

Review

Are Microbes Thermodynamically Optimised Self-Reproducing Machines?

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Abstract: To understand microbial growth with mathematical models has a long tradition that dates back to the pioneering work of Jacques Monod in the 1940s. Growth laws are simple mathematical expressions that aim at describing growth rates of microbes as functions of external parameters, in particular nutrient concentrations. These laws are now widely applied to construct, e.g., dynamic ecosystem models. However, to explain the growth laws from underlying (first) principles is extremely challenging. In the second half of the 20th century, numerous experimental approaches aimed at precisely measuring heat production during microbial growth to determine the entropy balance in a growing cell and to quantify the exported entropy. This has led to the development of thermodynamic theories of microbial growth, which have generated fundamental understanding and identified principle limitations of the growth process. Whereas these approaches considered a growing microbe as a black box, modern theories heavily rely on genomic resources to describe and model genome-scale networks to explain microbial growth. Interestingly, however, thermodynamic constraints are often included in modern modelling approaches only in a rather superficial fashion, and it appears that recent modelling approaches and classical theories are disconnected fields. In order to stimulate a closer interaction between these fields, we here review various theoretical approaches that aim at describing microbial growth based on thermodynamic principles. We start with classical black-box models of cellular growth, and continue with genome-scale modelling approaches that include thermodynamics, before we place these models in the context of fundamental considerations based on non-equilibrium statistical mechanics. We conclude by identifying conceptual overlaps between the fields and suggest how the various types of theories and models can be integrated. We outline how concepts from one approach may help to inform or constrain another, and we demonstrate how genome-scale models can be used to infer classical black-box parameters, which are experimentally accessible in growth experiments. Such integration will allow understanding to what extent microbes can be viewed as thermodynamic machines, and how close they operate to theoretical optima.

Keywords: Energy; Entropy; Anabolism; Catabolism; Microbial cultures; Biotechnology

1. Introduction

Life is certainly one of the greatest wonders on earth. The ability to grow and reproduce distinguishes living systems fundamentally from inanimate objects and engineered devices. However, of course any living organism is also subject to universal physical laws. As such, life is nothing special. No surprise, therefore, that generations of scientists have studied life, and in particular

microbial growth, using concepts and theories from physics. Whether microbes can be considered as sophisticated biochemical machines is a deeply philosophical and highly relevant question [1]. We do not wish to enter this discourse here, and rather adopt a pragmatic point of view. Certainly, every microbe uses available free energy gradients (either from chemicals in the environment, or – in the case of photosynthetic organisms – from sunlight) to convert chemicals from the environment into a copy of itself. Thus, microbes are not fundamentally different from machines, except that they do not convert free energy gradients into mechanical work that is useful for all kinds of purposes, but into chemical work, which is used to construct a copy of themselves. So why should not concepts, principles, and formulas from thermodynamics, which can so well describe steam engines, be applicable also to microbial growth? In fact, it was exactly this mindset that led generations of scientists to develop excellent thermodynamic theories of microbial growth in the second half of the 20th century. Notably, these researchers had only very limited knowledge of the detailed biochemical processes inside a cell, which are responsible for the macroscopic behaviour. Not surprisingly, the theories consider a microbe as a *black box*. This view matches well with the available experimental technologies at that time [2,3]. In bioreactors and fermenters, fluxes of matter and energy into and out of the cells could be precisely measured, and, together with observed growth rates, were ideal to challenge the developed theories [4–7].

With the advent of high-throughput technologies, we have now an increasingly detailed picture of the internal processes within a microbe. Especially metabolism, the biochemical networks converting chemicals into biomass, is described with increasing detail. Consequently, new modelling techniques and theories have emerged to describe and investigate *genome-scale metabolic networks*. Unfortunately, these modern approaches now form a separate research field that appears to be unconnected to existing thermodynamic theories of life. Scientists developing genome-scale network models and corresponding analysis techniques often seem to be unaware of the fundamental thermodynamic theories and thus ignore an invaluable heritage left behind by generations of great minds.

The purpose of this review is to provide a summary of essential developments of theory building, both from the times before high-throughput technologies and after. We aim at developing links and connections between modern modelling approaches and thermodynamic theories of microbial growth, with the hope that these existing and promising research fields can mutually benefit from each other. Genome-scale metabolic network models provide an unprecedented opportunity to peek inside the black boxes. In conjunction with recent developments of non-equilibrium thermodynamics to describe self-replication, we are convinced that an integration of the different approaches to describe microbial growth offers the chance to discover fundamental principles of life.

2. Historical Overview

The economy of energy flow in living matter is a fascinating and complex topic that is and was investigated by generations of great scientists. Starting with the pioneering work of Jaques Monod [8–10], this section aims to give a concise overview of the historical developments in the treatment of thermodynamic aspects in models describing microbial growth. By including not only pure theoretical work but also past biotechnological examples, this part illustrates the motivation of the researchers to focus on energetic aspects of microbial cultures, thereby uncovering numerous interesting relationships.

