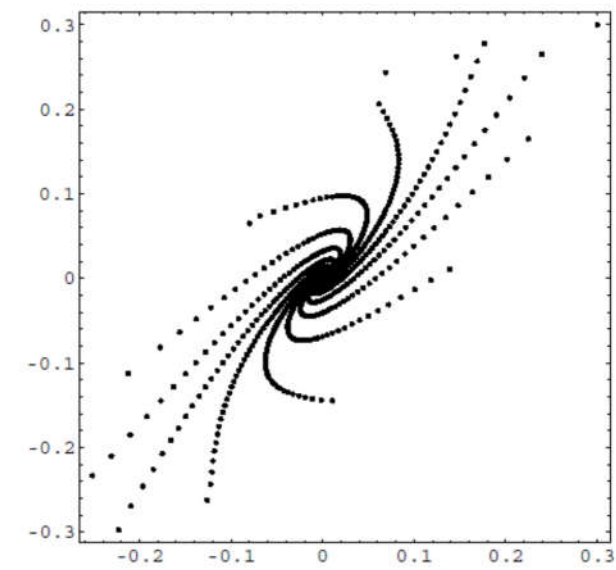


Figure 3.1. Trajectory for $\alpha = 1.4$, .

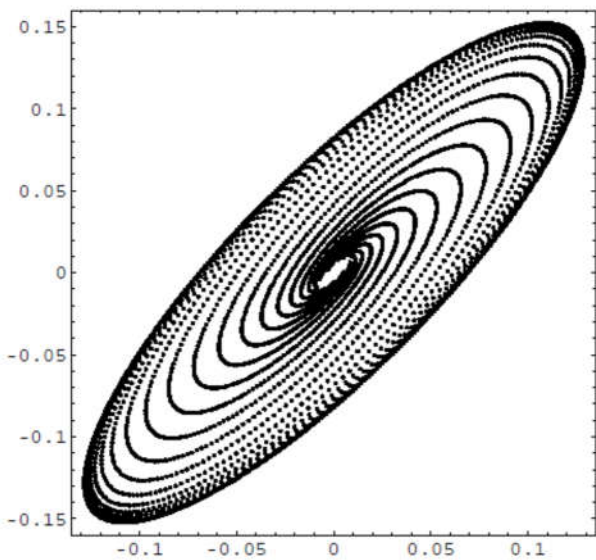


Figure 3.2. Trajectory for $\alpha = 1.5$, .

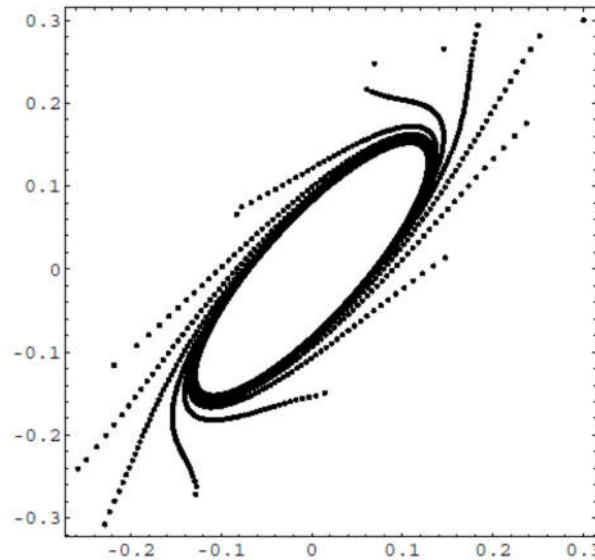


Figure 3.3. Trajectory for $\alpha = 1.5$.

The trajectories presented in Figure 3.2, Figure 3.3 are what we call orbits in classical sense.

This summary analysis does not reveal trajectories that are no longer classic curves. Even the asymptotically stable closed invariant curve, corresponding to the bifurcation value $\alpha = 1.5$, is a classical curve.

