

Selection of co-belonging ceramic fragments from archaeological excavations and their location in vase bodies from thermoremanent magnetization

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Abstract

Selection of co-belonging fragments from the numerous ceramic findings of an archaeological excavation remains a difficult process of questionable effectiveness, based exclusively on the experience and patience of the conservators. While the screening of the fragments is a central prerequisite and the most important stage of the process of vase reconstruction, established methods based on scientific criteria and guaranteed efficiency for the detection of co-belonging ceramic fragments suggested in the bibliography, do not exist. On the contrary, for methods dealing with the assembly of vases from co-belonging fragments, which is a secondary process that can be done more easily and effectively in an empirical way, there exist numerous studies based on fragment morphology. However, even these are also not implemented because of the time requirements, sheer volume and complexity of the proposed methods, in order for them to be applicable in practice. The proposed methods in this paper are based on thermoremanent magnetization (A/m), which is calculated from the weak magnetic field measurements by a fluxgate-sensor/magnet apparatus forming a three-dimensional orthogonal system. Experimental measurements from fragments of 6 vases show that the magnetization magnitude of co-belonging fragments display similar values, despite the magnetic anisotropy of the ceramic material, since these belong to vases that are made of the same clay and fired under the same conditions. This is the criterion for finding ceramic fragments of the same vase from archaeological excavations. The thermoremanent magnetism directionality of fragments, which is aligned along the geomagnetic field at the same place and time during the vase firing process, as it is configured by their rotational symmetry, defines the position of the fragments on the body of the 6 vases. The shape of the original vase can be reconstructed when only a few non adjacent fragments are available. The proposed measurement apparatus can be used for the construction of a useable portable magnetometer specialized for ceramic surface measurements to achieve the above objectives.

Keywords: *thermoremanent magnetization; fluxgate-sensor; selection of co-belonging ceramic fragments; vase reconstruction; vase surface grooves; vases with arbitrary rotational symmetry; archaeometry.*

1. Introduction

While the assembly of excavation fragments from their remaining thermal magnetism is suggested as a theoretically sound method in the bibliography [1], there are only a few older (1975) references [2] regarding corresponding research.

Research on thermoremanent magnetism in the archaeometry field since the 1980s mainly focused on the chronological dating of fired clay walls of kilns. At the same time, as reported in numerous publications [3-32], the reconstruction of vases is attempted almost solely using algorithms for optical data processing based on fragment morphology. Despite the intense scientific interest in resolving the matter using the capabilities of the new digital technology, it remains an exclusively empirical process due to the time requirements and sheer volume and complexity of the proposed methods, in order for them to be applicable in practice.

A central prerequisite before fragment assembly is the selection of the fragments that belong to the same vase from the numerous ceramic findings of an archaeological excavation. This is performed empirically and with great difficulty, especially in excavations conducted over long periods of time or ones with disturbed excavation layers.

However, the criterion for finding excavation fragments of the same vase should be the similarity of the values of the remanent magnetism, since these belong to vases that are made of the same clay, fired under the same conditions. Therefore, the natural way of assembling ceramic shards is not the processing of their imagery but rather the study of the directionality of their oriented thermoremanent magnetism along the geomagnetic field, which happened at the same place and time, and was configured by their rotational symmetry. In the present work, magnetic field measurements were performed using fluxgate sensors. All experiments were conducted in the Experimental Physics Laboratory, School of Electrical and Computer Engineering, Technical University of Crete. Application of the proposed method to six vases from the last two centuries fired in traditional kilns gave excellent results.

After a brief review of the clay magnetic compounds in Appendix A, which is suggested background reading material for non-experts in Mineralogy, we describe the detection methodology of the fragments' field by a fluxgate sensor in Sect. 2. The easy and unexpected assembling of base fragments and neighboring body fragments in 6 vases from their magnetic field directivity by a single fluxgate sensor, despite the irregular shape of fragments and the magnetic anisotropy of the fired clay, was the motivation for further investigation. For the systematic investigation and the comparison of the magnetic field, base and body fragment measurements were taken at the vertices of a lattice formed by the grooves and the vertical lines to them, with a specific orientation of the sensor.

The measurements in the above positions were repeated in an less tedious and more accurately way, by using a sensor/magnet apparatus forming a three-dimensional orthogonal system for the magnetic field, which is described in Sect. 3.

A series of experiments, mainly in base specimens of vases 4, 5 & 6 were carried out (see Appendices B.1-B.5) to investigate how the sensors on the fragments' surface were excited and the region that excites the sensors was identified, for the calculation of thermoremanent magnetism (Appendix B.6) from the sensor readings. We propose the experiments to be read in order to provide a clear focus in their results and the theoretical approach of the sensor excitation area, which were summarised in Sect. 4, and which were used for calculation of magnetization (Appendix B.7) in irregular fragments of the base vase 1,2&3, with the described methodology presented in Sect. 5.

The aggregate results and conclusions of calculating magnetization from magnetic field measurements on samples and fragments of the base of vases 1-6 in Sect. 6, were applied for the calculation of magnetization

(see Supplementary Files) in irregular fragments of the body vase 1-6, with the described methodology in Sect. 7.

In the 8th section, the mathematical relationships of the calculated magnetization changes in the fragments' reference systems in vases with cylindrical (Sect. 8.1) and arbitrary rotational (Sect. 8.2) symmetry are described. By using rotational transformations of magnetization components in the reference systems where field measurements were taken (Sect. 8.3), the location of fragments was identified in the bodies of the 6 vases. Finally, the Sect. 9 contains proposals for future research.

2. Detection of the weak magnetic field of ceramic fragments and its direction using a fluxgate sensor

The most important problem in detecting the weak field of fragments, on the order of a few tens of nT, is its overlay with the Earth magnetic field, which is about 1000 times stronger, as well as the lab magnetic background.

For the detection of the weak magnetic signal [65-67] of fragments within the geomagnetic field, a *fluxgate*¹ sensor [68] at a fixed horizontal position (**figure 1**) is oriented perpendicularly to the vertical line (\mathbf{B}_{vert}) and to the horizontal component (\mathbf{B}_{horz}) of the earth magnetic field ($\mathbf{B}_{\text{earth}}$), such that its reading is zero before the fragment placement. When the sensor distance from the fragment surface is greater than 1.5-2cm, its reading becomes zero.

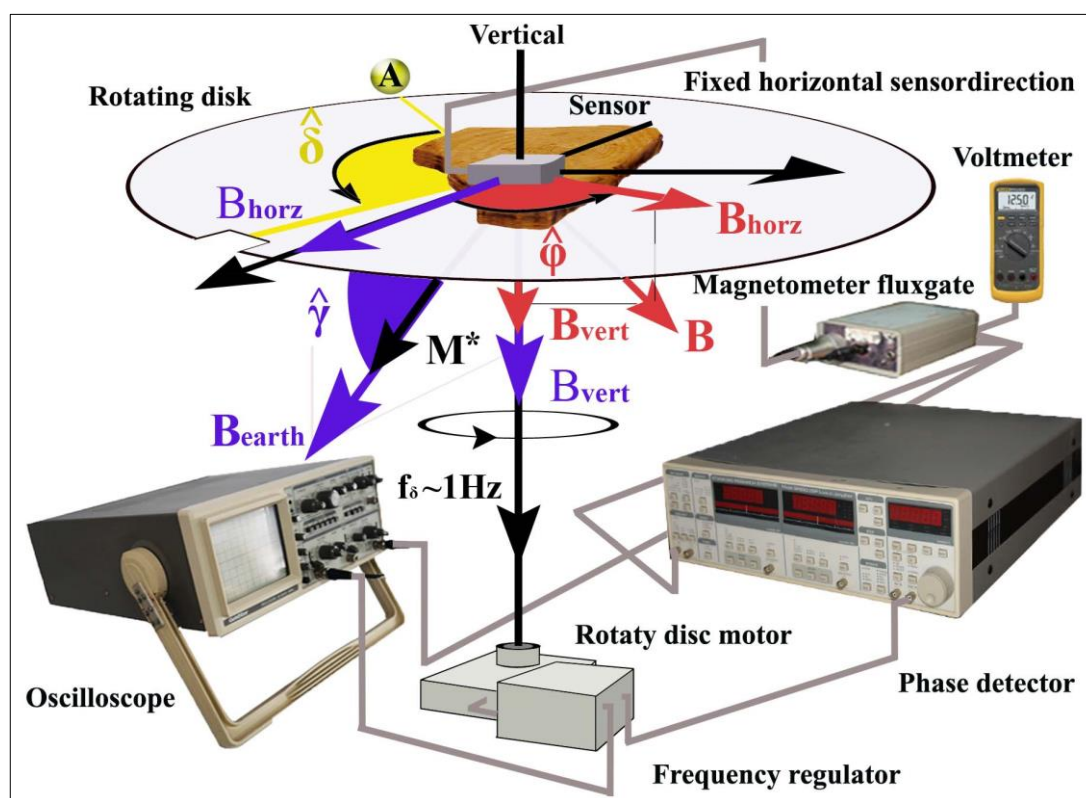


Figure 1. Illustration of the earth magnetic field ($\mathbf{B}_{\text{earth}}$) and of the detected projection of the horizontal component (\mathbf{B}_{horz}) of the fragment field onto the sensor direction during the initial placement of the fragment on the rotating disk. The initial phase angle, δ , of the disk is located between the fixed point, A, of the rotating system and the disk notch. The

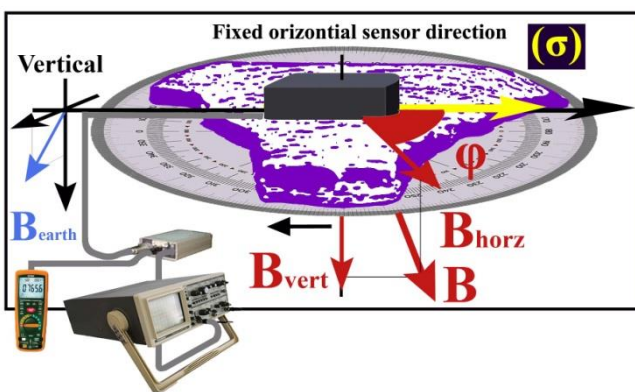
¹ Bartington 3 independent axes fluxgate magnetometer- model MAG 03 IE

initial phase angle, φ , of the magnetic field, \mathbf{B} , is located between the (fixed) vertical sensor direction and the \mathbf{B}_{horz} direction, which depends on the fragment placement on the disk.

- For the weak magnetic field extraction from the unwanted noise of lab magnetic background, the **Phase Sensitive Detection (PSD)** method [69] was used. The electric signal from the disk rotation, at a certain frequency, f , is used as a reference signal to the phase detector [70]. The direct electric current ($1\text{mV}/7\text{nT}$) from each fragment, detected by the magnetometer sensor, is converted during the disk rotation into alternating current and used as an input to the phase detector.

The electric signal from the horizontal rotating component, \mathbf{B}_{horz} , of each fragment is differentiated from the magnetic interferences because it is detected by the phase detector exclusively at the disk rotation frequency, f . If the magnetic field of each fragment displays constant directionality, then its magnitude and initial angular phase difference $\delta-\varphi$, are recorded, after the processing of the sinusoidal signals at the phase detector. During the disk rotation, the ceramic material induced magnetization, M^* , is oriented perpendicularly to the sensor along the direction of $\mathbf{B}_{\text{earth}}$.

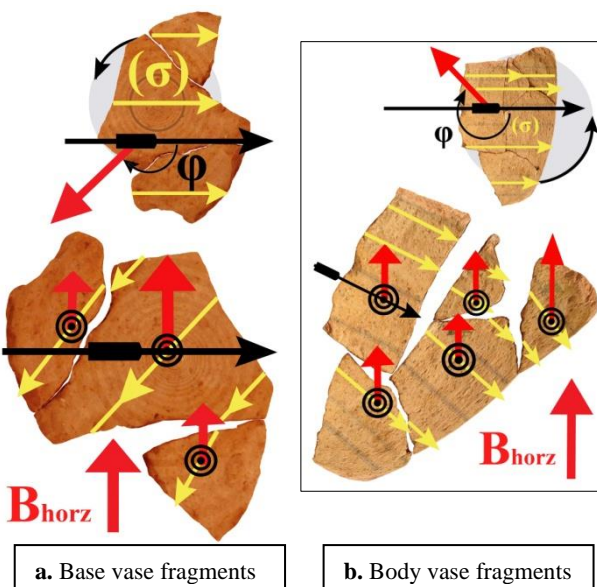
Using the PSD method, constant phase differences and weak magnetic field values with constant directionality of a few tens of nT can be found.



- After the detection of the fragments magnetic field (**figure 2**), the directionality of the base and body in co-belonging fragments is compared. The sensor remains in the same horizontal position, perpendicular to the vertical component and the horizontal component of the geomagnetic field.

Figure 2. Illustration of the horizontal component detection apparatus from vase fragments.

During the vase formation at the pottery wheel, circular grooves remain on the fragment surface. Parallel straight lines at an arbitrary direction (σ) are marked on the base fragments (**figure 3a**), which suggests their joining direction. Regarding the fragments from the vase bodies (**figure 3b**), their direction (σ) is selected along the surface grooves of the pottery wheel.



The fragments are positioned on a leveled angle-measuring apparatus, with the sensor axis aligned along the direction (σ), at a position where \mathbf{B}_{horz} diverges at an arbitrary angle, φ , from the sensor direction, while the disk is rotated manually until its reading is zeroed. At this position, where \mathbf{B}_{horz} is vertical to the sensor direction, its direction is marked.

Figure 3. If the horizontal components of the fragments are positioned such that their vectors are pointing in the same direction, then base fragments and only the neighboring body fragments are oriented such that they join. The distant body fragments are not oriented in a manner that they join.

Neighboring fragments are oriented in a joining manner because their magnetic field, acquired before their fracturing, is due to the thermoremanent magnetization along the direction geomagnetic field during their cooling in the kiln.

While finding the joining orientation of the base excavation fragments is not of essential practical value, because it is more easily performed empirically, it can be concluded from the above that not only does the magnetic anisotropy of the material not prevent their joining from the orientation of their magnetic field but also the irregular fragment shape; in fact this happens with remarkable precision. This is not the case for the distant body fragments, which come in a variety of shapes and are collected in large numbers. This is to be expected, because the thermomagnetic directionality varies depending on the orientation of the vase regions with respect to the geomagnetic field, and it is formed by the rotational symmetry of the ceramic vases.

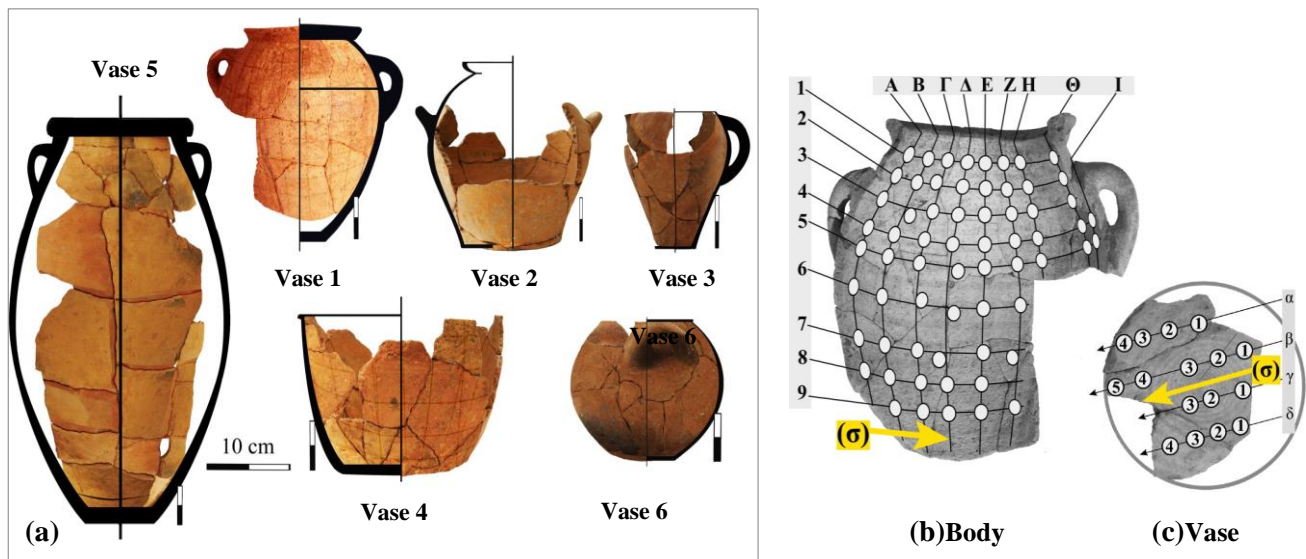


Figure 4. For the systematic study of the magnetic field directionality, single-sensor measurements from base and body fragments from 6 ceramic vases of the previous two centuries, fired in a traditional kiln, were taken (**figure a**).

The body fragments measurements (**figure b**) were taken at the vertices of a lattice formed by the grooves, which provide the reference directions (σ) , and vertical lines to them. The base fragments measurements (**figure c**) are taken on parallel lines, along the arbitrarily chosen joining direction (σ) of the fragments.

The fragments are placed on the horizontal angle-measuring apparatus disk, with the marked reference direction (σ) on their surface matching the fixed sensor direction.

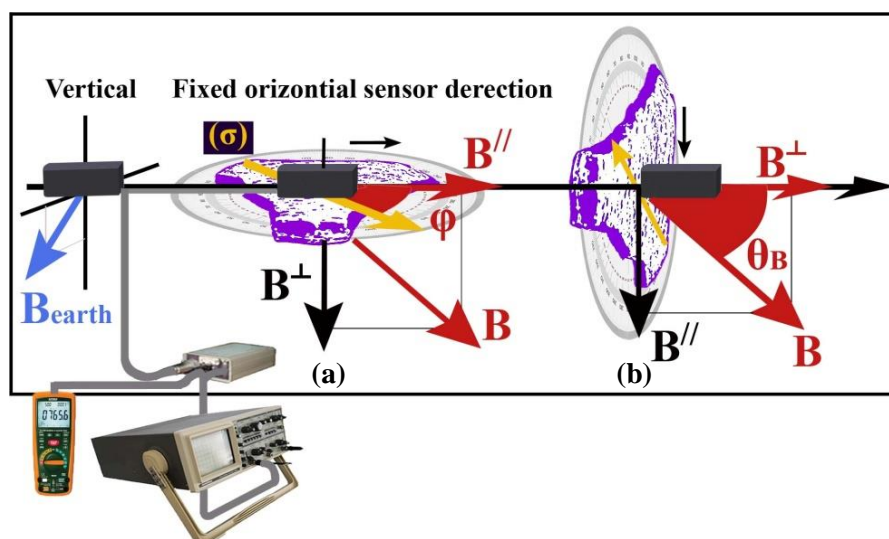


Figure 5. Illustration of the base and body fragment magnetic field measurement apparatus from the 6 vases, using a sensor that remains perpendicular to the vertical component and the horizontal component of the geomagnetic field, B_{earth} .

The component parallel to the fragment surfaces (B^{\parallel}) and the angle of divergence (φ) of the reference direction (σ) with the fixed sensor direction (**figure 5a**) are measured at the rotational position where the reading displays its maximal positive value.

The vertical component (B^{\perp}) and its angle (θ_B) of divergence from the fixed sensor direction (**figure 5b**) are measured at the same turning fragments position, albeit on the vertical disk

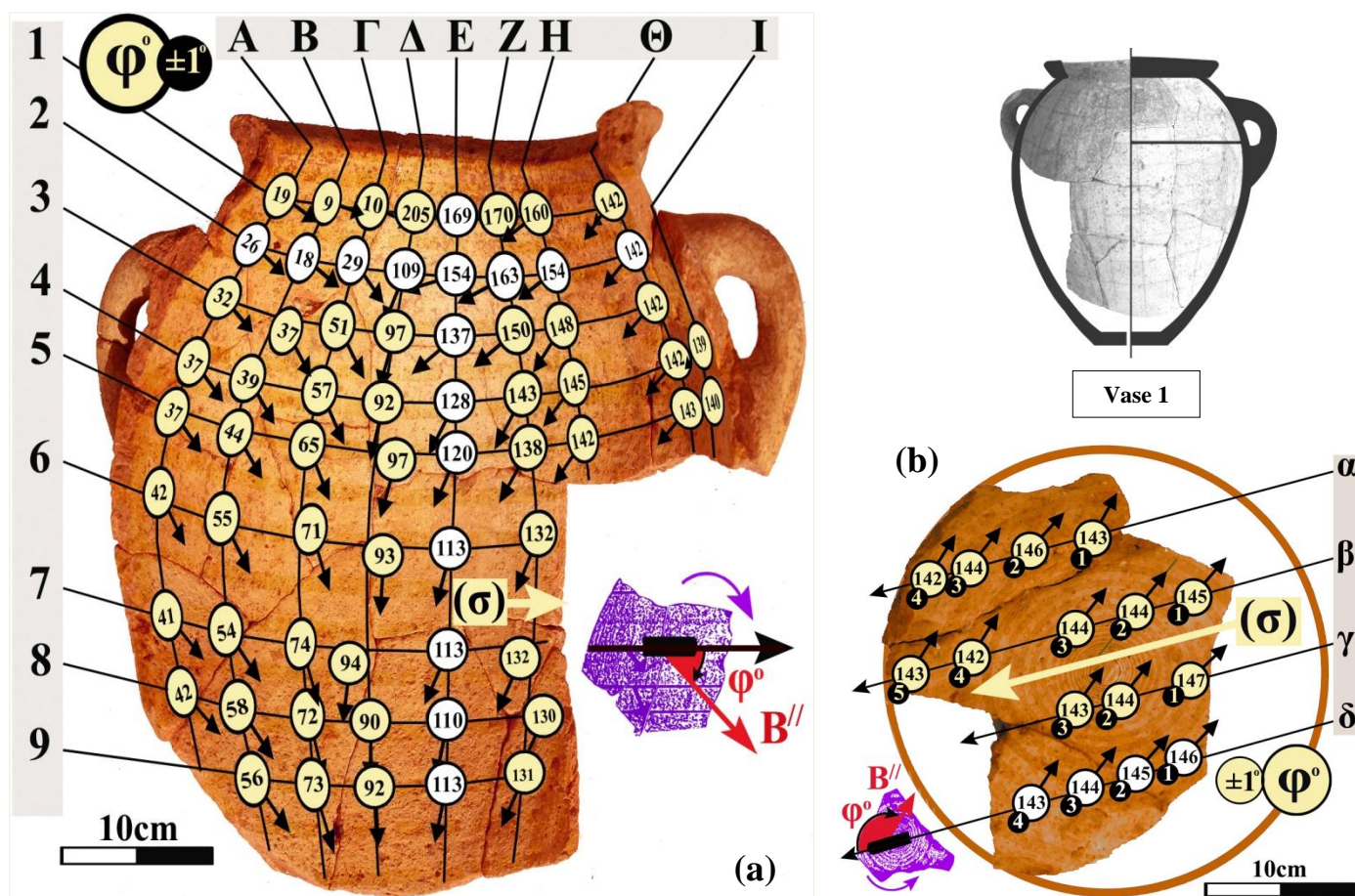


Figure 6. Representative measurements of the angle φ in body (**figure a**) and base fragments (**figure b**) of vase 1.

Diagram 1a (Body-column E)

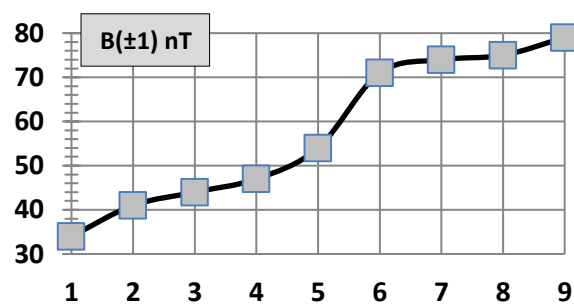


Diagram 1b (Body-line 2)

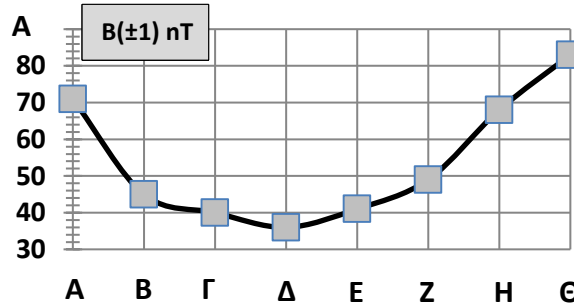


Diagram 1c (Vase-line delta)

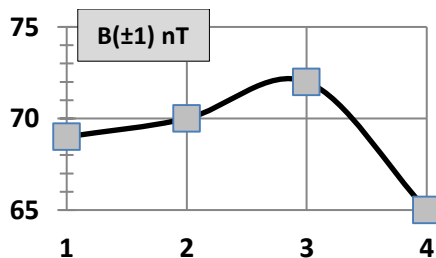


Diagram 1. Representative values of the magnetic field magnitude B computed at body positions (a,b) and base positions (c) of vase 1.

The computed magnitude, B , and the measured angle, ϕ , display similar values only in the co-belonging base fragments of each vase. For the co-belonging body fragments, the magnitude B differs, while the angles ϕ change in a systematic way.

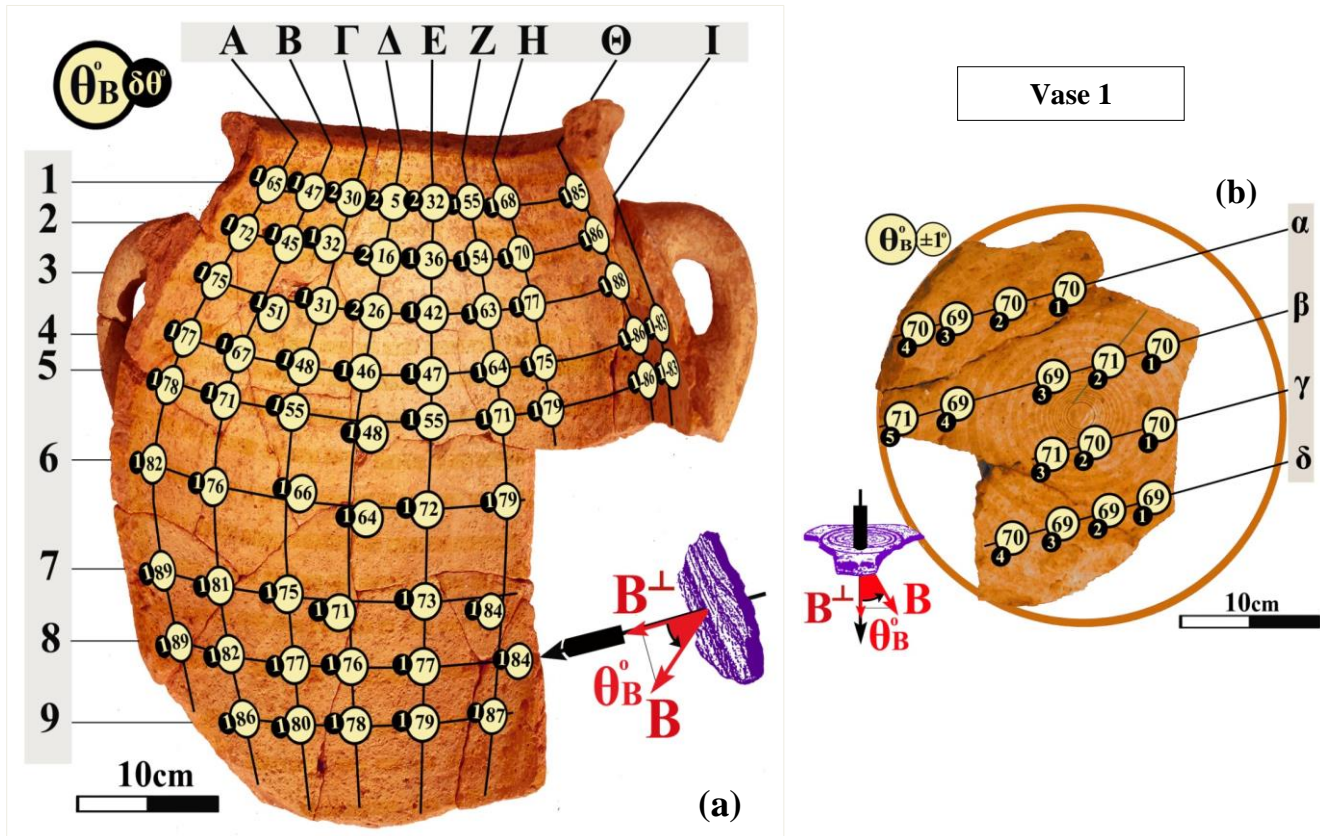


Figure 7. Representative measurements of the angle θ_B in body (figure a) and base fragments (figure b) of vase 1.

The computed angle θ_B displays similar values only in the base fragments of the vases, while it varies in a systematic way for the body fragments.

The measured magnetic field of the co-belonging base fragments by a *fluxgate* sensor shows similar magnitude and directionality. This is the result of thermoremanent magnetization, which despite the ceramic material magnetic anisotropy, maintains similar magnitude and directionality, because the co-belonging base fragments are composed of the same clay, were fired under the same conditions, at the same time and with a fixed orientation with respect to the earth magnetic field. In any case, local differences are not improbable, which can be due to the incomplete mixture homogenization of clays or eutectic additive materials, temperature variations or kiln aeration and later magnetizations.

The method by which measurements are taken, by fragments placed on the disc in horizontal and vertical positions, due to the time requirements and the required accuracy, is difficult to be used for wide application. The dependence of thermoremanent magnetization in vases shaped in pottery wheels was systematically studied using precise and less tedious magnetic field measurements, by the 3 *fluxgate* magnetometer sensors in a orthogonal coordinate system in three dimensions.

3. Magnetic field measurements of ceramic fragments by 3 fluxgate sensors in 3 orthogonal axes

The horizontal x -sensor is oriented vertically to the horizontal component and to the vertical component of the earth magnetic field (**figure 8**), measured by the y and z -sensors, respectively. The sensor readings ($1\text{mV}/7\text{nT}$) are zeroed with a maximum precision 0.01 mV or 0.1nT by adjusting their axial distance from the Al-Ni-Co type cylindrical magnets (400 KA/m), by turning their mounting screws. The excitation area of each sensor is located on the axis that passes through the center of its square cross section of side $2d=6,5\text{ mm}$ and at a distance of $a=6,5\text{ mm}$ from its curved edge.