Starting with Monod in the 1940s [9], who defined the biomass yield factor as the amount of biomass formed divided by the amount of limiting nutrient utilized (Y), the quantitative description of the production capabilities in microbial cultures evolved into a whole theory. The introduction of a simple hyperbolic relationship that connects the amount of limiting resources in the environment to growth rates of organisms opened the door for many different modeling approaches. Even today, the Monod growth law together with modern mathematical tools forms the basis for sophisticated theories of microbial growth. Although the work of Jaques Monod is undoubtedly one of the major achievements in the life sciences of the 20th century, it is worth mentioning that others have

81 proposed growth laws with slightly different functional dependencies of the growth rate on substrate
82 availability [11,12]. As Esener *et al.* show, all three models are able to realistically describe experimental
83 data of batch cultures [13].

84 While the description of batch cultivation could be adequately described by Monod's theory,
85 the growth of microbes in continuous cultures needed a different treatment. In 1956 Herbert *et al.*
86 presented a formal analysis of continuous cultures, introducing fundamental equations for describing
87 those systems efficiently [14]. In addition to this, Herbert realised in 1958 [13,15] that the yield factor
88 introduced in the 1940s by Monod is not constant but changes with the dilution rate. Both Herbert
89 (1958) and Pirt (1965) tried to explain this behaviour by introducing maintenance terms that describe
90 a growth-independent substrate requirement [15,16]. Although Herbert [15] and Pirt [16] provide
91 different explanations for the maintenance term, Esener *et al.* have argued that both explanations can
92 be used to essentially obtain the same results [13]. For a detailed discussion of different maintenance
93 parameters and their measurement see [17].

94 With these theories in hand, the question appeared how essential parameters of microbial systems
95 can be easily and relatively quickly estimated, given the diversity of organisms used in the laboratories
96 or in biotechnology, especially regarding the limited information about the metabolism available at
97 that time. The solution was the introduction of so called *black box models* that only need input and
98 output information, which can be readily obtained by controlling a bioreactor [2,3,13,18–22]. However,
99 since the applicability of black box models heavily depends on energy and mass balance, a profound
100 thermodynamic background is needed to understand and justify these approaches. Black box models
101 were successfully applied to reveal many interesting relationships between energy, biomass and yield.

102 As Mayberry *et al.* state in the 1960s, there was a considerable interest in investigating the
103 proportionality between the yield and the converted energy [23]. One important contribution was
104 made by Mayberry *et al.* in 1968, who observed that for bacteria growing on a single carbon compound
105 that serves as carbon and energy source, the amount of produced biomass per available electrons
106 ($\text{av } e^-$) is relatively stable ($3.14 \text{ gDW} / \text{av } e^-$) for a wide range of organic compounds. By redefining
107 the concept of the "degree of reduction", which was first proposed by Gunsalus and Shuster [24], as
108 "[number] of those electrons in a compound not involved in orbitals with oxygen [...]", and using the
109 above mentioned observation, Mayberry *et al.* stressed the usefulness of counting available electrons
110 for the analysis of growth data. A significant contribution was to introduce the degree of reduction in
111 the treatment of energetic aspects in biology as a measurement of reducing power of a compound. In
112 1973 Minkevich and Eroshin [25] derived a formula for calculating the degrees of reduction (γ) of the
113 electron donor, the biomass and metabolic products as well as a formula to obtain the dried biomass
114 per oxygen using an overall metabolic reaction. However, great care must be taken when calculating
115 and interpreting γ , since its derivation is dependent on the nitrogen source. Therefore, Roels later
116 defined a generalized formula for the degree of reduction that can be used for any available nitrogen
117 source [19].

118 By comparing numerous different organisms, carbon sources and literature, Minkevich and
119 Eroshin concluded that the degree of reduction of biomass (γ_b) for many species varies around 4.2
120 (for comparison, glucose and acetate have a degree of reduction of 4 per carbon, while methanol has
121 6, and oxalic acid 1). In addition to this, they used old thermochemical knowledge from the 19th
122 and early 20th century, namely that the molar combustion heats of organic compounds are nearly
123 proportional to the consumed oxygen (known as Thornton's rule [26]) to derive an efficiency parameter
124 of growth (η), which specifies the energy fraction (contained in the substrate) transferred to biomass.
125 This efficiency parameter sets a thermodynamic upper bound of the yield. Setting $\eta = 1$, one obtains
126 an upper bound of $Y < \frac{\gamma_s}{\gamma_b}$, where γ_b denotes the degree of reduction of the biomass and γ_s of the
127 electron donor [18,25]. Nonetheless, this relationship is only valid if both, the biomass and electron
128 donor, possess the same heat of combustion per available electrons. By studying different chemical
129 species, Minkevich and Eroshin confirmed that the heat of combustion per gram-equivalent of oxygen

130 consumed or per available electrons is for most organic substances around 26.94 kcal. Similar values
131 have been recalculated by numerous authors (for instance see [27,28]).