4. Case of a Discrete-Time Hopfield Neural Network of Two Neurons with a Two Delays and Self- Connections

In [69] an analysis of the dynamics of the neural network defined by

$$\begin{cases} x_{n+1} = ax_n + T_{11}g_1(x_{n-k_1}) + T_{12}g_2(y_{n-k_2}) \\ y_{n+1} = ay_n + T_{21}g_1(x_{n-k_1}) + T_{22}g_2(y_{n-k_2}) \end{cases} \quad \forall n \geq \max(k_1, k_2). \quad (4.1)$$

is undertaken. The external inputs are equal to zero and the system parameters are regarded as numerical expression of the state of network..The network state can change due to a disease for example.

Therefore the network voltage dynamics change too. In [69] this kind of change is analyzed .The conclusion in [69], is: "The results presented in this paper complete the bifurcation results obtained for discrete-time Hopfield neural networks with a single delay and self-connections presented in [68] and for networks with two delays and no self-connections presented in [66], with new results concerning the case of neural networks with two different delays and self-connections, revealing some resemblances (the existence of Neimark-Sacker, Fold, resonance 1:1, double Neimark-Sacker bifurcations) and some differences (the possible existence of Flip and Flip-Neimark-Sacker bifurcations), as well".

For illustrate what this means the following numerical example was considered:

$$\begin{cases} x_{n+1} = ax_n + \beta \tanh(x_{n-k_1}) - \sin(y_{n-k_2}) \\ y_{n+1} = ay_n + \gamma \tanh(x_{n-k_1}) + \beta \sin(y_{n-k_2}) \end{cases} \quad \forall n \geq \max(k_1, k_2). \quad (4.2)$$

with; $a = 0.5$, $k_1 = 4$, $k_2 = 10$, For $\gamma = 0.19$ the null solution of (4.2) is asymptotically stable and trajectories converges to the null solution.

Figure 4.1 illustrate that the trajectories of (4.2) converges to the null solution and the fact that these trajectories are what we call orbits in classic sense.

For $\gamma = 0.2$ the null solution is unstable and an asymptotically stable closed invariant curve is present.

The trajectories of (4.2) converges to the asymptotically stable invariant curve as is shown in Figure 4.2., Figure 4.3.

The trajectories presented in Figure 4.2, Figure 4.3 are what we call orbits in classical sense.

For $\beta = -0.82$ and $\gamma = -0.47$ the null solution of (4.2) is unstable and an asymptotically stable cycle of period 2 is present. The trajectory of (4.2) converges to the asymptotically stable cycle i.e. Flip (period-doubling) bifurcation occurs as is shown in Figure 4.4.. The trajectory presented in Figure 4.4, are what we call orbit in classical sense.

This summary analysis does not reveal trajectories that are no longer classic curves.

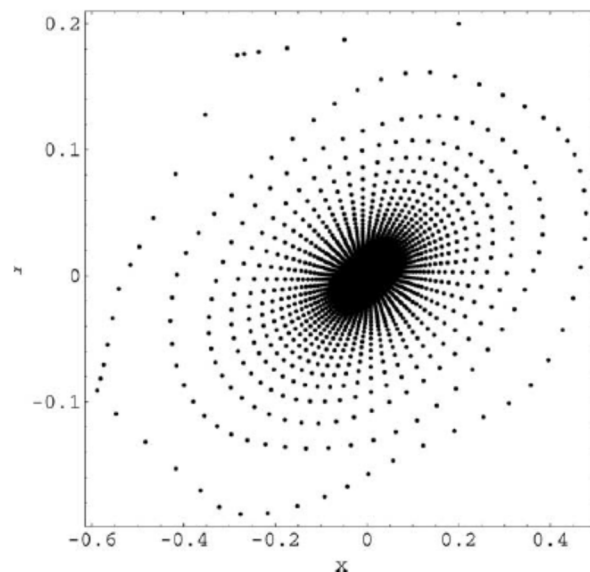


Figure 4.1. Trajectory for $\gamma = 0.19$.