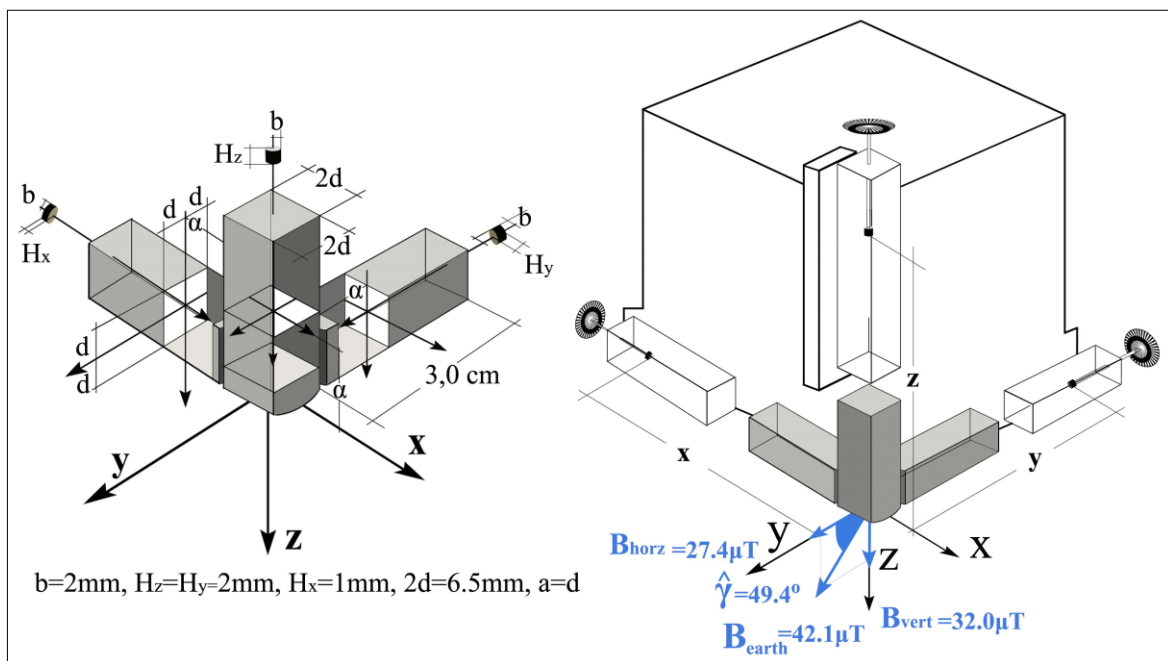


Figure 8. Illustration of the orthogonal coordinate system in three dimensions of the sensor/magnet apparatus for the compensation of the earth magnetic field and the zeroing of their readings.

Each magnet generates, besides its axial field at the sensor direction, also side fields, which excite the perpendicular sensors.

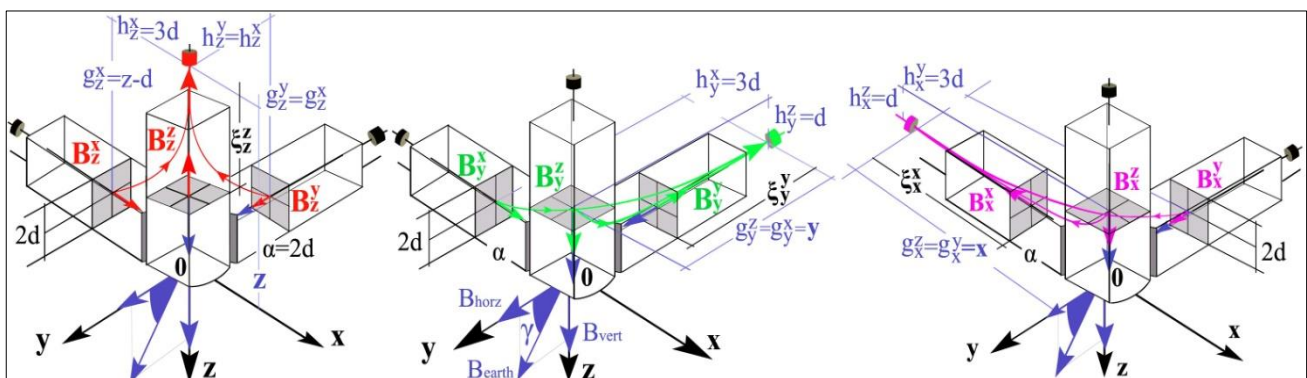


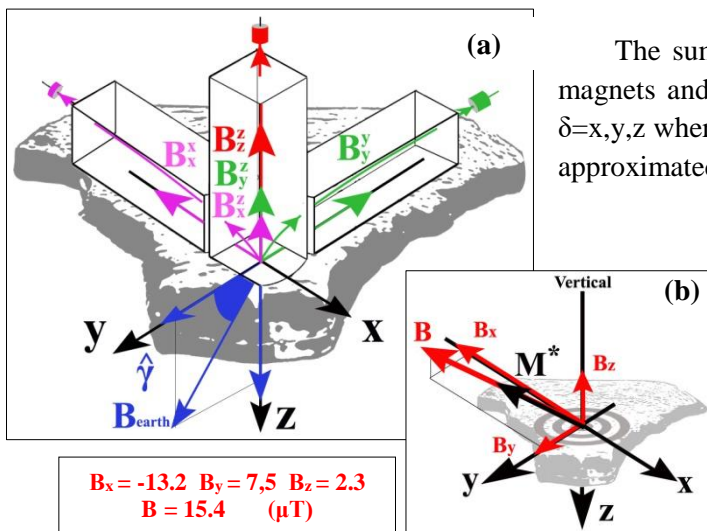
Figure 9. Illustration of the earth field, B_{earth} , and the fields B_{μ}^{α} of the magnets $\mu=x,y,z$ at the $\alpha=x,y,z$ sensors.

The measured axial fields B_z^z, B_y^y, B_x^x (**figure 9**) of the cylindrical magnets, with magnetization M , radius b , and height H_{μ} , with dipole moment $m = M \cdot \pi \cdot b^2 \cdot H_{\mu}$ and at axial distances ξ_{μ}^{α} from the sensors, are approximated (**table 1**) by the formula [71 pp.394-9]:

$$B_{\mu}^{\alpha=\mu} = \frac{\mu_0 \cdot M \cdot b^2 \cdot H_{\mu}}{2\xi_{\mu}^3}, (\mu=x,y,z) \quad (1)$$

The magnet side fields, $B_{\mu}^{\alpha \neq \mu}$, at axial distances $g_{\mu}^{\alpha \neq \mu}$ and side distances from the sensors, $h_{\mu}^{\alpha \neq \mu}$, are approximated by the formula:

$$B_{\mu}^{\alpha \neq \mu} = \frac{3\mu_0 \cdot M \cdot b^2 \cdot H_{\mu}}{4} \cdot \frac{g_{\mu}^{\alpha} \cdot h_{\mu}^{\alpha}}{[(g_{\mu}^{\alpha})^2 + (h_{\mu}^{\alpha})^2]^{5/2}}, (\mu, \alpha=x,y,z) \quad (2)$$



The summed components of the field B_{μ}^{δ} of the $\mu=x,y,z$ magnets and the earth field components at the axial distances $\delta=x,y,z$ where the sensor readings become zero (**figure 12**), are approximated [22], at the measurement area, by the formulas:

$$B_z^z = -\frac{\mu_0 \cdot M \cdot b^2 \cdot H_z}{2z^3} \quad (3a)$$

$$B_{\mu=x,y}^{\delta=\mu} = -\frac{\mu_0 \cdot M \cdot b^2 \cdot H_{\mu}}{4} \cdot \frac{2\delta^2 - d^2}{(\delta^2 + d^2)^{5/2}} \quad (3b)$$

$$B_{\mu=x,y}^{\delta=z} = -\frac{3\mu_0 \cdot M \cdot b^2 \cdot H_{\mu}}{4} \cdot \frac{\mu \cdot d}{(\mu^2 + d^2)^{5/2}} \quad (3c)$$

Figure 10. Illustration of the individual fields, B_{μ}^{δ} of the $\mu=x,y,z$ magnets and the earth's (**figure a**) and those of the total field (**figure b**) inside the measurement area, whose direction the induced magnetization M^* of the ceramic material is oriented.

Table 1. Computation of the fields at the sensors at the distances x,y,z that zero their readings. The magnetic field is computed at those distances, inside the measurement area.

Magnets		Fields (nT) at the sensors			Fields (μT) inside the measurement area			
Distances (mm)		z	y	X	z	y	x	
Z	z=44.91841	-61830.98			-22.2			
			60880.61					
				60880.61				
Y	ψ= 45.63295	17197.76			-4.3	-19.9		
			-112658.65					
				35669.95				
x	x= 41.33643	12633.18			-3.2			
			24378.01					
				-96550.59			-13.2	
Earth's field		32000			32.0	27.4		
Total fields		0.0	0.0	0.0	2,3	7,5	-13,2	

All the following measurements are taken at the same fixed position with respect to the geomagnetic field by the sensor/magnet apparatus forming a three-dimensional orthogonal system. The induced-magnetization contribution of the ceramic material in the sensor readings is checked in the following experiment by comparing the measurements (Annex B.2) taken inside a null magnetic field.

4. Experimental results: Theoretical approach of the sensors' excitation area for the calculation of magnetization by magnetic field measurements in base vase specimens

The remanent magnetization dependence and computation on the magnetic sensor measurements, is studied in Appendix B through a series of experiments (Appendix B.1-B.5) in irregular base fragments (Appendix B.1) and mainly in cylindrical specimens (Appendix B.2-B.5) from the base of vases 4-6.

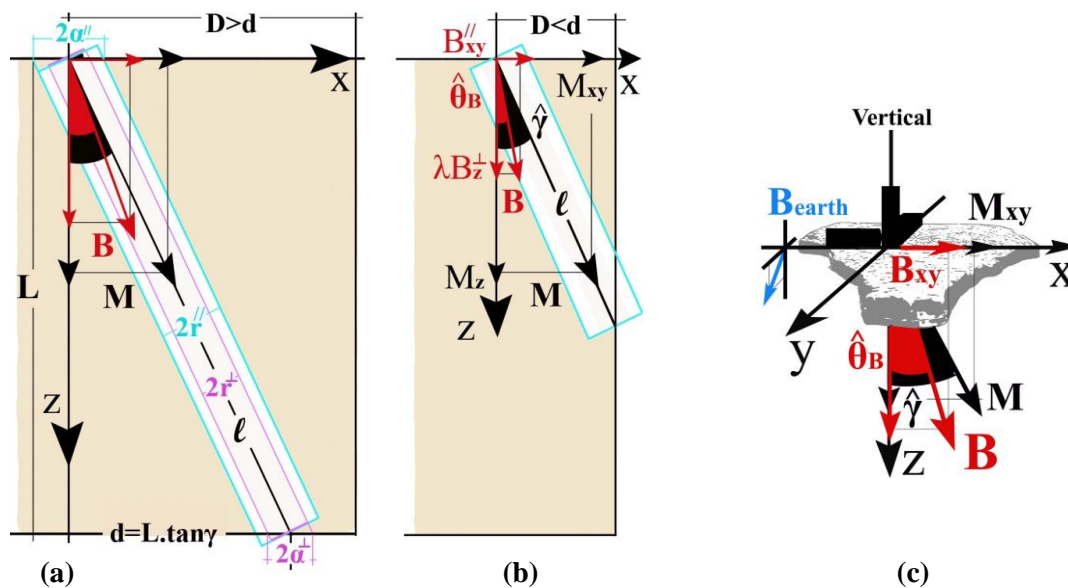
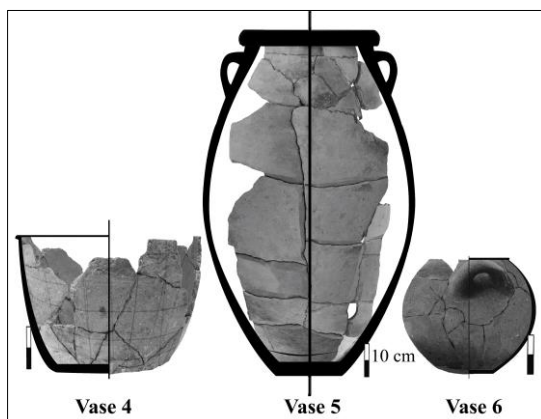


Figure 11. Illustration of the cylindrical excitation area dimensions of the vertical and parallel sensors in the ceramic material (figure a,b), depending on the distance, D , between the sensors position and the edge of the specimen. The x -sensor (figure c), remains perpendicular to the vertical and the horizontal components of the geomagnetic field. The measurements are taken along B_{xy} , which is oriented in the direction of the x -sensor. The angles, θ_B , $\gamma > \theta_B$, between the field, B , and the magnetization, M , increases as the distance to the measurement position from the edge of the fragment decreases.

From the experimental results of the measurements on specimens and fragments from the base vases 4-6, (figure 11) it is found that:



- The sensors are excited by the axial field B at the edge of a cylindrical region of ceramic material (appendices B.1-B.2, B.5) in the direction of the remanent magnetization, M .
- The measured magnetic field is only due to the remanent magnetization of the ceramic material (appendix B.2), and it is not altered by its induced magnetization.
- The height, ℓ , of the sensors' cylindrical area (appendices B.2-B.5, B.7) depends of the angle, γ , between the magnetization M and the vertical z -axis on specimens' surface, each specimen thickness, L , and the distance, D ,

between the sensors and the fragment edge along B_{xy} , when the vertical z -sensor reading is doubled.

When $D > d = L \cdot \tan \theta_B$ the height, $\ell = L / \cos \theta_B$, depends on the thickness, L , of the fragment.

When $D < d$ the height, $\ell = D / \sin \theta_B$ is dependent on the distance D .

- The radii (appendix B.5), $r^i = 2\alpha^i \cdot \cos \gamma$ ($i=, //$), of the excitation areas depends on the directionality (γ) of magnetization and of the parameters $\alpha^i, \alpha^{//}$, which, geometrically, are the sensitivity radius around measurement position of the vertical and the parallel sensors, respectively.

The field flow in the excitation coil of the vertical sensor B_i^{\perp} ($i=xy,z$), is smaller than that of the parallel ($B_i^{\parallel} \approx \lambda \cdot B_i^{\perp}$) sensor, for the measurement of the same component B_i of the field. The normalization factor $\lambda \approx 2$, for the correction of their readings (appendix B.5), has a similar value in non co-belonging fragments with different magnetization M due to the different orientation of the excitation coils on the specimen surface. The radii $\alpha^{\perp}, \alpha^{\parallel}$ increase with the magnetization magnitude, M , increase and display similar values in co-belonging shells.

- The sensor readings displays similar values when the ratios r/ℓ , of the cylindrical sensor excitation areas do not change, as does the measured axial field at the end of the solenoids with the same current and number of spirals, where the ratio of the radius, r , to their height, ℓ , is kept constant. The measured magnetic field B of the cylindrical sensors excitation areas is approximated by the magnetic field of a solenoid with the same dimensions.

- If $D \geq d = L \cdot \tan \gamma$ (**figure 11a**), the height, $\ell = L / \cos \gamma$, depends on the thickness, L , of the fragment and the components, M_i ($i=xy,z$), are approximated by the *sufficient length relationships*:

$$M_i = \frac{2B_i}{\mu_0} \cdot \sqrt{1 + \left(\frac{\alpha_i^{\perp//} \cdot \cos^2 \gamma}{L} \right)^2} \quad (4)$$

- If $D \leq d$ (**figure 11b**), the height $\ell = D / \sin \theta$ depends on the distance, D , of the measurement position from the fragment edge in the magnetization direction, M . The components M_i ($i=xy,z$) are approximated by

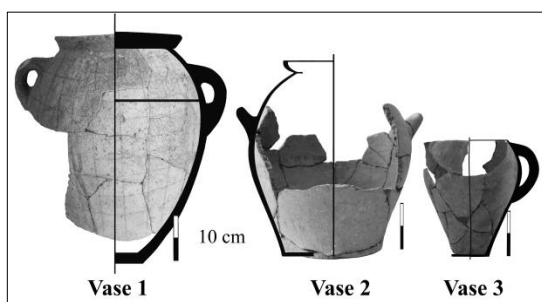
the *sufficient length relationships*:
$$M_i = \frac{2B_i}{\mu_0} \cdot \sqrt{1 + \left(\frac{\alpha_i^{\perp//} \cdot \sin^2 2\gamma}{2D} \right)^2} \quad (5)$$

Since $\alpha_{xy}^{\parallel} > \alpha_z^{\perp}$, the computed deviation, θ_B , of the field from the vertical (**figure 11c**) at the measurement is less than the angle γ of the magnetization M .

The horizontal components, B_{xy} , M_{xy} , of the measured field, B , and magnetization, M , diverge at the same angle φ from the horizontal direction of the x -sensor.

The quantities λ , α^{\perp} , α^{\parallel} , γ , M , M_i , ($i=xy,z$) were calculated with remarkable accuracy using the above relationships from the sensor readings and also by the thickness, L , and distance, D , measurements (appendix B.6) in base specimens of vases 4 and 6.

5. Computation of magnetization in irregular fragments of the vase base 1, 2 & 3



The components, M_z , M_{xy} , angle γ of the magnetization M , and sensitivity radii of the sensors, α^{\perp} , α^{\parallel} , in fragments from the bases of vases 1, 2 and 3 are computed by the same using way (appendix B.7) in specimens from base of vases 4,5 & 6, **without any need for cutting specimens**.

The three-dimensional orthogonal system apparatus of sensors/magnets remains in the same position with respect to the earth magnetic field (**figure 12**), with the x -sensor oriented perpendicularly ($B_y=0$, $B_x=B_{xy}>0$) to the horizontal and vertical component of the geomagnetic field.

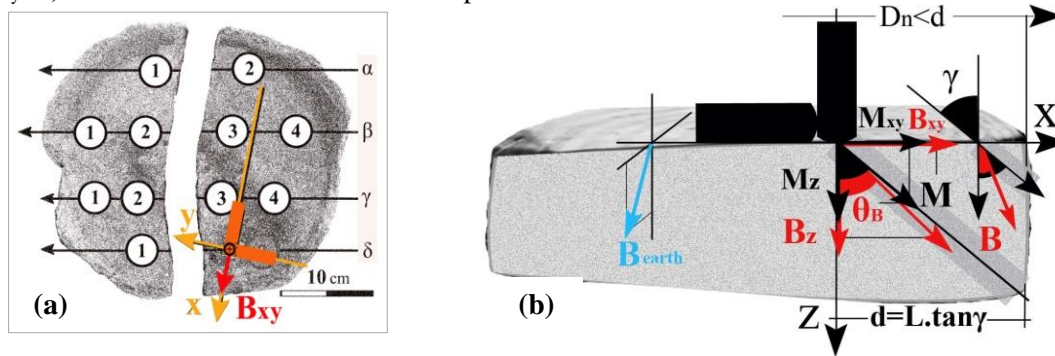


Figure 12. Illustration of the measurement methodology (b) of the magnetization components, M_i ($i=xy,z$), and the sensitivity radii of the vertical, α_z^{\perp} , and parallel sensor, α_{xy}^{\parallel} , in a base fragment (a) of vase 2.

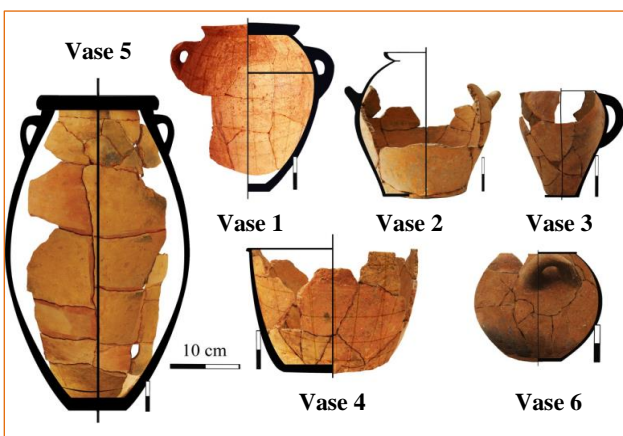
The measurements of B_{xy} , B_z are taken on B_{xy} in the direction of the x -sensor (**table 11**), from the measurement position where the systematic reduction in the magnetic field readings begins (**diagram 8**) at measured distances $D_n < d = L \cdot \tan \gamma$ from the fragments' edge.

The computation of sensor magnetization components, M_i ($i=xy,z$), and sensitivity radii, $\alpha_z^{\perp}, \alpha_{xy}^{\parallel}$, from the field measurements, B_i ($i=xy,z$), is performed using the least squares method (**annex A.7**), from the equation:

$$\left(\frac{\mu_0}{2B_i} \right)^2 = \frac{\alpha_i^2 \cdot \sin^2 2\gamma}{4 \cdot M_i^2} \cdot \frac{1}{D_n^2} + \frac{1}{M_i^2}, i=z,xy \quad (6)$$

The measurements' results are presented in appendix B.7.

6. Aggregate results and conclusions of calculating magnetization from magnetic field measurements on samples and fragments of the base of vases 1-6



The measured components of the magnetic field, B_i ($i=xy,z$), depend on the directionality (γ), magnetization magnitude M , and thickness, as well as the position and orientation of the sensors on the surface of the fragments. In the above perspective, the normalization constants, λ , α^{\perp} , α^{\parallel} , of the field flux, B_i , in the excitation coils of the vertical and parallel sensors and the magnetization (M_i , M , γ) in irregular base fragments of the vases 1, 2 & 3 (**annex B.6**) and in base samples of vases 4,5 & 6 (**annex B.7**) are computed.

Table 2. Computation of the magnitude, M , of the magnetization directionality (γ) and of the normalization constants, λ , α^{\perp} , α^{\parallel} , of the sensor readings in irregular base fragments of the vases 1,2 & 3 and in base specimens of the vases 4,5 & 6.

Vase	M(mA/m)	γ°	α^{\perp} (cm)	α^{\parallel} (cm)	λ
6	58.3±0.2	43.3±0.3	0,71 ±0.03	1.07±0.02	1.90±0.01
4	84.1±0.2	48.5±0.3	0.67±0.03	1.15±0.02	1.97±0.01
5	112.7±0.3	61.3±0.3	1.04±0.03	1.41±0.03	2.03±0.01

1	131.3±0.2	53.8±0.1	0.82±0.04	1.42±0.02	2.0±0.1
3	164.1±0.4	46.6±0.2	1.07±0.05	1.61±0.03	2.0±0.1
2	169.3±0.6	50.8±0.2	1.15±0.05	1.66±0.05	2.0±0.1

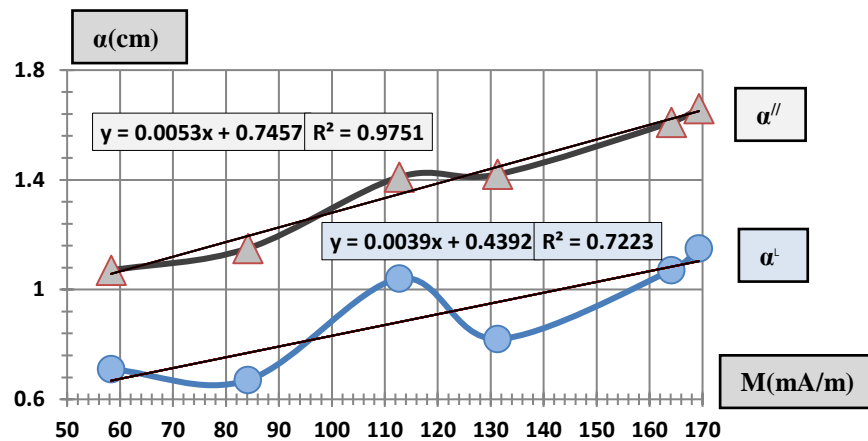
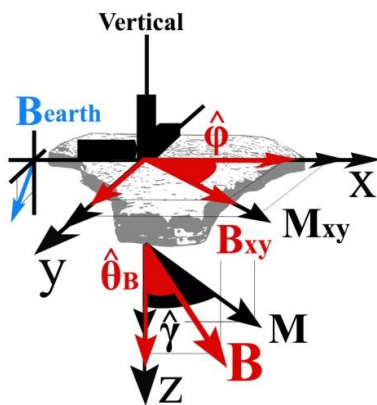


Diagram 2. Change of the sensitivity radii of the vertical α' and parallel sensor as a function of the computed magnetization magnitude, M , in the base fragments from the vases 1-6. While the number of points in the graph is not sufficient, their change appears linear, with a similar slope.

The sensitivity radii, α' , α'' , increase with the increase of the magnetization magnitude (**diagram 2**) and differ (**table 2**) in base fragments from different vases.



Since $\alpha_{xy''} > \alpha_z'$, according to the relationships (4,5), the computed deviation, θ_B , of the field from the vertical (**figure 13**) at the measurement is less than the angle γ of the magnetization M .

Since $\alpha_{xy''} = \alpha_{y''}$, the horizontal components, B_{xy} , M_{xy} , of the measured field, B , and magnetization, M , diverge at the same angle ϕ from the horizontal direction of the x -sensor. For this reason, the co-belonging base fragments or the adjacent body fragments of the vases are oriented in such a way that they join (sect. 2), when the B_{xy} are placed in parallel directions.

Figure 13. The magnetization, M , and the magnetic field, B , diverge at the same angle, ϕ , from the horizontal x -axis, while the angles, θ_B , $\gamma > \theta_B$, between the field, B , and the magnetization, M , deviate as the distance to the measurement position from the edge of the fragment decreases.

With the proposed perspective of the sensors' mode, the area of excitation and the normalization constants (λ , α' , α'') of their readings from measurements of the magnetic field in base vase fragments, the magnetization of co-belonging body fragments of vases 1-6 is computed.

7. Computation of magnetization in body fragments from vases 1-6

Measurements by the three-dimensional orthogonal system of sensors/magnet apparatus are taken in the same position, with the x -sensor oriented perpendicularly to the horizontal and vertical component of the geomagnetic field. The measurement positions are situated at the edges of a mesh formed from lines vertical to the traces of the grooves from the formation of the 6 vases (**figure 4b**) on a pottery wheel. **To exploit the rotational symmetry of vases (figure 14a), the measurements are taken with respect to a right-hand coordinate system, with the x -axis along the groove direction, the y -axis perpendicular to the surface and pointing towards the inside of the fragments, and the z -axis pointing towards the base of the vases.** The components M_i ($i=x,y,z$) are computed from the measurements B_i of the field, in the rotating position where the grooves are oriented in the direction of the x -sensor.

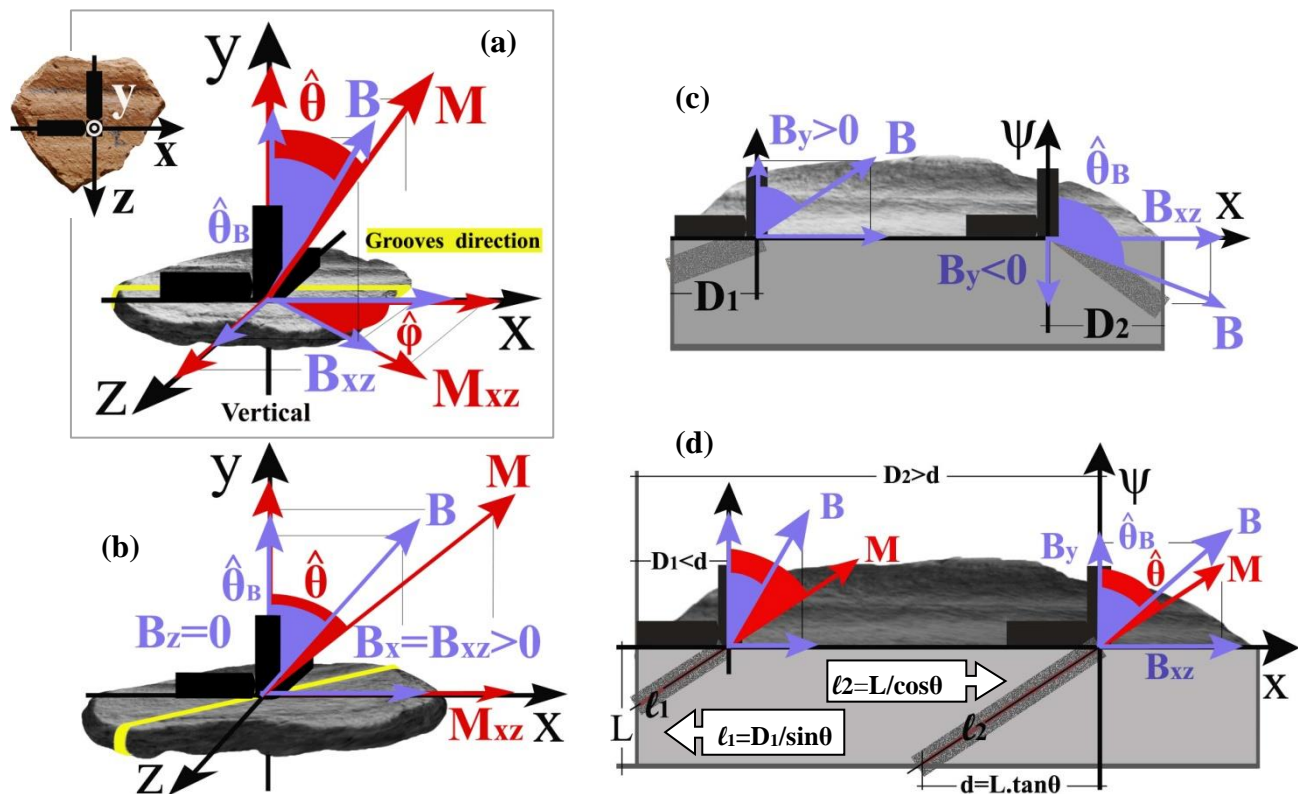


Figure 14. Methodology for calculating the components of the magnetization M from the measured components of the magnetic field B in the considered reference system.

The components of the measured field, B , and of the magnetization, M , in the considered coordinate system (**figure 14a**) are computed by the relations:

$$B_x = B \cdot \sin\theta_B \cdot \cos\phi, \quad B_y = B \cdot \sin\theta_B \cdot \sin\phi, \quad B_z = B \cdot \cos\theta_B, \quad M_x = M \cdot \sin\theta \cdot \cos\phi, \quad M_y = M \cdot \sin\theta \cdot \sin\phi, \quad M_z = M \cdot \cos\theta$$

To determine the deviation, θ , of the magnetization from the vertical (**figure 14b**), measurements of the B_{xz} , B_y are taken, by rotating the fragments at an angle $\phi_B = \phi$ and orienting B_{xz} in the x -direction, in the position where $B_z = 0$ and $B_x > 0$.

The distance, D , between the measurement position and the fragment edge in the common direction of the M_{xz} , B_{xz} (**Figure 14c**) is measured in the negative or positive x -semi-axis, when $B_y > 0$ or $B_y < 0$.

The distance D (**figure 14d**) may be sufficient ($D > d = L \cdot |\tan\theta|$) or insufficient ($D < d$), for the measured fragment thickness, L , in the measurement position. In any case, the magnetization angles $\theta = \theta_D$ or $\theta = \theta_L$ are approximated [72] by the solution of the 3rd degree equations (using the *Newton-Raphson* method):

$$D > d: \frac{B_{xz}}{B_y} = \tan\theta_B = \tan\theta_L \cdot \sqrt{\frac{1 + \left(\frac{\alpha_y \cdot \cos^2\theta_L}{L}\right)^2}{1 + \left(\frac{\alpha_{xz} \cdot \cos^2\theta_L}{L}\right)^2}} \quad D < d: \frac{B_{xz}}{B_y} = \tan\theta_B = \tan\theta_D \cdot \sqrt{\frac{1 + \left(\frac{\alpha_y \cdot \sin 2\theta_D}{2D}\right)^2}{1 + \left(\frac{\alpha_{xz} \cdot \sin 2\theta_D}{2D}\right)^2}} \quad (7)$$

The constants $\alpha_{xz}^{\parallel}, \alpha_y^{\perp}, \lambda$ are computed (**table 2**) from the field measurements in co-belonging base fragments of the vases.