132 Later Erickson together with Minkevich and Eroshin introduced two more efficiency measures:
133 The fraction of the energy in an organic compound transferred to products (ζ_p) and the fraction
134 that is evolved as heat (ϵ) during the growth process. Combined with the aforementioned energy
135 fraction transferred to biomass (η), Erickson *et al.* developed simple mathematical tests to check for the
136 consistency of various measured parameters, including growth and maintenance requirements [18].

137 However, as Roels states in his excellent report, those parameters are based on a not entirely
138 correct formulation of the second law of thermodynamics for open systems, because they assume that
139 cellular growth always releases heat to the environment [19]. However, there are examples of microbes
140 importing heat from the environment [2,5]. Roels provides a modified definition of maximum yield
141 on substrate that are consistent with the second law, ω_f , consistent with the heat transported to the
142 environment, ω_e , and consistent with the oxygen transport, ω_o . By showing that the values of ω_f and
143 ω_e most often only deviate by 10%, he basically confirmed that Minkevich's and Eroshin's η was a
144 useful parameter, at least under aerobic conditions.

145 Since ω_f is the maximum value of the yield that is consistent with the laws of thermodynamics,
146 Roels defined the thermodynamic efficiency of a growth process as $\eta_{th} = Y_{sx}/\omega_f$, where Y_{sx} is the
147 yield of biomass (x) given a certain substrate (s). By calculating the values of η_{th} for different organic
148 compounds as energy source, he observed that highly reduced as well as highly oxidized substrates
149 lead to a rather low efficiency. Roels concluded that substrates with a degree of reduction above 4.2
150 (the degree of reduction of biomass) contain enough energy to allow that all carbon could in principle
151 be converted into biomass. In contrast to this, for compounds with a degree of reduction below 4.2,
152 this is not possible due to energetic reasons.

153 How to incorporate these results into a mathematical model of bacterial growth is extensively
154 discussed by Esener in his report "Theory and Application of unstructured Models: Kinetic and
155 Energetic Aspects" [13]. Together with his paper published in 1982, in which Esener *et al.* describe
156 how to include aspects of varying biomass composition under a changing environment in a formal
157 description of bacterial growth [29], his work is a good example of an early attempt to develop
158 a consistent model explaining bacterial growth. These approaches are clearly related to modern
159 approaches to explain bacterial growth laws (see [30,31]). For a very well written discussion comparing
160 most of the efficiency measures discussed above (and much more), see [22,32].

161 In combination, these investigations and concepts of thermodynamic efficiency in bacterial
162 growth provided the basis for further developments, including methods to estimate the energy and
163 entropy of formation for biomass, which represents a parameter of fundamental importance for
164 energetic calculations concerning life [27,28,33,34]. Particularly valuable were Battley's contributions
165 for estimating the Gibbs free energy of formation of biomass, $\Delta_f G_b$ (-65.10 kJ/C-mol, $\gamma_b = 4.998$, N₂ as
166 nitrogen source), and the enthalpy of formation for biomass, $\Delta_f H_b$. By using the well known relation of
167 entropy, enthalpy and Gibbs free energy, $\Delta G = \Delta H - T\Delta S$, he even attempted to estimate the entropy
168 of formation of biomass, but concluded that this method is too prone for errors, because it highly
169 depends on the quality of the approximation of enthalpy and Gibbs free energy [28].

170 3. Recent applications of Thermodynamics in microbial growth

171 The introduction of novel experimental techniques, often referred to as high-throughput 'omics'
172 technologies, around the beginning of the 21st century, facilitated collecting extensive datasets and
173 information about important molecular components of cells, including metabolite, RNA and protein
174 levels. Falling cost and improved efficiency of genome sequencing gave rise to an ever growing
175 number of sequenced organisms [35]. The availability of omics data also resulted in comprehensive
176 biochemical databases that collect information about enzymatic reactions and their compounds [36,
177 37]. The construction of genome-scale metabolic models [38] is still a widely used approach to
178 investigate microbial growth rates, using flux balance analysis (FBA) as a method to predict metabolic

179 fluxes in optimised growth scenarios [39]. Well curated genome-scale metabolic models can provide
180 valuable information about flux distributions of cells growing on different media. In combination with
181 experimental measurements, this approach has successfully been applied to investigate uptake rates of
182 carbon sources and byproduct secretion rates for several microbial organisms [40–43]. However, there
183 are a number of phenomena associated with microbial growth which cannot be explained by these
184 type of models. For example, they do not explain the observation of fermentative metabolism and only
185 partial oxidation of organic substrates in aerobic conditions and high substrate availability (known as
186 the Crabtree effect in yeast or the Warburg effect in cancer cells) [44]. In general, flux balance analysis
187 is very limited when investigating fundamental principles of microbial growth.