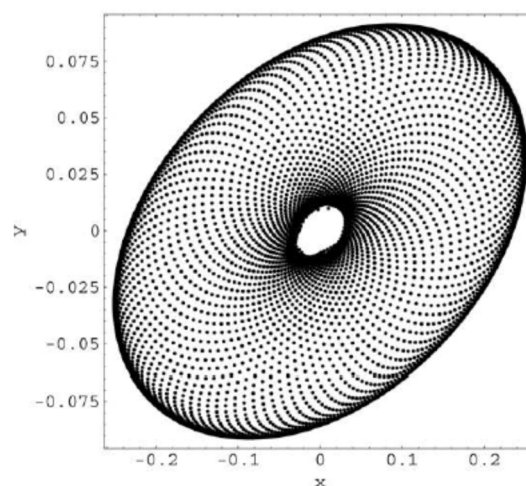


Figure 4.2. Trajectory for $\gamma = 0.2$.

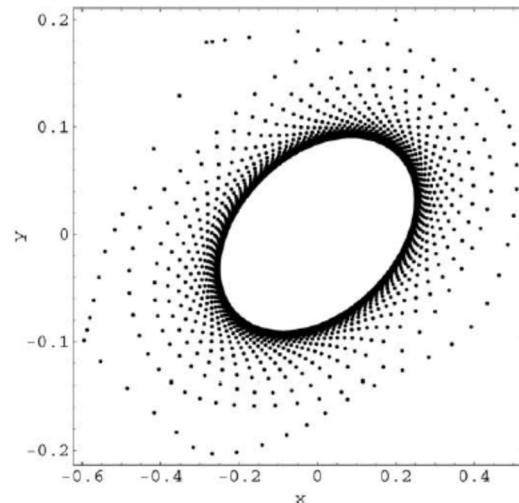


Figure 4.3. Trajectory for $\gamma = 0.2$.

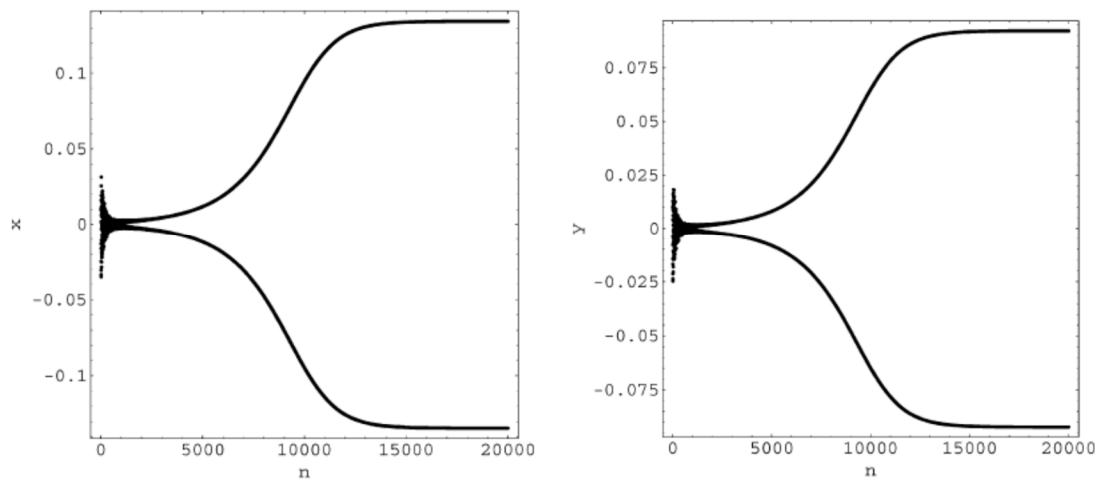


Figure 4.4. Trajectory for $\beta = -0.82, \gamma = -0.47$.