- If $\theta_L > \theta_D$, then $D > d = L \cdot \tan\theta_L$ and $\theta = \theta_L$. For the computation of the components of the magnetization, M, the sufficient length equations (4) are used.

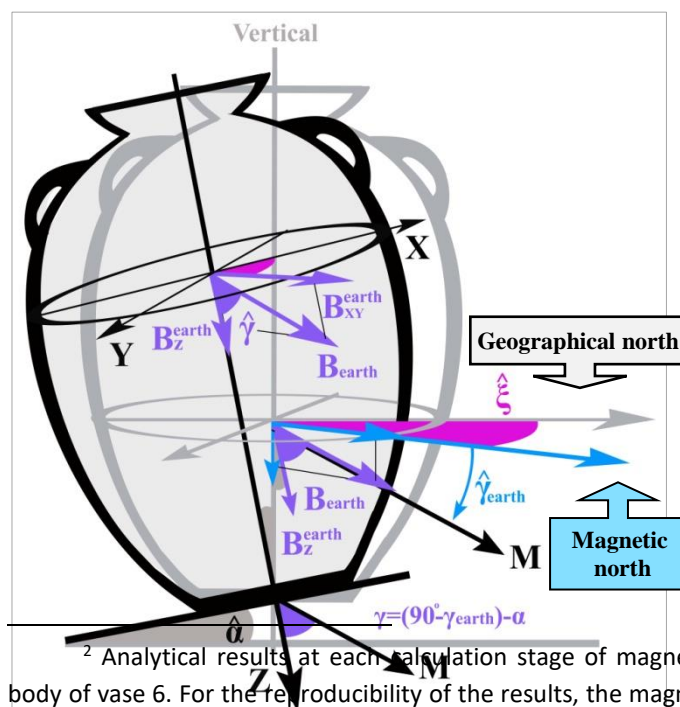
- If $\theta_L < \theta_D$, then $D < d = L \cdot \tan\theta_D$ and $\theta = \theta_D$. For the computation of the components of the magnetization, M, the sufficient length equations (5) are used.

Experimental results² show that the computed magnitude, M, **displays similar values in the base and body fragments, which is a criterion for finding ceramic fragments of the same vase from archaeological excavations.** The angles, θ, ϕ , change in a systematic way because their values are dictated by the rotational symmetry of the vase. This is confirmed in the next section, where the position of fragments in the bodies of 6 vases is identified by successive transformations of magnetization components.

8. Position localization of fragments in the body of a vase from the directivity of the remanent magnetization

If γ_{earth} is the inclination and ξ the declination of the local magnetic field of the earth (B_{earth}) during pottery firing (**figure 15**), then its components with respect to the right-hand axial coordinate system, XYZ, of the vase, where the Z-axis of rotational symmetry of the vase deviates from the vertical at α deviation angle, α , of the vase on the kiln floor and points towards its base, are computed by the equations:

$$B_X^{\text{earth}} = B_{\text{earth}} \cdot \sin\gamma \cdot \cos\xi, \quad B_Y^{\text{earth}} = B_{\text{earth}} \cdot \sin\gamma \cdot \sin\xi, \quad B_Z^{\text{earth}} = B_{\text{earth}} \cdot \cos\gamma, \quad \gamma = (90^\circ - \gamma_{\text{earth}}) \pm \alpha \quad (8)$$



The declination, ξ , of the geomagnetic field cannot be computed without knowing the vase orientation in the kiln.

The inclination, γ_{earth} , of the geomagnetic field cannot be computed without knowing the deviation angle, α , of the kiln with respect to the horizontal plane.

The circular grooves in the vase body, which are parallel to its base, deviate from the horizontal plane according to the same angle, α , due to the rotational symmetry of the base. In any vase with rotational symmetry, and for any considered axial reference system where the z-axis coincides with the axis of rotation, the X,Y axes are defined on

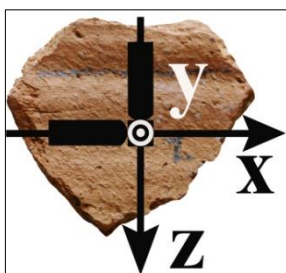
² Analytical results at each calculation stage of magnetization are presented in "Supplementary Files" for the body of vase 6. For the reproducibility of the results, the magnetic field measurements for the remaining vases in each position of their body are listed therein.

the plane of the same transverse groove on the vessel body.

Figure 15. Illustration of the geomagnetic field, B_{earth} , during firing of a vase with rotational symmetry with respect to the right-hand axial reference system XYZ where the Z -axis coincides with the axis of rotational symmetry and points towards the vase base.

The angle, γ , of inclination of the geomagnetic field from the vase axis of symmetry, is equal to the deviation angle of magnetization, \mathbf{M} , calculated from measurements of the magnetic field of base fragments (appendix B.7,8) since the ceramic material magnetization obtains the orientation of the geomagnetic field during firing.

8.1 Formation of remanent magnetization in vases with cylindrical symmetry



Measurements of the field, \mathbf{B} , are obtained in right-hand reference systems $(xyz)^*$, where the x^* -axis has the direction of the transverse grooves, the y^* -axis passes through the axis of rotational symmetry and the z^* -axis is tangential to the side walls, directed towards the vase base.

In the specific case of vases with cylindrical symmetry, because the z^* -axis in the measurement system of the \mathbf{B} field is parallel to the axis of symmetry, Z , (**figure 16a**) the position of each longitudinal axial section is defined by the rotation angle, u , of the right-hand axial reference system, XYZ , of the cylindrical vase (**figure 16b**), around the Z -axis of rotational symmetry.

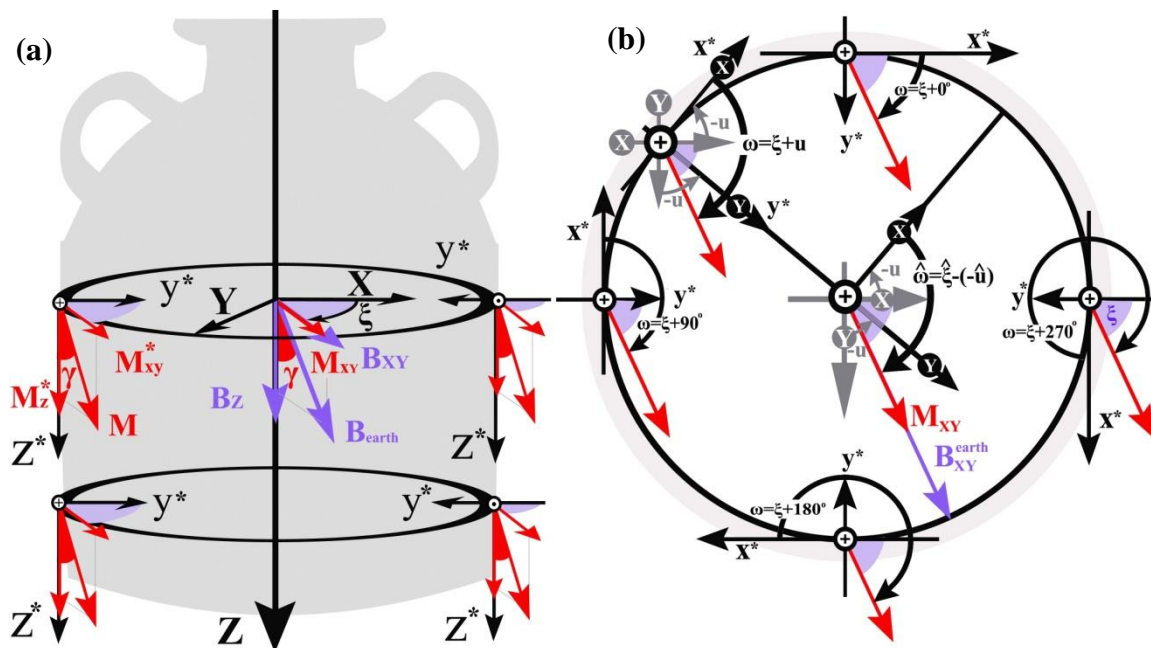


Figure 16. Illustration of the magnetization, \mathbf{M} , of ceramic material along the direction of the geomagnetic field, B_{earth} , during firing of a vase with rotational symmetry in right-hand reference systems, $(xyz)^*$, along longitudinal (**figure a**) and transverse cross-sections (**figure b**) of the vase body.

Since magnetization, M , obtains the orientation of B_{earth} in the pottery kiln, then for every rotation angle, u , of the axial reference system, XYZ , the M -magnetization components at the fringe of each transverse groove are determined by the rotation transformation of the peripheral reference systems $(xyz)^*$ at an angle $\omega = \xi - u$ around the z^* -axis:

$$M_x^* = M \cdot \sin\gamma \cdot \cos\omega = M_{xy}^* \cdot \cos\omega, \quad M_y^* = M \cdot \sin\gamma \cdot \sin\omega = M_{xy}^* \cdot \sin\omega, \quad M_z^* = M \cdot \cos\gamma \quad (9)$$

where: $\omega = \xi - u$, $\gamma = (90^\circ - \gamma_{earth}) \pm \alpha$

Although the angles u , ξ cannot be determined without knowing the orientation of the vase in the pottery kiln, the angle $\omega = \xi - u$ is determined by calculating the magnetization of the fragments of the vase body.

8.2 Formulation of remanent magnetization in vases with arbitrary rotational symmetry

In the general case of vases with arbitrary rotational symmetry (figure 18a), the side walls diverge at different angles, κ , from the axis of symmetry in each longitudinal section, $(\alpha, \beta, \gamma, \delta)$, of the vase and at the same angle κ in each transverse groove, (1,2,3).

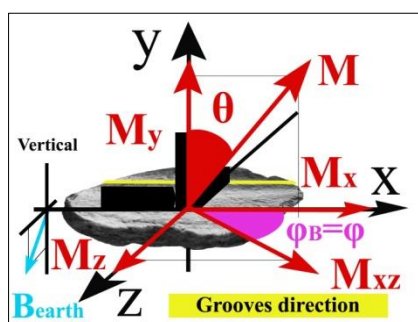


Figure 17. Illustration of magnetization components in the body fragment reference system of a vase whose magnetic field measurements are taken, and computed by the equations: $M_x = M \cdot \sin\theta \cdot \cos\phi$, $M_y = M \cdot \sin\theta \cdot \sin\phi$, $M_z = M \cdot \cos\theta$.

The magnetization, M , and the magnetic field, B , diverge at the same angle, $\phi = \phi_B$, from the horizontal x -axis

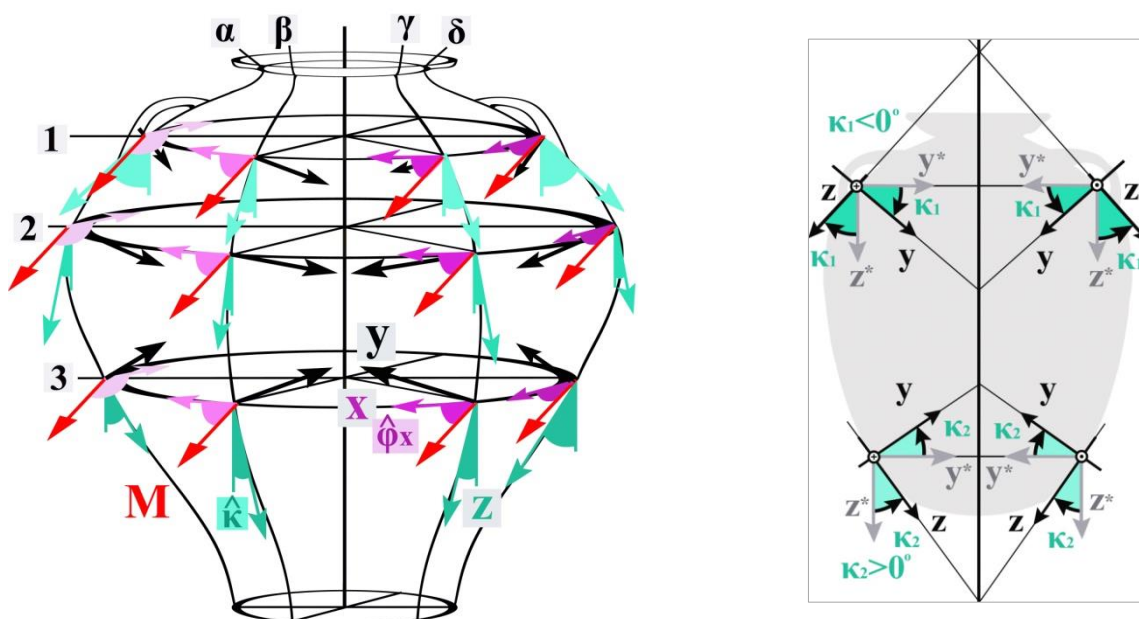


Figure 18a. Illustration of the reference systems (xyz) where the magnetic field measurements are taken, in transverse and longitudinal sections of a vase body with arbitrary rotational symmetry. The angles, ϕ_x , of magnetization, M , with the x -axes oriented in the direction of the grooves maintain a constant value along in each longitudinal section (1,2,3) and vary along in each cross-section $(\alpha, \beta, \gamma, \delta)$ of the vase. The angles, κ , between the vertical z -axes in the grooves and the axis of rotational symmetry maintain a constant value along in each transverse groove and vary along in each longitudinal section $(\alpha, \beta, \gamma, \delta)$ of the vase. The vertical y -axes on the surface of the fragments pass through the axis of rotational symmetry of the vase.

Figure 18b. The magnetization components in the reference systems (xyz) where the magnetic field measurements in vase fragments with arbitrary rotational symmetry are taken are derived from a rotational transformation of the reference systems (xyz)* of the cylindrical vase at an angle of inclination κ around the x^* -axes.

The components M_i ($i=x,y,z$) in the reference systems (xyz) of the B_i field measurements (**figure 17**) are derived from a rotational transformation (**figure 18b**) of the reference systems (xyz)* of the cylindrical vase at an inclination angle κ around the x^* -axes:

$$M_x = M_x^* = M \cdot \sin\gamma \cdot \cos\omega = M \cdot \sin\theta \cdot \cos\varphi = M \cdot \cos\varphi_x,$$

$$M_y = M_y^* \cdot \cos\kappa - M_z^* \cdot \sin\kappa = M \cdot \sin\gamma \cdot \sin\omega \cdot \cos\kappa - M \cdot \cos\gamma \cdot \sin\kappa = M \cdot \cos\theta \quad (10)$$

$$M_z = M_y^* \cdot \sin\kappa + M_z^* \cdot \cos\kappa = M \cdot \sin\gamma \cdot \sin\omega \cdot \sin\kappa + M \cdot \cos\gamma \cdot \cos\kappa = M \cdot \sin\theta \cdot \cos\varphi, \quad \omega = \xi - u \quad \kappa \alpha \quad \gamma = (90^\circ - \gamma_{\text{earth}}) \pm \alpha$$

❖ Based on the above transformation, the circumferential position of the fragments in each cross section of the vase is determined by the angle ω , which is computed by the formula: $\omega = \cos^{-1} \frac{\sin\theta \cdot \cos\varphi}{\sin\gamma}$ (11)

The angles $\tan\varphi = \frac{M_z}{M_x} = \frac{B_z}{B_x}$, $\cos\theta = \frac{M_y}{M}$ are calculated from the magnetic field measurements of the body fragments, while the angle γ is calculated from the magnetic field measurements in co-belonging fragments of the vase base.

❖ The height-position of the fragments in each longitudinal section of the vase is determined by the inclination κ , which is computed by the equation: $\kappa = \tan^{-1} \frac{M_z \cdot \sin\gamma \cdot \sin\omega - M_y \cdot \cos\gamma}{M_y \cdot \sin\gamma \cdot \sin\omega + M_z \cdot \cos\gamma}$ (12)

Determining the change in directionality of remanent magnetization, M , in the vase body in a manner dictated by its rotational symmetry, is performed from the calculated angles ω , κ at each measurement position, using the inverse rotational transform of the calculated magnetization components at an angle $-\kappa$ around the x -axis (**figure 19a**), from the reference system (xyz), where the magnetic field measurements are taken, to the (xyz)* reference system (**figure 19b**) of the cylindrical vase:

$$M_x^* = M_x = M \cdot \sin\gamma \cdot \cos\omega, \quad M_y^* = M_y \cdot \cos\kappa + M_z \cdot \sin\kappa = M \cdot \sin\gamma \cdot \sin\omega, \quad M_z^* = M_z \cdot \cos\kappa - M_y \cdot \sin\kappa = M \cdot \cos\gamma \quad (13)$$

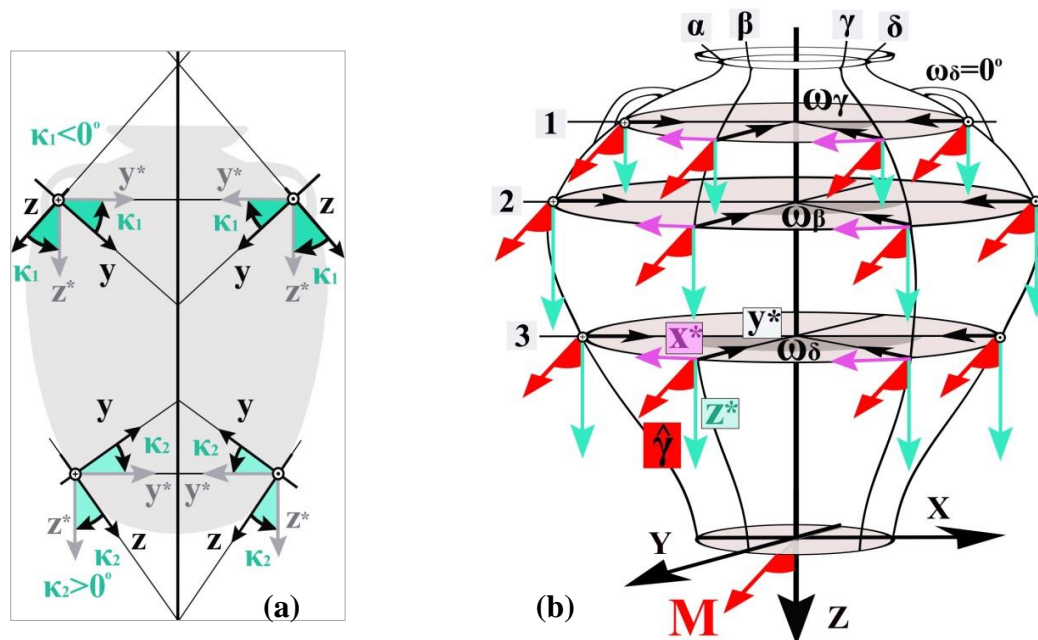
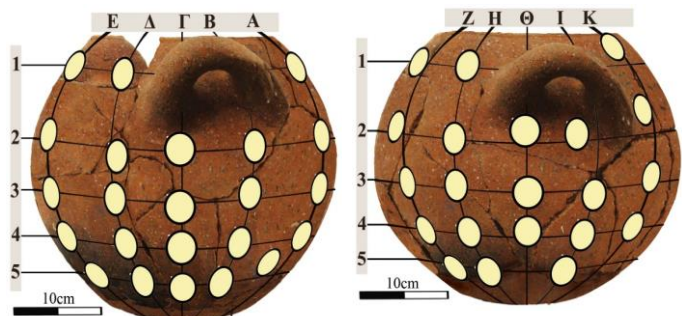


Figure 19. Illustration of the inverse transformation at an angle $-\kappa$ around the x -axes (**figure a**), from the reference systems (xyz), where the magnetic field measurements are taken, to the reference systems (xyz)* of the

cylindrical vase. Visualization of reference systems, $(xyz)^$, of the cylindrical vase (figure b), in a vase with arbitrary rotational symmetry.*

In the reverse transformation (**figure 33b**), the components M_x^*, M_y^* and the angle ω should display similar values in each longitudinal section of the body, based on the above equations, while M_z^* , which depends only on angle γ , should display similar values in each measurement position. Indicative experimental results with the reverse transformation are given from measurements of body fragments from **vase 6**.

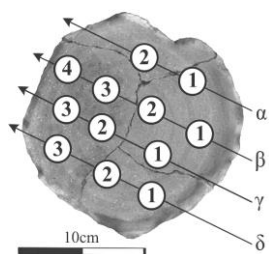
Table 3. Indicative values of the transformed components, M_i^* ($i=xyz$), of magnetization in the reference systems $(xyz)^*$ of the cylindrical vase, in fragments of the body from vase 6.



The components M_x^*, M_y^* and the angle ω display similar values in each longitudinal section of the body.

M_z^* displays similar values in each measurement position and **with that which computed for the base fragments of the vase.**

Inclination, κ , displays similar values in each transverse groove of the vase.



Vase 6	M_z	M_{xy}	M	γ°
(Base)	42.5	39.8	58.3	43.3
(mA/m)	± 0.2	± 0.2	± 0.2	± 0.3

Vase 6		A	B	Γ	Δ	E	Z	H	Θ	I	K
1	$M_y^*(\text{mA/m})$	-10±6			45±4	35±5	11±7	-16±5			-38±3
	$M_x^*(\text{mA/m})$	-34±2			-8±3	8±3	38±2	37±2			2±2
	$M_z^*(\text{mA/m})$	37±3			49±4	39±4	42±3	43±3			40±3
	ω°	197±8			100±10	77±12	16±9	336±5			272±2
	κ°	-40±8			-31±3	-29±4	-22±9	-33±5			-30±3
2	$M_y^*(\text{mA/m})$	-10±8	10±6	33±4	43±5	42±4	-7±9	-22±6	-31±6	-34±5	-44±4
	$M_x^*(\text{mA/m})$	-41±2	-35±2	-28±2	-9±2	9±2	40±2	38±2	24±2	10±2	-2±2
	$M_z^*(\text{mA/m})$	45±2	39±2	46±2	46±3	46±3	43±2	47±3	42±3	37±3	46±2
	ω°	194±7	163±8	131±4	101±3	78±3	350±10	331±3	308±3	286±3	268±2
	κ°	-16±8	-19±8	-12±3	-14±3	-6±2	-18±11	-18±5	-11±5	-16±5	-6±2
3	$M_y^*(\text{mA/m})$	3±18	15±6	26±5	40±5	41±5	-6±11	-12±6	-26±4	-43±4	-41±3
	$M_x^*(\text{mA/m})$	-43±2	-32±2	-21±2	-5±2	3±2	41±2	40±2	25±2	9±2	-9±2
	$M_z^*(\text{mA/m})$	46±2	37±3	35±3	43±3	43±3	45±2	44±2	38±2	46±2	44±2
	ω°	176±2	155±5	129±3	97±2	85±2	352±12	344±7	313±5	281±3	258±4
	κ°	24±23	13±7	19±5	17±4	21±4	9±13	14±7	16±4	14±2	18±2
4	$M_y^*(\text{mA/m})$	9±7	12±6	32±5	42±5	41±5	-10±9	-14±7	-24±5	-43±3	-40±3
	$M_x^*(\text{mA/m})$	-40±2	-41±2	-18±2	-2±2	8±2	37±2	35±2	33±2	2±2	-4±2
	$M_z^*(\text{mA/m})$	44±3	46±3	38±4	45±4	44±4	41±3	40±3	43±3	45±3	43±3
	ω°	168±9	164±6	119±3	92±2	79±2	345±11	338±10	324±8	273±20	265±14
	κ°	42±9	34±6	32±5	31±4	40±4	19±11	37±9	32±5	41±2	38±2
5	$M_y^*(\text{mA/m})$	-11±9	15±6	30±3	42±4	36±4	6±13	-19±6		-37±5	
	$M_x^*(\text{mA/m})$	-38±2	-31±2	-18±2	-13±2	11±2	39±2	32±2		7±2	
	$M_z^*(\text{mA/m})$	42±3	36±4	37±3	46±4	40±5	42±3	40±4		40±5	
	ω°	197±13	154±8	120±3	107±2	73±3	9±20	330±11		280±58	
	κ°	41±11	57±8	59±4	48±4	52±5	59±18	41±8		49±6	

8.3 Experimental results: Screening of co-belonging fragments and determination of their position in the body of a vase from their magnetization

The computed values for the inclination, κ , and the angle ω are displayed on the bodies of the 6 vases.

The magnetization magnitude M (mA/m) in body fragments displays similar values, and with those calculated from the base fragments of the vases.

The slope, κ , which displays similar values in each cross-section and determines the height-position of the fragments, is compared with the measured inclination of the walls with respect to the axis of symmetry in each vase. The computed values for the inclination, κ , from the sensor readings are similar to the measured ones.

The angle ω displays a similar value in each longitudinal section and varies at the transverse measurement positions, defining the peripheral position of the fragments on the body of each vase.

From the orientation (ω) of the transformed $M_{x\psi}^*$ in the vessel body and of $M_{x\psi}$ in the fragments of the base, the jointing position of the base and the body of each vase is determined. The shape of the original vase can be reconstructed when only a few non adjacent fragments are available.

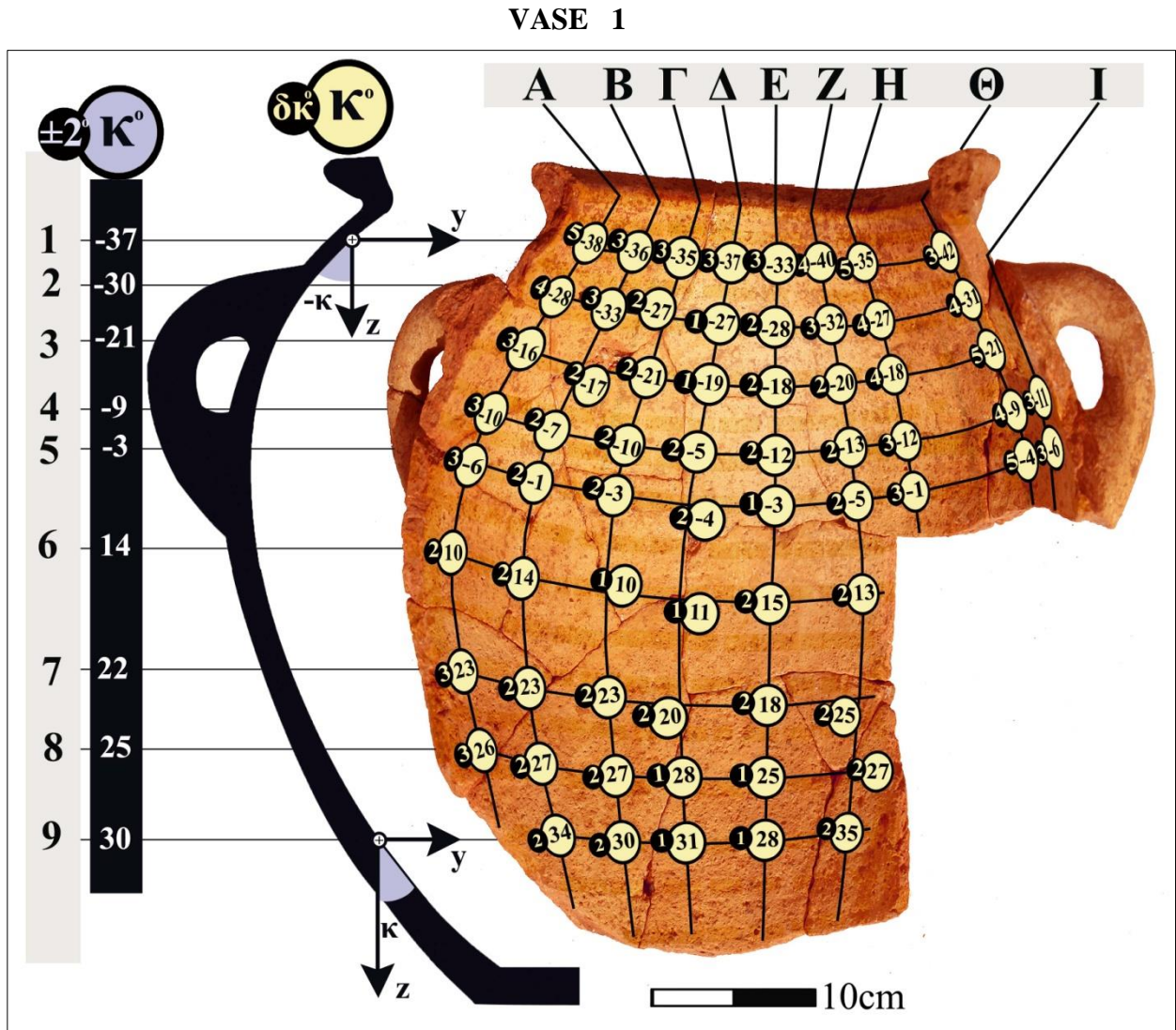


Figure 20. Angles, κ , display similar values in each transverse groove of the body. The computed values for the inclination, κ , from the sensor readings are similar to the measured ones.

Table 4. The magnetization magnitude in body and vase fragments of vase 2 display similar values. The values inside the colored frames computed by the sufficient length relationships.