Thermodynamic approaches are typically used in genome-scale metabolic models to avoid infeasible flux distributions in the solution space. Thermodynamics gives information about the correct direction of biochemical reactions [45]. Essentially any enzymatic reaction can be reversible, and the Gibbs free energy of reaction ($\Delta_R G$) defines the direction in which a reaction proceeds (always in the direction of negative $\Delta_R G$). The Gibbs free energy of reaction depends on the metabolite concentrations by

$$\Delta_R G = \sum_i n_i \Delta_f G_i^0 + RT \sum_i n_i \ln c_i, \quad (1)$$

188 where n_i is the stoichiometric coefficient of metabolite i , c_i its concentration, and $\Delta_f G_i^0$ the respective
189 standard Gibbs energy of formation. However, flux balance analysis does not incorporate metabolite
190 concentrations. Therefore, often physiological substrate concentration ranges are assumed [46],
191 and in combination with – often computationally derived – $\Delta_f G^0$ values [47], limitations on
192 reaction directionality can be calculated. This approach, if applied correctly, guarantees that only
193 thermodynamically feasible pathways [48] are found in the solution space. With incorrect assignment
194 of reaction directions, infeasible pathways can arise in genome-scale metabolic models and result
195 e.g. in pathways that produce ATP without external driving forces.

196 In recent years, resource allocation models and models including molecular crowding have
197 been developed to further investigate microbial growth limitations and overflow metabolism [49–51].
198 Whereas flux balance analysis neglects specific protein costs and benefits, resource allocation models are
199 focusing on incorporating costs for enzymes, ribosomes, available space in cells and other physiological
200 aspects. In fact, understanding microbial growth as an economic process, in which limited resources
201 (such as macromolecules or space) have to be efficiently allocated, dates back to the pioneering work
202 of Schaechter *et al.* [52,53] in the 1950s. Recently, comprehensive resource allocation models were
203 proposed, which allow calculating and thus further understanding balances and trade-offs during
204 microbial growth [49]. Molenaar *et al.* argue that it is beneficial to replace a simple biomass flux
205 objective with a more complex objective of self replication, because even the basic self-replicator model
206 introduced in their work exhibits a high degree of complex behaviours [44]. Noor *et al.* argue that
207 Flux-force relationships show that the logarithm of the ratio of forward and backward fluxes of a
208 reaction is proportional to the Gibbs energy dissipation of the reaction [54], and therefore conclude
209 that reactions with Gibbs energy dissipation close to zero require increased amounts of enzyme to
210 facilitate a given net flux. In most resource allocation models, these thermodynamic contributions are
211 not included. Extensive incorporation of thermodynamic principles has been applied by Niebel *et al.*
212 [55] to analyze growth of yeast and *E. coli*, including thermodynamic constraints. In their work,
213 it was suggested, based on experimental data, that an upper limit of the Gibbs energy dissipation
214 rate exists. This assumption allows the formulation of an equation for the Gibbs energy balance.
215 This balance constrains the sum of all Gibbs energy dissipation rates of all internal reactions to be
216 equal to the exchange of Gibbs energy with the environment. The resulting models were used to
217 predict oxygen uptake as well as byproduct and biomass production rates for different substrate
218 uptake rates. Further, some intracellular fluxes were predicted and then measured with ^{13}C -MFA.
219 In their work it was observed that with increasing substrate uptake rates, the model shifted flux
220 distributions from pathways with high dissipation rates, such as respiration, to pathways with lower

221 dissipation rates, such as fermentation. Thus they could provide a putative explanation, based on
 222 thermodynamic principles, for overflow metabolism phenomena, such as the Crabtree and Warburg
 223 effects. In contrast to standard flux balance analysis of genome-scale models, the model presented
 224 in their work also incorporates metabolite concentrations. This approach, while certainly increasing
 225 thermodynamic rigour and improving the predictive capabilities, comes at the cost of increased
 226 computational complexity. To find flux distributions that maximize growth, the authors used a
 227 mixed-integer non-linear program.