5. Case of a Discrete-Time Hopfield Neural Network with Delay and Ring Architecture

In [70] an extended analysis of the chaotic dynamics of the neural network with delay and ring architecture defined by

$$\begin{cases} x_1(n+1) = ax_1(n) + T_1 g_1(x_p(n-k_p)) \\ x_2(n+1) = ax_2(n) + T_2 g_2(x_1(n-k_1)) \\ \dots \\ x_p(n+1) = ax_p(n) + T_p g_p(x_{p-1}(n-k_{p-1})) \end{cases} \quad \forall n \geq \max(k_1, k_2, \dots, k_p). \quad (5.1)$$

was undertaken. The external inputs are equal to zero and the system parameters are regarded as numerical expression of the state of network. The network state can change due to a disease for example.

Therefore the network voltage dynamics change too. In [70] this kind of change is analyzed with respect to the network parameters. The stability domain of the null solution is found, the values of the characteristic parameter for which bifurcations occur at the origin are identified and the existence of Fold/Cusp,

Neimark-Sacker and Flip bifurcations is proved. These bifurcations were analyzed by applying the center manifold theorem and the normal form theory. It is proved that resonant 1:3 and 1:4 bifurcations may also be present. It is shown that the dynamics in a neighborhood of the null solution

become more and more complex as the characteristic parameter grows in magnitude and passes through the bifurcation values. A theoretical proof is given for the occurrence of Marotto's chaotic behavior, if the magnitudes of the interconnection coefficients are large enough and at least one of the activation functions has two simple real roots. For illustrate what this means the following numerical example was considered:

$$\begin{cases} x_1(n+1) = 0.5x_1(n) + T \tanh(x_5(n-2)) \\ x_2(n+1) = 0.5x_2(n) + \sin(x_1(n-1)) \\ x_3(n+1) = 0.5x_3(n) + T \tanh(x_2(n-2)) \\ x_4(n+1) = 0.5x_4(n) + \sin(x_3(n-1)) \\ x_5(n+1) = 0.5x_5(n) + T \tanh(x_4(n-2)) \end{cases} \quad \forall n \geq 2. \quad (5.2)$$

All numerical computations have been done using Mathematica. Denoting $b=T^3$, based on the theoretical result presented, we find that the critical values of b are $\{-7.593; -6.396; -3.784; -1.538; -0.433; -0.102; -0.036; 0.031; 0.055; 0.209; 0.851; 2.531; 5.148; 7.276\}$. The stability domain of the null solution, with respect to T is $D_s = (-0.331116; 0.31498)$.

At $T = 0.31498$, a Cusp bifurcation occurs at the origin, while at $T = -0.331116$ a supercritical Neimark-Sacker bifurcation takes place (as shown in Figure 5.1)

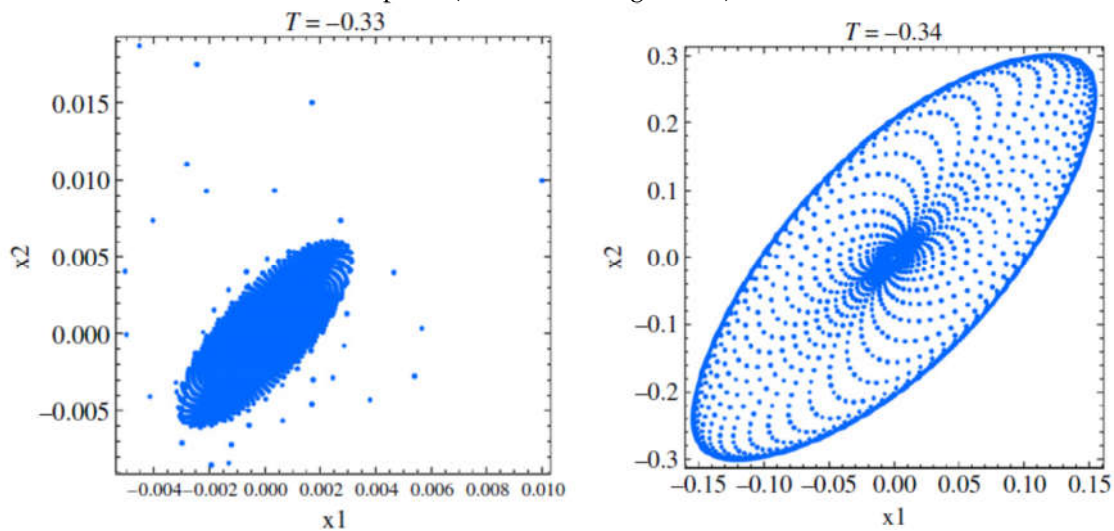


Figure 5.1. Supercritical Neimark-Sacker bifurcation at $T = -0.331116$.