Vase 1	M(Body)								
	A	B	Γ	Δ	E	Z	H	Θ	I
1	120±4	138±6	130±7	135±7	129±7	133±5	141±5	128±3	
2	135±3	126±5	136±6	132±6	133±5	132±5	132±3	127±3	
3	138±3	139±5	139±5	138±6	127±5	137±4	128±3	124±2	
4	130±3	125±4	130±4	131±5	135±5	121±4	130±3	131±3	123±3
5	119±3	137±3	140±5	132±5	138±4	128±3	129±3	138±3	131±2
6	125±3	128±3	131±4	131±4	129±4	123±3			
7	138±3	132±3	121±4	121±4	124±4	130±3			
8	133±3	141±3	123±3	132±4	135±3	134±3			
9		125±3	131±3	131±4	133±3	133±4			
	M(Base)		131.3±0.2(mA/m)						

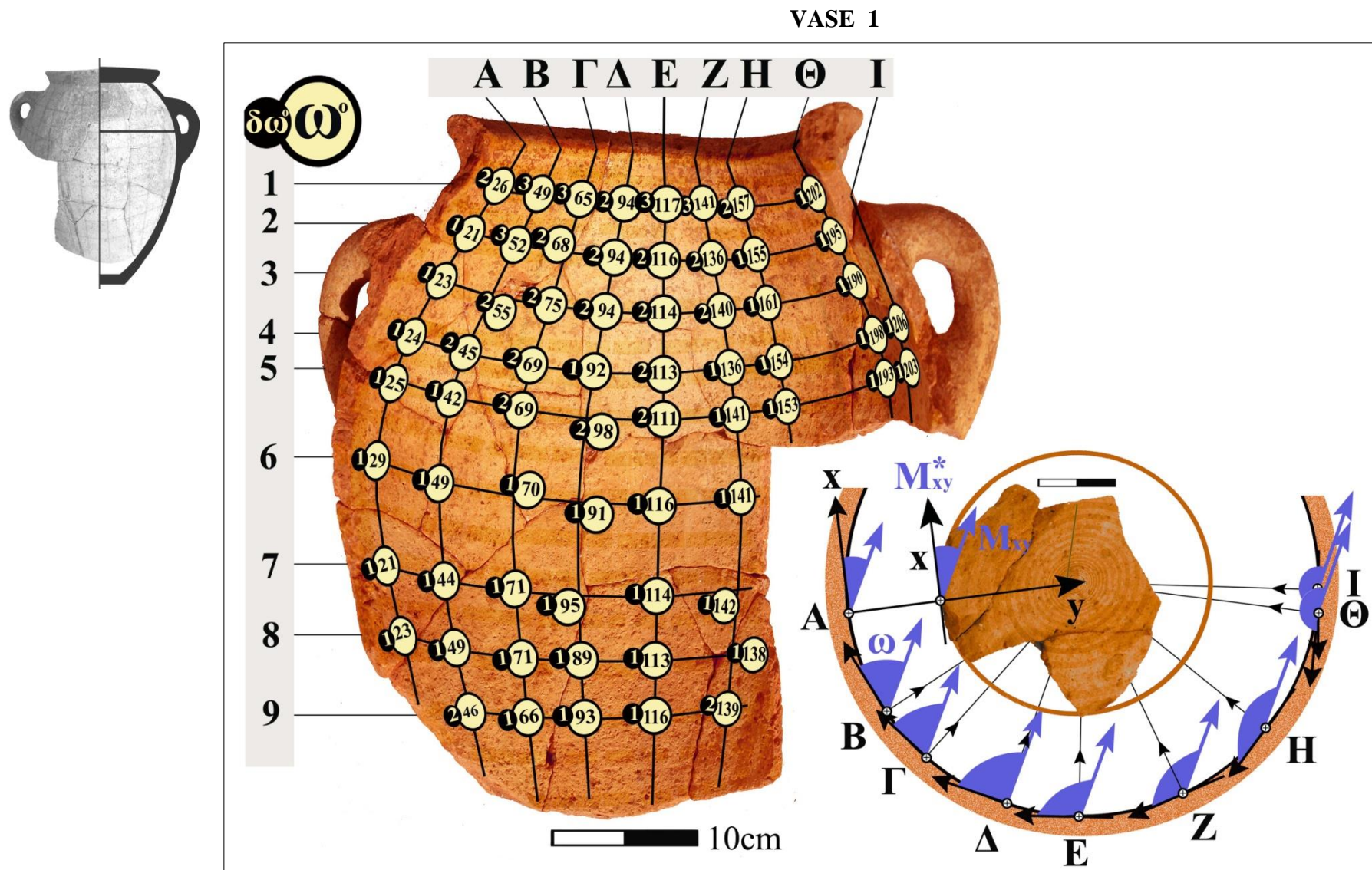


Figure 21. Angles, ω , display similar values in each longitudinal section of vase 1. If the common direction of the transformed components, M_{xy}^* , in each longitudinal section of body fragments aligned along the common direction of the components, M_{xy} , in base fragments, then the body and the base are oriented such that they join.

VASE 2

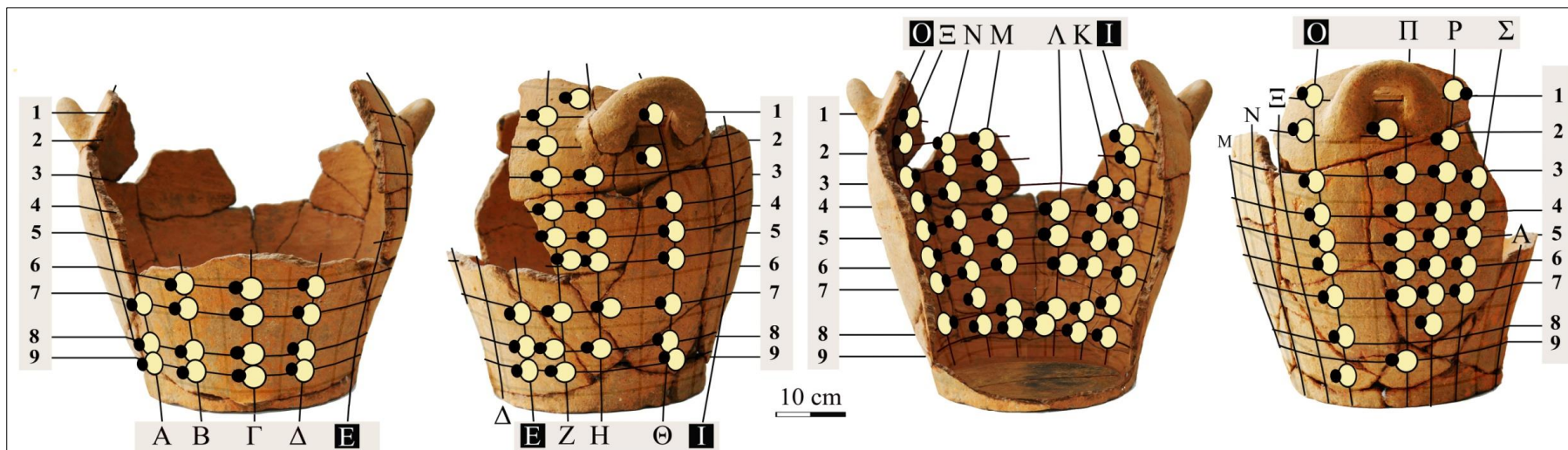


Table 5. The magnetization magnitude, M , in body and vase fragments of vase 2 display similar values. The values inside the colored frames computed by the sufficient length relationships.



M(Body)	A	B	Γ	Δ	E	Z	H	Θ	I	K	Λ	M	N	Ξ	O	Π	P	Σ	
1						182±12	153±12	179±14						169±3	171±3		168±4		
2						185±10			167±5				172±2	163±4	173±3	178±4	168±5	170±6	
3						181±7	184±7	181±8	167±3				167±3	174±4	164±3	159±4	170±6	170±6	163±5
4						185±5	169±4	164±4	173±4	162±3			170±3	169±4	174±4	166±4	174±6	173±5	170±5
5						186±4	175±5	168±4	163±3	164±3	161±2	168±2	169±3	169±4	173±5	178±7	174±7	167±6	
6		179±3	171±2	174±3		172±4	158±4	164±4	169±3	169±3	168±3	168±2	164±3	171±4	174±5	170±7	173±7	182±7	
7	156±6	175±3	170±2	162±3	169±3	168±4	161±4	167±4	170±3	168±3	173±3	164±3	176±4	173±5	177±6	185±8	170±7	178±8	
8	175±7	163±4	167±2	164±3	170±4	176±4	183±5	180±4	161±3	165±3	159±2	165±2	162±3	171±5	167±6		179±9		
9	175±7	163±4	167±2	164±3	170±4	176±4	183±5	180±4	161±3	159±3	169±2	171±3	166±3		173±7	178±10			
M(Base)	169.3±0.6 (mA/m)																		

VASE 2

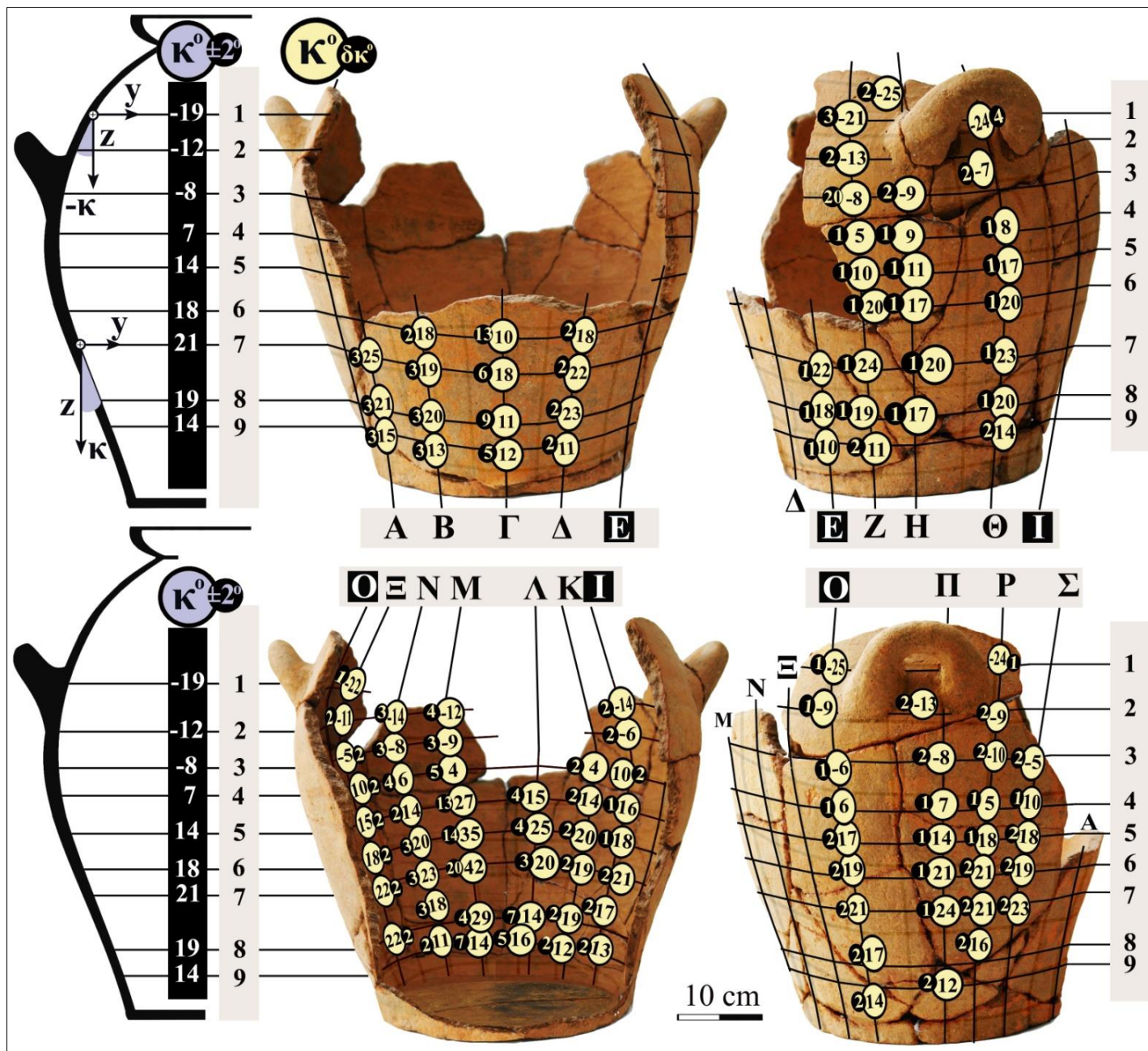


Figure 22. Angles, κ , display similar values in each transverse groove of body vase 2. The computed values for inclination, κ , from the sensor readings are similar to the measured ones.

VASE 3

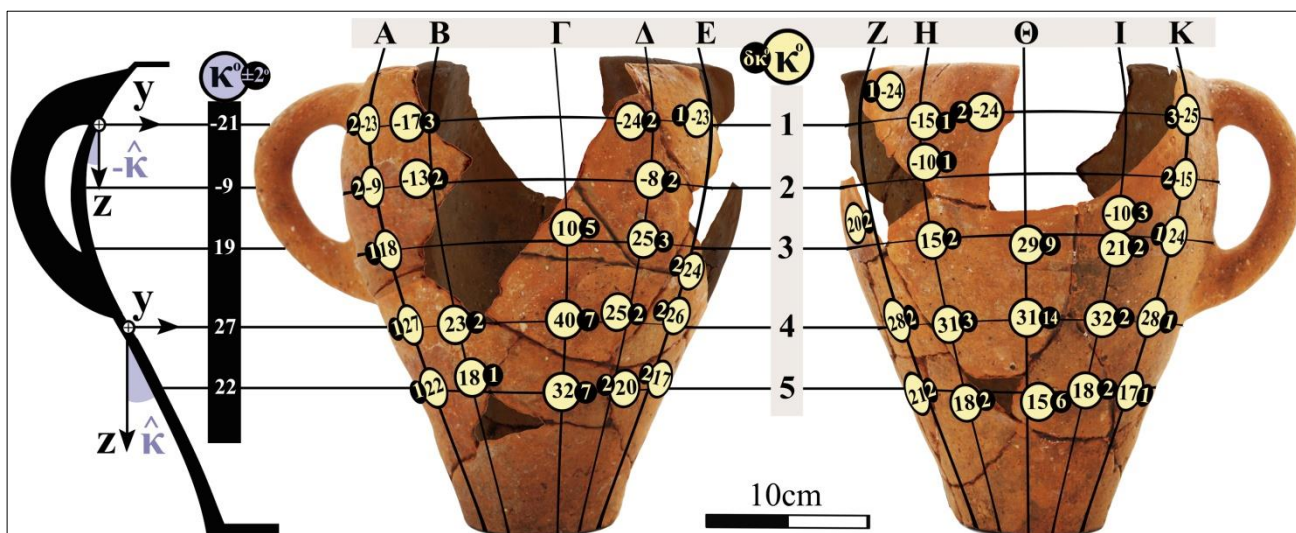


Figure 24. Angles, κ , display similar values in each transverse groove of the body. The computed values for inclination, κ , from the sensor readings are similar to the measured ones.

Table 6. The magnetization magnitude in body and base fragments of vase 3 display similar values. The values inside the colored frames are computed by the sufficient length relationships

M(Vase)	M(Body)	A	B	Γ	Δ	E	Z	H	Θ	I	K
164.1±0.4 (mA/m)	1	167±8	178±8		163±4	159±3	165±4	175±3	173±3		194±6
	2	166±5	154±4		153±3			177±3		169±3	169±5
	3	160±3		158±3	161±4	171±6	176±6	160±4	164±3	157±3	173±3
	4	168±3	161±3	162±3	162±5	161±5	181±7	162±5	166±5	167±4	155±3
	5	156±4	175±3	159±2	154±4	172±6	160±7	164±5	155±2	173±3	167±3

VASE 3

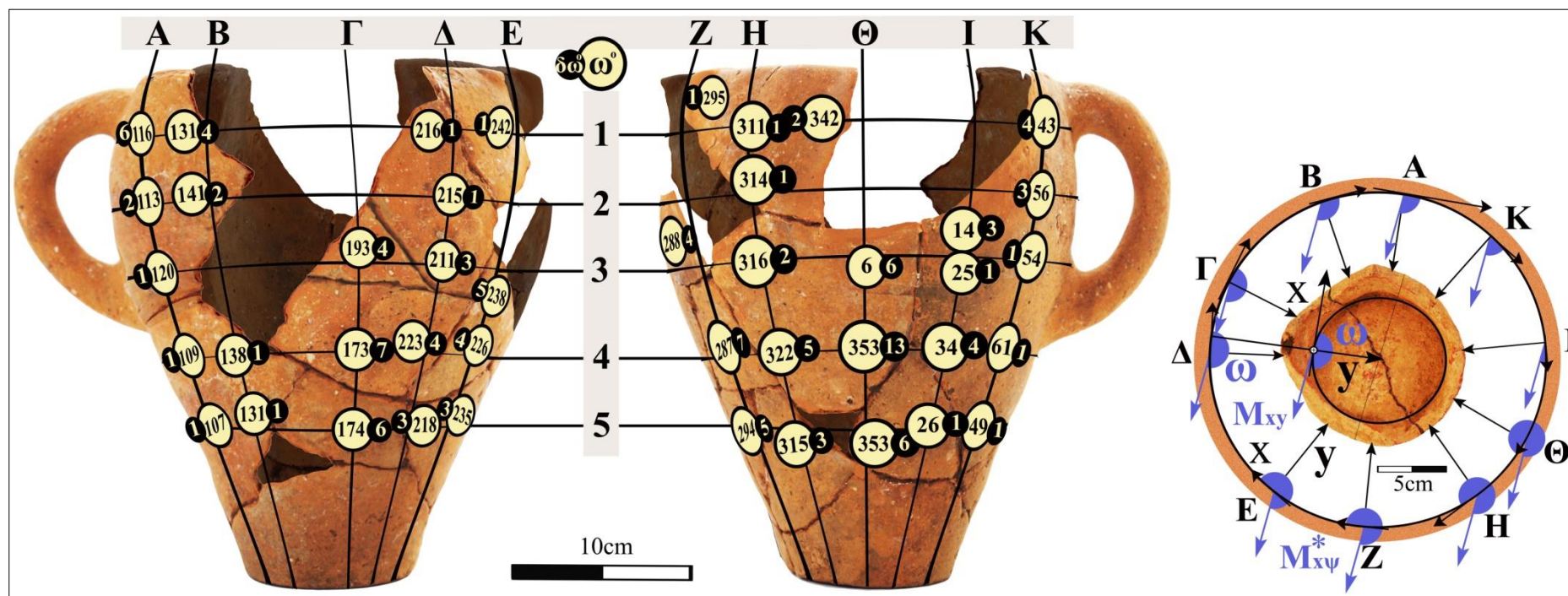


Figure 25. Angles, ω , display similar values in each longitudinal section of vase 3. If the common direction of the transformed components, M_{xy}^* , in each longitudinal section of body fragments aligned along the common direction of the components, M_{xy} , in base fragments, then the body and the base are oriented such that they join.

VASE 4

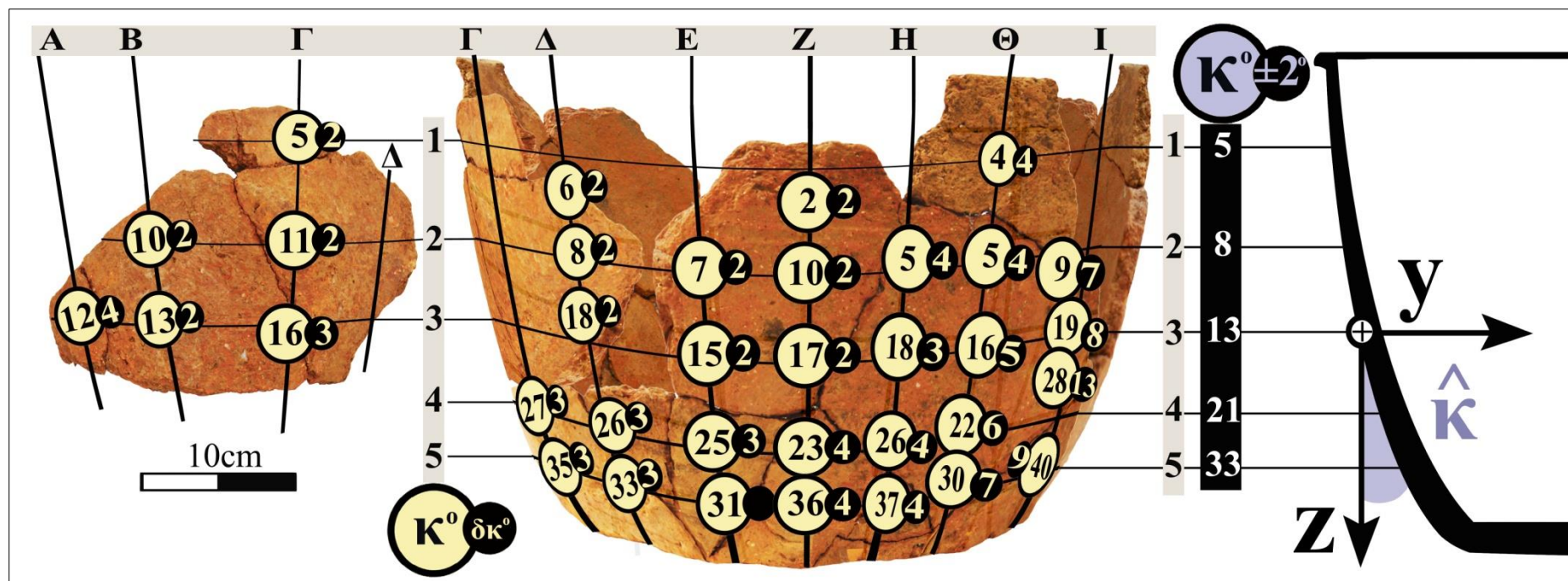


Figure 26. Angles, κ , display similar values in each transverse groove of the body. The computed values for inclination, κ , from the sensor readings are similar to the measured ones.

Table 7. The magnetization magnitude in body and vase fragments of vase 4 display similar values. The values inside the colored frames are computed by the sufficient length relationships.

M(Base)	M(Body)	A	B	Γ	Δ	E	Z	H	Θ	I
84.1±0.2 (mA/m)	1			81±3	78±3		93±3		89±2	
	2		97±3	86±3	85±3	92±3	72±3	76±3	76±3	77±2
	3	82±3	78±3	78±4	77±3	84±3	85±3	88±3	89±3	76±2
	4			87±4	76±4	90±4	80±4	78±4	78±3	82±3
	5			84±4	85±4	82±4	77±4	85±4	85±3	89±3

VASE 4

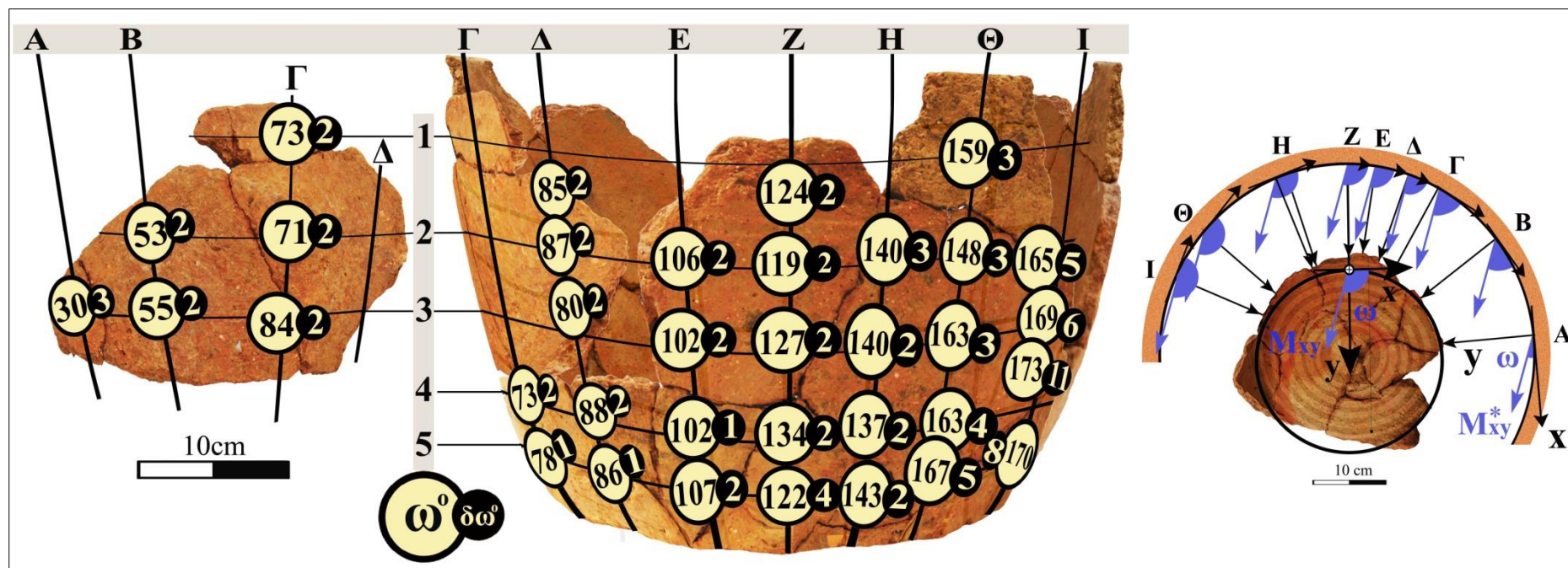


Figure 27. Angles, ω , display similar values in each longitudinal section of vase 4. If the common direction of the transformed components, M_{xy}^* , in each longitudinal section of body fragments aligned along the common direction of the components, M_{xy} , in base fragments, then the body and the base are oriented such that they join.

VASE 5

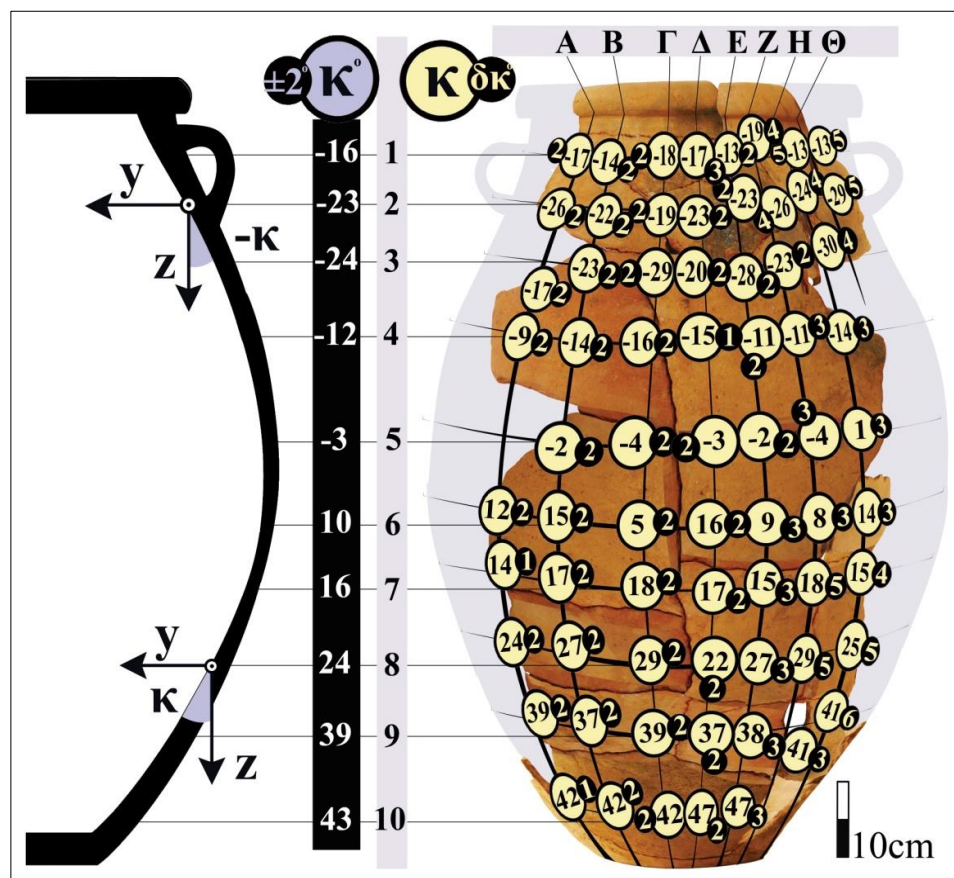


Table 8. The magnetization magnitude in body and vase fragments of vase 5 display similar values. The values inside the colored frames are computed by the sufficient length relationships

	M(Body)							
	A	B	Γ	Δ	E	Z	H	Θ
1	110±3	114±5	110±3	115±4	111±3	111±3	113±3	116±3
2	114±3	114±3	111±3	113±3	112±3	110±3	112±3	110±3
3	108±3	111±3	112±3	112±3	113±3	112±3	111±2	
4	113±4	112±3	111±3	113±3	111±3	115±3	118±3	
5		116±4	116±3	113±3	114±3	110±3	113±3	
6			114±4	116±3	110±3	115±3	111±3	
7			113±4	117±4	114±3	108±4	111±5	
8			119±4	113±3	112±3	113±4	117±4	
9			111±3	114±3	113±3	116±3	116±4	
10			119±3	116±3	115±3			
	M(Vase)	112.7±0.3 (mA/m)						

Figure 28. Angles, κ , display similar values in each traverse groove of the body. The computed values for inclination, κ , from the sensor readings are similar to the measured ones.

VASE 5

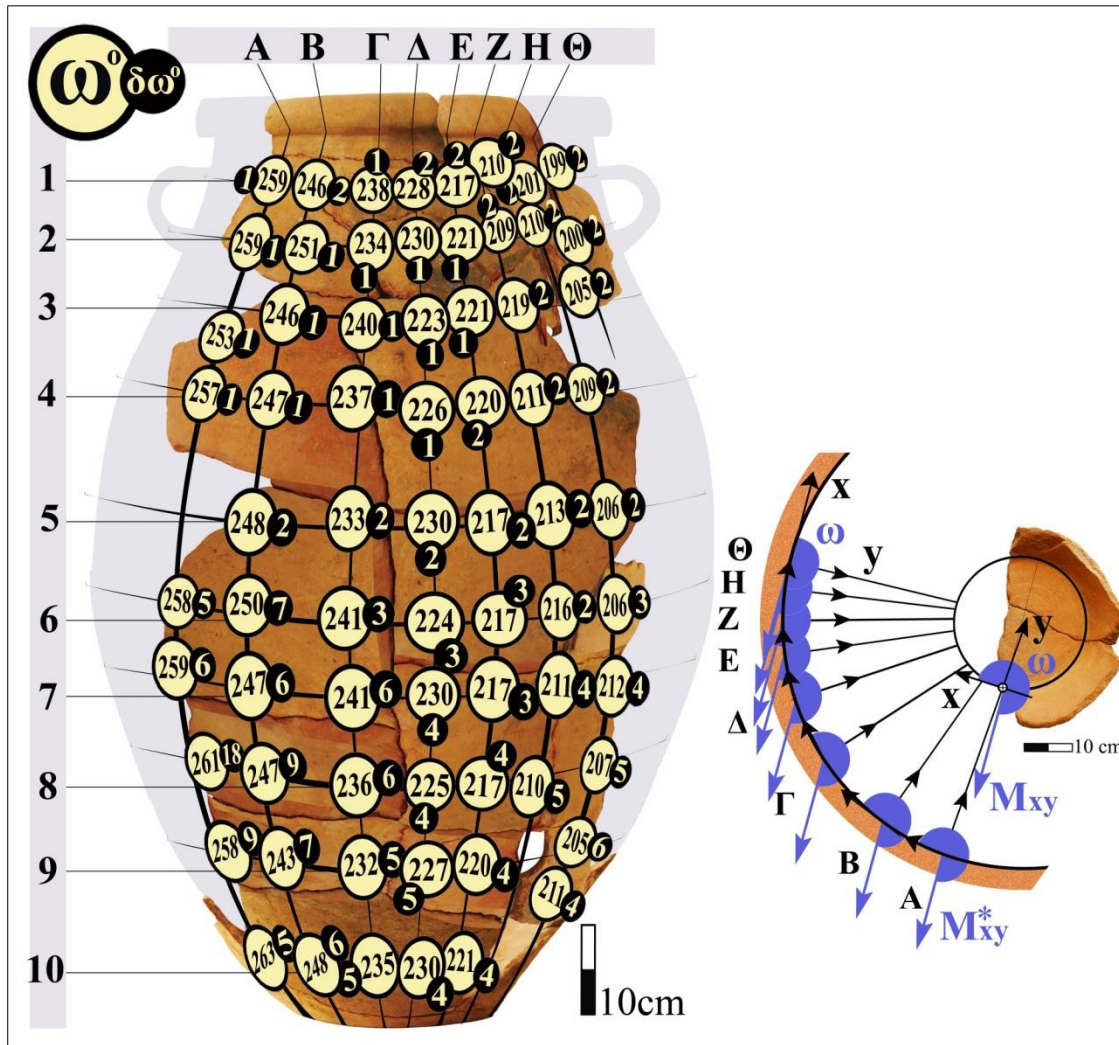


Figure 29. Angles, ω , display similar values in each longitudinal section of vase 5. If the common direction of the transformed components, M_{xy}^* , in each longitudinal section of body fragments aligned along the common direction of the components, M_{xy} , in base fragments, then the body and the base are oriented such that they join.