228 4. Thermodynamic approaches to self replication

In the 1990s, the development of the so-called fluctuation theorems [56–59] constituted a major advance in non-equilibrium statistical physics. Basically, they represent a generalisation of the second law of thermodynamics. They link the probability to observe an entropy increase σ during a time τ to the probability to observe an entropy decrease by the same amount in the same time. This class of theorems can be generally expressed by

$$\frac{\mathcal{P}(+\sigma)}{\mathcal{P}(-\sigma)} = \exp(\sigma\tau), \quad (2)$$

where $\mathcal{P}(+\sigma)$ and $\mathcal{P}(-\sigma)$ denote the probabilities to observe an entropy increase or decrease by σ during time τ , respectively. More recently, J. L. England [60] has proposed to apply this approach to self-replicating systems. Obviously, replication is a highly irreversible biological process and can be described in the language of statistical physics as a system that goes from a macrostate **I** (a single cell and the substrates in the surrounding medium), to a macrostate **II** (two daughter cells and the substrates), where each macrostate corresponds to an extremely high number of microstates. England's reasoning starts from the fact that particles obey classical mechanics at the microscopic scale, and therefore follow a reversible dynamics. This allows quantifying the reversibility of a microscopic transition by the associated change in entropy. Applying these microscopic considerations to the macroscopic scale, the author obtains a generalization of the second law of thermodynamics for macroscopic irreversible biological processes. While the classical second law of thermodynamics states that the increase of entropy of a closed system is always positive and obeys the inequality

$$\Delta Q_{\text{ex}} + T\Delta S_{\text{int}} \geq 0, \quad (3)$$

where ΔQ_{ex} is the amount of heat exchanged with the environment and ΔS_{int} the internal entropy increase of the system. England's derivation adds a new term to this relation,

$$\frac{\Delta Q_{\text{ex}}}{T} + \Delta S_{\text{int}} + \ln \frac{\pi(\mathbf{II} \rightarrow \mathbf{I})}{\pi(\mathbf{I} \rightarrow \mathbf{II})} \geq 0, \quad (4)$$

229 where $\pi(\mathbf{I} \rightarrow \mathbf{II})$ (respectively $\pi(\mathbf{II} \rightarrow \mathbf{I})$) stands for the probability that the system evolves from
 230 macrostate **I** to macrostate **II** (respectively from **II** to **I**). When a macroscopic transition is irreversible
 231 ($\pi(\mathbf{II} \rightarrow \mathbf{I}) \ll \pi(\mathbf{I} \rightarrow \mathbf{II})$), the logarithm becomes negative, increasing the lower bounds for heat
 232 dissipation and entropy increase.

Equation (4) allows to have a closer look at self-replication. England applies this relation to a population of exponentially growing cells. He denotes the growth rate as g and the reverse rate (highly unlikely to happen) as δ . It allows to express $\pi(\mathbf{I} \rightarrow \mathbf{II})$ as gdt and $\pi(\mathbf{II} \rightarrow \mathbf{I})$ as δdt . Hence, equation (4) becomes

$$\frac{\Delta q}{T} + \Delta s_{\text{int}} \geq \ln(g/\delta), \quad (5)$$

Employing now the formula to calculate the degree of reduction when N_2 is assumed to be the nitrogen source,

$$\gamma = 4n_C + n_H - 2n_O - 0n_N + 5n_P + 6n_S - n_{e^-}, \quad (7)$$

274 in combination with the relation between the energy of combustion for an organic compound and its
 275 degree of reduction (Battley assumes $-107.90 \text{ kJ/av } e^-$), he was able to estimate the energy of formation
 276 for biomass as -65.10 kJ/C-mol (not including ions). The n_X in equation (7) denotes the number of
 277 atoms of type X or charge (n_{e^-}), respectively. The energies of formation necessary for obtaining this
 278 value are listed in table 1.

279 Because all necessary information for repeating this calculation are contained in curated
 280 genome-scale models, it is rather straight-forward to transfer this old technique to modern modelling
 281 approaches. To demonstrate this, we determine the $\Delta_f G_b$ -values for all genome-scale models contained
 282 in the BIGG database [62], using the specified biomass reactions. Figure 1 shows the distribution of
 283 energies of formation as function of the degree of reduction.

Table 1. Standard free energy of formation for various organic compounds of interest necessary for estimating the energy of formation for biomass (see [28])

Substance	Formula	$\Delta_f G_i^0$ [kJ/mol]
Oxygen	O_2 (g)	0
Potassium hydroxide	KOH (c)	-379.11
Carbon dioxide	CO_2 (g)	-394.36
Nitrogen	N_2 (g)	0
Phosphorous decoxide	P_4O_{10} (c)	-2697.84
Potassium sulfate	K_2SO_4 (c)	-1321.43
Water	H_2O (lq)	-237.18

284 A comparison with the reported values ($\Delta_f G_b = -65.10 \text{ kJ/C-mol}$ and $\gamma = 4.998$ – see historical
 285 overview in Section 2) reveals that for most models energies of formation and degrees of reduction
 286 of the biomass are in agreement with former theoretical approximations. However, there are some
 287 values that deviate drastically from the mean (compare, for instance, upper right corner in figure 1).
 288 Possibly, the observed variation results from different biomass compositions that were assumed for
 289 the specific models and their particular research questions. However, this kind of calculation offers the
 290 opportunity to scrutinise the plausibility of model assumptions, in particular referring to the biomass
 291 functions. For example, an energy of formation of $+200 \text{ kJ/C-mol}$ seems highly unlikely.