For $T = -0.33$, the null solution is asymptotically stable, and the trajectory converges to the origin.

For $T = -0.34$, an asymptotically stable cycle (1-torus, drift ring) is present, and the trajectory converges to this cycle.

If modulus T is sufficiently large, chaotic behavior may be expected.

Figures 5.1–5.4 describe the route towards chaos occurring in system (5.2) in a neighborhood of the origin, as T decreases from 0 to -3 . The bifurcation diagram (Figure 5.2) and the computed four largest Lyapunov Characteristic Exponents shown in Figure 5.3, (computed using the Householder QR based method developed by Bremen, Udawadia, and Proskurowski (1997)), clearly describe this route, being consistent with the theoretical results presented.

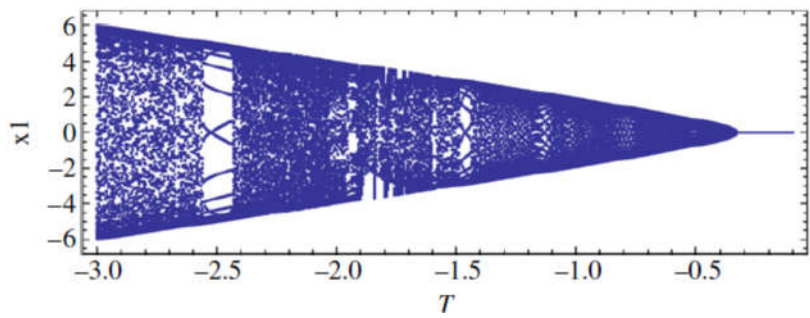


Figure 5.2. Bifurcation diagram for system (5.2), in the (T, x_1) -plane, for $-3 < T < 0$, with the step size of 0:005 for T .