VASE 6

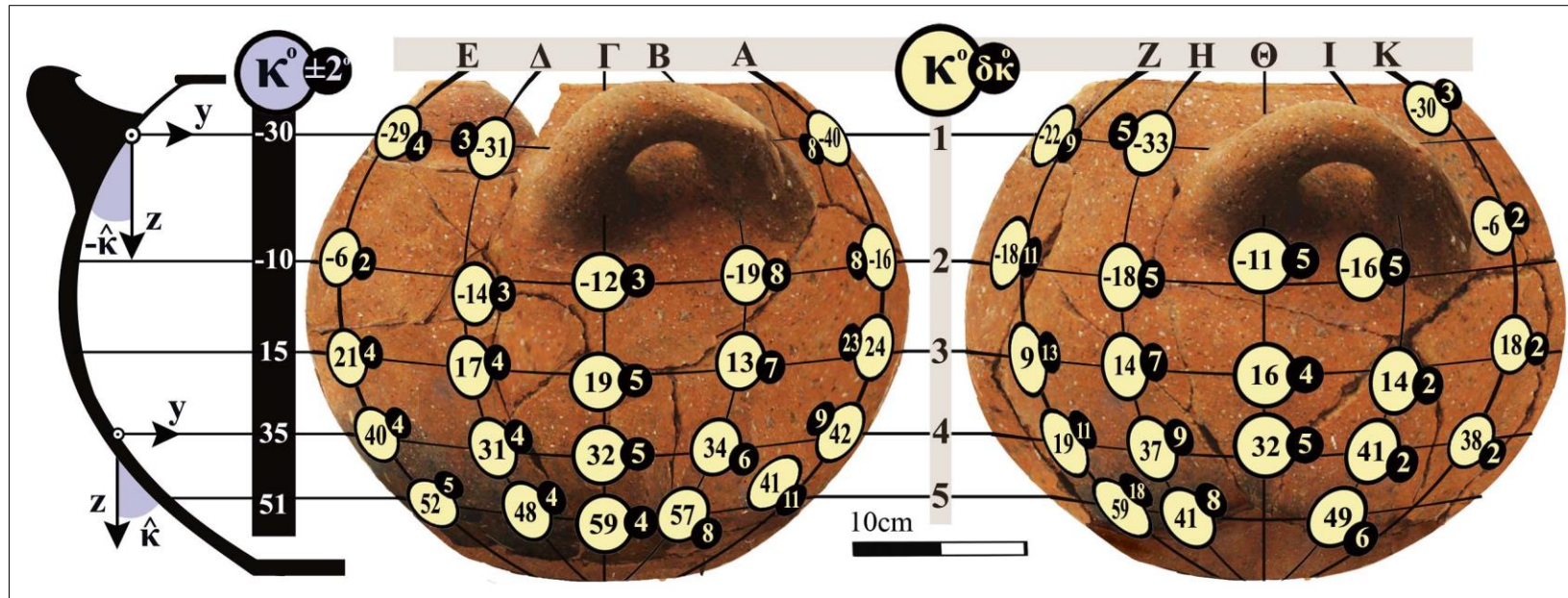


Figure 30. Angles, κ , display similar values in each transverse groove of the body. The computed values for inclination, κ , from the sensor readings are similar to the measured ones.

M(Base)	M(Body)	A	B	Γ	Δ	E	Z	H	Θ	I	K
58.3±0.2 (mA/m)	1	51±2			67±4	53±4	58±3	58±2			56±3
	2	62±2	53±2	64±3	63±4	63±3	59±2	64±3	57±4	51±4	64±3
	3	63±2	51±3	48±4	59±4	60±4	61±2	61±2	52±3	63±3	61±3
	4	60±3	63±2	53±4	62±4	61±4	56±3	55±3	59±3	62±3	59±3
	5	57±3	50±4	51±3	64±4	55±4	58±3	54±4		55±5	

Table 9. The magnetization magnitude in body and base fragments of vase 6 display similar values. The values inside the colored frames are computed by the sufficient length relationships.

VASE 6

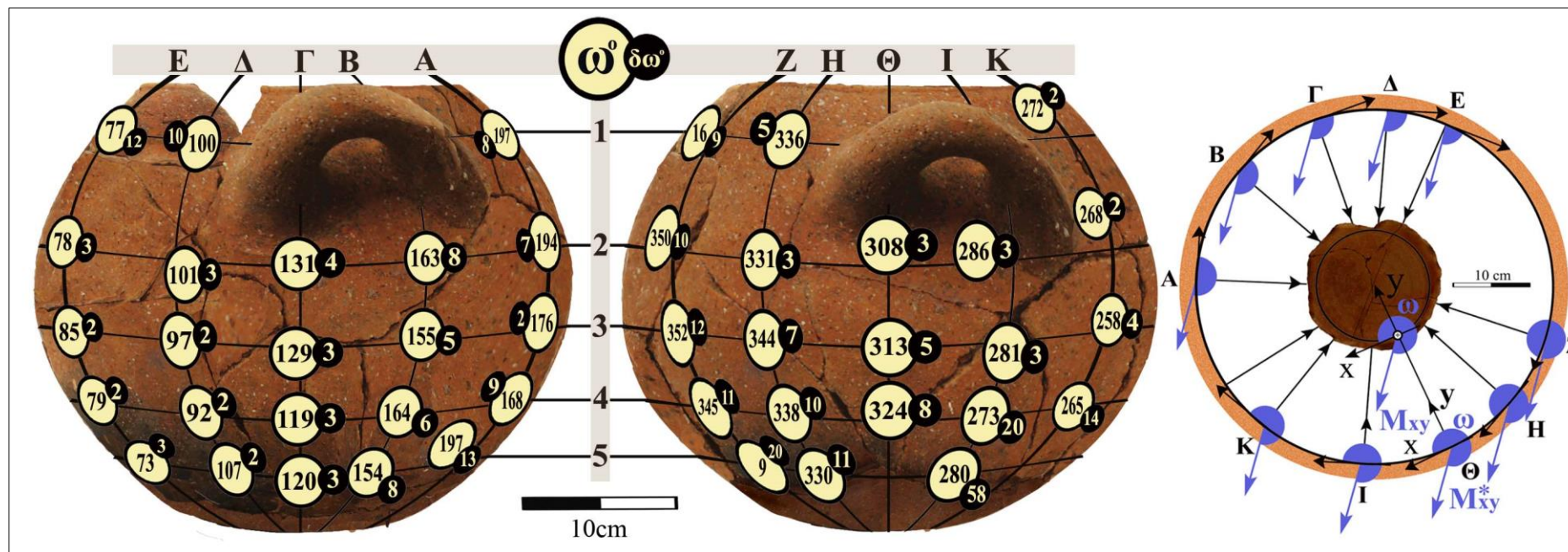


Figure 31. Angles, ω , display similar values in each longitudinal section of vase 6. If the common direction of the transformed components, M_{xy}^* , in each longitudinal section of body fragments aligned along the common direction of the components, M_{xy} , in base fragments, then the body and the base are oriented such that they join.

9. Conclusions and proposals for future research

From the experimental results it can be seen that despite the small range of the magnetization magnitude in the 6 vases, the screening of excavation fragments can be done on the basis of their similar magnetization values.

The position of co-belonging body fragments of the vases is determined with satisfactory precision by calculating the angles ω , κ . The methodology can be used to screen fragments from archaeological excavations under the following conditions:

- Finding co-belonging fragments of the vase base is essential for calculating the magnetization deviation angle, γ , and the parameters α^{\perp} , α^{\parallel} , λ from measurements of their magnetic field. This is not a problem, because in the usual process of screening co-belonging fragments, the base fragments, which are easily identified, are initially collected and the neighboring fragments of the vase body are then searched.

- In the fragment reference system, where measurements are taken, the x -axis is oriented in the direction of the vase grooves while the y -axis is perpendicular to the surface and directed towards the inside of the fragments. To take measurements on all fragments in relation to the same reference system, the z -axis must be oriented towards the base (or the aperture) of the vase. The orientation of each fragment with respect to the base or the aperture of the vase can be done easily, on the basis of its shape and the curvature of the grooves on its surface.

- The error in calculating the angles ω , κ increases as the sensor contact with the fragment surface becomes smaller. Determining the position of fragments with high curvature is hampered, at least with the ~3cm fluxgate sensors used in this research. However, in this case too, it is about a few fragments of small vases or shells with characteristic curvature of the shape of the vases, which are easier to categorize empirically.

Proposals for future research, based on additional preliminary investigations, are given below.

1. Experimental application of the method in fragments from archaeological excavations.

2. Construction of ceramic specimens from different raw materials and with different granulometry, fired in the same pottery kiln under different slopes of the kiln floor.

- Construction of ceramic specimens from the same raw material, fired in the same kiln, with the same slope, but under different oxidative conditions and at different temperatures.

- Analysis of the composition of the specimens and their content in magnetic oxides.

The above are necessary in order to study the dependence of remanent magnetization on the clay composition, the temperature/ventilation conditions in the pottery kiln, the orientation of the geomagnetic field, and the degree of magnetic anisotropy of the ceramic material. For the above research a traditional wood kiln is required. In ceramic specimens made from different commercially available clays fired in an electric kiln, the measured magnetic field was less than 10nT and without clear orientation, as a result of magnetic fields from the electric currents of the kiln.

3. Investigation of the sensitivity radii dependence of the vertical (α^{\perp}) and parallel (α^{\parallel}) sensors on the magnetization magnitude of the ceramic material.

- Investigation of the mode, the sensors' degree of excitation, and of the dependence of the normalization constant, λ , of the vertical and parallel sensor readings as well as of α^{\perp} , α^{\parallel} , on measurements of the same specimens at different distances from their surface.

A prerequisite for these are magnetic field measurements in ceramic materials with higher thermoremanent magnetization, so that their magnetic field is attenuated at longer distances from their surface.

- Investigation of the dependence of the normalization constant, λ , of the vertical and parallel sensor readings and their correlation with the dimensions and technical characteristics of the sensors from measurements of the same specimens with different types of *fluxgate* sensors.

4. Investigation of the feasibility of constructing a handy three-sensor fluxgate magnetometer for local measurements of weak surface fields (1-1000nT) in ceramic materials with a magnetization of 1-1000 mA/m for archaeological use.

5. Investigation of induced magnetization of ceramic materials in alternating magnetic fields as an additional criterion for the screening of co-belonging excavation fragments.

Experiments were carried out by the successive placement of cylindrical ceramic specimens from co-belonging vase fragments at specific locations inside a solenoid ($L \sim 3\text{mH}$) in an RLC circuit in a field ($\sim 50 \mu\text{T}$), and the induced magnetization was calculated by increasing the inductance of the coil at the resonance frequency [65 pp.21-34] of the circuit (*Null Detection* technique). In sequential addition of co-belonging tiles at different resonance frequencies ($f=40\text{-}95\text{KHz}$), similar increase in their induced magnetization was observed, which is a criterion for distinguishing samples from different vases.

-Investigation of the possibility of constructing a handy instrument for local surface measurements of induced magnetization of ceramic materials from an alternating magnetic field for archaeological use for the purpose of screening the co-belonging excavation fragments.

6. Investigation of the application of the method in rocks with remanent magnetization, for the localization of changes in local soil morphology, by comparing the orientation of the magnetic field of the parent rock and its detached parts. Experiments were conducted in fragments from serpentinite and were found to be oriented in a joining manner, from their thermoremanent magnetization.

These studies will be presented in follow up publications in the near future.

APPENTIX A

Remanent Magnetization in ceramic fractures

Since the 19th century it has already been observed [33,34] that rocks containing ferromagnetic materials carry the information about the direction of the Earth's magnetic field during their formation unaltered. Dated rocks are examined by paleomagnetic methods in order to investigate changes in the geomagnetic field. Similar archaeo-magnetic methods [35-37] are used for the extraction of technological information and the dating of ceramics that remain in the same position since the time of their last firing (fireplaces, kilns) using archaeomagnetic maps. Thermo-remnant magnetization in ceramics is mainly due to the formation of magnetite and hematite upon firing, at temperatures higher than [38, 39 pp.46-75] *Curie* temperature.

Clay does not have a particular mineral composition because its formation is due to the decomposition of fragmented pyrogenic, sedimentary and metamorphic rocks. Typical components of the raw material according to the "classical" triangular diagram [40-42] of *Levin* (1964) are various aluminosilicate compounds of SiO₂ (45-70%), Al₂O₃ (10-30%) and CaO (<20%) with lesser **admixtures** of iron oxide (<6%) and magnesium (2-3%). The final mineral composition of ceramics and their magnetic properties are formed during firing (at typical temperatures of 600°-1100° C) and depends on the composition of the clay and the heating conditions (heating/cooling rate) and ventilation (oxidative-reducing atmosphere) in the kiln.

The aluminosilicate minerals content in the firing clay composition are paramagnetic or diamagnetic [43] with a typical [44,45] magnetic susceptibility of 1000 times less than antiferromagnetic [43 pp.151-173] and 10000 times smaller than ferrimagnetic [43 pp.175-195] iron oxides. Magnetic iron oxides, according to the Fe-Ti triaxial phase (Butler 1992) diagram [46 pp.411], form three basic series of solid solutions with titanium dioxide (TiO₂), the *titanomagnetite*-ulvöspinel chain (Fe₃O₄-Fe₂TiO₄) the titanohematite/magemite-ilmenite series [(α,γ)Fe₂O₃-FeTiO₃] and the iron hydroxide series, ferro-pseudobrukite-pseudobrukite.

From the minerals of the three solid solutions of iron oxides, dominant role in the magnetic behavior of ceramics have the minerals of the titanomagnetite and the titanohematite, because the rest [47-57] are dehydrated or oxidized when firing the vases into minerals of low magnetic susceptibility or converted into minerals of titanomagnetite and titanohematite series.

The magnetic properties of the fired clay are determined [46 pp.409-412, 58] by the influence of the cooling temperature on the equilibrium compounds of solid solutions of titanomagnetite (*Nagata* 1961) and titanohematite (Robinson 2004), depending on their titanium content.

• Hematite- Fe_2O_3 (**figure A.1a**) forms a rhomboidal crystal [46 pp.29-31,61,64,81-87], with a corundum structure (**figure A.1b**). The metallic Fe^{+3} cations are placed in a hexagonal configuration, at equally-distanced levels, perpendicular to the C-axis.

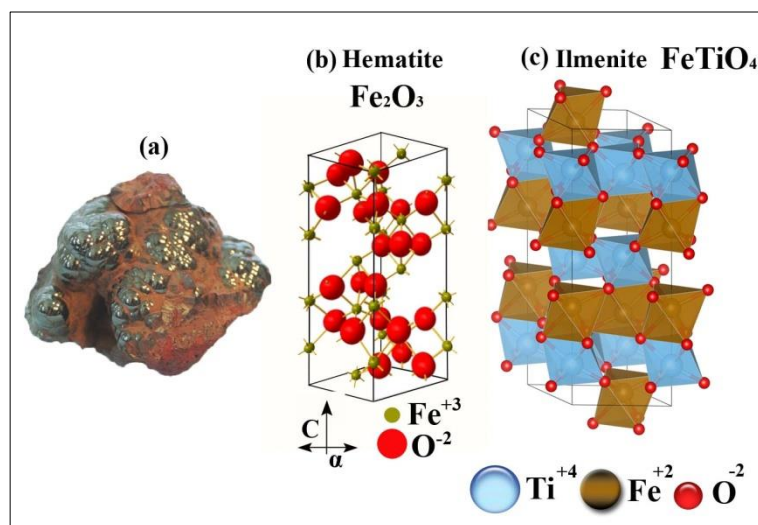


Figure A.1 Titanohematate series [hematite(Fe_2O_3)/ilmenite(FeTiO_3)]. (a) Hematite ore³. (b) Arrangement⁴ of ions of the rhombohedral hematite crystal with a corundum structure, (c) Hexagonal arrangement of the metal cations in ilmenite⁵.

Each ${}_{26}\text{Fe}^{+3}$ [$1s^2, 2s^2, 2p^6, 3s^2, 3p^6, 3d^5$] cation [46 pp.111-7] contains 5 unbound electrons (e^-) in the 3d orbital. Due to the charge compensation and the $2 \times 5 = 10$ single e^- with anti-parallel spins, according to the $|\text{Fe}^{+3}\text{Fe}^{+3}|\text{O}_3^{-2}$ scheme, hematite should in theory [46 pp.126-8] exhibit zero magnetization. Its weak antiferromagnetic behavior, with a magnetic susceptibility of $\chi_m = (1.19-1.69) \cdot 10^{-6} \text{ m}^3/\text{Kg}$ [45 pp.40] is due to grid defects, and increases with the involvement of other atoms in the crystal, which disturb the parallel alignment of opposing electron moments even more.

Substitution (**figure A.1c**) of Fe^{+3} by Ti^{+4} (Nagata 1961, Stacey and Banerjee 1974) increases [58] the saturation magnetization of the hematite ($\psi = 0$) - ilmenite ($\psi = 1$) series of minerals [$\text{Fe}_{2-\psi}\text{Ti}_\psi\text{O}_3$]. When the substitution is less than $\psi \approx 0.5$, the Ti^{+4} cations are distributed equally among the layers of Fe^{+3} cations and the crystal does not exhibit magnetization. For higher concentrations ($0.5 < \psi < 0.8$) of titanium, the ${}_{22}\text{Ti}^{+4}$ [$1s^2, 2s^2, 2p^6, 3s^2, 3p^6$] cations not containing unbound e^- are randomly distributed in alternating layers between the Fe^{+3} cations and the crystal obtains maximum antiferromagnetic properties for $\psi \approx 0.8$. For higher contents ($\psi > 0.8$), the titanium cations are distributed symmetrically between the layers of Fe^{+3} cations, with the result that ilmenite not exhibiting anti-magnetic, but rather paramagnetic behavior. The corresponding, Neel temperature of the Curie temperature for antiferromagnetic materials sharply decreases from the 685 °C, corresponding to hematite, with increasing Fe^{+3} substitution by Ti^{+4} .

Hematite is a primary compound of clay or is formed in fired clay [14] from the dehydration and recrystallization of the contained iron hydroxides and oxyhydroxides, and from the decomposition of the existing ferric (e.g., illite, chlorite) minerals. The thermoremanent magnetization is solely due to the

³This figure was taken from, <https://commons.wikimedia.org/wiki/File:Hematite.jpg> licensed under the Free Documentation License (GFDL).

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⁵This figure was taken from, <https://commons.wikimedia.org/wiki/File:Ilmenit-Struktur.png> licensed under the Creative Commons Attribution 4.0 International.

secondary hematite formation during the firing of the vases, since the hematite magnetic granules in the raw material do not retain specific orientation after mixing the clay.

High hematite concentrations (> 6%) are expected at high temperatures ($T > 950^{\circ}\text{C}$) and oxidative firing conditions in clays with a low CaO content (<5%), while in larger quantities, the hematite clay content is limited [59 pp.124-126] because Fe^{+3} participates in aluminum clinopyroxenes at the octahedral positions and in ghelinite at the octahedral and tetrahedral positions. The high percentage of CaO is due to the deliberate mixing of calcium aluminate clays or the addition of calcium carbonate to low-calcium clays, mainly due to the sintering of clay at lower temperatures [60] and the lesser need to control the temperature during firing. Moreover, the small variation of the thermal expansion-shrinkage coefficient obtained at temperatures of $850^{\circ}\text{-}1050^{\circ}\text{C}$ with the addition of calcium [61,62], reduces the possibility of forming micro-fractions when clay coatings are applied and gives bright colors to the color of the clay [46 pp.130-6], creating the necessary contrast to the usually dark-colored writing decoration.

• Magnetite (**figure A.2a**) crystallizes in holohedry according to a cubic system [46 pp. 32,61,64,87-92] with a reverse spinel structure. The oxygen anions (O^{2-}) form an internally-centered cubic lattice in each structural unit (**figure A.2b**), in which the Fe^{+2} and Fe^{+3} cations are placed in 4 tetrahedral positions (A) and 8 octahedral positions (B). To maintain the charge balance with the 4 O^{2-} anions, the Fe^{+2} occupy octahedral positions while the Fe^{+3} are equally distributed between the octahedral (B) and tetrahedral positions (A), according to the scheme: $\text{Fe}^{+3}|\text{Fe}^{+3}\text{Fe}^{+2}|\text{O}_4^{2-}$

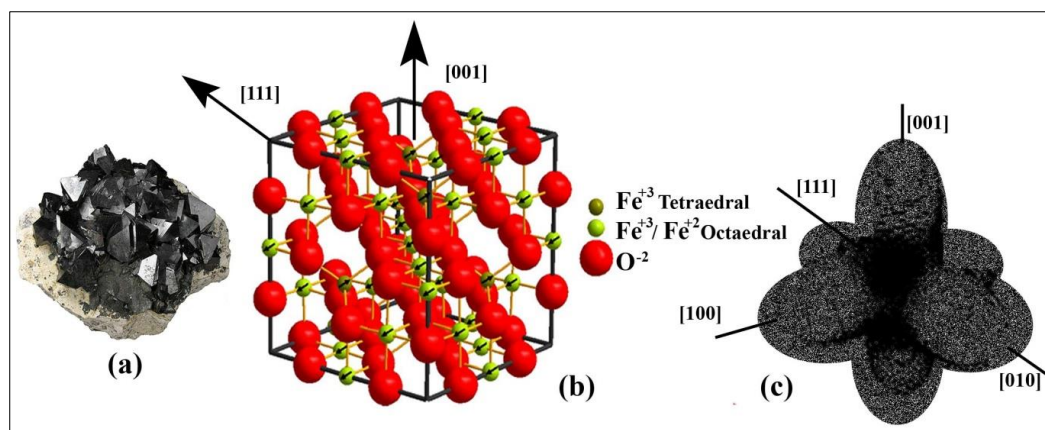


Figure A.2 Magnetite(Fe_3O_4): (a) Octahedral crystal of magnetite⁶. (b) Internal crystal structure⁷. The diagonal $[1,1,1]$ and the perpendicular direction $[0,0,1]$ to the cube are denoted by the arrows. (c) Directions of magnetic anisotropy of the crystal. The diagonal $[1,1,1]$ of the cube, with the smallest energy, is the "easy" magnetization direction.

The cations $_{26}\text{Fe}^{+2}$ [$1s^2, 2s^2, 2p^6, 3s^2, 3p^6, 3d^6$] contain 4 unbound e^- , while the $_{26}\text{Fe}^{+3}$ [$1s^2, 2s^2, 2p^6, 3s^2, 3p^6, 3d^5$] ones contain 5 unbound e^- in the 3d orbitals. Of the $5 + 4 = 9 e^-$ of the Fe^{+3} and Fe^{+2} in the octahedral positions (B), the $5e^-$ have antiparallel spins with the $5e^-$ of the Fe^{+3} in the tetrahedral positions (A) and their electron magnetic moments are mutually eliminated. Therefore, the ferrimagnetic behavior of magnetite [46 pp 128-130] corresponds to $9-5 = 4 e^-$ or 4 Bohr magnetons ($\mu_B=0,93.10^{-23}\text{J/T}$) per molecule (0°K).

Substitution of Fe^{+3} by Ti^{+4} , reduces the saturation magnetization of the magnetite ($z = 0$) - ulvospinel ($z = 1$) series of minerals $[\text{Fe}_{3-z}\text{Ti}_z\text{O}_4]$. Substitution of Fe^{+3} by Ti^{+4} in $\text{Fe}^{+3}|\text{Fe}^{+3}\text{Fe}^{+2}|\text{O}_4$ ($z=0$) causes Fe^{+3} to be converted to Fe^{+2} , to maintain the charge balance according to the scheme: $\text{Fe}^{+2}|\text{Fe}^{+2}\text{Ti}^{+4}|\text{O}_4$ ($z=1$, ulvospinel:

⁶ This figure was taken from <https://commons.wikimedia.org/wiki/File:Magnetite-244496.jpg>, licensed under the Creative Commons Attribution-Share Alike 3.0 Unported.

⁷ This figure was taken from <https://commons.wikimedia.org/wiki/File:Magnetit1.jpg>, licensed under the Creative Commons Attribution-Share Alike 3.0 Unported, 2.5 Generic, 2.0 Generic and 1.0 Generic license.

Fe₂TiO₄). Because the Ti⁴⁺ cations do not contain unbound e⁻, the four single e⁻ of the Fe²⁺ in the octahedral positions (B) have antiparallel spins with the 4 single e⁻ of the Fe²⁺ in the tetrahedral positions (A) and ulvospinel does not exhibit magnetic behavior. The gradual increase of Fe³⁺ substitution by Ti⁴⁺ [58] causes the reduction of magnetization (W.O'Reilly 1984), at Curie temperature increase of the size of the building blocks. The magnetic susceptibility of $\chi_m=(169-290).10^{-6} \text{ m}^3/\text{Kg}$ [45 pp.40] of the usual solid titanomagnetic solutions corresponds to 15% of the higher values of $\chi_m=(440-1116).10^{-6} \text{ m}^3/\text{Kg}$ of magnetite.

Magnetite is formed during firing, from the partial reduction of hematite under conditions of oxygen deficiency and at temperatures higher than 570°C inside the firing clay mass [63], which gets darker shades. Depending on the size of the magnetic grains, it displays the highest magnetic susceptibility values, compared to the other mineral components of firing clay. The magnetic susceptibility value of magnetite is about 1000 times greater than the magnetic susceptibility of the most powerful paramagnetic mineral and about 10000 times the values of diamagnetic αργίλοπυριτικών minerals.

Magnetite clay content is usually less than the lowest traceability limits (0.1%) by conventional methods of analysis. The total magnetic susceptibility of firing clay is estimated, theoretically (**table A.1**), by the sum of the values of the magnetic susceptibility of its most common diamagnetic, paramagnetic, antiferromagnetic and ferrimagnetic oxides. The typical clay content of basic metal oxides, other than magnetite, is determined by processing 50 analyzes using ICP-OSE (ENVIRO II ICP, Spectro Cirros ICP) method of ceramic material from fragments of different archaeological excavation vases.

Types of oxides	Basic oxides of firing clay	Typical clay content of oxides (C _i % w/w)	Mass magnetic susceptibility ⁸ $\chi_m(10^{-6} \text{ m}^3/\text{Kg})$	Reduced magnetic susceptibility $A_i=\chi_i.C_i/100$	Percentage content of mass magnetic susceptibility $A(\%)=(A_i/\Sigma A_i).100$
d	SiO ₂	59	-0.0058	-0.0034	-1.7
d	Al ₂ O ₃	16	-0.0046	-0.00074	-0.37
d	CaO	9	-0.0034	-0.00031	-0.16
a	αFe ₂ O ₃	7	+1.7	+0.12	+60
f	Fe ₃ O ₄	0.01	+800	+0.08	+40
d	MgO	3.2	-0.0032	-0.00010	-0.050
-	K ₂ O	3	-	-	-
d	Nα ₂ O	1.2	-0.0040	-0.000048	-0.0024
p	TiO ₂	1	+0.00093	+0.0000093	+0.0047
-	P ₂ O ₅	0.5	-	-	-
a	MnO	0.1	+0.86	+0.00086	+0.43
	Σύνολο	100		ΣA _i =0.20	98

Table 1. Estimation of the magnetic contribution of basic diamagnetic (d), antiferromagnetic (a), ferrimagnetic (f) and paramagnetic (p) oxides to the magnetic susceptibility of ceramic materials.

The magnetic properties of firing clay according to the above table are due to hematite and magnetite, for corresponding contents of 7% and 0.01%.

In conclusion, the thermoremanent magnetization of fragments is due to the formed magnetite and hematite oxides during firing [64], which are oriented during cooling in the direction of the earth's magnetic field. The paramagnetic behavior of firing clay in room temperature is due to the purely paramagnetic aluminosilicate

⁸ Landolt-Börnstein: «Numerical Data and Functional Relationships in Science and Technology», (New Series, II/2, II/8, II/10, II/11, II/12a II/16, III/19), Springer-Verlag: «Coordination and Organometallic Transition Metal Compounds», (Heidelberg, 1966-1984), Springer-Verlag: «Diamagnetic Susceptibility», Heidelberg 1986, Springer-Verlag: «Magnetic Properties of Metals», Heidelberg, 1986-1992), Masson: «Tables de Constantes et Données Numérique», Volume 7, Relaxation Paramagnetique, Paris, 1957.

compounds of firing clay and possibly [45 pp.43-8, 43 pp.359-408] due to the contribution of ferrimagnetic iron oxides and sulphides in small ($d < 0.03 \mu\text{m}$) grains.

APPENTIX B

Experimental investigation of the measured magnetic field dependence on the ceramic material remanent magnetization in base vase specimens

The remanent magnetization dependence and computation on the magnetic sensor measurements, is studied through a series of experiments answering the following research aspects:

B.1 The width of the region that excites the sensors and the degree of anisotropy of the ceramic material.

B.2 The induced magnetization contribution from the magnetic field ($<15.4 \mu\text{T}$) of the sensor/magnet apparatus, which forms a three-dimensional orthogonal system, and that of the earth on the sensor readings.

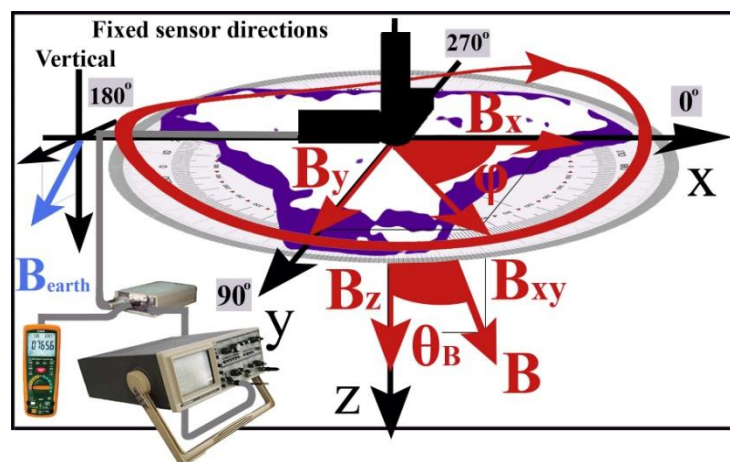
B.3 The dependence of the sensor excitation area on the thickness of the fragments.

B.4 The dependence of the sensor excitation area on the measurement position on the fragments surface and the directionality of their magnetization.

B.5 The comparison of sensor sensitivity degree of the vertical and the two parallel sensors to the ceramic surface, due to the differing positioning of the excitation solenoids.