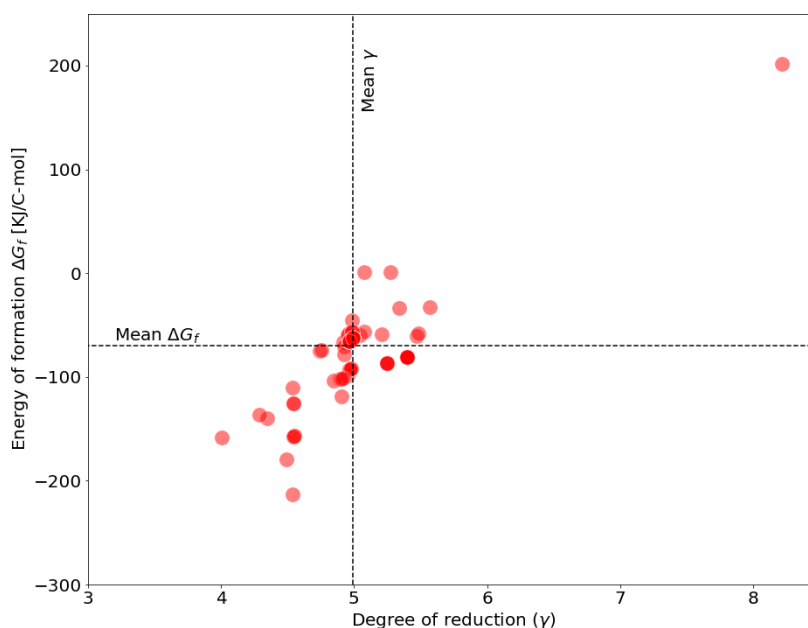


Figure 1. Energy of formation $\Delta_f G_b$ for biomass as encoded in genome-scale models of the BIGG database. Each point represents a biomass composition as described in the models. 85 models of the BIGG database containing 165 biomass functions were analysed. Mean $\Delta_f G_b = -70.079$ kJ/C-mol, Mean $\gamma = 4.996$.

292 Another possibility for a straight-forward approach to combine thermodynamic concepts from
 293 black box models with genome-scale models is a separate analysis of anabolism and catabolism.
 294 In order to calculate properties of anabolism, such as those predicted by Battley in 1993 [28], from
 295 genome-scale models, we pursue the following approach: Genome-scale models from the BIGG
 296 database were modified in two steps. First, all reactions that can produce ATP are disabled by
 297 introducing a dummy compound representing "unusable" ATP. Second, two strictly coupled reactions
 298 are introduced that import ATP into the cytosol, and export ADP and orthophosphate with the same
 299 rate. The strict coupling of import and export ensures that only energy but no matter is introduced into
 300 the system. Thus, the modified model is unable to produce ATP from any carbon source and instead
 301 must use the imported ATP as energy source. Therefore, this modification separates anabolism from
 302 catabolism by simulating an external "ATP battery" providing the organism with external chemical
 303 energy, replacing the usual catabolic pathways.

304 These modified models were used to simulate anabolism separate from catabolism. In particular,
 305 we calculated the minimum amount of ATP required to incorporate one carbon atom from the nutrients
 306 into biomass and the minimum number of CO_2 molecules that are released in this process. For this,
 307 the biomass production rate was fixed to the value obtained from the original model (in all cases
 308 the objective was maximisation of the biomass production rate) and subsequently minimizing all
 309 carbohydrate import fluxes. The results are shown in Figure 2. Every point represents one model
 310 from the BIGG database. The x-axis displays the ratio of carbon dioxide produced by anabolism only
 311 to that produced by the original metabolism, in which anabolism and catabolism are coupled. The
 312 y-axis shows how many moles ATP are minimally required to incorporate one mole carbon atoms into
 313 the biomass. The large number of data points sharing the same anabolism versus metabolism ratio
 314 of released CO_2 can be explained by the fact that a large proportion of models available in the BIGG
 315 database are for *E. coli*, and use the same biomass definition. For all *E. coli* models except the "core

316 model", the required ATP per biomass carbon is between 2 and 3.5. Interestingly, the anabolism versus
 317 metabolism ratio of released carbon dioxide for the *E. coli* models is very close to the ratio predicted by
 318 Battley [28] (indicated by the dotted black line).