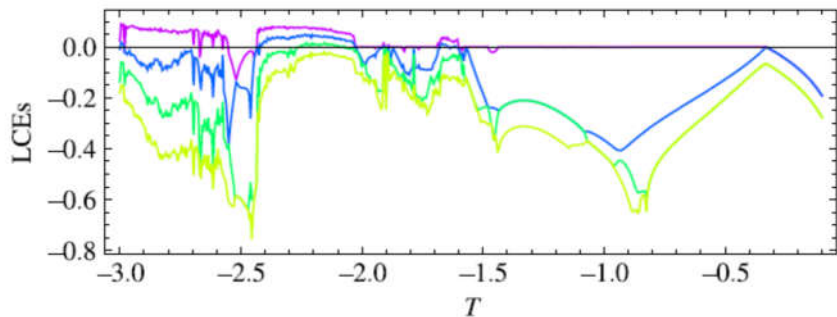
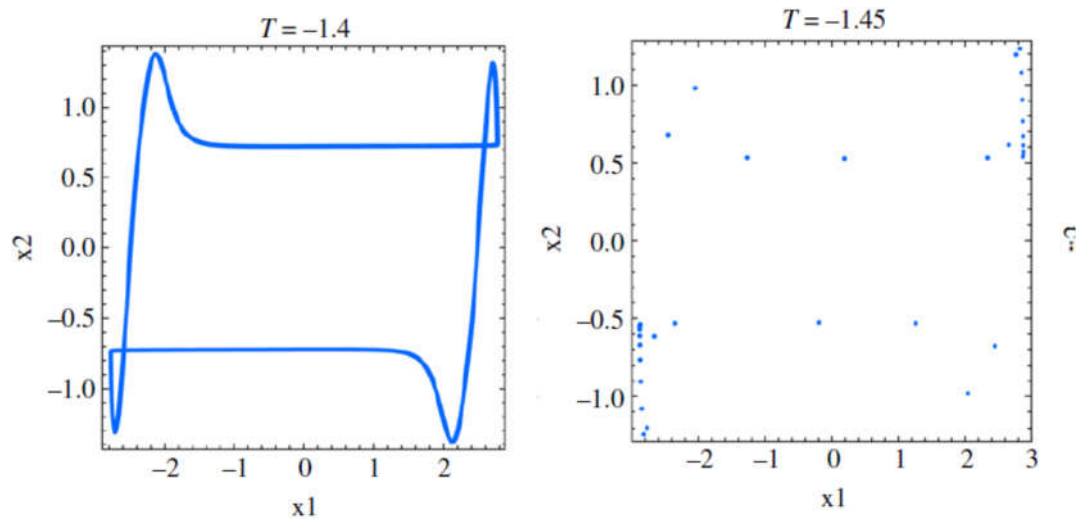
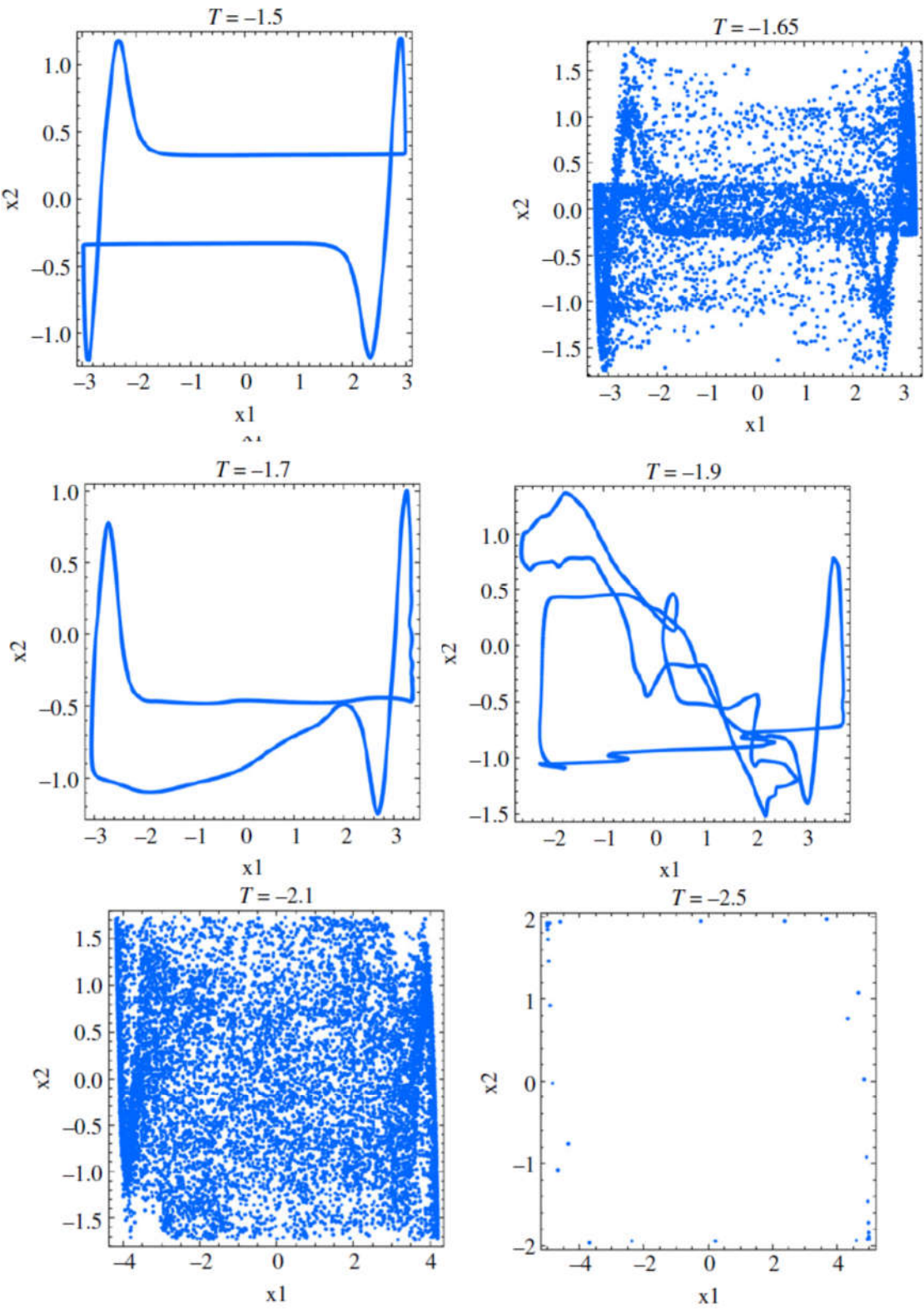