From the experimental results it can be concluded that the sensors are excited by a cylindrical area of the ceramic material, whose dimensions depend on the magnitude and the directionality of the remanent magnetization, the fragment thickness, and finally the sensor position. For the remanent magnetization computation from the magnetic field measurements, the readings of the vertical sensor are adjusted, because of its smaller sensitivity degree, compared with the parallel ones.

B.1 Investigation of the range and anisotropy degree of the sensor excitation area

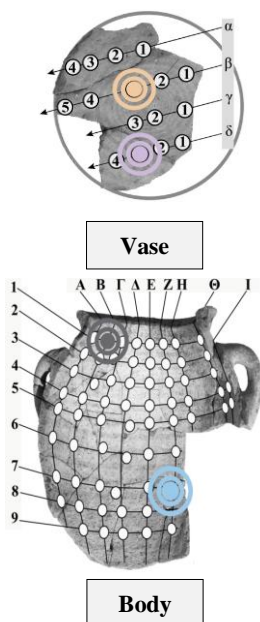


The three-dimensional orthogonal system layout of sensors/magnets (**figure B.1**) remains in the same position with respect to the geomagnetic field, with x -sensor ($B_x=0$) being perpendicular to the horizontal and vertical constituent of the geomagnetic field. The vase fragments are supported on a horizontal disk at a position where $B_y=0$ and $B_x=B_{xy}>0$. Measurements are taken with every 30° rotation of the disk, φ_δ , from 4 irregularly-shaped co-belonging fragments, 2 from the vase and 2 from body, for each of the 6 vases.

Figure B.1 Illustration of the experimental apparatus and measured components of the field B in the considered coordinate system by the three-dimensional orthogonal system layout of sensors/magnets.

With every rotation of the disk, the magnitude $B = \sqrt{B_x^2 + B_y^2 + B_z^2}$, angle $\theta_B = \cos^{-1} \frac{B_z}{B}$ and angle $\varphi = \tan^{-1} \frac{B_y}{B_x}$, which is compared with the disk rotation angle φ_δ , are computed from the sensor readings, B_i ($i=x,y,z$).

Table B.1 Change in the computed angle φ from the sensor readings with respect to the disk rotation angle, φ_δ , in 2 base fragments (β_3, δ_3) and 2 body fragments ($B1, Z7$) of vase I.



Vase 1				
φ_δ ($\pm 1^\circ$)	Base		Body	
	β_3	δ_3	B1	Z7
	φ ($\pm 1^\circ$)			
	$\pm 2^\circ$	$\pm 1^\circ$	$\pm 2^\circ$	$\pm 1^\circ$
0	0	0	0	0
30	31	33	25	32
60	63	60	63	63
90	89	88	92	91
120	121	117	122	117
150	148	152	155	150
180	178	179	180	181
210	209	213	215	212
240	241	243	245	243
270	271	271	272	271
300	299	301	299	300
330	329	334	330	329

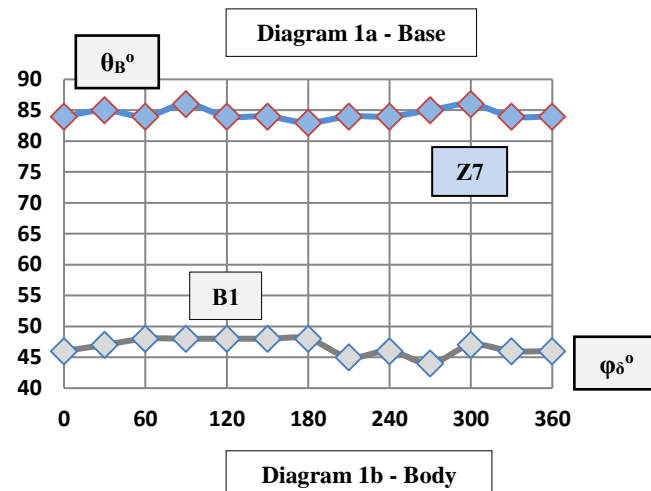
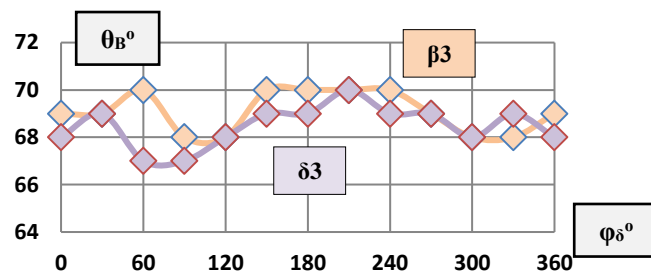


Diagram B.1 Change in the computed angle θ_B at disk rotation angles φ_δ in 2 base fragments (a) and 2 body fragments (b).

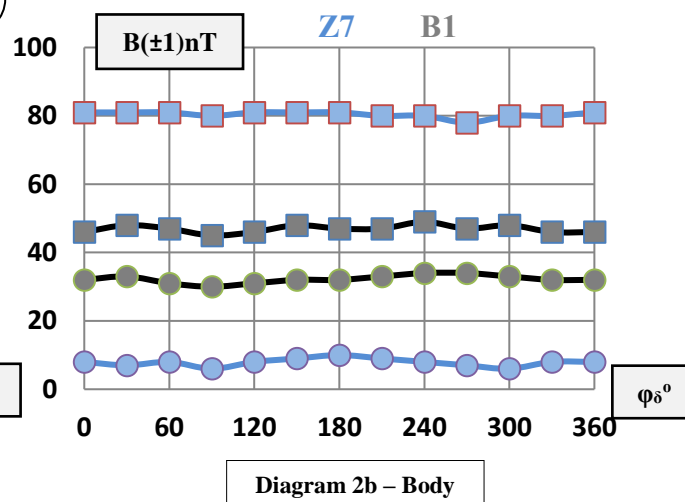
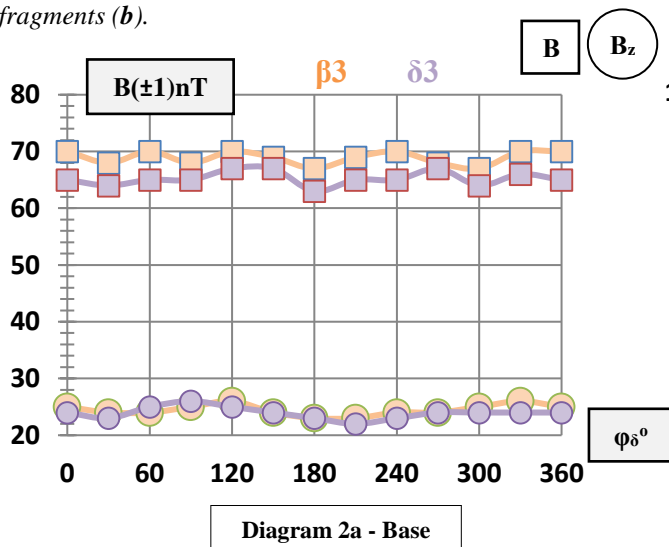
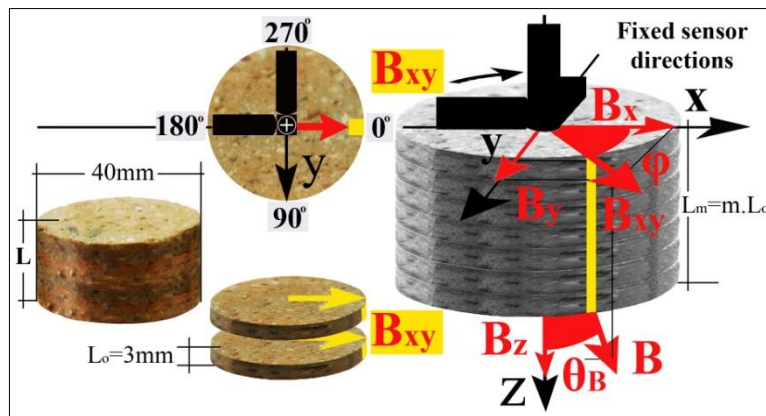


Diagram B.2 Change in the measured component B_z and computed magnitude, B , of the magnetic field at disk rotation angles φ_δ , in 2 fragments (β_3, δ_3) of the base (a) and 2 fragments ($B1, Z7$) of the body (b) of vase I.

At the rotation positions of the base and body fragments (table B.1), the computed angle φ is close to the measured disk rotation angle φ_δ . The sensors do not take measurements from the total shell area but rather from a restricted, rotationally symmetric area around the z-axis, which can be considered magnetically isotropic. For this reason irregularly shaped fragments from the vase and neighboring body fragments are oriented in a manner that they join.

In contrast to the co-belonging, and of similar thickness, base fragments (**diagrams B.1,B.2**), the values of B_z , B , θ_B vary in the vase body. The large deviations in the magnetic field magnitude of the co-belonging body fragments cannot be attributed to the differences in the thickness or the magnetic anisotropy of the ceramic material, since in the case of base fragments, of the same material and indeed of irregular shape, it displays similar values.

B.2 Investigation of the contribution of the ceramic material induced magnetization in the magnetic field measurements



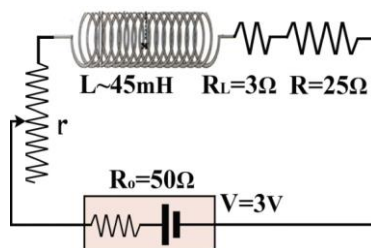
From the base fragments which display similar magnetic field, cylindrical specimens are constructed (**figure B.2**) and sliced into equilateral tiles. With B_{xy} oriented in the same direction, a tile column is assembled, with a magnetic field of equal magnitude and directionality to those of the parent specimen from the base of the vases **4,5 & 6**.

Figure B.2 Illustration of the experimental apparatus and measured components of the field B in the considered coordinate system by the three-dimensional orthogonal system layout of sensors.

Measurements are taken by the three-dimensional orthogonal system apparatus of sensors at the center of an oriented solenoid (**figure B.4**) in the direction of the geomagnetic field, where it is offset by the magnetic field of the coil. The axonal magnetic field in the coil center (**table B.2**) is set by a variable resistance, r , in an RL circuit (**figure B.3**), which is approximated [71 pp.223-6] by the formula:

$$B = \frac{\mu_o \cdot n \cdot V}{R_o + R + R_L + r} \cdot \left(1 + \frac{4b^2}{l^2}\right)^{-\frac{1}{2}} \quad (\text{B.1})$$

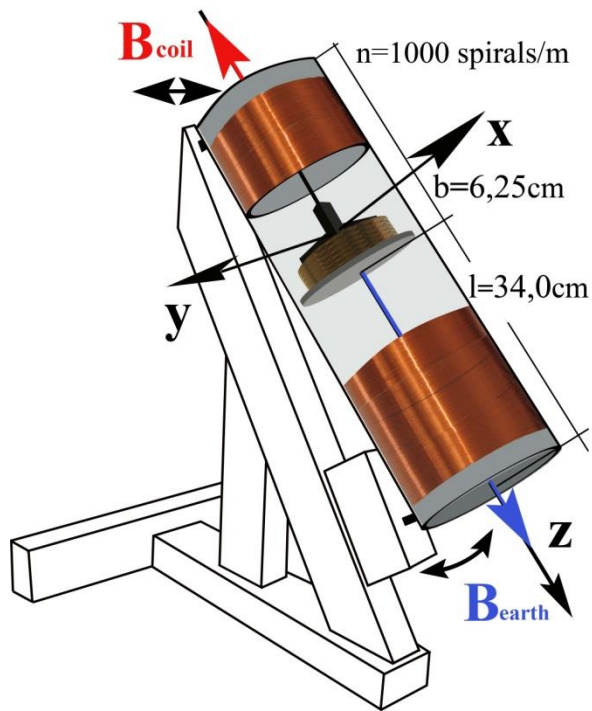
Table B.2 Offsetting ($\pm 1nT$) of the geomagnetic field by the axonal field of a solenoid, setting the resistance, r ($\pm m\Omega$), around the zeroing value of the total field in the center of the solenoid.



$r(\Omega)$	$R_{total}(\Omega)$	$B_{coil}(nT)$	$B_{earth}(nT)$	Total axonal field in the coil center (nT)
5.924	83.924	42161.4	42160.7	0.7
5.925	83.925	42160.9		0.2
5.926	83.926	42160.4		-0.3
5.927	83.927	42159.9		-0.8

Figure B.3. Elements of an RL circuit with variable resistance, r , functioning as a dimmer.

Measurements from the base specimen and the offspring tile column, at 30° rotation positions, are taken, from the position where B_{xy} is oriented ($B_{xy} > 0$, $B_y = 0$) in the x -sensor direction. Measurements are taken each time the thickness, $L_m = m \cdot L_o$, of the column is increased by one tile.



The measured magnetic field is only due to the remanent magnetization of the ceramic material, and it is not altered by its induced magnetization.

The measurements (**table B.3**) in a null field are similar to those of the next experiment (**diagram B.3**), taken from the same specimens by the three-dimensional orthogonal apparatus of sensors/magnets, in which the induced magnetization, M^* , of the ceramic material at the measurement area is oriented along the direction of the remanent magnetization ($15.4 \mu\text{T}$) of the earth and the magnets.

As the height of the column increases, the magnetic field increases and obtains a constant value after the addition of a variable number of tiles for each column, from the vase bases 4,5 &6.

The conclusions from the experiments of the dependence of the measured magnetic field on the shell thickness are given in the following experiment.

Figure B.4 Illustration of the measurement methodology of the cylindrical specimens magnetic field, inside a null field at the center of a solenoid oriented along the earth magnetic field.

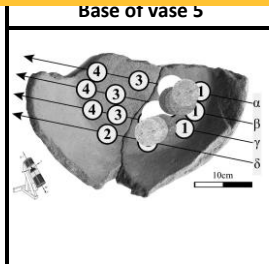


Table B.3 Indicative measurements of the magnetic field at the rotation positions of the parent specimen and of the column of thinner tiles, of increasing thickness, from the base of vase 5 inside a solenoid. The magnetic field obtains a constant value when the column thickness becomes greater than the shaded area of the table.



Base specimen		Magnetic field measurements from the cylindrical specimens at rotation positions $\phi_s (\pm 0,5^\circ)$ inside the coil														
L(± 1) mm		0°	30°	60°	90°	120°	150°	180°	210°	240°	270°	300°	330°			
	$B_x (\pm 1)$ nT	0	28	51	60	53	31	1	-31	-53	-60	-53	-32			
	$B_z (\pm 1)$ nT	16	16	17	16	17	16	17	17	16	17	17	17			
	$B (\pm 1)$ nT	63	61	62	62	62	62	62	62	62	62	62	62			
	$\theta_B (\pm 1^\circ)$	75	75	74	75	74	75	74	74	75	74	74	74			
$\phi (\pm 1^\circ)$	0	28	60	90	118	149	179	211	242	269	298	327				
Column of m tiles																
$L_m (\pm 0,2)$ mm	t	0°	30°	60°	90°	120°	150°	180°	210°	240°	270°	300°	330°			
m	3.0	$B_x (\pm 1)$ nT	43	37	21	1	-23	-39	-44	-36	-20	-2	21	38		
		$B_y (\pm 1)$ nT	0	22	39	43	36	22	-1	-23	-38	-45	-38	-21		
1	3.0	$B_z (\pm 1)$ nT	14	14	14	14	15	15	14	14	14	15	15			
		$B (\pm 1)$ nT	45	45	46	45	45	47	46	45	45	47	46			
		$\theta_B (\pm 1^\circ)$	72	72	72	72	71	71	72	72	72	73	71			
		$\phi (\pm 1^\circ)$	0	31	62	89	123	151	181	213	242	267	299	331		
		2	6.0	$B_x (\pm 1)$ nT	56	48	29	-1	-25	-48	-55	-46	-25	1	28	49
				$B_y (\pm 1)$ nT	0	27	47	55	48	27	1	-29	-49	-55	-47	-29
$B_z (\pm 1)$ nT	16			16	17	16	16	17	16	16	16	15	16	16		
$B (\pm 1)$ nT	58			57	58	57	56	58	57	57	57	57	57	59		
$\theta_B (\pm 1^\circ)$	74			74	73	74	74	73	74	74	74	75	74	74		
$\phi (\pm 1^\circ)$	0			29	58	91	118	151	179	212	243	271	301	329		
3	9.0	$B_x (\pm 1)$ nT	59	51	28	0	-29	-51	-59	-50	-31	1	28	53		
		$B_y (\pm 1)$ nT	0	28	52	59	50	30	-1	-31	-50	-59	-53	-28		
		$B_z (\pm 1)$ nT	16	16	16	16	16	16	16	16	17	17	16	16		
		$B (\pm 1)$ nT	61	60	61	61	60	61	61	61	61	61	62	62		
		$\theta_B (\pm 1^\circ)$	75	75	75	75	75	75	75	75	74	74	75	75		
		$\phi (\pm 1^\circ)$	0	29	62	90	120	150	181	212	238	271	298	332		
4	12.0	$B_x (\pm 1)$ nT	60	51	31	-1	-29	-50	-60	-52	-30	0	31	51		
		$B_\psi (\pm 1)$ nT	0	32	51	59	52	30	-2	-28	-51	-60	-50	-31		
		$B_z (\pm 1)$ nT	17	16	16	17	17	16	16	16	17	16	16	16		
		$B (\pm 1)$ nT	62	62	62	61	62	60	62	61	62	62	61	62		
		$\theta_B (\pm 1^\circ)$	74	75	75	74	74	75	75	75	74	75	75	75		
		$\phi (\pm 1^\circ)$	0	32	59	-89	119	149	182	208	240	270	302	329		
5	15.0	$B_x (\pm 1)$ nT	60	51	27	1	-30	-51	-51	-53	-30	1	28	52		
		$B_y (\pm 1)$ nT	0	31	53	59	52	31	-1	-29	-52	-61	-53	-29		
		$B_z (\pm 1)$ nT	16	16	17	16	16	16	17	17	17	17	17	16		
		$B (\pm 1)$ nT	62	62	62	61	62	62	54	63	62	63	62	62		
		$\theta_B (\pm 1^\circ)$	75	75	74	75	75	75	72	74	74	74	74	75		
		$\phi (\pm 1^\circ)$	0	31	63	89	120	149	181	209	240	271	298	331		
6	18.0	$B_x (\pm 1)$ nT	60	51	31	-1	-30	-52	-60	-51	-30	0	29	51		
		$B_y (\pm 1)$ nT	0	30	51	60	52	29	0	-31	-51	-60	-52	-32		
		$B_z (\pm 1)$ nT	16	16	16	17	16	17	16	16	16	17	17	17		
		$B (\pm 1)$ nT	62	61	62	62	62	62	62	62	61	62	62	63		
		$\theta_B (\pm 1^\circ)$	75	75	75	74	75	74	75	75	75	74	74	74		
		$\phi (\pm 1^\circ)$	0	30	59	91	120	151	180	211	240	270	299	328		
7	21.0	$B_x (\pm 1)$ nT	60	51	28	-1	-28	-51	-60	-52	-31	1	32	51		
		$B_y (\pm 1)$ nT	0	31	51	60	52	32	-1	-30	-51	-60	-51	-30		
		$B_z (\pm 1)$ nT	17	16	16	16	16	17	16	16	16	16	16	17		
		$B (\pm 1)$ nT	62	62	60	62	61	63	62	62	62	62	62	62		
		$\theta_B (\pm 1^\circ)$	74	75	75	75	75	74	75	75	75	75	75	74		
		$\phi (\pm 1^\circ)$	0	31	61	91	118	148	181	210	239	271	302	330		
8	24.0	$B_x (\pm 1)$ nT	60	52	29	1	-27	-53	-60	-51	-31	-1	28	51		
		$B_y (\pm 1)$ nT	0	30	52	60	53	28	-1	-32	-51	-59	-53	-32		
		$B_z (\pm 1)$ nT	16	16	17	16	16	17	16	16	17	17	17	16		
		$B (\pm 1)$ nT	62	62	62	62	62	62	62	62	62	61	62	62		
		$\theta_B (\pm 1^\circ)$	75	75	74	75	75	74	75	75	74	74	74	75		
		$\phi (\pm 1^\circ)$	0	30	61	269	117	152	181	212	239	269	298	328		

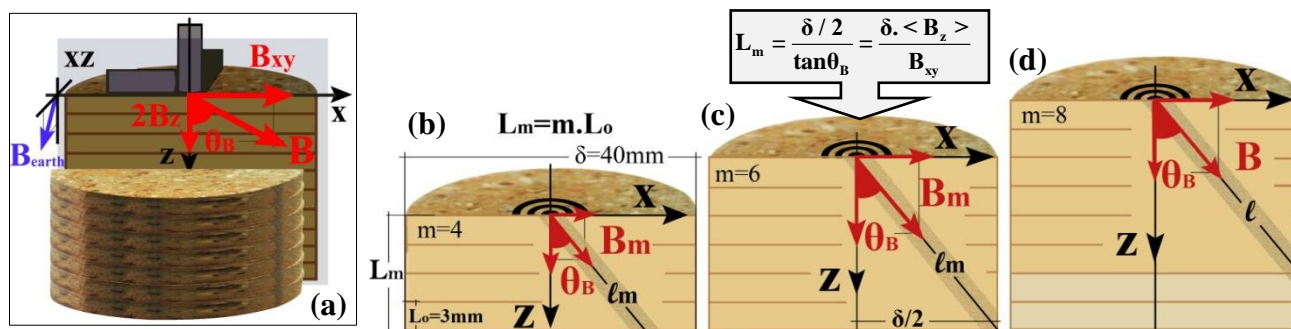


Figure B.6. For the geometric interpretation (b-d) of the results, B_{xy} is drawn (a) in the direction of the x-sensor.

The measured thickness, $L_m = m.L_o$, of each tile column from the vases 4,5 & 6 where the magnetic field magnitude B stabilizes (diagram B.4) is approximated geometrically by the column diameter and the sensor readings (table B.5), if it is assumed that the cylindrical sensor excitation area extends along the direction of the field B (figure B.6) and the reading of the vertical z -sensor doubles. The magnitude of the magnetic field increases with the increase of the column thickness, L_m , (figures B.6b,c) and stabilizes when the inclusion of additional tiles (figure B.6d) does not alter the height, l_m , of the cylindrical excitation area of the sensors.

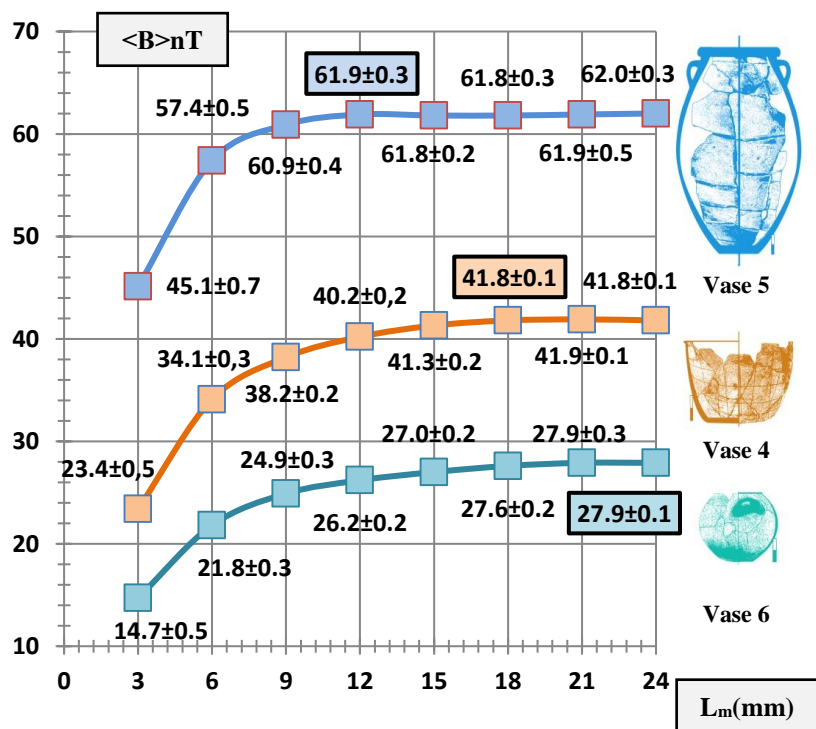


Diagram B.4 Change in average value, $\langle B \rangle$, of the magnitude of the magnetic field B with respect to the increase of the tile column thickness, $L_m = m.L_o$, from the vase bases 4,5 & 6. The average value, $\langle B \rangle$, is computed from the sensor readings at the 12 rotation positions of the tile column. The average value, $\langle B \rangle$, stabilizes, and does not change with the further addition of tiles, at the values indicated inside the colored frames.

With this view, sensor measurements are not only dependent on the thickness of the fragments, but also on the sensor locations, as confirmed in the next experiment.

Table B.5 Comparison of the measured and computed thickness $L_m = \frac{\delta/2}{\tan\theta_B} = \frac{\delta/2 \cdot \langle B_z \rangle}{2 \cdot B_{xy}}$ of the column from the geometry and the average value, $\langle B_z \rangle$, of the sensor readings at the 12 rotation positions of the column.

Vase	m	Measured thickness $L_m = m.L_o$ (± 0.5 mm)	B_{xy} (± 0.2 nT)	$\langle B_z \rangle$ (nT)	Computed thickness L_m (mm)
4	6	18.0	38.0	17.6 ± 0.1	18.5 ± 0.2
5	4	12.0	59.8	16.4 ± 0.1	11.1 ± 0.2
6	8	24.0	24.2	14.0 ± 0.2	23.1 ± 0.5

B.4 Investigation of the dependence of the sensor excitation area on their position on the surface of the shells and the directionality of their magnetization

Using the three-dimensional orthogonal sensor/magnets apparatus (**figure B.7a**), at the same position with respect to the geomagnetic field, measurements on the cylindrical specimens from the bases of the vases **4, 5 & 6** are taken, at distances D_n from their edges and along B_{xy} , which is oriented along the direction of the x -sensor, when $B_x = B_{xy} > 0$ and $B_y = 0$.

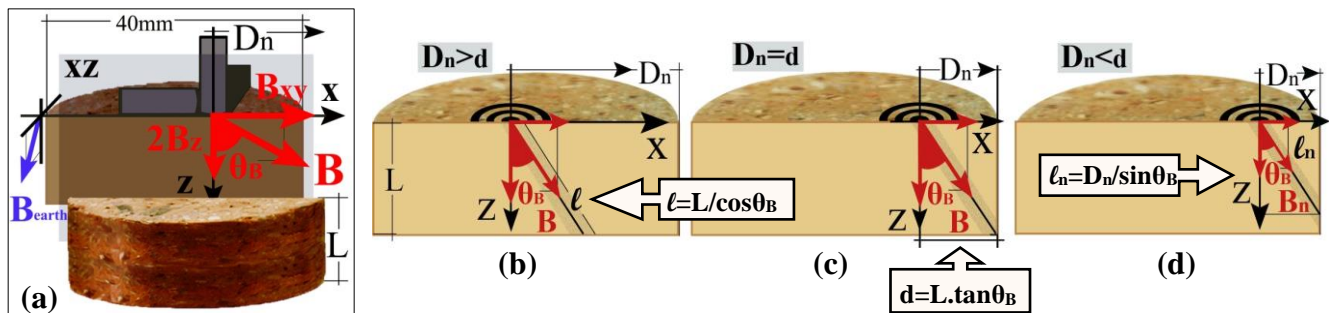
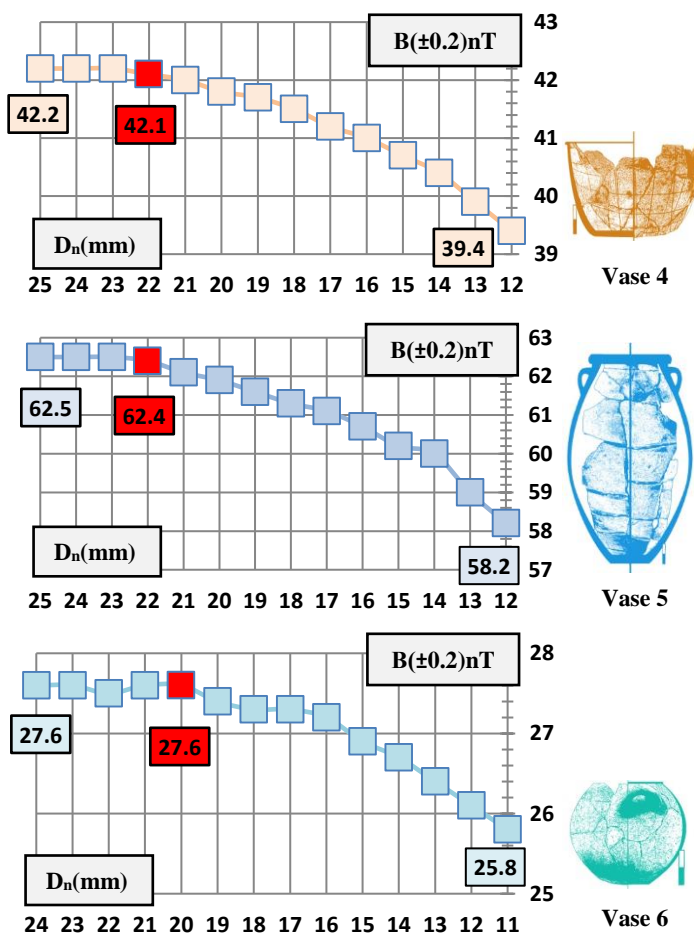


Figure B.7 Geometric interpretation (β - δ) of the experimental results.



At position $D_n=d$ (**figure B.7c**), the measured distance, D_n , for each specimen from the vase base is approximated by its thickness, L , and the readings of the magnetic field (**table B.6**), when the vertical z -sensor reading is doubled. The height, l_n , of the sensors' cylindrical area depends on the directivity (θ_B) of the magnetic field B and from their position (**diagram B.5**) on the surface:

-When $D_n > d = L \cdot \tan\theta_B$ (**figure B.7b**) the height depends on the thickness, L , of the specimen, and the measure B does not change, because the dimensions of the excitation area remain constant.

-When $D_n < d$ (**figure B.7d**) the height, l_n , is dependent on the distance D_n , and field B decreases as the sensors move towards the end of the sample, because the height of the stimulation are of the sensor is reduced.

Diagram B.5 Change of magnitude B versus distance D_n in base vases **4,5 & 6**. From the values in the red frame, B decreases as the distance D_n decreases.

Table B.6 Comparison of the measured and computed distances, $D_n = d = L \cdot \tan \theta_B = L \cdot B_{xy} / (2 \cdot B_z)$, from the specimen geometry and the sensor measurements.

Vase	L(±0.2)mm	$D_n = d$ (±0.2)mm (Measured)	B_{xy} (±0.2)nT	B_z (±0.2)nT	B (±0.2)nT	θ°	d(mm) (Computed)
4	20.0	22.0	38.2	17.7	42.1	65.1±0.3	21.6±0.3
5	12.0	22.0	60.1	16.6	62.4	74.6±0.2	21.7±0.3
6	21.0	20.0	24.2	13.2	27.6	61.4±0.5	19.3±0.3

B.5 Investigation of the dependence of the measured magnetic field on the sensor orientation on the fragments' surface. Theoretical approach of the sensor excitation area

From the co-belonging base fragments which display similar magnetic field, cylindrical tiles are constructed (**figure B.8**). With B_{xy} oriented at the same direction, a specimen is assembled, with a magnetic field of similar magnitude and directionality and with a diameter equal to its height.