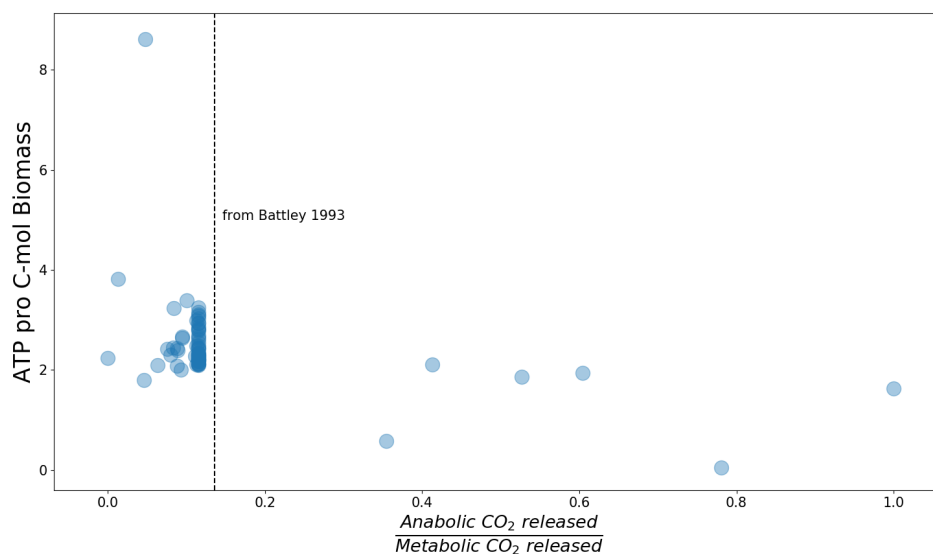


Figure 2. Anabolic properties of genome-scale models of the BIGG database. The y-axis indicates the minimum required amount of ATP per biomass carbon. The x-axis displays the ratio of carbon dioxide released by anabolism to carbon dioxide released by overall metabolism (including anabolism and catabolism).

319 6. Outlook & Conclusion

320 To increase fundamental understanding of microbial self-replication and its limitations,
 321 thermodynamic considerations and calculations can give valuable insights and expand the current
 322 knowledge about microbial growth. When attempting to answer the question whether prokaryotic
 323 organisms are thermodynamically optimised self-replicating machines, one is immediately confronted
 324 with a number of fundamental problems. There does not seem to be an agreement what exactly
 325 "thermodynamic optimality" means. The "black box" models, discussed in section 2, that arose in the
 326 20th century focussed on estimations of microbial growth efficiency parameters by only considering
 327 the exchange of matter and energy between cells and their environment. The parameters for aerobic
 328 microbial growth efficiency, proposed by Minkevich, Eroshin and Erickson [18,25], incorporate the
 329 degree of reduction of the biomass, substrates and electron donors to estimate a thermodynamic upper
 330 bound for the yield. This efficiency of biomass yield can be considered an important interpretation of
 331 thermodynamic optimality.

332 The upsurge of high-throughput data in the early 21st century gave rise to genome-scale
 333 metabolic models, opening the "black boxes" of microbial metabolism. These models provide excellent
 334 tools to investigate metabolic flux distributions. Some genome-scale metabolic models consider
 335 thermodynamics to constrain intracellular fluxes, but rarely to find constraints for self-replication
 336 and growth itself. The work of Niebel *et al.* [55], mentioned in section 3, is one of the most
 337 comprehensive approaches how to employ thermodynamics for the investigation of microbial
 338 growth using genome-scale metabolic models. Based on experimental data, the authors proposed
 339 an upper limit of Gibbs free energy dissipation during growth. They supported their hypothesis
 340 by implementing additional thermodynamic constraints reflecting this upper limit of Gibbs free
 341 energy dissipation. A flux distribution optimising growth, which obeys this upper limit can also
 342 be considered as an interpretation of "thermodynamic optimality". A redistribution of fluxes to
 343 fermentative pathways due to constraints by the upper limit of Gibbs energy dissipation described by
 344 Niebel *et al.* [55] only occurs at high glucose uptake rates in *E. coli* and *S. cerevisiae*. Such conditions

345 of excess nutrient availability are rarely found in nature and therefore an evolution towards such
346 thermodynamic optimality is questionable. However, a hypothetical environment of fluctuating
347 periods of excess nutrient availability could provide a fitness increase for microbial organisms adapted
348 to flux redistribution obeying the upper limit of Gibbs energy dissipation.