Figure 5.3. Four largest Lyapunov Characteristic Exponents for system (5.2).

For the computation of the Lyapunov spectrum, for each T value (step size 0:005 for T) the initial conditions were reset and 10^5 time steps were iterated before calculating the LCEs (which were computed over the next 10^5 time steps).

In Figure 5.4. Phase portraits for various values of T , $-3 < T < -1/4$ can be seen:





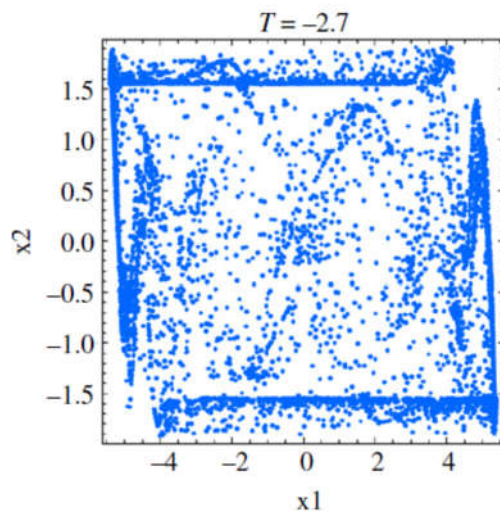


Figure 5.4. Phase portraits for: $T=-1.4$; $T=-1.45$; $T=-1.5$; $T=-1.65$; $T=-1.7$; $T=-1.9$; $T=-2.1$; $T=-2.5$; $T=-2.7$.

From the bifurcation point of view, the meaning of these figures is clear: $T=-1.4$, stable torus ($LE_1 = 0$); $T=-1.45$, stable period-32 orbit ($LE_1 < 0$); $T=-1.5$, stable torus ($LE_1 = 0$); $T=-1.65$, chaos ($LE_1 > 0$); $T=-1.7$; $T=-1.9$, stable tori ($LE_1 = 0$); $T=-2.1$, hyperchaos ($LE_3 > 0$); $T=-2.5$, stable period-38 orbit ($LE_1 < 0$); $T=-2.7$, chaos ($LE_1 > 0$).

The question is: are the phase portrait corresponding to $T=-1.45$, $T=-1.65$, $T=-2.1$, $T=-2.5$, $T=-2.7$ objects of Fractal Geometry? Have they real corresponding or they exist just in theory?

Our conjecture, based on the geometrical aspect, is that the phase portraits for $T=-1.45$, $T=-1.65$, $T=-2.1$, $T=-2.5$, $T=-2.7$ are fractals and they have real counterpart.

6. Results

6.1. In Case of a Discrete-Time Hopfield Neural Network of Two Neurons with Two Delays and No Self-Connections 23 Voltage trajectories are computed at bifurcation points. Among the 23 voltage trajectories we find 14 voltage trajectory which are not what we can call orbits in classic sense. The geometrical aspect of these trajectories suggest that they are fractals. Our conjecture is that the 14 trajectories in discussion are fractals and has real counterpart.