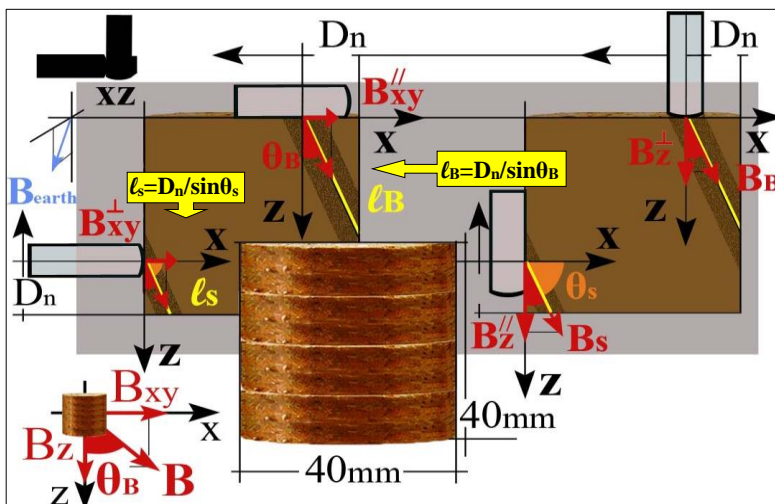


Figure B.8 Illustration of the experimental methodology for measuring the components (B_{xy} , B_z) of the field B by the vertical and horizontal sensors at equal distances, D_n , from the edges, from the base and on the side of specimens from the bases of vases 4, 5 & 6.

During measurements, the three-dimensional orthogonal system of sensors/magnets remains in a fixed position, with the x -sensor oriented perpendicularly to the vertical and horizontal component of the geomagnetic field.

The measurements are taken in positions at the base (B) and side (s) of each specimen and at equal measured distances D_n from its edge, on the oriented axial plane xz at the common direction of B_{xy} and the x -sensor, when $B_x = B_{xy} > 0$ and $B_y = 0$.

Since $D_n < d = L \cdot \tan \theta_B$, the heights, $l_B = \frac{D_n}{\sin \theta_B} \neq l_s = \frac{D_n}{\sin \theta_s}$, of the cylindrical sensor excitation areas

and the measured field B are increased at the equally-spaced measurement positions at the base and side of each specimen, with increasing distance D_n .

- If the directionality of the field is constant throughout the ceramic mass, as the result of its uniform magnetization, then the computed angles at the measurement positions of the base (θ_B) should display similar values. The same must be true for the angles (θ_s) at the specimen sides.

- If, in addition, the differently oriented sensors on the surface of the specimens are excited to the same extent, then due to the fact that the same component of the field is measured by the vertical and parallel sensors at the equally-distanced positions, by D_n , on the base and on the side of the specimens, the computed angles θ_B , θ_s must be complementary. In this case, the deviations in the value of the sum $\theta_B + \theta_s \approx 90^\circ$ can be computed, since they are due solely to the different height, $l_B \neq l_s$ of the cylindrical excitation areas of the sensors, which change in measurable way.

Under the above conditions, the directionality of the magnetic field in the ceramic material mass is tested and the degree of excitation of the vertical and parallel sensors from measurements of the same component of the field is compared (**table B.7**), at measuring locations equally-spaced from the edges of the specimen, in order to normalize their readings.

Table B.7 Indicative measurements of the same component, (B_{xy} , B_z), of the field by the vertical (\perp) and the parallel (\parallel) sensor at equal-distance locations, by D_n from the edges at the base (B) and sides (s) of the specimen from the base of vase

VASE 5		Measurements by the vertical and parallel sensors at the same position													
Table (a)															
L=40.4±0.2 (mm)		Vase(b)						Side(s)							
n	D_n (mm) (±0.2)	B_{xy}^{\parallel} (nT) (±0.2)	B_z^{\perp} (nT) (±0.2)	B_B (nT) (±0.2)	θ_B (±0.2°)	B_{xy}^{\perp} (nT) (±0.2)	B_z^{\parallel} (nT) (±0.2)	B_s (nT) (±0.2)	θ_s (±0.3°)	λ_{xy}	$\langle \lambda_{xy} \rangle$ (±0.01)	λ_z	$\langle \lambda_z \rangle$ (±0.01)		
1	14.0	56.9	15.9	59.1	74.4	29.2	31.4	42.9	47.1	1.95	1.98	1.97	1.99		
2	15.4	58.1	16.1	60.3	74.5	29.5	31.7	43.3	47.1	1.97		1.97			
3	16.8	58.6	16.2	60.8	74.5	29.7	32.0	43.7	47.1	1.97		1.98			
4	18.2	59.1	16.2	61.3	74.7	29.8	32.4	44.0	47.4	1.98		2.00			
5	19.6	59.2	16.3	61.4	74.6	29.9	32.5	44.2	47.4	1.98		1.99			
6	21.0	59.7	16.3	61.9	74.7	30.0	32.6	44.3	47.4	1.99		2.00			
7	22.4	59.9	16.4	62.1	74.7	30.0	32.8	44.5	47.6	2.00		2.00			
8	23.8	60.2	16.4	62.4	74.8	30.1	32.9	44.6	47.5	2.00		2.01			
9	25.2	60.4	16.5	62.6	74.7	30.2	33.2	44.9	47.7	2.00		2.01			
10	26.6	60.6	16.5	62.8	74.8	30.2	33.3	45.0	47.8	2.01		2.02			
Measurements by the same oriented sensors on the base(b) and on the side(s) of specimens															
Table (b)		(b)	(s)			(s)	(b)								
n	D_n (mm) (±0.2)	B_{xy}^{\parallel} (nT) (±0.2)	B_z^{\perp} (nT) (±0.2)	B^{\parallel} (nT) (±0.2)	θ^{\parallel} (±0.2°)	$\langle \theta^{\parallel} \rangle$ (±0.1°)	d^{\parallel} (mm) (±0.4)	B_{xy}^{\perp} (nT) (±0.2)	B_z^{\parallel} (nT) (±0.2)	B^{\perp} (nT) (±0.2)	θ^{\perp} (±0.3°)	$\langle \theta^{\perp} \rangle$ (±0.0°)	d^{\perp} (mm) (±0.4)		
1	14.0	56.9	31.4	65.0	61.1	61.3°	73.7	29.2	15.9	33.2	61.4	61.4°	74.1		
2	15.4	58.1	31.7	66.2	61.4			29.5	16.1	33.6	61.4				
3	16.8	58.6	32.0	66.8	61.4			29.7	16.2	33.8	61.4				
4	18.2	59.1	32.4	67.4	61.3			29.8	16.2	33.9	61.5				
5	19.6	59.2	32.5	67.5	61.2			29.9	16.3	34.1	61.4				
6	21.0	59.7	32.6	68.0	61.4			30.0	16.3	34.1	61.5				
7	22.4	59.9	32.8	68.3	61.3			30.0	16.4	34.2	61.3				
8	23.8	60.2	32.9	68.6	61.3			30.1	16.4	34.3	61.4				
9	25.2	60.4	33.2	68.9	61.2			30.2	16.5	34.4	61.3				
10	26.6	60.6	33.3	69.1	61.2			30.2	16.5	34.4	61.3				
Corrected measurements															
Table (c)		Vase(b)					Side(s)								
n	D_n (mm) (±0.2)	B_{xy}^{\parallel} (nT) (±0.2)	B_z^{\perp} (nT) (±0.6)	B_B (nT) (±0.3)	θ_B	ℓ_B (mm)	B_{xy}^{\perp} (nT) (±0.9)	B_z^{\parallel} (nT) (±0.2)	B_s (nT) (±0.8)	θ_s (±0.2°)	ℓ_s (mm)				
1	14.0	56.9	31.7	65.1	60.9±0.5	16.0±0.2	58.0	31.4	65.9	28.4	29.4±0.5				
2	15.4	58.1	32.1	66.4	61.1±0.5	17.6±0.2	58.6	31.7	66.6	28.4	32.3±0.5				
3	16.8	58.6	32.3	66.9	61.1±0.4	19.2±0.2	58.9	32.0	67.1	28.5	35.2±0.5				
4	18.2	59.1	32.3	67.4	61.3±0.4	20.7±0.2	59.1	32.4	67.4	28.7	37.9±0.5				
5	19.6	59.2	32.5	67.5	61.2±0.4	22.4±0.2	59.3	32.5	67.7	28.7	40.8±0.5				
6	21.0	59.7	32.5	68.0	61.4±0.4	23.9±0.2	59.5	32.6	67.9	28.7	43.7±0.5				
7	22.4	59.9	32.7	68.3	61.4±0.4	25.5±0.3	59.5	32.8	68.0	28.8	46.4±0.5				
8	23.8	60.2	32.7	68.5	61.5±0.4	27.1±0.3	59.7	32.9	68.2	28.8	49.3±0.5				
9	25.2	60.4	32.9	68.8	61.4±0.4	28.7±0.3	59.9	33.2	68.5	29.0	52.0±0.5				
10	26.6	60.6	32.9	69.0	61.5±0.4	30.3±0.3	59.9	33.3	68.6	29.1	54.8±0.5				
Table (d)		n	D_n (mm)	B_{xy}^{\perp} (nT) (±0.2)	B_z^{\parallel} (nT) (±0.2)	M_z (mA/m)	M_{xy} (mA/m)	θ°							
		10	26.6±0.2	30.2	16.5	26.3±0.3	48.1±0.3	61.3±0.3							

The results of the measurements in the above table show:

❖ The values of the computed angles (**table B.7a**) at the measurement positions, at the equally spaced positions by D_n , from the base (θ_B) and the side (θ_s) of the specimen show similar values. The directionality of the field \mathbf{B} is constant but the degrees of sensor excitation ($\theta_B + \theta_s \neq 90^\circ$) are different.

The parallel sensor ($B_{xy}^//, B_z^//$) readings, for the measurement of the same component of the field, are approximately double of those of the vertical (B_{xy}^\perp, B_z^\perp) sensor, which displays lower growth rate and range of values **for each specimen from the vases 4,5& 6**. For the correction of the readings, the ratios $\lambda_{xy} = B_{xy}^// / B_{xy}^\perp, \lambda_z = B_z^// / B_z^\perp$ and the average values $\langle \lambda_{xy} \rangle \approx \langle \lambda_z \rangle \approx \lambda = 2$ are computed.

• When, in the equally-distanced measuring positions at the base and side (**table B.7b**), the magnitude ($B^//, B^\perp$) and directionality ($\theta^//, \theta^\perp$) of the field are computed from only the readings of the parallel ($//$) or only the vertical (\perp) sensor, it can be seen that $B^// \neq B^\perp$ and $\theta^// \approx \theta^\perp$. The parallel and vertical sensor readings for the measurement of the same field component are differentiated by the same factor, λ , such that:

$$\tan^{-1} \frac{B_{xy}^//}{B_z^//} = \theta^// \approx \theta^\perp = \tan^{-1} \frac{\lambda \cdot B_{xy}^\perp}{\lambda \cdot B_z^\perp}$$

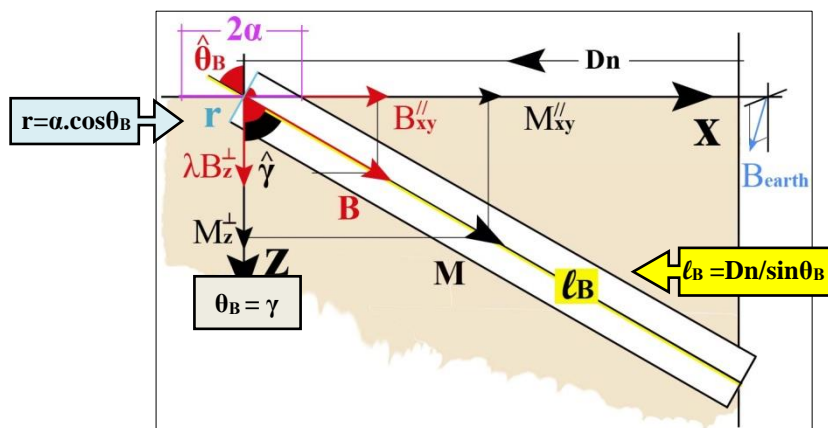
From the average values $\langle \theta^// \rangle, \langle \theta^\perp \rangle$, the maximum distances, $d^{//\perp} = L \cdot \tan(\langle \theta^{//\perp} \rangle)$, of the sensor positions from the edge of the fragments are computed such that the excitation range extends over the entire thickness of the specimen. Since $D_n < d^{//\perp}$, the magnitude of the magnetic field increases with the distance D_n .

• After the correction ($*B_z^\perp = \langle \lambda_z \rangle \cdot B_z^\perp, *B_{xy}^\perp = \langle \lambda_{xy} \rangle \cdot B_{xy}^\perp$) of the readings B_z^\perp, B_{xy}^\perp of the vertical sensor in the equally-spaced positions on the base and side of the specimen (**table B.7c**), it can be seen that $\theta_B + \theta_s \approx 90^\circ$.

• However, the magnitude of the field at the same positions displays similar values ($B_B \approx B_s$) while the heights, $l_B = \frac{D_n}{\sin \theta_B}, l_s = \frac{D_n}{\sin \theta_s}$, of the cylindrical sensor excitation areas are different.

This is because the radii, $r_B = \alpha \cdot \cos \theta_B, r_s = \alpha \cdot \cos \theta_s$, of the cylindrical sensor excitation areas on the base and the side of the specimen, also depend on the directionality, (θ_B, θ_s), of the sensor and also because

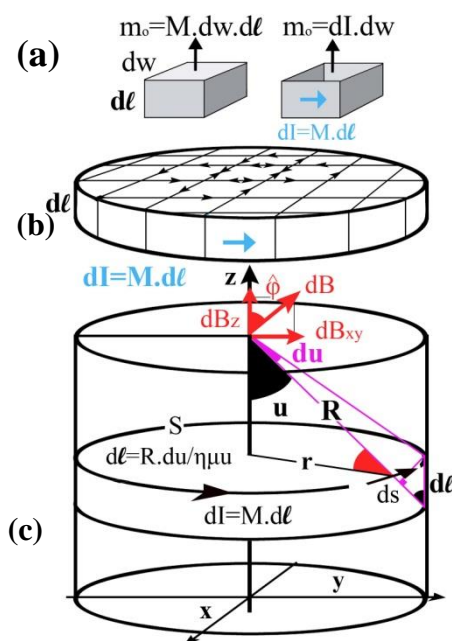
$$\frac{r_B}{l_B} = \frac{\alpha \cdot \sin \theta_B \cdot \cos \theta_B}{D_n} = \frac{\alpha \cdot \sin \theta_s \cdot \cos \theta_s}{D_n} = \frac{r_s}{l_s}, \text{ since } \theta_B + \theta_s \approx 90^\circ, \text{ also holds.}$$



❖ The magnetic field at the equally-distanced positions on the base and sides of the specimen displays similar values because the ratios $r_B/l_B, r_s/l_s$ of the cylindrical sensor excitation areas do not change, as does the measured axial field at the end of the solenoids with the same current and number of spirals, where the ratio of the radius, r , to their height is kept constant.

Figure B.9 Illustration of the cylindrical area of sensor excitation at measurement locations on the base of specimens and along the B_{xy} , which is oriented in the x -sensor direction. The diameter, $2r = 2\alpha \cdot \cos \theta_B$, depends on the directionality, θ_B , of the field \mathbf{B} and the parameter α , which, geometrically, is the sensitivity radius of the sensors around the measurement position. As can be seen below, the sensitivity radii of the vertical and parallel sensors are different and the angle, γ , of the magnetization, \mathbf{M} , is not the same as the angle θ_B of the field \mathbf{B} .

The cylindrical sensor [71 pp.413-8] excitation area (**figure B.9**), of radius r and height $\ell = \ell_B$, can be divided into elemental pieces of area dw and thickness $d\ell$ or equivalent loops (**figure B.10a**) with magnetic bipolar moments $m_o = M \cdot dw \cdot d\ell = dI \cdot dw$ along the axial direction, with uniform magnetization $M = dI / d\ell$.



Due to the canceling out of opposite currents from adjacent tiles (**figure B.10b**), each thin disk, of radius r and height $d\ell$, with magnetic bipolar moment $m = dI \cdot (\pi \cdot r^2) = M \cdot (\pi \cdot r^2) \cdot d\ell$, is equivalent to a current loop with an axial magnetic field approximated by the *Biot-Savart law*.

At every point of the z -axis (**figure B.10c**), the dB_{xy} components cancel each other out due to the rotational symmetry, and the \mathbf{B} field has only the axial z -component, which is approximated [71 pp. 223-6] by the individual contributions, dB , of the elementary curve segments, ds , in each tile of thickness $d\ell$:

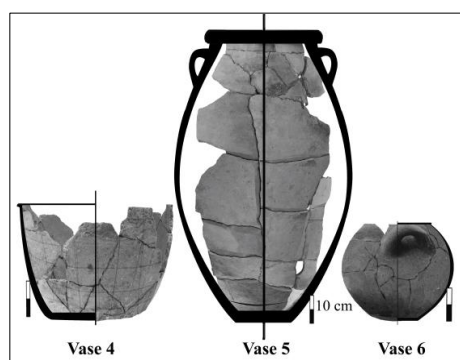
$$dB_z = dB \cdot \cos\phi = \frac{\mu_o \cdot dI \cdot d\vec{s} \times \hat{R}}{4\pi R^2} \cdot \frac{r}{R} = \frac{\mu_o \cdot (M \cdot d\ell) \cdot r}{4\pi R^3} \cdot ds \quad (\mathbf{B.2})$$

Figure B.10 The measured magnetic field \mathbf{B} of the cylindrical sensor excitation is approximated by the magnetic field of a solenoid with the same dimensions.

The axial magnetic field B at the end of a cylinder of radius $r = R \sin u$ and height ℓ is the result of contributions from individual tiles of thickness $d\ell = R \cdot du / \sin u$ and it is approximated by the integral:

$$B = B_z = \frac{\mu_o \cdot M \cdot r}{4\pi R^3} \cdot \int_{s=2\pi} ds \int d\ell = \frac{\mu_o \cdot M \cdot r^2}{2R^3} \cdot \int \frac{R \cdot du}{\sin u} = \frac{\mu_o \cdot M}{2} \sin u \cdot du = \frac{\mu_o \cdot M}{2\sqrt{1 + (\frac{r}{\ell})^2}} \quad (\mathbf{B.3})$$

where $r = \alpha \cdot \cos \gamma$ και $\ell = D_n / \sin \gamma$



B.6 Computation of the magnetization in specimens from the base of the vases 4,5 & 6

For the calculation of the magnetization components according to the above considerations (**figure B.11**), knowledge of the deviation, γ , of magnetization from the z -axis, which does not necessarily coincide with the angle of deviation, θ_B , of the measured magnetic field \mathbf{B} is necessary.

The sensitivity radius, α'' , of the parallel sensor is larger than that of the vertical sensor, α' (**table B.7a**), as evidenced by its readings, $(B_z^\perp, B_{xy}^\perp)$, which show smaller range and growth rate.

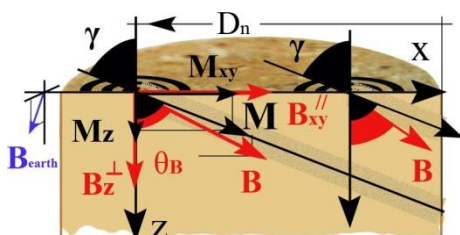


Figure B.11 Illustration of the cylindrical sensor stimulation area in measurement locations at the base of the specimens and along B_{xy} , which is oriented in the x -sensor direction. Since the sensitivity radius, α_z'' , of the vertical sensor is smaller than the radius of the parallel sensor, α_{xy}'' , the angle, γ , of magnetization M with the z -axis, deviates from the angle $\theta_B < \gamma$ of the measured field \mathbf{B} , as the distance, D_n , of the measurement position from the edge of the specimen decreases.

This is because, according to the above considerations, the measured components of the field, B_i , are approximated by the components of the magnetization, M_i , by the relation:

$$B_i = \frac{\mu_o \cdot M_i}{2 \sqrt{1 + \left(\frac{\alpha_i \cdot \sin 2\gamma}{2D_n} \right)^2}}, i=xy,z \quad (\text{B.4})$$

Since $\alpha^t < \alpha^//$, the readings, (B_z^t, B_{xy}^t) , of the vertical sensor tend to the limit values $B_i^t = \mu_o \cdot M_i / 2$ ($i=xy,z$), more quickly (**table B.7d**) the ratio $\alpha^t \cdot \sin 2\gamma / D_n$ in the previous relation is zeroing faster, since the heights, $\ell_B = D_n / \sin \gamma$ and $\ell_s = D_n / \sin(90^\circ - \gamma)$, of the sensor excitation areas on the base and sides of the specimen increase with the distance D_n , and hold:

$$\text{Vase (B): } \tan \theta_B = \frac{B_{xy}^//}{\lambda B_z^t} < \tan \gamma \cdot \sqrt{\frac{1 + \left(\frac{\alpha_z^t \cdot \sin 2\gamma}{2D_n} \right)^2}{1 + \left(\frac{\alpha_{xy}^// \cdot \sin 2\gamma}{2D_n} \right)^2}}, \text{ Side (s): } \tan \theta_s = \frac{B_z^//}{\lambda B_{xy}^t} < \tan(90^\circ - \gamma) \cdot \sqrt{\frac{1 + \left[\frac{\alpha_{xy}^t \cdot \sin 2(90^\circ - \gamma)}{2D_n} \right]^2}{1 + \left[\frac{\alpha_z^// \cdot \sin 2(90^\circ - \gamma)}{2D_n} \right]^2}} \quad (\text{B.5})$$

The value of the angle $\gamma \approx \theta^t = \tan^{-1}(B_{xy}^t / B_z^t)$, which, in the case of the greatest distance, $D_{n=10}$, is approximated by the readings of the vertical sensor only (**table B.7b**), and the formulas for computing the magnetization components from the magnetic field measurements, are used in the least squares method, for the measurements ($i=z,xy$), by the vertical, B_i^t , and the parallel, $B_i^//$, sensor at the equally-spaced positions from the edge by D_n , at the base and side of the specimens:

$$\text{Measurements at the base of the specimens: } \left(\frac{\mu_o}{2B_i^//} \right)^2 = \frac{(\alpha_i^//)^2 \cdot \sin^2 2\gamma}{4 \cdot (M_i^//)^2} \cdot \frac{1}{D_n^2} + \frac{1}{(M_i^//)^2}, (i=z,xy) \quad (\text{B.9})$$

$$\text{Measurements at the side of the specimens: } \left(\frac{\mu_o}{2B_i^t} \right)^2 = \frac{(\alpha_i^t)^2 \cdot \sin^2 2\gamma}{4 \cdot (M_i^t)^2} \cdot \frac{1}{D_n^2} + \frac{1}{(M_i^t)^2}, (i=z,xy) \quad (\text{B.10})$$

Table B.8 Indicative measurements at the base ($B_{xy}^//, B_z^t$) and side ($B_{xy}^t, B_z^//$) of the specimen from the base of vase 5, from the vertical and parallel sensor.

Vase 5	D_n	B_z^t	B_{xy}^t	$B_z^//$	$B_{xy}^//$
n	($\pm 0.2\text{mm}$)	($\pm 0.2\text{nT}$)			
1	14.0	15.9	29.2	31.4	56.9
2	15.4	16.1	29.5	31.7	58.1
3	16.8	16.2	29.7	32.0	58.6
4	18.2	16.2	29.8	32.4	59.1
5	19.6	16.3	29.9	32.5	59.2
6	21.0	16.3	30.0	32.6	59.7
7	22.4	16.4	30.0	32.8	59.9
8	23.8	16.4	30.1	32.9	60.2
9	25.2	16.5	30.2	33.2	60.4
10	26.6	16.5	30.2	33.3	60.6



Vase 5

The above view is confirmed by the linearity of relationships, and the sensitivity constants, $\alpha^t, \alpha^//$, the components of the magnetization, M_i , ($i=xy,z$), and the sensor normalization λ , are computed from the field measurements (**table B.7a, B.8**), with remarkable accuracy.

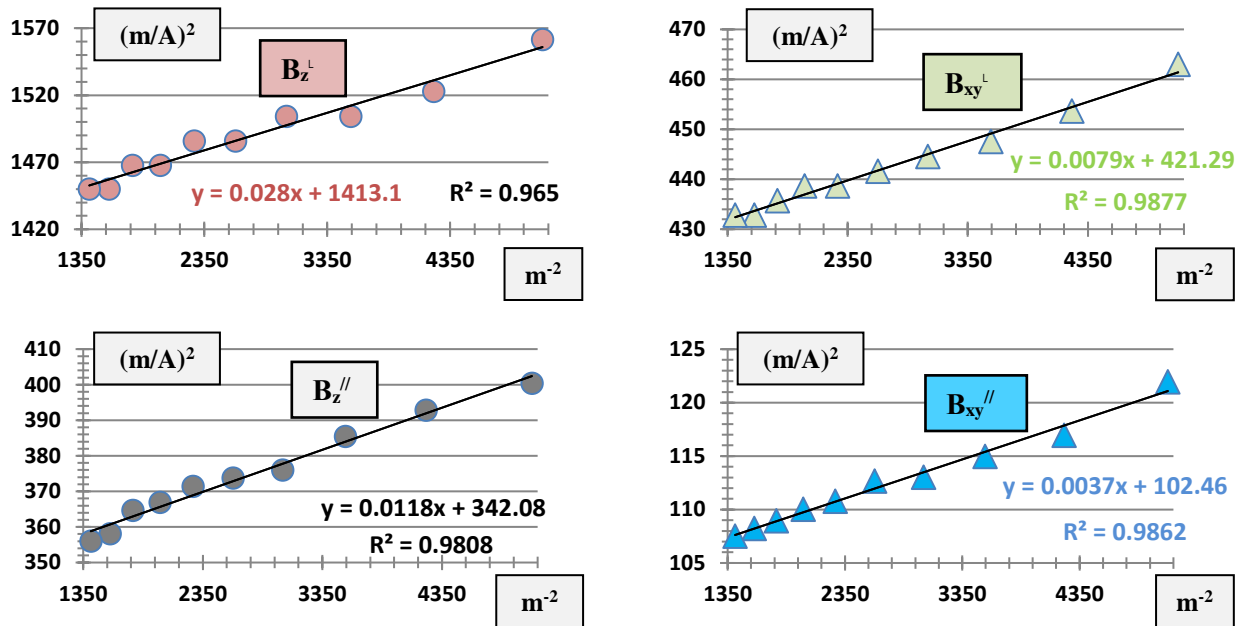
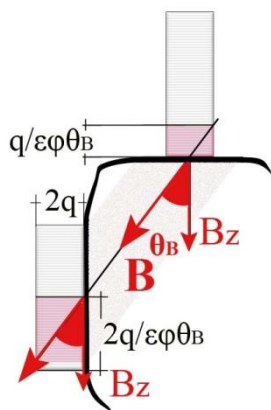


Diagram B.6 The proposed perspective of the sensors' mode and area of excitation is ascertained, using the least squares method, from the linearity of the measurements of the same-field component with the vertical and horizontal sensor at the equally-distant positions by D_n on the base and on the side of the specimen from the vase base 5.

Table B.9 The components $M_i^⊥, M_i^∥$ ($i=xy,z$) of the field are estimated using the least squares method from the vertical and horizontal sensor readings, where it is found that $M_{xy}^∥/M_{xy}^⊥ \approx M_z^∥/M_z^⊥ \approx 2 = \lambda$ and $\alpha_z^⊥ \approx \alpha_{xy}^⊥, \alpha_z^∥ \approx \alpha_{xy}^∥$, and the components, $M_z = 1/2 \cdot (M_z^∥ + \lambda M_z^⊥), M_{xy} = 1/2 \cdot (M_{xy}^∥ + \lambda M_{xy}^⊥)$, magnitude of magnetization, M , and sensitivity radii, $\alpha^⊥ = 1/2 \cdot (\alpha_{xy}^⊥ + \alpha_z^⊥), \alpha^∥ = 1/2 \cdot (\alpha_{xy}^∥ + \alpha_z^∥)$, of the vertical and parallel sensors are computed.

Vase 5	mA/m	M_z	M_{xy}	M	γ°	$\alpha_z^⊥$	1.06±0.04	$\alpha^⊥$ (cm)	1.04±0.03
$M_z^∥$	54.1±0.1	54.0	98.9	112.7	61.3	$\alpha_{xy}^⊥$	1.03±0.02		
$M_z^⊥$	26.6±0.01					$\alpha_z^∥$	1.40±0.04		
$M_{xy}^∥$	98.8±0.2	±0.2	±0.3	±0.3	±0.3	$\alpha_{xy}^∥$	1.42±0.04	$\alpha^∥$ (cm)	1.41±0.03
$M_{xy}^⊥$	48.7±0.1					λ	2.03±0.01		



The normalization factor $\lambda \approx 2$ displays similar values in base fragments from different vases and does not depend on the magnetization, M , of the ceramic material (**figure B.12**), but only on the different orientation of the excitation coils of the sensors vertical and parallel to the fragments' surface.

Figure B.12 The flow of the measured axial field B from the cylindrical excitation area is double at the parallel sensor, for the measurement of the same field component.

The quantities $\lambda, \alpha^⊥, \alpha^∥, M, M_i$ ($i=xy,z$) are also calculated (**tables B.10, B.11**) from the base specimens of vases 4 and 6 in the same way.

Table B.10a Indicative measurements of the same component, (B_{xy} , B_z), of the field by the vertical (\perp) and the parallel (\parallel) sensor at equal-distance locations, by D_n , from the edges at the base (B) and sides (s) of the specimen from the base of vase 4.