349 A promising complementary approach attempts to describe microbial self-replication by first
350 principles from physics. In stark contrast to genome-scale models, which heavily depend on
351 high-throughput data and computational power, first principle concepts require only a minimum
352 amount of data. The work by England, mentioned in section 4, aims at understanding microbial growth
353 from a thermodynamic perspective. By applying fluctuation theorems to the non-equilibrium process
354 of self-replication, a generalization of the second law of thermodynamics for irreversible transitions
355 between two macroscopic states has been derived. This allowed calculating a lower bound for the
356 produced heat during self-replication as a function of the internal entropy, growth and decaying rates.
357 A consequence of these calculations is that a self-replicating microbe that dissipates heat with a rate
358 close to the thermodynamic minimum is optimal in the sense that energy loss is minimised. However,
359 the maximal rate of self-replication increases with increased heat dissipation. The finding that the
360 heat dissipation of *E. coli* is not far from the calculated minimum needed for self-replication hints at
361 evolution towards thermodynamic efficiency. However, the calculations imply that replication rates
362 are increased with higher internal entropy and an increased rate of spontaneous self-decay. Both
363 properties are not commonly found in microbial organisms. These properties may be beneficial to
364 increase growth rate from a thermodynamic perspective, but are probably disadvantageous regarding
365 other evolutionary pressures.

366 Evidently, there is no simple and unique answer to the question whether microbes are
367 thermodynamically optimised self-replicating machines. Although all three concepts described here
368 are concerned with the same phenomenon, each represents a different perception and viewpoint on
369 thermodynamic optimality of microbial growth. In our opinion these three concepts, as different as
370 they may be, host an enormous potential to complement each other into an extended understanding of
371 thermodynamic limitations and optimality of microbial growth.

372 The minimum amount of heat dissipation, and the upper limit of Gibbs free energy dissipation
373 define fundamental thermodynamic limitations of microbial growth. The lower bound is a consequence
374 of the extended second law of thermodynamics: It is impossible to replicate with less dissipated heat.
375 The upper bound is an empiric observation which so far has not experienced a theoretical explanation.
376 It implies that there is a principle upper limit for microbial growth rates. In addition, black box models
377 allow calculating upper bounds of the yield based on physical-chemical properties of substrate and
378 biomass, most importantly the degree of reduction. It is not yet clear how these three limitations can
379 be considered simultaneously, whether they are in agreement, and how combining them might result
380 in a more confined thermodynamic space. An interesting observation was made by von Stockar [5],
381 who reviewed thermodynamic data on microbial growth. This review includes examples of microbes
382 that import heat from their environment, and compensate this by exporting chemical entropy [2].
383 Other microbes can do the opposite and even reduce the chemical entropy of their environment, at the
384 expense of increased heat dissipation [63].

385 After pointing out that by combining the different interpretations of thermodynamics of microbial
386 growth their potential could be drastically enhanced, we suggest that attempts to merge different
387 thermodynamic theories should be further intensified. One example of such a combined approach is
388 the analysis of genome-scale models using constraints and concepts from black box models of microbial
389 metabolism. Section 5 shows that in a relatively simple fashion, the combination of genome-scale
390 models and black box models provides promising strategies to further understand thermodynamics
391 of microbial growth and metabolism. However, to incorporate black box model concepts, metabolic
392 networks need to guarantee to obey fundamental conservation laws. For example, calculating the upper
393 thermodynamic limit of the yield ($Y < \frac{\gamma_s}{\gamma_b}$), requires the determination of the degrees of reduction
394 for the biomass and all substrates. This heavily depends on a fully mass- and charge-balanced

395 metabolic network. During our analysis, we noticed that only few models can actually ensure this.
 396 This is understandable if one considers that most models were constructed for a completely different
 397 purpose. Finding flux distributions that optimise the incorporation of carbon into biomass also work
 398 without a perfectly charge-balanced network, and even an incorrect mass balance, which only concerns
 399 hydrogen atoms, will not affect the results. Nonetheless, in order to be truly reusable, in particular
 400 for thermodynamic calculations, any genome-scale model that is published should adhere to these
 401 fundamental chemical principles. We hope that this review encourages further activities to integrate
 402 different thermodynamic concepts and motivates the introduction of stricter standards to ensure
 403 reusability of genome-scale metabolic models.

404 **Supplementary Materials:** All code for analysing published genome-scale metabolic networks in Section 5 is
 405 available at <https://gitlab.com/qt-b-hhu/thermodynamics-in-genome-scale-models>.

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415 Abbreviations

416 The following abbreviations are used in this manuscript:

417 Y	Biomass yield
γ	Degree of reduction
$\Delta_R G$	Gibbs free energy of reaction
$\Delta_f G_b$	Gibbs free energy of formation for biomass
$\Delta_f G_i^0$	Standard Gibbs free energy of formation for a metabolite i
σ	Entropy production
418 τ	Time
T	Temperature in Joule, k_B set to 1
$\mathcal{P}(\pm\sigma)$	Probability to observe an entropy production of $\pm\sigma$
$\pi(\mathbf{II} \rightarrow \mathbf{I})$	Probability to observe a transition from macrostate \mathbf{II} to \mathbf{I}
$\pi(\mathbf{I} \rightarrow \mathbf{II})$	Probability to observe a transition from macrostate \mathbf{I} to \mathbf{II}
g	Duplication rate
δ	Decaying rate

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