6.2. In case of a discrete-time Hopfield neural network of five neuron with delay and ring architecture 9 voltage trajectories appear at bifurcation points. Among the 9 voltage trajectories we find 5 voltage trajectory which are not what we can call orbits in classic sense. The geometrical aspect of these trajectories suggest that they are fractals. Our conjecture is that the 5 trajectories in discussion are fractals having real counterpart..

7. Discussion

The scientific literature concerning fractals in the description of nervous system is extensive.

In [21] for example it is shown that, Currents through Ion Channels has fractal properties in time, Electrical Activity of Auditory Nerve Cells is described in terms of fractal.

In [22] authors sustain that “fractal and conventional morphometry may represent complementary analytical/ quantitative tools to elucidate the diversity of morphological patterns and functional parameters which characterize neural cells and brain structures”.

In [23] the author writ: “the present survey provides experimental data confirming that biological processes including growth, proliferation, apoptosis, epigenetic and genetic mechanism, morphologic/ultrastructural and functional organization occurring in living shaped elements and complex structured tissues may follow fractal rules. The large agreement with the fractal nature of the brain and nervous cell system sustained by theoretical, experimental and heuristic foundations is nowadays consolidated and intervenes more than thirty years after the publication of the Fractal Geometry of Nature, in which Mandelbrot recognized that “the notion that neurons are fractals

remains conjectural” Its relevance and contribution to the cultural development of mankind (as comprehensive of humanistic and scientific thinking) is keen underlined by the observation of some years ago arguing that the fractal geometry could be considered as a biological design principle for living organisms.

Paper [24] use fractal time series analysis (detrended fluctuation analysis; DFA) to examine the spontaneous activity of single neurons in an anesthetized animal model, specifically, the mitral cells in the rat main olfactory bulb. DFA bolstered previous research in suggesting two subclasses of mitral cells. Although there was no difference in the fractal scaling of the interspike interval series at the shorter timescales, there was a significant difference at longer timescales.

Neurons in Group B exhibited fractal, power-law scaled interspike intervals, whereas neurons in Group A exhibited random variation. These results raise questions about the role of these different cells within the olfactory bulb and potential explanations of their dynamics. Specifically, self-organized criticality has been proposed as an explanation of fractal scaling in many natural systems, including neural systems. However, this theory is based on certain assumptions that do not clearly hold in the case of spontaneous neural activity, which likely reflects intrinsic cell dynamics rather than activity driven by external stimulation. Moreover, it is unclear how self-organized criticality might account for the random dynamics observed in Group A, and how these random dynamics might serve some functional role when embedded in the typical activity of the olfactory bulb.

In the paper [25] the authors conclude:” Given the central role of brain’s “wiring”, our previous research focused on the importance of fractal scaling in establishing connectivity between neurons. Diagnostic Analysis (DA) was show to relate to the optimization of competing functional constraints—the ability of dendrites to reach out and connect to other neurons versus the costs associated with doing so. Within this model, different neuron types were predicted to different DA values depending on the relative importance of connectivity and material cost with higher DA values indicating a greater emphasis on connectivity. In the current investigation, we hypothesize that pathological state of neurons might also affect this fractal optimization and consider whether changes in DA might therefore, be used as a diagnostic tool. This analysis represents an appealing development because it relates form to function rather than relying purely on pattern characterization.”

Our approach in this paper is different from those presented in the above refered literature. We start considering the discrete-time Hopfield neural network model which claim to describe the electrical activity of a neural system. The bifurcation analysis of some low dimensional networks show that for certain values of the bifurcation parameter the geometry of the voltage orbits of neurons are what we call classic orbits and there exists other bifurcation values for which the geometrical aspect of trajectories are not at all what we call classic orbits. Their geometrical aspects suggest that they are moreover fractals. In the present paper present to the reader both type of trajectories putting the natural question: Are some of the presented voltage state orbits, fractals? Exists these orbits in reality or they exists just in theory i.e. has these orbits a real counterpart?

Our conjecture is that some of the voltage trajectories presented here are fractals and have real counterpart.

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