VASE 4		Measurements by the vertical and parallel sensors at the same position											
Table (a)													
L=40.5±0.2 (mm)													
		Vase(b)				Side(s)							
n	D_n (mm) (±0.2)	B_{xy}^{\parallel} (nT) (±0.2)	B_z^{\parallel} (nT) (±0.2)	B_B (nT) (±0.2)	θ_B (±0.1°)	B_{xy}^{\perp} (nT) (±0.2)	B_z^{\perp} (nT) (±0.2)	B_s (nT) (±0.2)	$(\theta_s)^{\circ}$	λ_{xy}	$\langle \lambda_{xy} \rangle$ (±0.02)	λ_z	$\langle \lambda_z \rangle$ (±0.02)
1	13.0	36.2	17.3	40.1	64.5	19.5	32.0	37.5	58.6±0.2	1.86	1.91	1.85	1.91
2	14.6	36.8	17.4	40.7	64.7	19.6	32.9	38.3	59.2±0.2	1.88		1.89	
3	16.2	37.1	17.5	41.0	64.7	19.6	33.2	38.6	59.4±0.2	1.89		1.90	
4	17.8	37.6	17.6	41.5	64.9	19.7	33.5	38.9	59.5±0.1	1.91		1.90	
5	19.4	37.9	17.6	41.8	65.1	19.8	33.7	39.1	59.6±0.1	1.91		1.91	
6	21.0	38.1	17.6	42.0	65.2	19.8	33.9	39.3	59.7±0.1	1.92		1.93	
7	22.6	38.3	17.7	42.2	65.2	19.8	34.0	39.3	59.8±0.1	1.93		1.92	
8	242	38.4	17.7	42.3	65.3	19.9	34.1	39.5	59.7±0.1	1.93		1.93	
9	25.8	38.7	17.7	42.6	65.4	20.0	34.2	39.6	59.7±0.1	1.94		1.93	
10	27.4	38.8	17.7	42.6	65.5	20.0	34.3	39.7	59.8±0.1	1.94		1.94	
Measurements by the same oriented sensors on the base(b) and on the side(s) of specimens													
Table (b)		(b)	(s)					(s)	(b)				
n	D_n (mm) (±0.2)	B_{xy}^{\parallel} (nT) (±0.2)	B_z^{\parallel} (nT) (±0.2)	B^{\parallel} (nT) (±0.2)	$(\theta^{\parallel})^{\circ}$	$\langle \theta^{\parallel} \rangle$ (±0.1°)	d^{\parallel} (mm) (±0.3)	B_{xy}^{\perp} (nT) (±0.2)	B_z^{\perp} (nT) (±0.2)	B^{\perp} (nT) (±0.2)	$(\theta^{\perp})^{\circ}$ (±0.3°)	$\langle \theta^{\perp} \rangle$ (±0.1°)	d^{\perp} (mm) (±0.3)
1	13.0	36.2	32.0	48.3	48.5±0.2	48.4	45.6	19.5	17.3	26.1	48.4	48.4	45.5
2	14.6	36.8	32.9	49.4	48.2±0.2			19.6	17.4	26.2	48.4		
3	16.2	37.1	33.2	49.8	48.2±0.2			19.6	17.5	26.3	48.2		
4	17.8	37.6	33.5	50.4	48.3±0.2			19.7	17.6	26.4	48.2		
5	19.4	37.9	33.7	50.7	48.4±0.2			19.8	17.6	26.5	48.4		
6	21.0	38.1	33.9	51.0	48.3±0.1			19.8	17.6	26.5	48.4		
7	22.6	38.3	34.0	51.2	48.4±0.1			19.8	17.7	26.6	48.2		
8	242	38.4	34.1	51.4	48.4±0.1			19.9	17.7	26.6	48.3		
9	25.8	38.7	34.2	51.6	48.5±0.1			20.0	17.7	26.7	48.5		
10	27.4	38.8	34.3	51.8	48.5±0.1			20.0	17.7	26.7	48.5		
Corrected measurements													
Table (c)		Vase(b)					Side(s)						
n	D_n (mm) (±0.2)	B_{xy}^{\parallel} (nT) (±0.2)	B_z^{\parallel} (nT) (±0.5)	B_B (nT) (±0.4)	θ_B (±0.5°)	ℓ_B (mm)	B_{xy}^{\perp} (nT) (±0.6)	B_z^{\perp} (nT) (±0.2)	B_s (nT) (±0.5)	θ_s (±0.5°)	ℓ_s (mm)		
1	13.0	36.2	33.0	49.0	47.6	17.6±0.3	37.3	32.0	49.1	40.6	20.0±0.4		
2	14.6	36.8	33.2	49.6	47.9	19.7±0.3	37.5	32.9	49.9	41.3	22.1±0.4		
3	16.2	37.1	33.4	49.9	48.0	21.8±0.3	37.5	33.2	50.1	41.5	24.4±0.4		
4	17.8	37.6	33.6	50.4	48.2	23.9±0.3	37.7	33.5	50.4	41.7	26.8±0.4		
5	19.4	37.9	33.6	50.7	48.4	25.9±0.3	37.8	33.7	50.7	41.7	29.2±0.4		
6	21.0	38.1	33.6	50.8	48.6	28.0±0.3	37.8	33.9	50.8	41.9	31.5±0.4		
7	22.6	38.3	33.8	51.1	48.6	30.1±0.3	37.8	34.0	50.9	41.9	33.8±0.4		
8	242	38.4	33.8	51.2	48.6	32.2±0.4	38.0	34.1	51.1	41.9	36.3±0.4		
9	25.8	38.7	33.8	51.4	48.9	34.3±0.4	38.2	34.2	51.3	41.8	38.7±0.5		
10	27.4	38.8	33.8	51.5	48.9	36.3±0.4	38.2	34.3	51.4	41.9	41.0±0.5		
Table (d)		n	D_n (mm)	B_{xy}^{\perp} (±0.2nT)	B_z^{\perp} (±0.2nT)	M_z (mA/m)	M_{xy} (mA/m)	θ°					
		10	27.4±0.2	20.0	17.7	28.2±0.3	31.8±0.3	48.5±0.3					

Table B.10b The quantities λ , α^{\perp} , α^{\parallel} , M , M_i ($i=xy,z$) are estimated using the least squares method from the above sensors readings.

Vase 4	mA/m	M_z	M_{xy}	M	γ°	α_z^{\perp}	0.66±0.02	α^{\perp} (cm)	0.67±0.03
M_z^{\parallel}	55.8±0.1	±0.2	±0.2	±0.2	±0.3	α_{xy}^{\perp}	0.68±0.04		
M_z^{\perp}	28.41±0.03					α_z^{\parallel}	1.13±0.03		
M_{xy}^{\parallel}	63.0±0.1					α_{xy}^{\parallel}	1.17±0.01		
M_{xy}^{\perp}	32.0±0.1					λ	1.97±0.01		

Table B.11a Indicative measurements of the same component, (B_{xy} , B_z), of the field by the vertical (\perp) and the parallel (\parallel) sensor at equal-distance locations, by D_n , from the edges at the base (B) and sides (s) of the specimen from the base of vase 6.

VASE 6		Measurements by the vertical and parallel sensors at the same position												
Table (a)														
L=41.2±0.2 (mm)														
		Vase(b)					Side(s)							
n	D_n (mm) (± 0.2)	B_{xy}^{\parallel} (nT) (± 0.2)	B_z^{\parallel} (nT) (± 0.2)	B_B (nT) (± 0.2)	θ_B ($\pm 0.2^\circ$)	B_{xy}^{\perp} (nT) (± 0.2)	B_z^{\perp} (nT) (± 0.2)	B_s (nT) (± 0.2)	θ_s ($\pm 0.2^\circ$)	λ_{xy}	$\langle \lambda_{xy} \rangle$ (± 0.02)	λ_z	$\langle \lambda_z \rangle$ (± 0.03)	
1	11.0	22.9	13.3	26.5	59.9	12.7	23.9	27.1	62.0	1.80		1.80		
2	13.0	23.4	13.6	27.1	59.8	12.8	24.2	27.4	62.1	1.83		1.78		
3	15.0	23.6	13.7	27.3	59.9	12.9	25.1	28.2	62.8	1.83		1.83		
4	17.0	24.0	13.7	27.6	60.3	12.9	25.4	28.5	63.1	1.86	1.85	1.85	1.85	
5	19.0	24.1	13.8	27.8	60.2	13.0	25.6	28.7	63.1	1.85		1.86		
6	21.0	24.3	13.8	27.9	60.4	13.0	25.8	28.9	63.3	1.87		1.87		
7	23.0	24.5	13.9	28.2	60.4	13.1	25.9	29.0	63.2	1.87		1.86		
8	25.0	24.6	13.9	28.3	60.5	13.1	26.0	29.1	63.3	1.88		1.87		
9	27.0	24.7	14.0	28.4	60.5	13.2	26.2	29.3	63.3	1.87		1.87		
10	29.0	24.8	14.0	28.5	60.6	13.2	26.3	29.4	63.3	1.88		1.88		

Measurements by the same oriented sensors on the base(b) and on the side(s) of specimens													
Table (b)		(b)	(s)				(s)	(b)				(s)	(b)
n	D_n (mm) (± 0.2)	B_{xy}^{\parallel} (nT) (± 0.2)	B_z^{\parallel} (nT) (± 0.2)	B^{\parallel} (nT) (± 0.2)	θ^{\parallel} ($\pm 0.2^\circ$)	$\langle \theta^{\parallel} \rangle$ ($\pm 0.2^\circ$)	d^{\parallel} (mm) (± 0.3)	B_{xy}^{\perp} (nT) (± 0.2)	B_z^{\perp} (nT) (± 0.2)	B^{\perp} (nT) (± 0.2)	θ^{\perp} ($\pm 0.3^\circ$)	$\langle \theta^{\perp} \rangle$ ($\pm 0.1^\circ$)	d^{\perp} (mm) (± 0.2)
1	11.0	22.9	23.9	33.1	43.8			12.7	13.3	18.4	43.7		
2	13.0	23.4	24.2	33.7	44.0			12.8	13.6	18.7	43.3		
3	15.0	23.6	25.1	34.5	43.2			12.9	13.7	18.8	43.3		
4	17.0	24.0	25.4	34.9	43.4	43,4	39,0	12.9	13.7	18.8	43.3	43.3	38.9
5	19.0	24.1	25.6	35.2	43.3			13.0	13.8	19.0	43.3		
6	21.0	24.3	25.8	35.4	43.3			13.0	13.8	19.0	43.3		
7	23.0	24.5	25.9	35.7	43.4			13.1	13.9	19.1	43.3		
8	25.0	24.6	26.0	35.8	43.4			13.1	13.9	19.1	43.3		
9	27.0	24.7	26.2	36.0	43.3			13.2	14.0	19.2	43.3		
10	29.0	24.8	26.3	36.1	43.3			13.2	14.0	19.2	43.3		

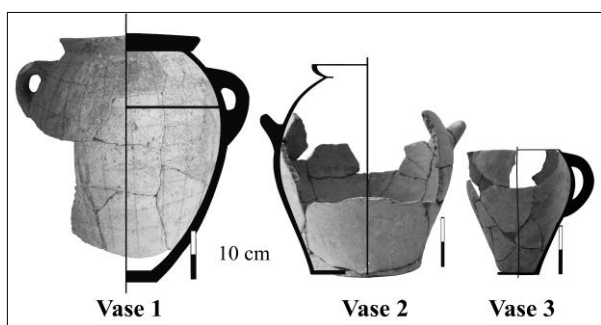
Corrected measurements												
Table (c)		Vase(b)					Side(s)					
n	D_n (mm) (± 0.2)	B_{xy}^{\parallel} (nT) (± 0.2)	B_z^{\parallel} (nT) (± 0.6)	B_B (nT) (± 0.3)	θ_B ($\pm 0.8^\circ$)	ℓ_B (mm)	B_{xy}^{\perp} (nT) (± 0.9)	B_z^{\perp} (nT) (± 0.2)	B_s (nT) (± 0.8)	θ_s ($\pm 0.2^\circ$)	ℓ_s (mm) (± 0.3)	
1	11.0	22.9	24.6	33.6	43.0	16.1±0.4	23.5	23.9	33.6	45.4	15.4	
2	13.0	23.4	25.1	34.3	43.0	19.1±0.4	23.7	24.2	33.9	45.6	18.2	
3	15.0	23.6	25.3	34.6	43.0	22.0±0.4	23.9	25.1	34.7	46.4	20.7	
4	17.0	24.0	25.3	34.9	43.5	24.7±0.5	23.9	25.4	34.9	46.7	23.4	
5	19.0	24.1	25.5	35.1	43.4	27.7±0.5	24.1	25.6	35.2	46.7	26.1	
6	21.0	24.3	25.5	35.2	43.6	30.4±0.5	24.1	25.8	35.3	46.9	28.7	
7	23.0	24.5	25.7	35.5	43.7	33.3±0.6	24.3	25.9	35.5	46.8	31.5	
8	25.0	24.6	25.7	35.6	43.8	36.1±0.6	24.3	26.0	35.6	46.9	34.2	
9	27.0	24.7	25.9	35.8	43.7	39.1±0.6	24.5	26.2	35.9	46.9	36.9	
10	29.0	24.8	25.9	35.8	43.8	41.9±0.7	24.5	26.3	35.9	47.1	39.6	

Table (d)		n	D_n (mm)	B_{xy}^{\perp} (± 0.2 nT)	B_z^{\perp} (± 0.2 nT)	M_z (mA/m)	M_{xy} (mA/m)	θ°
		10	29.0±0.2	13.2	14.0	42.6±0.2	40.1±0.2	43.3±0.3

Table B.11b The quantities λ , α^\perp , α^\parallel , M , M_i , ($i=xy,z$) are estimated are estimated using the least squares method from the above sensors readings.

Vase 6	mA/m	M_z	M_{xy}	M	γ°	α_z^\perp	0.75±0.03	α^\perp (cm)	0.71±0.03
M_z^{\parallel}	42.5±0.2	±0.2	±0.2	±0.2	±0.3	α_{xy}^\perp	0.67±0.04		
M_z^\perp	22.41±0.05					α_z^{\parallel}	1.13±0.03		
M_{xy}^{\parallel}	39.8±0.1					α_{xy}^{\parallel}	1.00±0.01		
M_{xy}^\perp	21.0±0.1					λ	1.90±0.01		

B.7 Computation of magnetization in irregular fragments of the vase base 1,2& 3



The components, M_z , M_{xy} , angle γ of the magnetization M , and sensitivity radii of the sensors, α^{\perp} , α^{\parallel} , in fragments from the bases of vase 2 are computed the same way, **without any need for cutting specimens**, at measured distances $D_n < d = L \cdot \tan \gamma$ from the fragments' edge, using the least squares method ($i=z,xy$) from the equation (6):

$$\left(\frac{\mu_0}{2B_i}\right)^2 = \frac{\alpha_i^2 \cdot \sin^2 2\gamma}{4 \cdot M_i^2} \cdot \frac{1}{D_n^2} + \frac{1}{M_i^2} \Leftrightarrow y_n = w_i \cdot x_n + b_i$$

Table B.12a Measurements of the magnetic field components B_{xy} , B_z at distances D_n from its edge along B_{xy} in base fragments of vase 2. The normalization factor, $\lambda \approx 2$, displays a similar value in fragments from different vases. For the normalization of B_{xy} the value $\lambda = 2 \pm 0.1$ is used. $i = z, xy$

Vase 2			
	D_n	$\lambda \cdot B_z^{\perp}$	B_{xy}^{\parallel}
n	(± 0.2 mm)	(± 0.2 nT)	(± 0.2 nT)
1	26.0	66.0	78.0
2	25.0	65.6	78.6
3	24.0	65.2	78.4
4	23.0	65.2	77.9
5	22.0	65.0	77.4
6	21.0	64.8	77.0
7	20.0	64.6	76.4
8	19.0	64.4	75.8
9	18.0	64.0	75.2
10	17.0	64.0	74.2

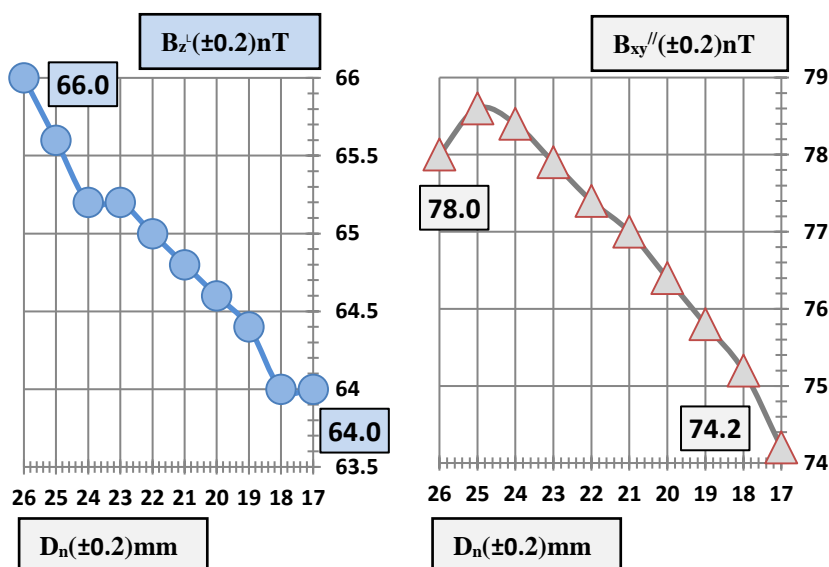


Diagram B.7a The sensor readings B_{xy} , B_z decrease as the distance D_n decreases.

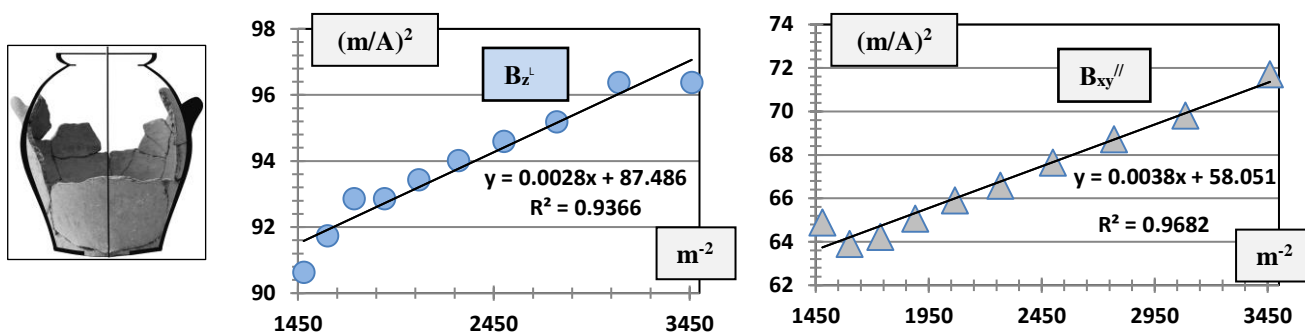


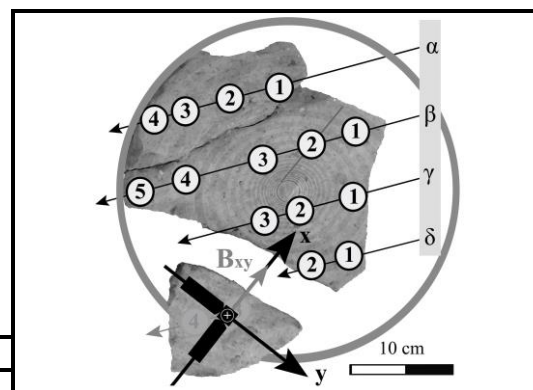
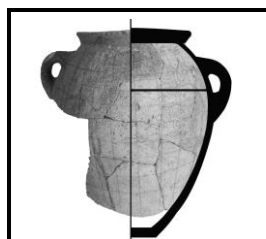
Diagram B.7b The quantities α^{\perp} , α^{\parallel} , M , M_i , ($i=xy,z$) are estimated using the least squares method from the above sensors readings.

Table B.12b Computation of the quantities α^{\perp} , α^{\parallel} , γ , M , M_i , ($i=xy,z$) of the base fragments of vase 2.

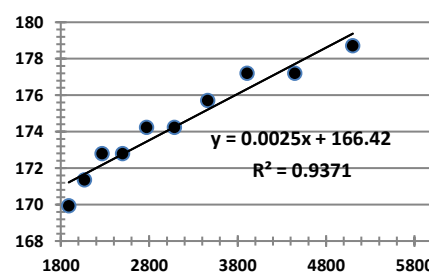
M_z^{\perp} (mA/m)	M_{xy}^{\parallel} (mA/m)	M (mA/m)	γ°	α_z^{\perp} (cm)	α_{xy}^{\parallel} (cm)	λ
106.9 \pm 0.4	131.2 \pm 0.7	169.3 \pm 0.6	50.8 \pm 0.2	1.15 \pm 0.05	1.66 \pm 0.05	2.0 \pm 0.1

The same quantities are computed (tables B.13, B.14) for the fragments of the vases 1 and 3, in the same way.

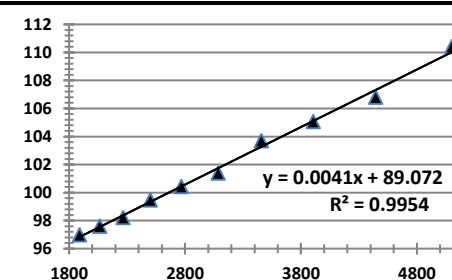
Table B.13 Computation of the quantities α' , α'' , γ , M , M_i , ($i=xy,z$) from measurements in irregular fragment of base vase 1.



Vase 1		Measurements in irregular fragment of base vase 1					
$L \pm \Delta L$ (mm)		Measurements along B_{xy} , which is oriented along the direction of the x-sensor					
21.2±0.5		Measurements of B_z from the vertical z- sensor on the fragment surface					
n	$\lambda \cdot B_z$ (±0.2nT)	D_n (±0,2mm)	x_n	ψ_n	$(x_n - \langle x \rangle)^2$	$\psi_n \cdot (x_n - \langle x \rangle)$	$(\psi_n - \langle \psi \rangle)^2$
1	48.2	23.0	1890.4	169.9	1585056.5	-213938.0	20.1
2	48.0	22.0	2066.1	171.3	1173395.3	-185609.2	9.4
3	47.8	21.0	2267.6	172.8	777528.4	-152356.9	2.7
4	47.8	20.0	2500.0	172.8	421654.7	-112197.3	2.7
5	47.6	19.0	2770.1	174.2	143842.9	-66083.1	0.0
6	47.6	18.0	3086.4	174.2	3960.1	-10964.8	0.0
7	47.4	17.0	3460.2	175.7	96632.8	54621.7	1.7
8	47.2	16.0	3906.3	177.2	572898.5	134126.5	7.8
9	47.2	15.0	4444.4	177.2	1677271.1	229497.2	7.8
10	47.0	14.0	5102.0	178.7	3813003.7	348977.7	18.5
		$\langle x \rangle$	$\langle \psi \rangle$	$\Sigma(x_n - \langle x \rangle)^2$	$\Sigma \psi_n \cdot (x_n - \langle x \rangle)$	$\Sigma(\psi_n - \langle \psi \rangle)^2$	ΣR_i^2
		3149.3	174.4	10265244.0	26073.8	70.7	4.4
		b_z	δb_z	w_z	δw_z		
		0.0025	0.0002	166.4	0.8		
		M_z (A/m)	δM_z	α_z' (m)	$\delta \alpha_z'$		
		0.0775	0.0002	0.0082	0.0004		

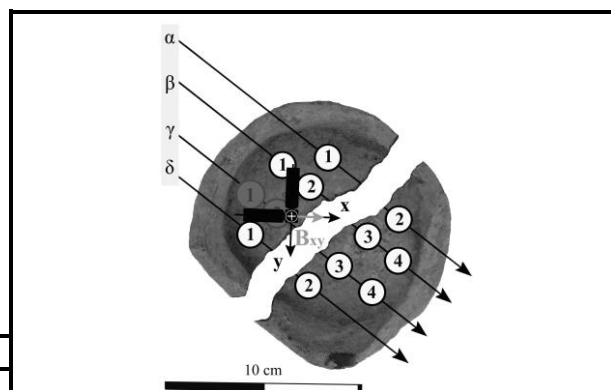


Vase 1		Measurements of B_{xy} from the parallel x- sensor on the fragment surface					
n	B_{xy} (±0.2nT)	D_n (±0,2mm)	x_n	ψ_n	$(x_n - \langle x \rangle)^2$	$\psi_n \cdot (x_n - \langle x \rangle)$	$(\psi_n - \langle \psi \rangle)^2$
1	63.8	23.0	1890.4	97.0	1585056.5	-122107.1	25.2
2	63.6	22.0	2066.1	97.6	1173395.3	-105722.6	19.4
3	63.4	21.0	2267.6	98.2	777528.4	-86604.3	14.3
4	63.0	20.0	2500.0	99.5	421654.7	-64588.8	6.4
5	62.7	19.0	2770.1	100.4	143842.9	-38086.3	2.5
6	62.4	18.0	3086.4	101.4	3960.1	-6380.4	0.4
7	61.7	17.0	3460.2	103.7	96632.8	32236.8	2.9
8	61.3	16.0	3906.3	105.1	572898.5	79520.2	9.3
9	60.8	15.0	4444.4	106.8	1677271.1	138310.2	23.0
10	59.8	14.0	5102.0	110.4	3813003.7	215571.3	70.5
		$\langle x \rangle$	$\langle \psi \rangle$	$\Sigma(x_n - \langle x \rangle)^2$	$\Sigma \psi_n \cdot (x_n - \langle x \rangle)$	$\Sigma(\psi_n - \langle \psi \rangle)^2$	ΣR_i^2
		3149.3	102.0	10265244.0	42149.2	173.9	0.8
		b_{xy}	δb_{xy}	w_{xy}	δw_{xy}		
		0.0041	0.0001	89.1	0.3		
		M_{xy} (A/m)	δM_{xy}	α_{xy}'' (m)	$\delta \alpha_{xy}''$		
		0.1060	0.0002	0.0142	0.0002		

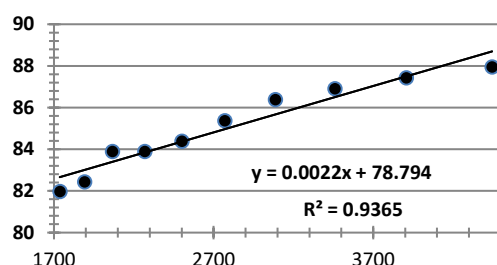


Vase 1							Computation of the quantities α' , α'' , γ , M , M_i , ($i=xy,z$)	
$M_z \pm \delta M_z$ (mA/m)	$M_{xy} \pm \delta M_{xy}$ (mA/m)	$M \pm \delta M$ (mA/m)	$\gamma^\circ \pm \delta \gamma^\circ$	$\alpha' \pm \delta \alpha'$ (cm)	$\alpha'' \pm \delta \alpha''$ (cm)	$\lambda \pm \delta \lambda$		
77.5±0.2	106.0±0.2	131.3±0.2	53.8±0.1	0.82±0.04	1.42±0.02	2.0±0.1		

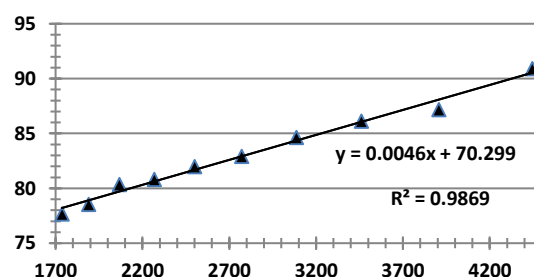
Table B.14. Computation of the quantities $\alpha^I, \alpha^{II}, \gamma, M, M_i, (i=xy,z)$ from measurements in irregular fragment of base vase 3.



Vase 3		Measurements in irregular fragment of base vase 3						
$L \pm \Delta L$ (mm)		Measurements along B_{xy} , which is oriented along the direction of the x-sensor						
22.3 ± 0.5		Measurements of B_z from the vertical z- sensor on the fragment surface						
n	$\lambda \cdot B_z$ ($\pm 0.2nT$)	D_n ($\pm 0.2mm$)	x_n	ψ_n	$(x_n - \langle x \rangle)^2$	$\psi_n \cdot (x_n - \langle x \rangle)$	$(\psi_n - \langle \psi \rangle)^2$	
1	69.4	24.0	1736.1	82.0	1159165.2	-88249.7	9.57	
2	69.2	23.0	1890.4	82.4	850816.8	-76044.1	6.86	
3	68.6	22.0	2066.1	83.9	557472.4	-62635.9	1.37	
4	68.6	21.0	2267.6	83.9	297224.2	-45735.5	1.37	
5	68.4	20.0	2500.0	84.4	97816.6	-26390.9	0.46	
6	68.0	19.0	2770.1	85.4	1821.0	-3643.3	0.10	
7	67.6	18.0	3086.4	86.4	74891.6	23641.9	1.77	
8	67.4	17.0	3460.2	86.9	419193.0	56266.1	3.40	
9	67.2	16.0	3906.3	87.4	1195728.1	95595.5	5.57	
10	67.0	15.0	4444.4	87.9	2662405.7	143498.5	8.32	
		$\langle x \rangle$	$\langle \psi \rangle$	$\Sigma(x_n - \langle x \rangle)^2$	$\Sigma \psi_n \cdot (x_n - \langle x \rangle)$	$\Sigma(\psi_n - \langle \psi \rangle)^2$	ΣR_i^2	
		2812.8	85.1	7316534.7	16302.5	38.8	2.46	
		b_z	δb_z	w_z	δw_z			
		0.0022	0.0002	78.8	0.6			
		M_z (A/m)	δM_z	α_z^I (m)	$\delta \alpha_z^I$			
		0.1127	0.0004	0.0107	0.0005			



Vase 3		Measurements of B_{xy} from the parallel x- sensor on the fragment surface						
n	B_{xy} ($\pm 0.2nT$)	D_n ($\pm 0.2mm$)	x_n	ψ_n	$(x_n - \langle x \rangle)^2$	$\psi_n \cdot (x_n - \langle x \rangle)$	$(\psi_n - \langle \psi \rangle)^2$	
1	71.3	24.0	1736.1	77.7	1159165.2	-83609.1	29.68	
2	70.9	23.0	1890.4	78.5	850816.8	-72441.1	20.88	
3	70.1	22.0	2066.1	80.3	557472.4	-59984.0	7.65	
4	69.9	21.0	2267.6	80.8	297224.2	-44050.2	5.32	
5	69.4	20.0	2500.0	82.0	97816.6	-25635.8	1.29	
6	69.0	19.0	2770.1	82.9	1821.0	-3538.5	0.03	
7	68.3	18.0	3086.4	84.6	74891.6	23159.8	2.32	
8	67.7	17.0	3460.2	86.1	419193.0	55768.6	9.18	
9	67.3	16.0	3906.3	87.2	1195728.1	95311.6	16.46	
10	65.9	15.0	4444.4	90.9	2662405.7	148329.0	60.84	
		$\langle x \rangle$	$\langle \psi \rangle$	$\Sigma(x_n - \langle x \rangle)^2$	$\Sigma \psi_n \cdot (x_n - \langle x \rangle)$	$\Sigma(\psi_n - \langle \psi \rangle)^2$	ΣR_i^2	
		2812.8	77.7	7316534.7	33310.3	153.7	2.02	
		b_{xy}	δb_{xy}	w_{xy}	δw_{xy}			
		0.00455	0.00019	70.3	0.55			
		M_{xy} (A/m)	δM_{xy}	α_{xy}^{II} (m)	$\delta \alpha_{xy}^{II}$			
		0.1193	0.0005	0.0161	0.0003			



Vase 3						
Computation of the quantities $\alpha^I, \alpha^{II}, \gamma, M, M_i, (i=xy,z)$						
$M_z \pm \delta M_z$ (mA/m)	$M_{xy} \pm \delta M_{xy}$ (mA/m)	$M \pm \delta M$ (mA/m)	$\gamma^0 \pm \delta \gamma^0$	$\alpha^I \pm \delta \alpha^I$ (cm)	$\alpha^{II} \pm \delta \alpha^{II}$ (cm)	$\lambda \pm \delta \lambda$
112.7±0.4	119.3±0.5	164.1±0.4	46.6±0.2	1.07±0.05	1.61±0.03	2.0±0.1

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